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Theoretical Design of n-Type Organic Semiconducting Materials **Containing Thiazole and Oxazole Frameworks**

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Supporting Information

ABSTRACT: The characteristics of molecular structure and charge transport of some new n-type organic semiconductors containing thiazole 1a-6a and oxazole 1b-6b frameworks and trifluoromethylphenyl as terminal groups were predicted using density functional theory (DFT) methods. The energy levels of HOMO and LUMO of these compounds are decreased when thiophene and furan units are replaced by thiazole and oxazole units, respectively. The same trend was observed when benzo[1,2-d:4,5-d']bisthiazole groups were replaced with benzo[1,2-d:4,5-d']bisthiazole-4,8-diones. The reorganization energies for electron of compounds are computed in a range of 0.21-0.37 eV, which is comparable to the value of 0.25 eV of well-known n-type semiconductors such as perfluoropentacene. Some important trends can be pointed out as follows: (i) replacing the core thiazolothiazole unit of compounds 1a and 2a by the larger core benzo[1,2-d:4,5-d']bisthiazole units of 3a and 4a decreases both

reorganization energies for electron (λ_e) ; (ii) the λ_e values of compounds containing thiazole 2a, 4a, and 6a are smaller than those of compounds containing thiophene 1a, 3a, and 5a, respectively; (iii) there is no clear trend when replacing benzene rings of compounds 3a and 4a by quinone rings of 5a and 6a. The λ_e values of 5 and 6 are only somewhat larger. The same trend is also found for compounds containing oxazole 1b-6b. The intermolecular charge transports in solid state of these compounds mainly occur among molecules in the same molecular layer, whereas intermolecular interactions between molecules in different molecular layers are very small. Generally, beside some experimentally reported molecules 1a-4a, the remaining molecules designed here are good candidates for n-type organic semiconducting materials with small reorganization energies for electron and low energy levels of LUMO.

1. INTRODUCTION

The π -conjugated organic semiconducting materials are of considerable current interest owing to their key role in the fabrication of high performance organic electronic devices such as organic field-effect transistors (OFETs), organic light emitting diodes (OLEDs), and organic photovoltaic cells (OPVs). 1-3. A good organic semiconductor is required to exhibit high ambient stability and efficient charge transport characteristics. In this context, some p-type organic semi-conductors such as pentacene and rubrene^{4,5} have emerged as promising materials with high FET mobility that exceeds that of amorphous silicon field effect transistors (FETs, α -Si:H, mobility $\sim 1.0 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$). However, similarly effective ntype organic semiconductors are rather scarce although they play a primordial role in fabricating complete and complementary circuits.^{7,8} As a consequence, the search for n-type semiconductors constitutes a stimulating and actual challenge.

Despite the fact that numerous π -conjugated organic semiconducting materials were experimentally investigated over the past four decades, an understanding about the relationship between molecular structure and charge transport properties of materials remains quite limited. One of the most important advantages of organic semiconductors is that their molecular structures can easily be modified by either combining several different molecular units to form cooligomers and copolymers or substituting different functional groups into parent compounds. 9,10 This allows us to design tailor-made materials with considerably improved physicochemical properties. However, we are also facing many challenges as the performance of organic electronic devices not only depends on the nature of materials but also is affected by several different factors such as the methods of production of device, purity of materials, type of electrodes, and so on. Fortunately, these complex effects can be avoided in theoretical investigations using quantum chemical computations that are now able to provide us with realistic structure-property correlations.

In this context, we set out in the present work to design and investigate the characteristics of molecular structure and charge transport properties of some new n-type organic semiconductors containing thiazole 1a-6a and oxazole 1b-6b frameworks (Figure 1) and having trifluoromethylphenyls as terminal groups. Thiazole and its derivatives such as benzo [1,2d:4,5-d'] bisthiazole, bithiazole, and thiazolothiazole are wellknown as electron-accepting planar heterocyclic compounds

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1a:
$$X = S$$
; $Y = CH$; $Z = CH$ 1b: $X = O$; $Y = CH$; $Z = CH$
2a: $X = S$; $Y = N$; $Z = CH$
2a-iso: $X = S$; $Y = CH$; $Z = CH$
2a-iso: $X = S$; $Y = CH$; $Z = CH$
2a-iso: $X = CH$
2a-iso: $X = CH$
2a-iso: $X = CH$
2a-iso: $X = CH$
3b: $X = CH$
4a: $X = CH$
4a-iso: $X = CH$

Figure 1. Chemical structures of molecules 1a-6a and 1b-6b.

