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# Dissolution Potential of SO<sub>2</sub> Co-Injected with CO<sub>2</sub> in Geologic Sequestration

LAUREN E. CRANDELL, BRIAN R. ELLIS, AND CATHERINE A. PETERS\*

Department of Civil & Environmental Engineering, Princeton University, Princeton, New Jersey 08540

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Sulfur dioxide is a possible co-injectant with carbon dioxide in the context of geologic sequestration. Because of the potential of SO<sub>2</sub> to acidify formation brines, the extent of SO<sub>2</sub> dissolution from the CO<sub>2</sub> phase will determine the viability of co-injection. Pressure-, temperature-, and salinity-adjusted values of the SO<sub>2</sub> Henry's Law constant and fugacity coefficient were determined. They are predicted to decrease with depth, such that the solubility of SO<sub>2</sub> is a factor of 0.04 smaller than would be predicted without these adjustments. To explore the potential effects of transport limitations, a nonsteady-state model of SO<sub>2</sub> diffusion through a stationary cone-shaped plume of supercritical CO<sub>2</sub> was developed. This model represents an end-member scenario of diffusion-controlled dissolution of SO2, to contrast with models of complete phase equilibrium. Simulations for conditions corresponding to storage depths of 0.8-2.4 km revealed that after 1000 years, 65-75% of the SO<sub>2</sub> remains in the CO<sub>2</sub> phase. This slow release of SO<sub>2</sub> would largely mitigate its impact on brine pH. Furthermore, small amounts of SO<sub>2</sub> are predicted to have a negligible effect on the critical point of CO<sub>2</sub> but will increase phase density by as much as 12% for mixtures containing 5% SO<sub>2</sub>.

## Introduction

Geologic sequestration of carbon dioxide is emerging as a promising carbon mitigation strategy (1-3). Electric power plants, which are the largest point sources from which CO<sub>2</sub> will be captured, emit other important pollutants such as sulfur dioxide, and there is the possibility that such pollutants could be disposed of by co-injection with CO<sub>2</sub> into deep geologic formations. Every year, the electric power industry collectively spends more than \$5 billion on permits for the right to emit SO<sub>2</sub> (4). Given the environmental and human health benefits of controlling SO<sub>2</sub> emissions (5), it may be economically advantageous to dispose of SO<sub>2</sub> with CO<sub>2</sub>. However, the effects that impurities in the CO<sub>2</sub> stream will have on injection and long-term storage are largely unknown. Regulations governing the injection of CO2 have been proposed by the U.S. EPA; however, there has been no specification of the required purity of the injected stream except that it must not meet the definition of a hazardous waste (6). John Gale, of the IEA, recently commented on the need to understand the effects of impurities on capture, transport, and storage before setting purity regulations (7).

In aqueous solution, SO<sub>2</sub> may form sulfurous acid, sulfuric acid, and even H<sub>2</sub>S (8). With the exception of H<sub>2</sub>S, these are acids stronger than carbonic acid, formed from aqueous CO<sub>2</sub>. Acidity in the context of geologic CO2 sequestration will accelerate mineral dissolution and precipitation, which may affect formation porosity and permeability, integrity of wellbore cements, and caprock integrity (9-13). Gunter et al. (11) examined water-rock reactions for injection streams of CO<sub>2</sub>, H<sub>2</sub>SO<sub>4</sub>, and H<sub>2</sub>S and predicted that these acids could be effectively buffered by the formation mineralogy, thus reducing the potential for long-term brine acidification. However, other geochemical modeling studies (14-16) predict that storage of sulfur co-injectants with CO2 would create persistent, highly acidic conditions. Knauss et al. (15) used a reactive transport simulation and predicted that even small amounts of SO<sub>2</sub>, such as 10<sup>-6</sup> bar partial pressure, would create pH conditions near unity. Under these acidic conditions, porosity will be increased because of mineral dissolution and carbonate mineral precipitation will be inhibited. Similarly, Xu et al. (16) predicted that co-injection of SO<sub>2</sub> would create a larger and more acidic zone than injection of CO<sub>2</sub> alone and would increase porosity in the acidic zone due to mineral dissolution and decrease porosity at the acid front due to sulfate precipitation.

