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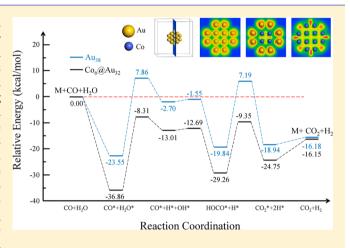
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Theoretical Investigation of the Mechanism of the Water-Gas Shift Reaction on Cobalt@Gold Core-Shell Nanocluster

Sheng-Ke Wu,^{†,‡} Ren-Jie Lin,^{†,‡} Soonming Jang,[§] Hui-Lung Chen, [∥] Shih-Min Wang, [⊥] and Feng-Yin Li*, [†]

Supporting Information

ABSTRACT: We studied the mechanism of the water-gas shift reaction (WGSR; CO + $H_2O \rightarrow CO_2 + H_2$) catalyzed by Co₆@Au₃₂ core-shell nanoalloy using density-functional theory (DFT) calculations to investigate the bimetallic effects on the catalytic activation. The molecular structures and adsorbate/substrate interaction energies were predicted, along with the potential energy surface constructed using the nudged elastic band (NEB) method. Our results indicated that the energetic barriers of the two hydrogen dissociation reactions are lower on the core-shell nanoalloy than on Au₃₈. Furthermore, all of the related chemical species of the WGSR can adsorb stably on Co₆@Au₃₂ to allow the reactions to take place under ambient pressure. To gain insight into the synergistic effect in the catalytic activity of the Co6@Au32 nanoalloy, the nature of the interaction between the adsorbate and substrate was analyzed by detailed electronic local densities of states (LDOS) as well as molecular structures.



INTRODUCTION

Since Haruta et al.'s pioneering findings on the catalysis of CO oxidation with small gold nanoparticles, the use of gold-based catalysts has become a topic of considerable interest mostly because of their potential usage in various industrial applications, such as carbon monoxide removal in hydrogen purification for polymer electrolyte fuel cells and some selective hydrogenation reactions in the synthesis of fine chemicals.^{2–10} To explore more active catalysts with superior performance and stability, considerable research has been devoted to goldrelated bimetallic nanoparticles, which can provide more tunability than their single-metal counterparts because their properties depend not only on the composition but also on the atomic ordering. Furthermore, bringing one more metals together can lead to completely new phenomena, so-called synergistic effects, 22,23 which have no counterpart in bulk materials or single-metal nanoparticles. However, these advantages simultaneously present a great challenge for actual implementation because there are so many possible potential conformer structures and each of them could exhibit certain unpredictable properties. Among the various types of bimetallic nanoparticles, core-shell structures have recently become one of the increasingly important systems because, with defined composition and structure, they can provide predictable control

over structural, geometric, and electronic effects on the catalytic properties; moreover, a generalized synthetic procedure was reported to create core-shell nanoparticles utilizing a redox transmetalation process.²⁴ The catalytic applications of coreshell bimetallic nanoparticles were reviewed recently by Jiang and Xu.²⁵ To evaluate the probable influence of a second metal in gold-based bimetallic nanoparticles, the water-gas shift reaction (WGSR) can serve as a good study system because this reaction has been studied extensively and its reaction mechanism is well-understood.^{26–28}

The WGSR, CO + $H_2O \rightarrow CO_2 + H_2$ is an important industrial reaction in providing high-purity hydrogen for fuel cells and numerous industrial applications. ^{26–28} In addition, it is essential for any practical device that generates hydrogen from a hydrocarbon source. The interest in the WGSR has grown significantly because of its potential capability for the mass production of H_2 , as required by the hydrogen economy and fuel cell applications. ^{29–35} Industrially, the WGSR is catalyzed by two different materials according to the operating temperature: FeCr-based catalysts for high temperature

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(350-450 °C) and CuZn-based catalysts for low temperature (180-250 °C). However, in addition to requiring lengthy preactivation to prevent agglomeration of the metal, Cu-based catalysts are sensitive to condensed water and become pyrophoric if exposed to air. 36,37 These features make Cubased catalysts unsuitable for fuel cell applications. Therefore, a renewed search has been sparked for high-efficiency catalysts to improve the WGSR rate that are also highly stable, especially operating at temperatures lower than 250 °C. Noble metals (e.g., Au and Pt) supported on ceria oxide carriers exhibit high catalytic activity for the low-temperature WGSR without preactivation and are nonpyrophoric. 37,38 The main reason for this high catalytic activity, as recently reviewed by Rodriguez,³⁹ is that the metal nanoparticle and the oxide support can work cooperatively to facilitate the different reaction steps: the adsorption and dissociation of water take place on the oxide, CO adsorption occurs on Au nanoparticles, and all subsequent reaction steps occur at oxide-metal interfaces. However, the detailed reaction mechanism of the WGSR catalyzed by Au-ceria catalysts still remains unsettled. 40,41 Two main alternatives have been proposed in the literature, namely, redox and associated mechanisms, where the latter can be further divided into formate and carboxyl mechanisms. 42 In the redox mechanism, the CO adsorbed on the metal surface is oxidized by a labile oxygen atom of the support to form an oxygen vacancy site, and the oxygen vacancy site is then reoxidized by water, leading to the formation of a hydrogen molecule. 43-47 According to the associated mechanism, the adsorbed CO reacts with terminal hydroxyl groups of the oxidic support to form carbon-containing intermediates, namely, formate and carboxyl, which further react or decompose to form CO₂ and H₂.⁴⁸⁻⁵⁰ Recently, using steadystate isotopic transient kinetic analysis (SSTIKA), Kalamaras et al. concluded that formate is mainly a spectator species. 43,44 Therefore, the associated mechanisms are mainly through carboxyl pathway.