6b:

that are widely used in both n-type and p-type organic semiconductors. ^{11,12} Introduction of these moieties into novel organic semiconductors based on thiophene is expected to downshift the energy levels of their HOMO and LUMO, which results in better ambient stability. In addition, the lowering of the LUMO levels of n-type semiconductors tends to decrease the energy barrier between their |LUMO| and work function of gold electrode (~5.1 eV) that subsequently improves their transport characteristics. In fact, the syntheses of some compounds such as 1a-4a shown in Figure 1 were previously reported.¹³ However, the effect of molecular structures on their FET mobility has not been well understood yet due to the different conditions used in fabricating FET devices. For instance, the FET mobility of compound 2a was reported to be 0.64 cm² V⁻¹ s⁻¹ by Mamada et al. 14 which is higher than the value of 0.30 cm² V⁻¹ s⁻¹ of 1a reported by Ando et al., 15 but lower than the value of 1.2 cm² V⁻¹ s⁻¹ of 1a reported in ref 16. Although the compounds 2a and 2a-iso were found to have similar electrochemical properties and crystal structures, the FET characteristics of bottom contact devices made from these materials turn out to be much different from each other. While 2a revealed FET mobility of 0.64 cm² V⁻¹ s⁻¹, that of 2a-iso is very low $(0.3 \times 10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1})$. ¹⁴ Similar observations were also found for compounds 4a and 4b-iso. 17 Thus, a theoretical investigation on these compounds provides us with more insight into these interesting phenomena.

In order to design new n-type semiconductors, we replace benzo [1,2-d:4,5-d'] bisthiazole moieties in compounds 3a and 4a by benzo [1,2-d:4,5-d'] bisthiazole-4,8-dione units to form compounds 5a and 6a, respectively, whose levels of LUMO are expected to be decreased. Finally, replacement of thiazole units in 1a-6a by oxazole counterparts also leads to different n-type semiconducting materials 1b-6b.

2. COMPUTATIONAL METHODS

2.1. Charge Transport Properties. In order to probe the charge transport properties, the hopping model is used. $^{18-23}$ In this model, the hole and electron carriers can be jumped between the adjacent molecules of the organic crystals. The rate of charge hopping (K) can be evaluated by using the Marcus—Hush equation (eq 1): 24

$$K = (t^2/\hbar)\exp(-\lambda/4k_{\rm B}T)\sqrt{\pi/\lambda k_{\rm B}T}$$
 (1)

where t is the transfer integral between two adjacent molecules, λ the reorganization energy, $k_{\rm B}$ the Boltzmann constant, and T the working temperature (being 298 K in our calculations). Accordingly to eq 1, a high rate of charge hopping K is attained when the transfer integral (t) between two molecules is high, and the monomers have low reorganization energy (λ) .

The transfer integral is defined by eq 2,25-27

$$t = \langle \varphi_i^{0, \text{site1}} | F | \varphi_j^{0, \text{site2}} \rangle \tag{2}$$

where $\phi_i^{0,\text{site1}}$ and $\phi_j^{0,\text{site2}}$ present the HOMO (or LUMO) of isolated molecules 1 and 2, respectively. F is the Fock operator $(F = SC\varepsilon C^{-1})$ of the dimer with a density matrix from non-interacting dimer, where S is the intermolecular overlap matrix, and C and ε are the molecular orbital coefficients and energies from one-step diagonalization without iteration, respectively. The transfer integral for hole (t_h) is calculated as half of the energy difference between the HOMO and HOMO-1 of the dimer, while the transfer integral for electron (t_e) is the half of the energy difference between the LUMO and LUMO+1 of the dimer. These orbital energy levels are obtained using density functional theory (DFT) with the PW91 functional 28,29 in conjunction with the 6-31G(d,p) basis set. This approach has been proven to be suitable and effective in calculations of the transfer integrals of many organic compounds. $^{30-38}$

2.2. Reduction and Oxidization Potentials in Solution. The reduction and oxidization potentials of compounds in solution are calculated by using a protocol recently developed by Davis and Fry.³⁹ First, optimizations of geometries and calculations of harmonic vibrational frequencies of each compound in the neutral, anionic, and cationic charge states are performed at the B3LYP/6-31+G(d) level. Their single point electronic energies are obtained using a larger basis set, B3LYP/6-311++G(2df,2p). Addition of a set of diffuse functions (++) in the basis set is necessary to describe the structures of the ions involved in the evaluation. The SMD/ IEF-PCM solvation model⁴⁰ is used to probe the solvent effects. These methods were proven to be effective to obtain the reduction and oxidation potentials of PHA compounds in solution.^{39,41} In the present case, we use acetonitrile (CH₃CN) as the working solvent.