The magnitude of brine acidification will depend on the extent and rate of dissolution of SO<sub>2</sub> out of the injected CO<sub>2</sub>. One of the objectives of this work was to determine a means of predicting the solubility of SO<sub>2</sub> in formation brines under geologic sequestration conditions. Solubility of CO2 under these conditions has been fairly well examined (17-19); however, little is known of SO<sub>2</sub> solubility at high temperature, pressure, and salinity conditions. At injection depths below 800 m, where CO<sub>2</sub> would exist as a supercritical fluid (20), temperatures are greater than 30 °C, and pressures are greater than 74 bar. The solubility of SO<sub>2</sub> in pure water has been determined at temperatures up to 130 °C, but at pressures up to only 25 bar (21-23). Furthermore, at depths relevant for geologic sequestration, brine salinity in sedimentary basins is as high as 225 g/L (approximately 3.8 M) (20). SO<sub>2</sub> solubility has been measured in 0.1-6 M NaCl solutions (24-26), and in mixed electrolyte solutions (27-29) but not at the pressures relevant for geologic sequestration.

In this work, we present a means of determining the phase partitioning of  $SO_2$  from supercritical  $CO_2$  (scCO<sub>2</sub>) by using pressure-, temperature-, and salinity-adjusted parameters including the Henry's Law constant and the  $CO_2$ -phase fugacity coefficient. We examine how these parameters and the resulting  $SO_2$  brine concentrations would vary with depth, in comparison to  $CO_2$  concentrations. Additionally, densities of  $CO_2$ – $SO_2$  mixtures under relevant pressure and temperature (PT) conditions, along with mixture critical points, are calculated. These properties are important to demonstrate the effects of  $SO_2$  additives on  $CO_2$  injection depth and migration potential. Finally, the diffusivity of  $SO_2$  in sc $CO_2$  is estimated, and its variation with depth and PT conditions is examined.

In addition to the need to quantify the potential for equilibrium phase partitioning of  $SO_2$ , there is a need to estimate the flux of  $SO_2$  from the injected  $CO_2$  to the bulk brine phase. Previous modeling studies (15, 16) of co-injection of  $SO_2$  with  $CO_2$  assume sustained phase equilibrium between all the  $scCO_2$  and brine. This is an extreme case scenario in which there is no limitation on contact of  $SO_2$  with the brine. The opposite extreme case is a scenario in which  $SO_2$  is limited by diffusion through a stationary  $scCO_2$  phase. These two extremes bound reality in which the actual rate of  $SO_2$  contact

 $<sup>^{\</sup>ast}$  Corresponding author. Tel: 609-258-5645; E-mail: cap@princeton.edu.

with brine is controlled by process complexities and system heterogeneities.

To address the need to understand the diffusion-limited bounding scenario, the second major objective of this work was to describe diffusion-limited dissolution behavior of  $SO_2$  into brines in the context of co-injection with  $CO_2$  in geologic sequestration. A nonsteady-state model of  $SO_2$  diffusion through a stationary cone-shaped plume of  $scCO_2$  was developed and used to simulate fluxes of  $SO_2$  into the brine outside the  $scCO_2$  plume. Pressure- and temperature-adjusted binary diffusion coefficients of  $SO_2$  in  $scCO_2$  were estimated. A variety of depths were considered to determine how changes in temperature and pressure affect the properties. Two different mixture compositions, 1 and 5%  $SO_2$ , were considered, corresponding to moderate-to-high ratios of  $SO_2$  to  $CO_2$  in emissions from electric power plants.

#### Methods

 ${
m CO_2-SO_2}$  Mixture Properties. Critical points for mixtures of  ${
m CO_2}$  and  ${
m SO_2}$  were estimated by using the corresponding states method as given by Lee and Kesler (30). Yang et al. (31) examined the accuracy of this method in describing  ${
m CO_2}$  mixtures and concluded that Lee–Kesler type equations of state are as accurate, sometimes more so, than more complex cubic equations of state. This estimation used values of the pure component critical molar volumes, critical temperatures, and acentric factors from reported data (31–34) (Supporting Information, Table S1).

The densities of mixtures of  $CO_2$  and  $SO_2$  were estimated by using the method given in Lee and Kesler (30). For consistency, this method was also used to estimate the densities of pure  $CO_2$  despite the availability of  $CO_2$  equations of state (35). Densities were calculated at PT conditions corresponding to a variety of depths, based on gradients from reported data (36, 37) as described in the Supporting Information (Section S-1).