Although Au—ceria catalysts are particularly active and stable, their potential in fuel cell applications is severely hindered by their rapid deactivation due to the buildup of the adsorbed carbonate species on ceria surface, the irreversible over-reduction products generated by the ceria support. These carbonate species act as reaction-inhibiting spectators by either blocking the reaction active sites or preventing the reaction intermediates from accessing these active sites. Even though these carbonate species can be removed by annealing in O_2 environment at 400 °C to regain 95% of the initial activity on the fresh catalyst, the rapid deactivation makes the half-lives of Au—ceria catalysts too short for some fuel cell applications.

Recently, Au nanoparticles smaller than 1.4 nm supported on inert materials were reported to be efficient and robust catalysts for selective oxidation by O₂. Without help from the reducible support, the energetic barriers for hydrogen dissociation of the hydroxyl groups, an important reaction step in the redox mechanism, are rather high, and therefore, the lowest-energy pathway for the WGSR over Au nanoparticles with an inert support is an associated mechanism through a carboxyl intermediate, as suggested by density-functional theory (DFT) studies. Previously, we showed that gold nanoparticle catalysts meet this purpose well with their high surface-to-volume ratios, which not only reduce the amount of catalytic material required but also enhance the catalytic activity compared with that of their bulk counterparts. To further

reduce the usage of the precious metal and simultaneously increase the catalytic activity and stability of the catalysts, gold-based bimetallic nanoparticles 2 are promising candidates. For example, several experimental studies have demonstrated that gold-based bimetallic nanoparticles on different oxide supports can act as catalysts at lower temperature than regular catalysts or even below room temperature. $^{11-21}$ However, there are few theoretical reports on the catalytic effects of bimetallic nanoparticles, and all of them focus on the CO oxidation reaction. $^{56-58}$

In this study, we applied periodic density-functional calculations to investigate the WGSR process on the Co6@ Au₃₂ core-shell nanostructure. The reasons for choosing the Co₆@Au₃₂ core-shell nanostructure as a low-temperature WGSR catalyst are the following: First, there are two main disadvantages of using Au₃₈ for the WGSR, as found in our previous study.⁵⁵ For the WGSR on Au₃₈, the transition states for two O-H bond dissociation reactions have adsorption energies larger than the total energy of the bare Au₃₈ and the gas-phase CO₂ and H₂O. This fact requires these reactions to be performed under high pressure to prevent the reactive complexes from flying away from the surface. Furthermore, the energetic barriers for these bond dissociation reactions on Au₃₈ are relatively high (31.41 and 27.03 kcal/mol), and any energy reduction in these two barriers should enhance the catalytic performance. As predicted by Huang et al. through a systematic reaction mechanism analysis among nine different transitionmetal surfaces covering rows 3d-5d with groups 9 (Co, Rh, and Ir), 10 (Ni, Pd, and Pt), and 11 (Cu, Ag, and Au), 59 the Co metal surface has relatively low energetic barriers for O-H bond dissociation of water molecules and hydroxyl radicals but a rather high barrier for C-O bond formation, whereas the Au metal surface shows the opposite trend, that is, it has a high barrier for O-H bond dissociation and a low barrier for C-O bond formation. Because the Cu metal surface has mild barriers in both O-H bond dissociation and C-O bond formation, it makes the Cu surface have better catalytic activity than the Co and Au surfaces.⁵⁹ Furthermore, the Co surface has higher adsorption energies for both CO and H2O than the Au surface. 59 We surmised that, with the synergetic effects of the combination of Co and Au metals on the WGSR, Co₆@Au₃₂ core—shell nanoparticles might overcome these disadvantages. Second, Co@Au core-shell nanoparticles have been synthesized and characterized successfully by several research groups. 24,60-62 So far, the research on Co@Au core-shell nanoparticles has mainly focused on the potential application of their magnetic and optical properties, along with their synthesis and characterization methods. 31,61,63 To the best of our knowledge, this study is the first regarding the catalytic activity of Co@Au core-shell nanoparticles on the WGSR. The closest related study is perhaps that of Venugopal et al., 11 who compared the WGSR activity catalyzed by a Co-Au nanoalloy against those of a series of Au-based bimetallic nanoclusters, but these nanoclusters were not core-shell structures because their synthesis method was the deposition coprecipitation method. Other core-shell bimetallic catalysts related to Co@Au nanoparticles have been investigated, for example, Ag@Au for glucose oxidation, 64 Pd@Au for the oxidation of alcohols, 65 and Co@Pt for the dehydrogenation of aqueous NH3BH3.66 Finally, Co nanoparticles smaller than 4.4 nm are subject to water-induced oxidation to become cobalt oxides because of their surface free energies, but this type of oxidation will not take place on the bulk Co surface. $^{67-69}$ Therefore, we excluded Co nanoparticles from the energetic calculations because of their instability in water environments. Furthermore, as reported by Ban et al., Co@Au core—shell nanoparticles are stable in water and under acidic conditions because of the protection of the Au shell. The operating cost and reactivity of the WGSR can be significantly improved if the core Co atoms enhance the adsorption stability of CO and H₂O and, at the same time, the surface Au atoms catalyze the reactions enhanced by the core Co atoms. We believe that this study is vital for understanding the factors governing the WGSR catalyzed synergistically by nanosized gold bimetallic clusters.