The absolute potentials at 298 K are obtained using eq 3,

$$E^{298}_{\text{Abs}} = E^{298}_{\text{calc}} - 0.03766 \tag{3}$$

where E^{298}_{calc} is calculated from the difference in free energy (G) at 298 K between the couples neutral/cation and anion/

Table 1. Oxidation and Reduction Potentials (eV) and Energy Levels of HOMO and LUMO (eV) of Compounds Considered in Acetonitrile Solution (B3LYP/6-311++ $G(2df_{2}p)$ //B3LYP/6-31+G(d) Using the SMD/IEF-PCM Solvation Model)

| | $E_{1/2}^{ m red}$ | $E_{1/2}^{\mathrm{oxd}}$ | LUMO | НОМО | gap | | $E_{1/2}^{ m red}$ | $E_{1/2}^{\mathrm{oxd}}$ | LUMO | НОМО | gap |
|----|---------------------|--------------------------|-------|-------|------|----|--------------------|--------------------------|-------|-------|------|
| 1a | -1.62 $(-1.88)^a$ | 1.01 | -3.18 | -5.81 | 2.62 | 1b | -1.78 | 0.92 | -3.02 | -5.72 | 2.70 |
| 2a | -1.43 $(-1.48)^a$ | 1.28 | -3.37 | -5.81 | 2.71 | 2b | -1.56 | 1.21 | -3.24 | -6.01 | 2.77 |
| 3a | -1.57 $(-1.80)^a$ | 1.16 | -3.23 | -5.96 | 2.73 | 3b | -1.87 | 1.08 | -2.93 | -5.88 | 2.95 |
| 4a | -1.49 $(-1.39)^a$ | 1.37 | -3.31 | -6.17 | 2.86 | 4b | -1.70 | 1.37 | -3.10 | -6.17 | 3.07 |
| 5a | -0.60 | 1.42 | -4.20 | -6.22 | 2.02 | 5b | -0.53 | 1.41 | -4.27 | -6.21 | 1.94 |
| 6a | -0.54 | 1.71 | -4.26 | -6.51 | 2.25 | 6b | -0.52 | 1.67 | -4.32 | -6.47 | 2.15 |

^aThe values in parentheses were obtained from available experimental reports (1a, ref 15; 2a, ref 14; 3a, ref 17; 4a, ref 17).

neutral. The value of 0.03766 accounts to small thermal correction for a free electron at 298 $\rm K.^{42}$

The reduction potential is predicted in CH_3CN solution for the ferrocene—ferrocenium (Fc/Fc⁺) pair, and the standard redox couples are calculated by using expressions 4 and 5:

reduction:
$$E_{1/2}^{298} = 1.056E_{Abs}^{298} - 4.90$$
 (4)

oxidation:
$$E_{1/2}^{298} = 0.932E_{Abs}^{298} - 3.94$$
 (5)

2.3. Electronic Structure Calculations. Optimizations of the relevant geometries and calculations of their harmonic vibrational frequencies are fully performed using DFT methods. The functionals including the hybrid B3LYP⁴³⁻⁴⁵ are used in conjunction with the 6-31G(d,p), 6-31+G(d,p), and 6-311++G(2df,2p) basis sets. 46,47 For open-shell systems including the cations and anions, the unrestricted formalism (UB3LYP) is used. This functional has again been demonstrated to be reliable in predicting geometrical parameters and charge transport properties of many π -conjugated systems. 35-37 For each species considered, the corresponding neutral, anionic, and cationic states are characterized. All electronic structure calculations are carried out using the Gaussian 09⁴⁸ suite of programs.

3. RESULTS AND DISCUSSION

3.1. Geometries of Compounds Considered. Geometry optimizations and calculations of harmonic vibrational frequencies of all compounds examined are performed using the hybrid B3LYP functional in conjunction with the polarized 6-31G(d,p) basis set. Reliable geometries and charge transport properties of the organic semiconductors have been obtained using this approach in many previous studies. The chemical structures of all compounds considered are shown in Figure 1, whereas their shapes and coordinates are given in the Supporting Information.

A good agreement between our computed geometries and the available experimental parameters can be noted. For instance, the computed C–C distance (d_1) between the thiazolethiazole core and two thiophene moieties of 1a is 1.440 Å, which is close to the experimental values of 1.439–1.443 Å. Similarly, the computed C–C bonding length (d_1) of 1.440 Å between the thiazolethiazole core and two thiazole moieties of 2a also agrees well with the experimental values of 1.444–1.449 Å. Interestingly, our predictions show that compounds 2a, 4a, and 6a containing thiazole have planar geometries, whereas compounds 1a, 3a, and 5a containing thiophene are slightly twisted out of the molecular plane. In

agreement with experimental results, 14 the compound 2a has a planar structure, while compound 2a-iso is twisted. Interestingly, all compounds 1b-6b are found to have planar structures.