**Partitioning into Brine.** For the case of phase equilibrium partitioning of  $SO_2$  between  $scCO_2$  and brine, the convention of the infinite-dilution reference state for the aqueous phase is adopted. This results in an expression relating the  $SO_2$  concentration in the brine phase, C[M], to the partial pressure in the  $scCO_2$  phase,  $P_{SO2}$  [atm]:

$$C = \phi_{\text{SO2}} P_{\text{SO2}} K_{\text{H}} \tag{1}$$

where  $\phi_{SO2}$  is the fugacity coefficient of  $SO_2$  in the  $scCO_2$  phase and  $K_H$  is the Henry's Law constant in [M/atm].

The adjustment of the Henry's Law constant for high pressures was done by using the Krichevsky–Ilinskaya equation, which also adjusts for variation in activity coefficient, important for highly soluble gases (38). This equation relates the Henry's law constant,  $K_{\rm H,P}$ , at system pressure, P, to the Henry's law constant at a lower reference pressure,  $K_{\rm H,P}$ , through

$$\ln K_{H,P} = \ln K_{H,P^*} - \frac{A}{RT}(x_j^2 - 1) - \frac{\bar{v}_{SO2}^{*}(P - P^*)}{RT} + \ln\left(\frac{\bar{V}_{P^*}}{\bar{V}_P}\right)$$
(2)

where A is the Margules constant,  $x_j$  is the mole fraction of water in the brine phase,  $\bar{V}_{SO2}^*$  is the partial molar volume of  $SO_2$  at infinite dilution,  $P^*$  is the reference pressure, R is the universal gas constant, T is the temperature, and  $\bar{V}_P$  is the molar volume of water. In this work, the last term is ignored because changes in the molar volume of water with pressure are assumed negligible. Henry's Law constants for  $SO_2$  at  $P^*$  of 1 bar were taken from Rabe and Harris (21) for several temperatures representative of geologic sequestration conditions (given in Table S2 in the Supporting Information).

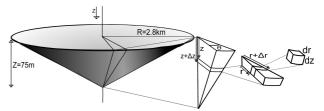


FIGURE 1. Diagram of model injection plume showing an example differential volume element.

These values of  $K_{\text{H,P}}^*$  refer to aqueous  $SO_2$  in the unreacted form, as opposed to the total  $SO_2$  in solution. Partial molar volumes of  $SO_2$  at infinite dilution (Table S1 in the Supporting Information) were taken from Brelvi and O'Connell (39). The Margules constants were estimated via regression of measurements of  $SO_2$  solubility in pure water at pressures up to 11 bar (22).

The Henry's Law constants were adjusted for saline conditions by using the Schumpe model (27) for mixed electrolyte solutions with an adjusted gas-specific parameter given by Rodriguez-Sevilla et al. (24). A 1 M NaCl solution was selected to represent the brine.

The fugacity coefficient for  $SO_2$  in the  $scCO_2$  was determined from the second virial coefficients, determined from the reduced-virial-coefficient method in Tarakad and Danner (40)

 $SO_2$  **Diffusivity in scCO<sub>2</sub>.** The diffusion coefficient for  $SO_2$  in scCO<sub>2</sub>,  $D_{SO2,CO2}$  at high pressures was estimated by using the Takahashi correlation (41) as described in the Supporting Information (Section S-4). Values of  $D_{SO2,CO2}$  at relevant conditions are also given in the Supporting Information (Section S-4).

Model System and Simplifications. A model system was conceptualized to represent an extreme-end-member scenario in which SO<sub>2</sub> contact with the bulk brine is entirely diffusion-limited within the scCO2 phase. By "bulk brine", we refer to the brine phase outside the scCO<sub>2</sub> plume. The system is a scCO2 phase within a geologic formation after the injection period (Figure 1). We selected a cone as simplified geometry to mimic the shape of a CO2 plume trapped underneath a caprock seal. For simplicity, the scCO2 is assumed to be stationary. During the injection process, pressure will force the CO<sub>2</sub> to flow into the formation (see e.g refs 42-44). Post-injection, however, these pressures will dissipate, and CO2 flow will be slower, driven only by buoyancy override and hydrodynamic flow. Furthermore, in deep aquifers, the flow of water is slow, 1-10 cm/yr (45), which means that hydrodynamic forces are minimal.