■ COMPUTATIONAL METHODS

All of the reaction species of the WGSR, including reactants, intermediates, transition states, and products, as well as the nanoparticles, were calculated using spin-polarized densityfunctional theory (DFT)⁷¹ with a plane-wave basis set, as implemented in the Vienna ab initio simulation package (VASP) program.^{71–76} The generalized gradient approximation (GGA)⁷⁷ with the PW91 functional⁷⁸ was employed in conjunction with the projector augmented wave (PAW) potential to describe the interaction between core and valence electrons. The open-shell spin states of the radical intermediates, such as OH and H radicals, were set as ISPIN = 2 in the calculations. For the supercell, we chose a cube size of 25³ Å³, which is large enough to ensure that the integration of nanoparticles with their images is negligible. A 400 eV cutoff energy, which was adequate for convergence to a total energy of 1×10^{-4} eV, was used together with the Γ point for the summation in the Brillouin zone. In this work, we calculated the adsorption energies according to the expression

$$\Delta E_{\rm ads} = E_{\rm total} - E_{\rm Co_6@Au_{32}} - E_{\rm adsorbate} \tag{1}$$

where E_{total} , $E_{\text{Co}_6 @ \text{Au}_{32}}$, and $E_{\text{adsorbate}}$ correspond to the electronic energies of the adsorbed species on the $\text{Co}_6 @ \text{Au}_{32}$ nanoparticle, the bare $\text{Co}_6 @ \text{Au}_{32}$ nanoparticle, and a free adsorbate, respectively. Furthermore, the coadsorption energies of CO with H_2O molecule were calculated as

$$\Delta E_{\text{coads}} = E_{\text{total}} - E_{\text{Co}_6 @ \text{Au}_{32}} - E_{\text{CO}} - E_{\text{H},\text{O}}$$
 (2)

The nudged elastic band (NEB) method^{79,80} was applied to construct the minimum-energy paths (MEPs) by connecting reactants, intermediates, transition states, and products. All chemical species were verified by the vibrational frequencies of the adsorbed structures, which were analyzed by diagonalizing the Hessian matrix of selected atoms within the VASP approach. Atomic charges of the selected structures were calculated by utilizing Bader charge^{81,82} analysis with a program designed by Henkelman et al.⁸³

To calculate the rate constants, we employed harmonic transition state theory (hTST)⁸⁴

$$k = v^* \exp\left(-\frac{E_a}{k_B T}\right) \tag{3}$$

where $E_{\rm a}$ is the activation energy, $k_{\rm B}$ is the Boltzmann constant, T is the temperature in Kelvin, and

$$\nu^* = \frac{k_{\rm B}T}{h} \frac{\prod_{i=1}^{D} \left[1 - \exp\left(-\frac{\hbar\omega_i^{\rm R}}{k_{\rm B}T} \right) \right]}{\prod_{i=1}^{D} \left[1 - \exp\left(-\frac{\hbar\omega_i^{\rm S}}{k_{\rm B}T} \right) \right]}$$
(4)

is the quantum mechanically modified form of the prefactors, where D, ω_i^R , and ω_i^S are the number of normal modes and the nonimaginary vibrational frequencies of the *i*th mode in the reactants and saddle-point states, respectively.

■ RESULTS AND DISCUSSION

Comparison between $Co_6@Au_{32}$ and Au_{38} Nanoparticles. Similarly to Au_{38} , so the $Co_6@Au_{32}$ nanoparticle is modeled as an tetrakaidecahedral structure with O_h symmetry within a 0.1-Å tolerance through structural optimization. As shown in Figure 1, there are eight equivalent hexagonal

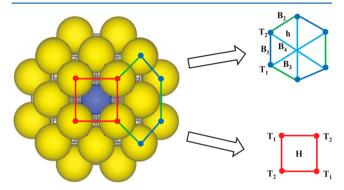


Figure 1. Schematic presentation of a $Co_6@Au_{32}$ nanoparticle. High-symmetry adsorption sites on a hexagonal face of the $Co_6@Au_{32}$ nanoparticle. T_1 and T_2 represent the top sites, B_1-B_4 represent the bridge sites, H represents the hole site, and h represents the hcp site.

fcc(111)-like faces and six equivalent square fcc(100)-like faces, similar to those found in the Au₃₈ cluster, characterized as top (T), bridge (B), hollow (H), and hexagonal-close-packed (hcp) sites. However, because the radius of the Au atom is slightly larger than that of the Co atom, the nanoparticles with cobalt atoms as the shell layer require more than 32 atoms to form stable structures. This is why gold-cobalt core-shell nanoparticles are not included in this study. Several characteristic differences between Co₆@Au₃₂ and Au₃₈ nanoparticles, such as geometry, structure, and adsorption information, are listed in Table 1, along with their charge information. For example, the calculated average nearest-neighbor bond distance between the shell gold atoms is 2.74 Å, and that between the core cobalt atoms is 2.42 Å. These two distances are slightly shorter than those of their single-metal counterparts as reported in the literature (2.78 Å for Au₃₈ nanoparticles⁷⁹ and 2.50 Å for Co₃₈ nanoparticles⁸⁶). However, the average bond distance between a shell gold atom and a core cobalt atom is 2.93 Å, slightly larger than the distance between their counterparts in Au₃₈ nanoparticles (2.83 Å). This Co-Co distance shrinkage in Co@Au nanoparticles was also found by Wang et al. 87 They attributed this shrinkage to the decrease of the total magnetic moment compared with that found in bare Co nanoparticle due to the antiparallel ordering of the spin moments of the shell Au atoms with respect to the spin moments of the Co core atoms. 87 The distance shrinkage between the shell Au atoms is because the size of the Co core is smaller than that of the Au core in Au₃₈ nanoparticles, making the shell Au atoms move slightly toward the center of the core-shell nanoparticles. This is why the size and surface area of a Co₆@Au₃₂ nanoparticle are slightly smaller than those of a Au₃₈ nanoparticle. However, the size shrinkage of the Co core is slightly more than the inward movement of the Au shell, which is why the average bond

Table 1. Geometrical and Structural Comparison between Co_6 @ Au_{32} and Au_{38} Nanoparticles, along with Their Adsorption Information with CO and H_2O Adsorbed onto Their Surfaces