3.2. Electrochemical Properties. Frontier molecular orbitals (FMO) are primordial descriptors in probing organic semiconductors. For n-type organic semiconductors, earlier studies showed that their LUMO level should be in the range of -3.0 to -4.5 eV. This not only makes the materials kinetically more stable with respect to ambient oxidation conditions but also reduces energy barriers between the n-type semiconductors and work function of gold electrode (\sim 5.1 eV in vacuum). Although some other metals such as aluminum, magnesium, and so on were also used as electrodes with lower work functions, they are not preferred because of their lower stability in ambient conditions.

Since experimental measurements are usually carried out in solution, theoretical predictions performed using continuum solvent models allow us to obtain more comparable and useful results. In this work, reduction and oxidation potentials and subsequently energy levels of frontier orbitals of the molecules are calculated using the SMD/IEF-PCM solvent model. So Our computed results summarized in Table 1 are in good agreement with the available experimental values. The computed reduction potentials of 2a and 4a are $-1.43~\rm eV$ and $-1.49~\rm eV$, which are close to the experimental values of $-1.48~\rm eV^{14}$ and $-1.39~\rm eV$, respectively. The $E^{1/2}_{\rm red}$ values of compounds 1a and 3a are predicted to be $-1.62~\rm eV$ and $-1.57~\rm eV$, respectively, which are somewhat higher than the experimental values of $-1.88~\rm eV^{15}$ and $-1.80~\rm eV$.

Interestingly, the HOMO and LUMO values listed in Table 1 reveal that of all compounds are in a range of -3.18 to -4.26 eV, which turns out to be very good for n-type organic semiconductors. More importantly, some variable trends for the energy levels of their HOMO and LUMO emerged as follows:

- (i) First, replacement of thiophene units by thiazole units results in decreasing energy levels of both HOMO and LUMO. The HOMO and LUMO values of 2a, 4a, and 6a are systematically lower than those of 1a, 3a, and 5a, respectively.
- (ii) Second, small differences in the HOMO and LUMO levels of energy are observed when the core thiazolelothiazole units in 1a and 2a are replaced by benzo[1,2-d:4,5-d']bisthiazole units in 3a and 4a, respectively.
- (iii) Our predictions consistently reveal that replacement of benzo[1,2-d:4,5-d']bisthiazoles in 3a and 4a by benzo-[1,2-d:4,5-d']bisthiazole-4,8-diones in 5a and 6a consid-

erably reduces their HOMO and LUMO energy levels. The LUMO energy levels of **5a** and **6a** are located at -4.20 and -4.26 eV, respectively, which are very good for n-type organic semiconductors.

(iv) Finally, the effect of dione on LUMO energy levels of oxazole analogues is larger than on that of thiazole analogues. The LUMO energy levels of compounds 5b and 6b are systemically lower than those of compounds 5a and 6a, respectively. The LUMO values of 1b-4b are higher than those of 1a-4a, respectively.

Similar trends can be found for the compounds 1b-6b containing oxazole units. The HOMO and LUMO values of 2b, 4a, and 6b are in fact lower than those of 1b, 3b, and 5b, respectively.

3.3. Reorganization Energy. The reorganization energy is a key parameter in the evaluation of both electron and hole mobilities of a molecule. It is defined as the sum of geometrical relaxation energies when the species goes from the neutral state geometry to a charged state geometry, and vice versa. A description of these terms on a potential energy surface is schematically shown in Figure 2. Accordingly, a molecule

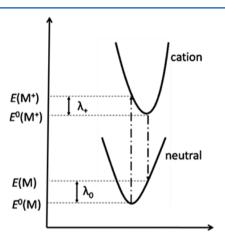


Figure 2. A schematic definition of the reorganization energy for hole.

exhibits high charge mobilities when it has low reorganization energies that can be defined as follows:

$$\lambda_{\rm h} = \lambda_0 + \lambda^+ = [E(M) - E^0(M)] + [E(M^+) - E^0(M^+)]$$
(6)

$$\lambda_{e} = \lambda_{0} + \lambda^{-} = [E(M) - E^{0}(M)] + [E(M^{-}) - E^{0}(M^{-})]$$
(7)

where λ_h and λ_e are reorganization energies for hole and electron, respectively; $E(M^+)$ and $E(M^-)$ are the total energies of the cationic and anionic states with the optimized geometries of the neutral ground state, respectively; E(M) is the total energy of the neutral state with the optimized geometries of the cationic species (for eq 6) and of the anionic species (for eq 7); and $E^0(M)$, $E^0(M^+)$, and $E^0(M^-)$ are the total energies of the neutral, cationic, and anionic ground state species, respectively.