The other type of CO<sub>2</sub> flow, which is not considered here, is the flow driven by the new density gradients that will be created because of spatial differences in SO<sub>2</sub> concentrations within the scCO<sub>2</sub> phase. A circulation pattern may emerge in which SO<sub>2</sub>-rich CO<sub>2</sub> near the top, center of the cone will exchange with the less dense CO<sub>2</sub> that is depleted of SO<sub>2</sub>. (Note that this does not affect mixing between the brine phase and the scCO<sub>2</sub> phase because the density of the SO<sub>2</sub>-rich scCO<sub>2</sub> phase is still significantly less than the density of the brine phase.) This advective transport is one of the processes that will enhance SO<sub>2</sub> flux relative to the extreme case of diffusion in a stagnant fluid, modeled here.

Another process not considered is partitioning of  $SO_2$  into residual brine trapped within the volume of the cone. Partitioning to residual brine may be significant. The effect, relative to what is modeled here, is that the diffusive flux to the bulk brine would be diminished because of the reduced  $SO_2$  concentrations in the  $scCO_2$  phase.

A cone volume of  $6.1 \times 10^8 \,\mathrm{m}^3$  was calculated on the basis of a CO<sub>2</sub> density of 750 kg/m<sup>3</sup>, an aquifer porosity of 20%, and a 50 year injection period at a rate of 1.83 Mton/year

(46). By assuming a formation thickness of 75 m and by assuming that the height of the cone equals this thickness, a cone with the selected injection volume would extend to a radius of 2.8 km.

 $SO_2$  **Diffusion Modeling.** The cone radius was discretized in radial coordinates by first dividing it into wedges with angle  $\theta$ . The wedges with total height Z were divided into slices with thickness  $\Delta z$ . Slices were then divided into segments with width  $\Delta r$ , as shown in Figure 1. The derivation of the equation for diffusive mass transport of  $SO_2$  in  $scCO_2$  in a differential volume element is presented in detail in the Supporting Information (Section S-5). The resulting model equation is

$$\frac{\partial C}{\partial t} = D_{\text{eff}} \left( \frac{1}{r} \frac{\partial C}{\partial r} + \frac{\partial^2 C}{\partial r^2} + \frac{\partial^2 C}{\partial z^2} \right) \tag{3}$$

where C is the concentration of SO<sub>2</sub> and  $D_{\rm eff}$  is the effective diffusion coefficient. The effective diffusion coefficient is calculated from  $D_{\rm SO2,CO2}$ , formation porosity, n, and tortuosity,  $\sigma$  (47).

$$D_{\text{eff}} = D_{\text{SO2,CO2}} \frac{n}{\sigma} \tag{4}$$

For this work, tortuosity was selected as  $\sigma = 2$  to represent high diffusive flux.

The initial condition assumes a spatially uniform concentration of 1 or 5% SO<sub>2</sub>, within the scCO<sub>2</sub>. A no-flux boundary was assumed at the top of the cone, representing an aquifer bounded by an impermeable caprock. Diffusion within the cone is driven by a concentration gradient created by preferential dissolution of SO<sub>2</sub> into the bulk brine at the cone boundary. SO<sub>2</sub> dissolution was predicted by using eq 1 and the estimated Henry's Law constant and fugacity coefficient. The model does not account for dissolution of CO<sub>2</sub>. In Ellis et al. (8), we describe this process and examine the change in the scCO<sub>2</sub> plume size by considering different scenarios for brine-phase transport. To create a maximum driving force for dissolution, at each time step, the brine at the cone boundary was reset to have a zero concentration of SO<sub>2</sub>. This is a realistic condition because dissolved SO<sub>2</sub> will quickly convert to acid reaction products (see Ellis et al. (8)). The diffusion equation was numerically solved by using a time-split explicit difference method as described in the Supporting Information (Section S-5).