	Au ₃₈	Co ₆ @Au ₃₂
volume (ų)	256.70	239.42
area (Ų)	$20.79 (8.04)^a$	20.06 (7.51)
$d_{ m shell-shell}$ (Å)	2.80	2.74
$d_{\text{core-core}}$ (Å)	2.80	2.42
$d_{ m core-shell}$ (Å)	2.83	2.93
charge on the shell (lel)	-0.380	-2.146
charge on the core (lel)	0.380	2.146
$d_{\mathrm{H_2O}}$ (Å)	$0.97/0.97^b$	0.97/0.98
$d_{ m Au-O}$ (Å)	2.56 ^c	2.72
$E_{\rm ads}({\rm H_2O})$ (kcal mol ⁻¹)	-6.38^{d}	-4.04
d_{CO} (Å)	1.17^{e}	1.17
$d_{\mathrm{Au-C}}$ (Å)	$2.14/2.14^f$	2.12/2.13
$E_{\rm ads}({\rm CO})~({\rm kcal~mol^{-1}})$	-20.15^{g}	-31.81

"Area of hexagon (area of square). b Two O-H bond lengths of one H_2O adsorbed on a T_2 site. Distance between the oxygen atom of the adsorbed H_2O and the gold atom of a T_2 site. Adsorption energy of one H_2O adsorbed on a T_2 site. C-O bond length of one CO adsorbed on a T_2 site. Adsorbed between the carbon atom of the adsorbed CO and the two gold atoms of a T_2 site. Adsorption energy of one CO adsorbed on a T_2 site.

distance between a shell gold atom and a core cobalt atom is slightly larger than the distance between their counterparts in Au₃₈ nanoparticles.

Because the work function⁸⁸ of pure Co metal (5 eV) is slightly smaller than that of pure gold metal (5.10–5.47 eV), a negative charge distribution is built up in the shell gold atoms. This charge separation along with the area shrinking consequently influences the adsorption behavior of H_2O and CO molecules on the surface of this bimetallic nanoparticle. For example, the distance between the oxygen atom of H_2O and its adsorbed gold atom of a $Co_6@Au_{32}$ nanoparticle is 2.72 Å, slightly larger than that found in a Au_{38} nanoparticle (2.56 Å). The adsorption energy of CO on a B_1 bridge site of $Co_6@Au_{32}$ nanoparticle is -31.81 kcal mol^{-1} , significantly different from that found on a B_1 site of a Au_{38} nanoparticle (-20.15 kcal mol^{-1}).

Adsorption of CO, H₂O, OH, and H on a Co₆@Au₃₂ Nanoparticle. Understanding the adsorption structures and related energetics of the WGSR intermediates is significant for exploring the WGSR process on the nanoparticle surface. To locate the probable stable adsorption configurations of WGSR intermediates, such as Co₆@Au₃₂-CO, Co₆@Au₃₂-OH, Co₆@

Au₃₂-H₂O, and Co₆@Au₃₂-H, we placed CO, OH, H₂O, and H species, respectively, at those sites on the surface of a Co₆@ Au₃₂ nanoparticle, as shown in Figures S1-S4, respectively, of the Supporting Information, and the corresponding adsorption energies of these isomeric adsorption structures are listed in Table 2. For CO adsorption, the molecule can adsorb on the Co₆@Au₃₂ nanoparticle in various isomeric structures (Figure S1, Supporting Information), all with the "end-on" configuration, that is, using the C atom to bind to one of the shell Au atoms. Among all of the calculated CO adsorption structures on the Co₆@Au₃₂ surface, only the adsorption energy of OC-B₁ (-31.18 kcal/mol) was larger than that found on the Au₃₈ surface (-20.15 kcal/mol), and the adsorption energies for the rest of adsorption configurations were all smaller than their counterparts on the Au₃₈ surface. Compared with their counterparts on the Au_{38} surface, the most stable COadsorption structures were found to be OC-T₂ and OC-B₁ in the same end-on configuration, with smaller adsorption energies as -20.09 and -20.15 kcal/mol, respectively. A similar pattern was also found in the OH radical adsorption configurations (Figure S2, Supporting Information); that is, only the adsorption energy of OH-B₁ (-77.34 kcal/mol) was larger than that found on the Au₃₈ surface (-66.41 kcal/mol), and the adsorption energies for the rest of the adsorption configurations were all smaller than their counterparts on the Au₃₈ surface. Interestingly, the most stable adsorption configuration of Co₆@Au₃₂-HO is the HO-B₁ configuration, but the most stable adsorption configuration of Au₃₈-HO is the HO-B₂ configuration. For Co₆@Au₃₂-H₂O adsorption (Figure S3, Supporting Information), the adsorption energies for all structures found in this study ranged from -0.88 to −6.07 kcal/mol, slightly smaller than their specific counterparts on the Au₃₈ surface. Interestingly, the H₂O-T₂ configuration, with the oxygen atom adsorbed on a T2 site of Co6@Au32-H₂O, was found to be energetically the most stable among all of the adsorption structures, but for the Au₃₈-H₂O adsorption structures, the most stable configuration was the H₂O-T₁ configuration. However, the results were quite different for the adsorption of the hydrogen atom because all of the adsorption configurations on the Co₆@Au₃₂ surface (Figure S4, Supporting Information), except the $H-B_2$ configuration, were more stable than their specific counterparts on the Au₃₈ surface. Obviously, the main reason for these differences is the influence of the core Co atoms, but the role played by these Co atoms is rather subtle because there is no direct interaction between the adsorbate and the core Co atoms. Most likely, these core Co atoms exert their influence first by changing the properties of the shell Au atoms and then the shell Au atoms interact with