Although the values of reorganization energies for hole (λ_h) of all compounds considered are close to their λ_e values, they are rather large as compared to the typical values of p-type semiconductors previously reported.^{35,51} Thus, these molecules cannot be suggested for p-type organic semiconductors.

The reorganization energies for electron of compounds 1a–6a are computed to be in a range of 0.21–0.34 eV. These values

are comparable to that of 0.25 eV of perfluoropentacene (PF-PEN)³⁵ which is one of the best known n-type organic semiconductors. Additionally, **1a** was also found to be a good candidate for n-type organic semiconductors with a high FET mobility of 1.2 cm² V⁻¹ s⁻¹. Because **1a** is characterized by the largest λ_e value as compared to those of the remaining molecules, the compounds designed above are thus expected to exhibit good charge transport character as well.

More importantly, some remarkable observations can also be pointed out. First, replacement of a core thiazolothiazole unit of compounds 1a and 2a by larger core units of 3a and 4a decreases their reorganization energies for electron (λ_e) . Second, the λ_e values of compounds containing thiazole 2a, 4a, and 6a are smaller than those of compounds containing thiophene 1a, 3a, and 5a, respectively. However, no clear trend emerges when replacing benzene rings of compounds 3a and 4a by quinone rings of compounds 5a and 6a. The λ_e values of 5a and 6a are only marginally larger.

The computed results listed in Table 2 also point out that the λ_e values of compounds **2a-iso** and **4a-iso** are much larger than

Table 2. Reorganization Energies for Electron and Hole of Compounds Obtained at the B3LYP/6-31G(d,p) Level of Theory

| compd | $\lambda_{\rm e}~({\rm eV})$ | $\lambda_{\rm h}~({\rm eV})$ | compd | $\lambda_{\rm e}~({\rm eV})$ | $\lambda_{\rm h}~({ m eV})$ |
|--------|------------------------------|------------------------------|--------|------------------------------|-----------------------------|
| 1a | 0.34 | 0.31 | 1b | 0.31 | 0.32 |
| 2a | 0.25 | 0.33 | 2b | 0.29 | 0.36 |
| 2a-iso | 0.37 | 0.33 | 2b-iso | 0.34 | 0.32 |
| 3a | 0.29 | 0.27 | 3b | 0.24 | 0.25 |
| 4a | 0.21 | 0.29 | 4b | 0.23 | 0.30 |
| 4a-iso | 0.31 | 0.28 | 4b-iso | 0.27 | 0.26 |
| 5a | 0.29 | 0.20 | 5b | 0.35 | 0.16 |
| 6a | 0.26 | 0.20 | 6b | 0.35 | 0.22 |

those of their isomers 2a and 4a, respectively $(\lambda_e(2a) = 0.25 \text{ eV}, \lambda_e(2a\text{-iso}) = 0.37 \text{ eV}, \lambda_e(4a) = 0.21 \text{ eV}, \lambda_e(4a\text{-iso}) = 0.31 \text{ eV})$. This is consistent with the previous experimental result that the FET mobility of 2a is much higher than that of 2a-iso $(2a, \mu = 0.64 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1})$, and 2a-iso, $\mu = 0.3 \times 10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1})$. Our computed results show that, following addition of one excess electron, all anions 2a, 2a-iso, 4a, and 4a-iso have planar geometries (see Table S3, Supporting Information). While the neutrals 2a and 4a also exhibit planar structures, the 2a-iso and 4a-iso have twisted geometries. The large differences in geometrical structures between neutral and anionic states are the main reason for the much larger λ_e values of 2a-iso and 4a-iso as compared to those of 2a and 4a, respectively.

Similar trends are also found for the series of compounds 1b-6b containing oxazole units (Table 2). Moreover, the reorganization energies for electrons of 1b-6b are comparable to those of 1a-6a, respectively. As shown above, their HOMO and LUMO levels of energy are very close together. As a consequence, all of the compounds considered here are expected to be good materials for n-type organic semiconductors.