### **Results and Discussion**

 ${
m CO_2-SO_2}$  Mixtures. The critical points estimated for mixtures of  ${
m CO_2}$  and  ${
m SO_2}$  are shown in the phase diagram in Figure 2. For reference, gas—liquid phase boundaries for pure  ${
m CO_2}$  and  ${
m SO_2}$  are also shown. (Published correlations were used for the phase boundaries for  ${
m CO_2}$  (48) and  ${
m SO_2}$  (34).) Whereas the critical temperature of  ${
m SO_2}$  is much higher than that of  ${
m CO_2}$ , pure  ${
m SO_2}$  and pure  ${
m CO_2}$  have similar critical pressures. Also, critical points for mixtures of  ${
m SO_2}$  and  ${
m CO_2}$  are predicted to be fairly constant in pressure. This implies that, in the context of injection, to ensure a supercritical mixture, less than an additional 1 m depth is required for  ${
m 1\% SO_2}$  mixtures and an additional 3.5 m for  ${
m 5\% SO_2}$  mixtures.

Estimated densities are shown in Figure 3, along a trajectory corresponding to a surface temperature of 10 °C. Uncertainties are 5.3% (see Supporting Information). The density of pure  $\rm CO_2$  ranges from 754 to 759 kg/m³ from depths of 0.8 to 2.4 km (Figure 3a). This variation is small because, coincidentally, the PT depth trajectory is closely aligned with an iso-density (isopycnic) line in the supercritical regime. Densities of mixtures of  $\rm CO_2$  and  $\rm SO_2$  (Figure 3b,c) are larger than those of pure  $\rm CO_2$ . For the conditions of interest, densities are up to 3% larger for 1%  $\rm SO_2$  mixtures and up to 12% larger for 5%  $\rm SO_2$  mixtures. This increase is expected

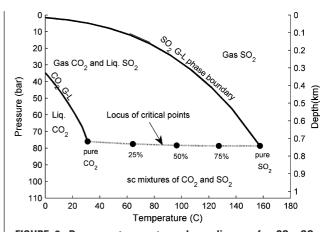


FIGURE 2. Pressure—temperature phase diagram for  $\text{CO}_2$ ,  $\text{SO}_2$ , and their mixtures. Gas—liquid (G-L) phase boundaries are shown for the pure substances, along with the locations of critical points for mixtures of  $\text{CO}_2$  and  $\text{SO}_2$ . The diagram is inverted, showing pressure decreasing on the left ordinate axis to correspond to variation with depth, depicted on the right ordinate axis.

given the larger density of pure  $SO_2$ . For pressures and temperatures that correspond to depths between 0.8 and 2.4 km, the density of  $SO_2$ , as measured by Ihmels (49), ranges from 1373 to 1300 kg/m³. Thus, small amounts of  $SO_2$  have a large effect on the mixture density. In the context of geologic sequestration, this may be beneficial. The densities of brines vary from 950 to 1200 kg/m³ (50). A decrease in the density difference between the two phases would decrease the buoyancy of  $scCO_2$  and may decrease leakage potential.

**Equilibrium Partitioning of SO<sub>2</sub> into Brine.** Salinity-adjusted values of the Henry's Law constant for SO<sub>2</sub> were determined at 20, 40, 60, and 70 °C, for pressures from 1 to 260 bar, corresponding to depths up to 2.4 km. These values have estimated errors of 14% (see Supporting Information). For example, at 1.2 km depth, the Henry's Law constant is predicted to be  $0.58 \pm 0.08$  M/atm. Interpolations of the computed values are shown as contours in a PT plot in the Supporting Information (Section S-6). Values of  $K_{\rm H,P^*}$  decrease with temperature and pressure. Therefore, the Henry's Law constant decreases with increasing injection depth such that SO<sub>2</sub> solubility at geologic sequestration conditions is less than that under conditions at the land surface (value given in Table S1 in the Supporting Information).