Table 2. Calculated Adsorption Energies (kcal/mol) of Adsorbed CO, H₂O, OH, and H Species on Various Sites of a Co₆@ Au₃₂ Nanoparticle

site	CO	H_2O	ОН	Н
T_1	$-17.14 \ (-17.16)^a$	-4.04 (-6.38)	$-64.39 \text{ (NA}^b)$	$-65.51 (-57.24)^a$
T_2	-17.66 (-20.09)	-5.14 (-6.24)	-57.10 (-63.51)	-64.05 (-57.24)
B_1	-31.81 (-20.15)	-2.63 (-3.42)	-77.34 (-66.41)	-65.19 (-55.17)
B_2	-14.79 (-18.24)	-3.12 (-6.09)	-64.48 (-72.06)	-53.04 (-56.69)
B_3	-10.89 (-17.93)	NA	NA (-61.57)	NA (-56.87)
B_4	-8.53 (-17.92)	NA	NA (-61.83)	NA (-56.69)
Н	-3.11 (-20.09)	-0.75 (-1.64)	-53.93 (-62.82)	NA
hcp	NA	NA	NA (-65.17)	NA (-54.53)

 $[^]a$ Values in parentheses are results on Au $_{38}$. 55 b Not available

the adsorbates. To resolve the above phenomena, we investigated the effects of the core Co atoms on the WGSR through analysis of their electronic properties, such as the electronic local density of states (LDOS), in addition to the geometric and charge differences listed in Table 1. Before performing LDOS analysis, we investigated the probable reaction profile for the WGSR on the $\rm Co_6 @Au_{32}$ surface. In this way, we focused on only the critical conformations of the chemical species on the $\rm Co_6 @Au_{32}$ surface without going through all possible conformations.

Reaction Profile for WGSR on a Co₆@Au₃₂ Nanoparticle. The generally accepted reaction mechanism with the minimum energy barriers for the WGSR is the associative mechanism, ^{89,90} which involves the following steps

$$CO(g) + S \leftrightarrow CO(ads)$$
 (i)

$$H_2O(g) + S \leftrightarrow H_2O(ads)$$
 (ii)

$$H_2O(ads) \rightarrow H(ads) + OH(ads)$$
 (iii)

$$CO(ads) + OH(ads) \rightarrow OCOH(ads)$$
 (iv)

$$OCOH(ads) \rightarrow CO_2(ads) + H(ads)$$
 (v)

$$CO_2(ads) \rightarrow CO_2(g)$$
 (vi)

$$2H(ads) \rightarrow H_2(g)$$
 (vii)

where S represents the nanoparticle. We also tested the two main reactions related to the redox mechanism, $OH(ads) \rightarrow H(ads) + O(ads)$ and $OH(ads) + OH(ads) \rightarrow H_2O(ads) + O(ads)$, and found their energetic barriers to be 33.34 and 57.43 kcal/mol, respectively (Figure S6, Supporting Information), higher than the barrier of reaction v in the associative mechanism on Au_{32} . Therefore, we focused on only the associative mechanism, particularly the hydroxyl pathway.

Based on our adsorption results, we constructed the potential energy surface (PES) for the WGSR on a $Co_6@Au_{32}$ nanoparticle according to the above mechanism, as depicted in Figure 2. The important geometric illustrations of the intermediates, transition states, and products along the WGSR PES are presented in Figure S5 (Supporting Information), and the adsorption sites of these chemical species on a $Co_6@Au_{32}$

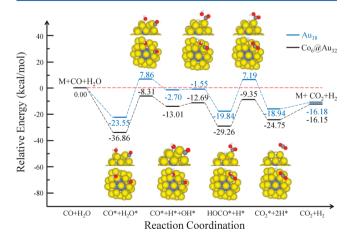


Figure 2. Calculated probable potential energy diagram for the reaction of H_2O and CO on both the $Co_6@Au_{32}$ and Au_{38} nanoparticles, where the numbers are the energies related to the reactants in kcal/mol.