3.4. Charge Transport Characteristics. An examination of the intermolecular transport behaviors of charge carriers requires parameters of the dimers of the molecules in different packing motifs. Due to lack of relevant experimental results, only molecules **1a**, **2a**, **2a-iso**, and **3a**, whose crystal structures were experimentally known, will be examined. Additionally, we

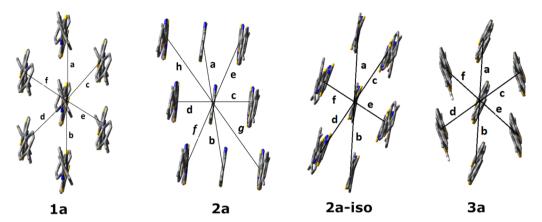


Figure 3. X-ray structures (refs 14, 15, and 17) and charge transport paths of compounds 1a, 2a, 2a-iso, and 3a. Hydrogen and fluorine atoms are removed. The X-ray structure of 3a was refined as compared to that reported in ref 17.

only emphasized the transport characteristics for electron because, as mentioned above, these compounds are probed to be good materials for n-type semiconductors. The corresponding dimers are constructed using experimental parameters of their crystalline lattices. Several charge transport pathways of compounds are shown in Figure 3 and Figure S2 of the Supporting Information.

Our computed results point out that intermolecular charge transports of these compounds mainly occur among molecules in the same molecular layer. Intermolecular interactions between molecules in different molecular layers are in fact very small (Figure S2, Supporting Information).

It is known that the computed intermolecular transfer integrals are very sensitive to DFT functionals. Recently, Sutton et al. 56 showed that increasing amount of HF exchange in DFT functionals will increase the predicted transfer integral values in a nearly linear fashion. In the present work, we performed calculations of transfer integrals for various transport pathways of above compounds by using several functionals (PW91, B3LYP, M06,⁵⁷ BhandHLYP,⁵⁸ and M06HF⁵⁹ that contain amount of HF exchange of 0%, 20%, 27%, 50%, and 100%, respectively). Our results in Table 3 are consistent with the previous reports that the transfer integral values of compounds are largest when functional M06-HF (with percent of HF exchange of 100%) is used. The smallest values are obtained when pure functional PW91 is used. However, we would note that the effect of HF exchange (the actual fraction) also depends on the compounds considered and is not always in linear motif as previously shown. Interestingly, it can be observed that the transfer integrals of pathways c and e in compound 1 are much larger than those of d and f, respectively, although their intermolecular distances are approximate. This can be understood from the fact that arrangements of thiophene groups in these dimers differ from each other. As shown in Figure S4 (Supporting Information), in dimers c and e the two thiophene groups in two monomers can be regarded as having a head-to-tail configuration, whereas they are rather in a head-to-head configuration in both dimers d and f.

While geometries of two isomers 2a and 2a-iso are similar, their packing motifs in X-ray structures are quite different from each other (Figure 3). For both isomers, the intermolecular electron transports are mainly carried out along four pathways, c, d, e, and f, while the remaining directions are much less effective. Interestingly, the values of transfer integrals of the isomer 2a-iso are even larger than those of 2a. The maximum

Table 3. Distance between Two Monomers (d, Å) and Transfer Integral (t_e, meV) in Several Charge Transport Paths of Compounds 1a, 2a, 2a-iso, and 3a

| | | | trans | fer integral | (meV) | |
|------------|-------|--------|--------|--------------|----------------|--------|
| pathways | d (Å) | PW91 | B3LYP | M06 | BH and HLYP | M06HF |
| | | | 1a | | | |
| a, b | 6.05 | 5.03 | 5.58 | 6.67 | 6.12 | 4.90 |
| c | 5.11 | 43.13 | 48.16 | 48.16 | 53.06 | 57.82 |
| d | 5.68 | 1.09 | 1.77 | 0.54 | 8.30 | 10.88 |
| e | 4.26 | 70.34 | 83.54 | 84.08 | 103.68 | 128.57 |
| f | 5.32 | 4.76 | 7.35 | 7.48 | 13.88 | 20.41 |
| | | | 2a | | | |
| a, b | 5.914 | 8.98 | 10.61 | 11.16 | 12.93 | 13.20 |
| c, d | 3.982 | 26.80 | 33.74 | 30.34 | 42.04 | 42.72 |
| e, f | 6.333 | 14.42 | 16.60 | 16.74 | 20.54 | 23.40 |
| g, h | 7.769 | 4.49 | 7.21 | 7.48 | 11.16 | 14.42 |
| | | | 2a-iso | | | |
| a, b | 8.43 | 1.63 | 1.90 | 1.63 | 2.18 | 2.18 |
| c, d | 5.86 | 26.67 | 29.39 | 30.07 | 32.79 | 34.15 |
| e, f | 4.665 | 51.02 | 60.95 | 59.05 | 76.06 | 98.51 |
| | | | 3a | | | |
| a, b | 6 | 27.48 | 32.93 | 32.52 | 40.68 | 46.12 |
| c, d, e, f | 4.728 | 120.82 | 119.59 | 110.48 | 126.13 | 143.00 |

value of transfer integral for **2a-iso** attains a value of 98.51 meV (at level of theory M06-HF), which is significantly larger than that of 42.72 meV found for **2a**.