The depth variations of all the factors that govern SO<sub>2</sub> equilibrium phase partitioning into brine (eq 1) are shown in Figure 4. Partial pressure increases linearly with depth, but  $\phi_{SO2}$  and  $K_H$  decrease nonlinearly. At shallow depths,  $SO_2$ would behave ideally in the  $CO_2$  phase, with  $\phi_{SO2}$  approximately equal to unity, and at greater depths, substantial negative deviations from ideality are predicted. Also shown in Figure 4 is the resulting molar concentration of SO2 in brine and its variation with depth. This was computed with the simplifying assumptions that, over the entire depth, the aqueous phase is a 1 M NaCl brine, and the brine is in equilibrium with scCO<sub>2</sub> containing 1 or 5% of SO<sub>2</sub>. Also, this calculation does not account for the fact that phase partitioning of a finite mass of SO<sub>2</sub> would deplete it from the scCO<sub>2</sub> phase, and the resulting concentration of SO<sub>2</sub> in the brine would be much smaller than what is computed here. (This more realistic case is considered in Ellis et al. 2009 (8).) However, this simplified case is examined for the sake of illustration. At land-surface conditions of 1 bar and 10 °C, the solubility of SO<sub>2</sub> would be 0.02 M for 1% SO<sub>2</sub> and 0.09 M for 5% SO<sub>2</sub>. As depth increases, the increase in partial pressure causes phase partitioning to increase by one-totwo orders of magnitude to a maximum of 0.45 M at 0.8 km

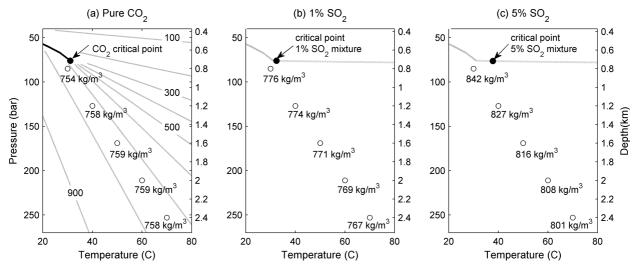


FIGURE 3. Densities of pure  $CO_2$  (a) and mixtures of 1% (b) and 5% (c)  $SO_2$  with  $CO_2$  along a *PT* trajectory corresponding to a pressure gradient of 105 bar/km and a temperature gradient of 25 °C/km, with a surface temperature of 10 °C. Contours in (a) correspond to constant densities of pure  $CO_2$  in kg/m<sup>3</sup>.

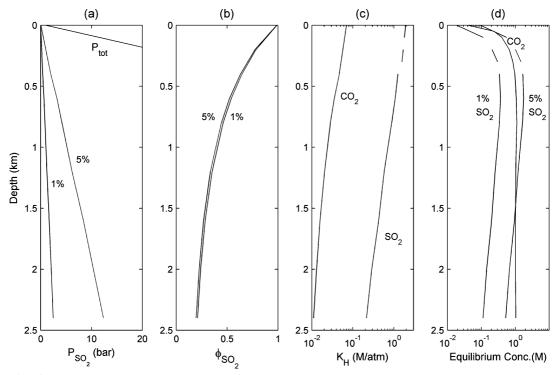


FIGURE 4. (a-c) Variation with depth of the factors that determine phase partitioning between  $scCO_2$  and 1 M NaCl brine and (d) resulting brine-phase concentrations.

for 1% SO<sub>2</sub> and 2.2 M at 0.6 km for 5% SO<sub>2</sub>. At greater depths, the effects of increasing nonideality in the scCO<sub>2</sub> phase and decreasing Henry's Law constant reverse the trend, causing the phase partitioning of SO<sub>2</sub> to decrease with depth.

 $SO_2$ -phase partitioning can be compared with the solubility of  $CO_2$  by using the values calculated by Duan and Sun (18). A major difference between the two cases is that  $SO_2$  is a minor mixture component with significant nonideal solution behavior and  $CO_2$  is a nearly pure component with less extreme deviations from ideality. Consequently, the value of  $\phi$  is not as important for  $CO_2$  as it is for  $SO_2$ . At PT conditions representative of the land surface, the solubility of  $CO_2$  in 1 M NaCl is 0.04 M, and it increases by two orders of magnitude to a value of 1.06 M at a depth of 1.2 km (Figure 4d). This similarity to our findings for  $SO_2$  at shallow depths is expected