nanoparticle are summarized in Table S1 (Supporting Information), along with the corresponding sites for a Au₃₈ nanoparticle. At first, CO and H₂O are coadsorbed to form IM1 on the nanoparticle, where CO adsorbs on a B₁ site and H₂O adsorbs on a T₂ site for both nanoparticles. The coadsorption energy was calculated to be -36.86 kcal/mol on the Co6@Au32 surface, lower than that found on the Au38 surface (-23.55 kcal/mol), and the coadsorption energy was roughly the same as the sum of the corresponding adsorption energies of CO(ads) and H₂O(ads) individually for Co₆@Au₃₂ (-36.95 kcal/mol) but slightly larger than that found on Au₃₈ (-26.58 kcal/mol). For IM1 of a Co₆@Au₃₂ nanoparticle, the distance between the O atom of H₂O and Au is 2.587 Å, but this distance for a Au₃₈ nanoparticle is 2.652 Å. Although H₂O is closer to the gold atoms in a $Co_6@Au_{32}$ nanoparticle than in a Au₃₈ nanoparticle, there is only small charge transfer between them according to the Bader charge analysis (Table 3). H₂O dissociation from IM1 into adsorbed OH and H producing IM2 for a Co₆@Au₃₂ nanoparticle requires an energetic barrier of 28.55 kcal/mol (TS1) with an endothermicity of 23.85 kcal/ mol, but the corresponding barrier and reaction heat for a Au₃₈ nanoparticle are 31.41 and 20.85 kcal/mol, respectively. Compared to Cu-related materials in terms of the reaction barrier of the water dissociation reaction, a Co6@Au32 nanoparticle is more reactive than a Cu(111) surface (31.36 kcal/mol) but still less reactive than a Cu(100) surface (26.06 kcal/mol) or a Cu_{29} nanoparticle (21.44 kcal/mol). ^{29,31,91} In the structure of TS1 for a Co6@Au32 nanoparticle, H2O moves to an hcp site with one of the two hydrogen atoms pointing toward a B2 site and the associated OH bond stretching from 0.98 to 1.26 Å; however, in the TS1 structure for a Au₃₈ nanoparticle, H2O remains on a T2 site, and the leaving hydrogen atom points toward a B2 site, with the to-be-broken OH bond extended from 1.00 to 1.35 Å. This indicates that a Co₆@Au₃₂ nanoparticle can stabilize the WGSR-related species more than a Au₃₈ nanoparticle. Interestingly, OH(ads) of IM2 for a Co₆@Au₃₂ nanoparticle moves from an hcp site to a T₂ site to react with CO(ads) and produce an OCOH(ads) species by passing through TS2 with a low energy barrier of 0.32 kcal/ mol and an exothermicity of 16.25 kcal/mol; the corresponding energy barrier and reaction heat for a Au₃₈ nanoparticle are 1.15 and -17.14 kcal/mol, respectively. In TS2 for a Co₆@Au₃₂ nanoparticle, the distance between the hydroxyl oxygen atom and the original T₂-site gold atom of IM1 is 4.11 Å, and the forming C-O bond is 1.85 Å; at the same time, for a Au₃₈ nanoparticle, the corresponding distance and C-O bond length are 2.53 and 1.89 Å, respectively. The final decomposition of OCOH(ads) occurs by OH bond cleavage to produce CO₂(ads) and H(ads) on T₁ and B₂ sites, respectively, for both nanoparticles. For a Co₆@Au₃₂ nanoparticle, this process requires a barrier of 19.91 kcal/mol with an endothermicity of 4.51 kcal/mol, but for a Au₃₈ nanoparticle, it needs 27.03 kcal/ mol to overcome the corresponding barrier with an endothermicity of 0.90 kcal/mol. The broken H-O bond of TS3 was calculated to be 1.24 and 1.37 Å for Co₆@Au₃₂ and Au₃₈ nanoparticles, respectively. The adsorption energy of CO₂ was found to be 0.01 kcal/mol for a Co₆@Au₃₂ nanoparticle and -0.92 kcal/mol for a Au₃₈ nanoparticle, suggesting that the CO₂ molecule is only physisorbed on both nanoparticles. Before H₂ desorption from a Co₆@Au₃₂ nanoparticle, the two adsorbed hydrogen atoms that are generated in reaction iii on a B₁ site and in reaction v on a B₂ site, with the sum of their adsorption energies being 5.69 kcal/mol, recombine to form

Table 3. Bader Charges of Metal Atoms; C and O1 of CO; and H1, H2, and O2 of H2O for the Stationary States along the Minimum-Energy Pathway of the WGS Reaction on a Co₆@Au₃₂ Nanoparticle

IM4	4.000 (4.000)	-2.011 (-2.028)	-2.014 (-2.003)	-0.118 (-0.150)	-0.115 (-0.076)		-0.015 (-0.056)	-0.007 (-0.033)	0.034 (0.030)	0.374	0.319
IM3	2.608 (2.489)	-1.910 (-1.881)	-1.928 (-1.891)	-0.108 (-0.109)	1.000 (1.000)		-0.022 (0.139)	0.080 (-0.008)	-0.026 (0.025)	0.381	0.314
IM2	1.763 (1.836)	-1.918 (-1.968)	-1.469 (-1.470)	-0.098 (-0.089)	1.000 (1.000)		-0.001 (0.059)	0.369 (0.030)	0.103 (0.257)	0.379	0.318
IMI	CO ^a 1.734 (1.737)	-1.942 (-1.942) H.O ^a	-1.950 (-1.969)	0.999 (1.000)	1.000 (1.000)	$Co_6 @Au_{32}^b$	0.026 (0.068)	0.074 (0.068)	0.020 (0.030)	0.369	0.336
$H_2O-Co_6@Au_{32}$		μ	-1.962 (-1.969)	1.000 (1.000)	1.000 (1.000)	Co			-0.048 (0.0455)		0.344
OC-Co ₆ @Au ₃₂	1.705 (1.707)	-1.913 (-1.960)					0.050 (0.082)	0.055 (0.072)		0.370	
H ₂ O, CO, Co ₆ (@Au ₃₂	1.844	-1.844	-1.999	1.000	1.000		-0.065 (-0.010)	-0.066 (-0.006)	-0.013 (-0.025)	0.377	0.334
	O	01	02	H1	H2		Au(C1)	Au(C2)	Au(O2)	Co(Au1,Au2)	Co(Au3)

²OI represents the O atom of CO, O2 represents the O atom of H₂O, H1 represents the first dissociated hydrogen atom of H₂O, and H2 represents the second one. Values in parentheses are those corresponding to a Au₃₈ nanoparticle. ^bAu(C1), Au(C2), and Au(O2) represent the Au atom bonding with the C atom of CO and the O atom of H₂O, respectively, whereas Co(Au1,Au2) and Co(Au3) represent the nearest Co atoms to Au(C1,C2) and Au(O2), respectively $H_2(ads)$ on a T_1 site. The adsorption energy of $H_2(ads)$ is 9.51 kcal/mol, implying that $H_2(ads)$ can desorb from the $Co_6@$ Au_{32} nanoparticle surface to become $H_2(g)$ at room temperature.