On the other hand, the compound 3a exhibits the packing structure of herringbone type, rather than π -stacking structures of compounds 1a, 2a, and 2a-iso, whereas the compounds 4a and 4a-iso were suggested as amorphous or distorted structures. Our computed results show that the charge transport for electron of 3a is mainly carried out along transport pathways c, d, e, and f in face-to-edge motif ($t_e = 143.0 \text{ meV}$). The transfer integral along pathways a and a in face-to-face motif is much lower (a).

The hopping drift mobility (μ) is calculated by using the Einstein relation (eq 8):

$$\mu = eD/k_{\rm B}T\tag{8}$$

where e is the electronic charge and D the diffusion coefficient that is calculated from eq 9:^{60,61}

Table 4. Hopping Rate (K, s^{-1}) for Electron in Several Charge Transport Pathways and Mobility $(\mu, cm^2 V^{-1} s^{-1})$ of Compounds 1a, 2a, 2a-iso, and 3a

| | PW91 | B3LYP | M06 | BH and HLYP | M06-HF |
|---|-----------------------|-----------------------|-----------------------|-----------------------|----------------------|
| | | Compound 1a | | | |
| hopping rate $K \times 10^{-12} \text{ (s}^{-1}\text{)}$ | | - | | | |
| a | 0.027 | 0.033 | 0.047 | 0.039 | 0.025 |
| b | 0.027 | 0.033 | 0.047 | 0.039 | 0.025 |
| c | 0.001 | 0.003 | 0.000 | 0.072 | 0.125 |
| d | 1.958 | 2.441 | 2.441 | 2.963 | 3.519 |
| e | 0.024 | 0.057 | 0.059 | 0.203 | 0.438 |
| f | 5.207 | 7.344 | 7.440 | 11.311 | 17.397 |
| mobility $(cm^2 V^{-1} s^{-1})$ | 0.53×10^{-1} | 0.74×10^{-1} | 0.75×10^{-1} | 1.13×10^{-1} | $1.76 \times 10^{-}$ |
| | | Compound 2a | | | |
| hopping rate $K \times 10^{-12} \text{ (s}^{-1})$ | | | | | |
| a, b | 0.238 | 0.332 | 0.367 | 0.492 | 0.513 |
| c, d | 2.117 | 3.356 | 2.713 | 5.210 | 5.380 |
| e, f | 0.613 | 0.812 | 0.825 | 1.244 | 1.614 |
| g, h | 0.059 | 0.153 | 0.165 | 0.367 | 0.613 |
| mobility (cm 2 V $^{-1}$ s $^{-1}$) | 0.19×10^{-1} | 0.29×10^{-1} | 0.24×10^{-1} | 0.45×10^{-1} | $0.48 \times 10^{-}$ |
| | | Compound 2a-i | so | | |
| hopping rate $K \times 10^{-12} \text{ (s}^{-1}\text{)}$ | | | | | |
| a, b | 0.002 | 0.003 | 0.002 | 0.004 | 0.004 |
| c, d | 0.536 | 0.651 | 0.681 | 0.810 | 0.879 |
| e, f | 1.961 | 2.799 | 2.627 | 4.357 | 7.309 |
| mobility (cm ² V ⁻¹ s ⁻¹) | 0.24×10^{-1} | 0.35×10^{-1} | 0.33×10^{-1} | 0.55×10^{-1} | $0.94 \times 10^{-}$ |
| | | Compound 3a | | | |
| hopping rate $K \times 10^{-12} \text{ (s}^{-1})$ | | | | | |
| a, b | 1.400 | 2.010 | 1.960 | 3.068 | 3.944 |
| c, d | 27.063 | 26.517 | 22.628 | 29.492 | 37.910 |
| e, f | 27.063 | 26.517 | 22.628 | 29.492 | 37.910 |
| mobility (cm ² V ⁻¹ s ⁻¹) | 3.84×10^{-1} | 3.72×10^{-1} | 3.17×10^{-1} | 4.10×10^{-1} | $5.27 \times 10^{-}$ |

$$D = (1/2n) \sum_{i} d_i^2 K_i P_i \tag{9}$$

where n is the dimensionality (n = 3), d_i the center of mass distance to neighbor i, K_i the hopping rate which is calculated using eq 1, and P_i the relative probability for charge carrier to an ith neighbor $(P_i = K_i / \sum_i K_i)$.