because the determining factor is the change in pressure. Beyond a depth of 1.2 km, the solubility of  $\mathrm{CO}_2$  levels off (Figure 4d), which is explained mostly by the increasing importance of the Henry's Law constant. By using the  $K_{\mathrm{H}}$  correlations presented by Bachu and Adams (17) and correcting them for salinity by using the Schumpe model (27), the pressure-adjusted Henry's Law constants for  $\mathrm{CO}_2$  were calculated. These values were found to decrease from 0.053 M/atm at the surface to 0.009 M/atm at a depth of 2.4 km (Figure 4c). In comparison with  $\mathrm{SO}_2$ , values of  $K_{\mathrm{H,P}}$  for  $\mathrm{CO}_2$  are always smaller, reflecting the higher aqueous solubility of  $\mathrm{SO}_2$ . In fact, the solubility of  $\mathrm{SO}_2$  is so much higher that its concentration in the brine phase is comparable to that of  $\mathrm{CO}_2$  despite the fact the mole fraction of  $\mathrm{SO}_2$  is so small. At 1.2 km, the equilibrium concentration in the brine

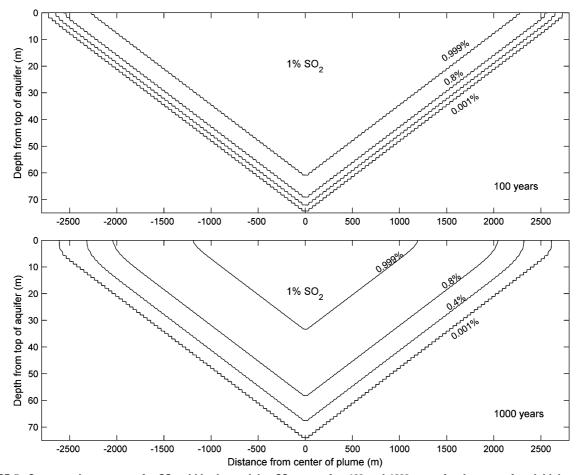


FIGURE 5. Concentration contours for  $SO_2$  within the model  $scCO_2$  cone after 100 and 1000 years for the case of an initial condition of 1%  $SO_2$ , temperature of 40 °C, and pressure of 127 bar corresponding to a depth of 1.2 km. The phase boundary between  $scCO_2$  and brine is coincident with the 0.001% line.

is 0.38 and 1.8 M for 1 and 5% SO<sub>2</sub>, respectively and 1.03 M for  $\rm CO_2$  at 1.2 km (Figure 4d).

SO<sub>2</sub> Diffusion through scCO<sub>2</sub>. The predicted diffusion of SO<sub>2</sub> through the model scCO<sub>2</sub> cone is shown in Figure 5, for the case of 1% SO<sub>2</sub>, as concentration profiles through a vertical section of the cone. The concentration of SO<sub>2</sub> near the cone boundary decreases early on because of the relatively high solubility of SO<sub>2</sub> in the bulk brine. The resulting depletion of SO<sub>2</sub> at this boundary creates a concentration gradient within the scCO<sub>2</sub> that causes diffusive flux of SO<sub>2</sub> from the center of the plume. The simulations predict that the concentration contours are roughly parallel with the cone boundary. The parallel contours establish a "zone of depletion" of fairly uniform thickness at the boundary of the cone. If we arbitrarily define the zone of depletion to be bounded by where the SO<sub>2</sub> concentration equals 0.999%, after 100 years, it is 15 m thick, and after 1000 years, it is approximately 41 m thick. Even after 1000 years, the concentration of SO<sub>2</sub> in a large portion of the center of the cone remains unchanged.

The effective flux of  $SO_2$  from the entire cone over time is shown in Figure 6a for depths of 0.8 and 2.4 km. Initially, the flux of  $SO_2$  into the bulk brine is limited mostly by solubility, and later, the thick zone of depletion acts as a barrier for diffusion. For both PT conditions, after the first 200 years, the flux is two orders of magnitude less than the flux during the initial years. If the model allowed for  $SO_2$  to accumulate in the bulk brine, the flux of  $SO_2$  from the cone would be even slower (see Ellis et al. (8)). Slower flux would also be predicted in the case of a model that described  $SO_2$  partitioning to a residual brine phase within the  $SCO_2$  cone.

Figure 6b shows the total amount of  $SO_2$  that remains in the cone over time. Even though the effective flux in the case of 5%  $SO_2$  is consistently five times larger than the effective flux for the 1% case (Figure 6a), this difference has a negligible effect on the percent of  $SO_2$  remaining. In the first few years, nearly 5% of the  $SO_2$  leaves the cone. After this, there is a more gradual change in the percent of  $SO_2$  removed because of a slower flux of  $SO_2$  from the cone during these years. After 1000 years, 64-75% of  $SO_2$  still remains in the cone for all of the PT conditions simulated.