According to the reaction profile on a Co6@Au₂₂ nanoparticle, the relative energies of all of the WGSR-related chemical species are less than 0 compared with the energies of the free adsorbates and clear Co₆@Au₃₂ nanoparticle; therefore, the reaction should be able take place without a high-pressure environment. On the other hand, for the reactions on a Au₃₈ nanoparticle, the relative energies of TS1 and TS3 are slightly greater than 0. Moreover, the three energy barriers of the WGSR on a Co6@Au32 nanoparticle are lower than those found on a Au₃₈ nanoparticle (2.86, 0.83, and 7.12 kcal/mol lower for reactions iii-v, respectively). Thus, a Co₆@Au₃₂ nanoparticle is more reactive than a Au₃₈ nanoparticle, which was predicted to exhibit better performance than various copper and gold surfaces, as shown in our previous work.⁵⁵ We calculated the rate constant of the rate-limiting step $[H_2O(ads) \rightarrow H(ads) +$ OH(ads)] with the quantum-mechanical modified hTST and found that hydrogen dissociation of a water molecule can take place with rate constants ranging from 6.78×10^{-2} to 0.495 s^{-1} within the WGSR low-temperature range (180-250 °C) and from 1.26×10^2 to 1.09×10^3 s⁻¹ for the WGSR hightemperature range (350-450 °C).

Interestingly, the activation energies and reaction heats (both in kcal/mol) for reactions iii-v on Co6@Au32 are (28.55, 23.85), (0.32, -16.25), and (19.19, 4.51), respectively, whereas the corresponding values on Au₃₈ are (31.41, 20.85), (1.15, -17.14), and (27.03, 0.90), respectively. The activation energies of these three reactions on Co6@Au32 are lower than the corresponding ones on Au₃₈, but all of the reaction heats on Co6@Au32 are more endothermic than the corresponding values on Au₃₈. Howvere, these observations seem to contradict the prediction of the general Brønsted–Evans–Polanyi (BEP) relation, 92–97 a linear relationship between the transition state energy and reaction energy of a reaction. To resolve this issue, we resorted to the molecular and electronic structures along the reaction pathway. Before doing that, we investigated the LDOS of the related nanoparticles in this study to provide general information about the electronic structures to facilitate the discussion.

Analysis of the Electronic State of the Nanoparticles. To understand the bimetallic effect, we calculated the LDOS of pure and core-shell nanoparticles as shown in Figure 3. The dband centers for Au₃₈, Co₆@Au₃₂, and Co₃₈ were found to be -3.495, -3.188, and -2.329 eV, respectively. According to the d-band model of chemisorption developed by Hammer and Nørskov, 98 the metal surface with its d-band center close to the Fermi energy tends to produce a strong interaction between the metal surface and the adsorbed molecule. This partially explains why the relative energies of all of the WGSR-related chemical species on Co₆@Au₃₂ are more stable than their counterparts on Au₃₈. In addition, the d-band distribution of Co₆@Au₃₂ around the Fermi energy is similar to that of Au₃₈ because both of them have an energy gap (~1.0 eV) between the bands around the Fermi energy, although this energy gap for Co₃₈ is less than 0.1 eV, if only the d bands are considered. This indicates that the reactivity of Co6@Au32 should be similar to that of Au₃₈ and is quite different from that of Co₃₈. This energy gap found in Co₆@Au₃₂ is strongly influenced by the dband distribution of the core Co atoms. For example, the upward shifting in the two d bands of Au₃₈ (around −1 and 0

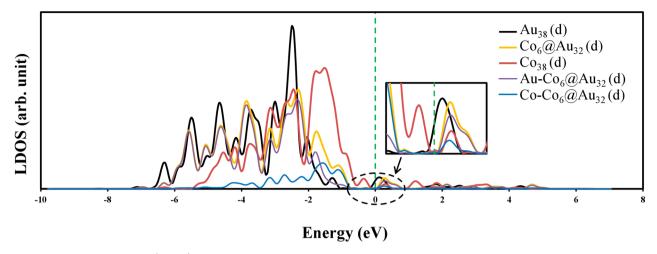


Figure 3. Local densities of states (LDOS) of Au_{38} , $Co_6@Au_{32}$, and Co_{38} nanoparticles along with the contributions from the shell Au atoms and core Co atoms in $Co_6@Au_{32}$. The green dashed line represents the Fermi level.

eV) and the downward shifting in the d bands of Co₃₈ below the Fermi energy match with the corresponding d bands of the core Co atoms. However, the d band of Co₆@Au₃₂ just above the Fermi energy remains almost the same as those of Co₃₈, except that there is a shoulder band around 0.8 eV, which is not found in either Co38 or Au38, contributed exclusively from the shell Au atoms of Co₆@Au₃₂. The main reason for the above observation is the hybridization between the high-lying occupied d bands of the core Co atoms and the unoccupied d bands just above the Fermi energy in the shell Au atoms. Because of this hybridization, those high-lying occupied d bands of the core Co atoms shift downward, and those unoccupied bands of the shell Au atoms near the Fermi energy shift upward with a shoulder band. Similarly, the reverse of the above interaction, that is, the occupied Au d bands and unoccupied Co d bands near the Fermi energy, should take place. However, because of the small portion of the Au d band on the Fermi energy being occupied, only a small upwardshifted shell Au d band is found (approximately -1.0 eV), and almost no changes are found for the unoccupied Co d band (\sim 0.2 eV). Hence, by summing the above two processes, d electrons are transferred from the core Co atoms to the shell Au atoms, as predicted by their work functions. This phenomenon can be easily observed through an electron localization function (ELF) analysis (Figure 4). The electron density distribution of Co₆@Au₃₂ can be divided into two parts,

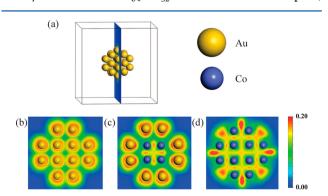


Figure 4. Electron localization function (ELF) diagrams (a) sliced along the (200) plane in the bounding box of (b) Au₃₈, (c) Co₆@Au₃₂, and (d) Co₃₈ nanoparticles.

the shell and the core. The shell electron density closely resembles the electron density of Au_{38} , but more electron density localizes on the shell Au atoms. The core electron density shares the pattern found in Co_{38} . The catalytic reactions, on the other hand, is mainly controlled by the shell Au atoms, and therefore, the WGSR on $\mathrm{Co}_6 @ \mathrm{Au}_{32}$ is very similar to that found on Au_{38} , but the electron density of the shell Au atoms near the Fermi level was shifted upward due to the Co core atoms and enhanced the catalytic activity of $\mathrm{Co}_6 @ \mathrm{Au}_{32}$ without breaking the stability. Therefore, a better Au -based core—shell catalyst has the occupied Au d bands shifted upward toward the Fermi level to enhance its catalytic activity but preserves the d-band pattern of noble metals near the Fermi level to maintain the stability.