The electron hopping rate for several transport pathways of compounds and their drift mobility are given in Table 4. The mobility of 1a varies in the range of 0.53×10^{-1} to 1.76×10^{-1} cm² V⁻¹ s⁻¹. These values are in good agreement with the experimental value of $1.2 \times 10^{-1} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ (obtained at substrate temperature of 25 °C). 15 The computed mobility of compound 3a is in the range of 3.84×10^{-1} to 5.27×10^{-1} cm² V⁻¹ s⁻¹, which are somewhat higher than the experimental values of 1.0×10^{-1} to 2.4×10^{-1} cm² V⁻¹ s^{-1.14} At substrate temperature of 20 °C, the earlier experimental study showed that the electron mobility of compound ${\bf 2a}$ is 0.03×10^{-1} to 3.8×10^{-1} cm² V⁻¹ s⁻¹, ¹⁷ which are comparable to our theoretical values of 0.19×10^{-1} to 0.48×10^{-1} cm² V⁻¹ s⁻¹. Although the electron mobility of **2a-iso** is equal to 0.003×10^{-1} cm² V⁻¹ s⁻¹ in a previous experimental report, 14 our theoretical value is much larger and varies in the range of 0.24×10^{-1} to $0.94 \times$ 10^{-1} cm² V⁻¹ s⁻¹. This value is even larger than the corresponding value of compound 2a. The reason for this discrepancy comes from the much larger transfer integrals of 2a-iso as compared to those of 2a. We expect that a reexamination of X-ray structures of two these compounds could be carried out in either theoretical or experimental studies in the near future.

4. CONCLUDING REMARKS

In the present work, we designed and investigated the characteristics of molecular structure and charge transport properties of some new n-type organic semiconductors containing thiazole (1a-6a) and oxazole (1b-6b) frameworks and terminal groups of trifluoromethylphenyl using density functional theory (DFT) methods. The main points emerge as follows:

- (i) The HOMO and LUMO energy levels of compounds considered are decreased when thiophene and furan moieties of compounds 1a-2a and 1b-2b are replaced by thiazole (of compounds 3a-4a) and oxazole (of compounds 3b-4b) moieties, respectively. The same trend was observed when benzo[1,2-d:4,5-d']bisthiazole groups are replaced with benzo[1,2-d:4,5-d']bisthiazole-4,8-diones. The effect of the dione on LUMO energy levels of oxazole analogues is larger than that of the thiazole counterparts. The LUMO energy levels of compounds 5b and 6b are systemically lower than those of 5a and 6a, respectively. The LUMO values of 1b-4b turn out to be higher than those of 1a-4a, respectively.
- (ii) The reorganization energies for electron of compounds are in the range of 0.21–0.37 eV, which is comparable to the value of 0.25 eV of well-known n-type semi-conductors such as perfluoropentacene.
- (iii) Replacement of the core thiazolelothiazole unit of compounds 1a and 2a by the larger core units of 3a and 4a invariably decreases their reorganization energies for electron (λ_e) .

- (iv) The λ_e values of compounds containing thiazole (2a, 4a, and 6a) are predicted to be smaller than those of compounds containing thiophene (1a, 3a, and 5a), respectively.
- (v) No clear trend can be emphasized when replacing benzene rings of compounds 3a and 4a by quinone rings of compounds 5a and 6a. The λ_e values of 5a and 6a are only slightly larger. A similar trend is also found for compounds containing oxazole 1b-6b.
- (vi) The intermolecular charge transports in solid state of these compounds mainly occur among molecules in the same molecular layer, whereas intermolecular interactions between molecules in different molecular layers are very small.
- (vii) The transfer integral of dimers considerably depends on the amount of HF exchange of the functionals used, but there is no linear correlation.

In summary, besides the molecules 1a-4a that were experimentally reported, the remaining molecules designed in the present work can be proposed as very good candidates for n-type organic semiconducting materials with small reorganization energies for electron and low LUMO energy levels.

ASSOCIATED CONTENT

S Supporting Information

Shapes and Cartesian coordinates of the optimized geometries of compounds 1a-6a and 1b-6b. The charge transport paths in solid state of compounds 1a, 2a, and 2a-iso. Shapes of molecules 2a, 2a-iso, 4a, and 4a-iso in both neutral and anionic states. Direction of dimers a, b, c, and d of compound 1a. This material is available free of charge via the Internet at http://pubs.acs.org.

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Notes

The authors declare no competing financial interest.

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