The variation in flux with depth seen in Figure 6 is due to the variation in the diffusion coefficient of SO<sub>2</sub> in scCO<sub>2</sub>. For 1% SO<sub>2</sub>, at 0.8 km, the diffusion coefficient is  $4.76 \times 10^{-8}$  $m^2/sec$ , compared with  $1.9 \times 10^{-8}$   $m^2/sec$  for the 2.4 km depth, both with error of  $\pm 7\%$  (see Supporting Information). As shown in Figure S-1 in the Supporting Information, at shallower depths, there is little change in  $D_{SO2,CO2}$  with temperature, and values are quite sensitive to changes in pressure. A pressure increase from 80 bar to only 100 bar produces a decrease in  $D_{\rm SO2,CO2}$  from 6 imes 10<sup>-8</sup> to 3 imes 10<sup>-8</sup> m<sup>2</sup>/sec. At greater depths, D<sub>SO2,CO2</sub> values decrease slightly with pressure, and the increase with temperature is more substantial than at shallower depths. The balance of these two effects means that, for depths of 1.2–2.4 km, the  $D_{\rm SO2,CO2}$ contours align with the PT gradient such that there is little change in  $D_{SO2,CO2}$  values.

Implications for Geologic  $CO_2$  Sequestration. As mentioned in the Introduction, prior co-injection modeling studies predicted low pH conditions because of the formation of sulfur-containing acids and that these conditions would be sustained for decades. It is certainly the case that  $SO_2$  is

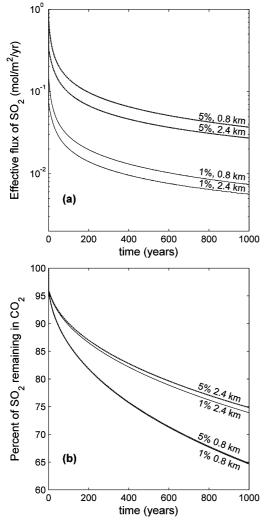


FIGURE 6. (a) Effective flux of  $SO_2$  from  $scCO_2$  cone and (b) percent of  $SO_2$  remaining in the cone for 1 and 5%  $SO_2$  at depths of 0.8 and 2.4 km.

very soluble and much more soluble than CO2. However, our results show that SO<sub>2</sub> solubility decreases substantially with depth. By not accounting for changes in  $K_H$  and  $\phi$  with depth, the partitioning of SO<sub>2</sub> would be overestimated by more than an order of magnitude at depths of 2.4 km. Furthermore, if the fluids are fairly stationary for long time periods, conditions that would not favor mixing of the brine and scCO<sub>2</sub>, diffusion limitations may significantly limit the extent of SO<sub>2</sub> mass transfer to the bulk brine. This would limit the formation of sulfur-containing acids and their impact on brine pH. This finding is very different from the conclusions presented in other studies. Of course, reality is represented neither by the perfectly mixed phase-equilibrium scenario of previous studies nor by the perfectly stagnant diffusion-limited scenario modeled here. Reality lies somewhere between these two extremes, in which some degree of mixing will occur and some degree of transport limitations will exist. This work has demonstrated the important role of such transport limitations in governing the fate of co-injected SO<sub>2</sub>. Future reactive transport studies should account for the possibility of diffusion-limited release of SO<sub>2</sub> from the scCO<sub>2</sub> phase. Some of these effects are discussed in Ellis et al. 2009 (8) who considered a variety of reaction and transport scenarios to estimate the potential for brine acidification from SO<sub>2</sub> fluxes.

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# **Supporting Information Available**

Supporting Information includes sedimentary basin *PT* gradients, data used to estimate Henry's Law constants and mixture properties, *PT* contour plot of adjusted Henry's Law constants for SO<sub>2</sub>, derivation of the diffusion model equations, diffusion coefficients for SO<sub>2</sub> in scCO<sub>2</sub> and error propagation for calculations of Henry's law constant, mixture density, and the diffusion coefficient. This material is available free of charge via the Internet at http://pubs.acs.org.

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