Comparison between Co₆@Au₃₂ and Au₃₈ during the WGSR. We applied detailed LDOS analyses to the adsorbed constructs of H₂O and CO species, as well as with the dprojected electron density of the bonded Au atoms and nearby Co atoms, as depicted in Figure 5, along the reaction pathway, where Figure 5a shows the LDOS of the isolated H₂O, CO, and Co₆@Au₃₂ nanoparticle and panels b-h of Figure 5 correspond to the LDOS of the IM1, TS1, IM2, TS2, IM3, TS3, and IM4 configurations, respectively. For comparison, a similar LDOS analysis of a Au₃₈ nanoparticle is shown in Figure 6. In this section, we investigate the different behaviors in the electronic states between Co₆@Au₃₂ and Au₃₈ nanoparticles along the WGSR pathway to resolve the energetic issue regarding the BEP relationship as found in the previous section, combined with the corresponding molecular structures and Bader charge analysis (Table 3).

The higher adsorption energy of IM1 on $Co_6@Au_{32}$ compared to Au_{38} is because the small O1 band of $Co_6@Au_{32}$ around the Fermi energy is completely in the bonding state whereas that of Au_{38} still has some antibonding characteristics (Figures 5b2 and 6b2). It is mostly due to π back-donation of Au d electrons and the charge transfer between the unoccupied $1\pi^*$ orbital of CO and the Au d states with Au(C1) and Au(C2) changing from negative to positive charge. The main difference between TS1 on these two nanoparticles is the H_2O adsorption site, hcp for $Co_6@Au_{32}$ and T_2 for Au_{38} , mainly because of the combination of the CO coadsorption and the distance shrinkage between the shell Au atoms in $Co_6@Au_{32}$. Without the Au distance shrinkage, this

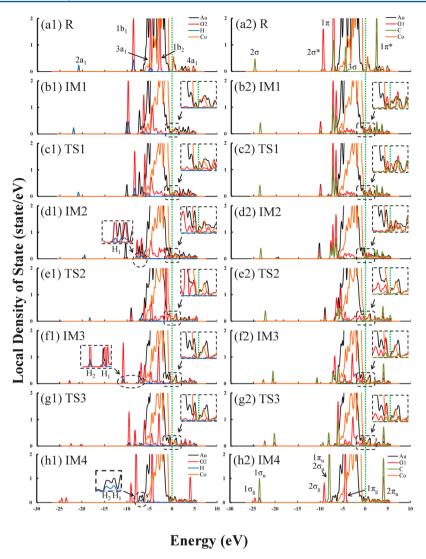


Figure 5. Calculated electronic LDOS of the system projected onto the orbitals for the adsorbed constructs of H_2O (left panel) and CO (right panel) species on $Co_6@Au_{32}$, as well as the d-projected electron density of the bound Au (black) and Co (purple) atoms. The blue, red, olive green, orange, and black lines represent the LDOS of hydrogen, oxygen, carbon, Co (d band), and Au (d band), respectively. The green dashed line represents the Fermi level.

site on Au_{38} is still not a stable H_2O adsorption site even with CO coadsorption. Although CO coadsorption plays an important role in the stabilization of H_2O adsorption on an hcp site, the orbitals of CO(ads) between **IM1** and **TS1** remain almost the same. Therefore the stabilization of H_2O adsorption is through the interaction with the Au atom [Au(C2) in Table 3], which bonds with both CO and H_2O . Because H_2O resides on an hcp site for $Co_6@Au_{32}$, the **TS1** conformation requires the to-be-broken OH bond to be stretched from 0.98 to 1.26 Å toward a B_2 site, which is less than that for Au_{32} (from 1.00 to 1.35 Å), with H_2O locating at a T_2 site and the dissociating H atom also pointing toward a B_2 site. The shorter stretching bond length is the main reason why the energetic barrier to water dissociation on $Co_6@Au_{32}$ is smaller than that on Au_{38} .

After the first hydrogen dissociation, $H_2O(ads)$ becomes HO(ads) with one extra unpaired electron in addition the original nonbonding $1b_2$ paired electrons, and the adsorption changes into chemisorption. The configuration of IM2 has OH adsorbed on an hcp site and the dissociated H atom (H1 in Table 3) on a B2 site on $Co_6@Au_{32}$, but the adsorbed OH radical on Au_{38} resides on a T_2 site. A new peak in the Au_{38}

LDOS (approximately -0.9 eV) indicates the existence of the extra unpaired electron on O2 and stabilization by the metal (Figure 6d1), but this peak is barely noticeable on the Co₆@ Au₃₂ LDOS (Figure 5d1). It is because the short Au-Au distance in Co₆@Au₃₂ and the occupation of an hcp site brings OH(ads) close to both Au(C2) (2.46 Å) and H1(ads) (3.32 Å) to allow an interaction with them. The interaction between OH(ads) and Au(C2) causes the 1π orbital of CO splitting and increases the intensity of the O1 antibonding states at the same time. Meanwhile, on Co6@Au32, the interaction between OH(ads) and H1(ads) reduces the radical characteristics of OH(ads) to cause H1(ads) to still overlap with the O2 orbital and suppress the peak of the unpaired electron in O2 (approximately -0.9 eV) as found for OH(ads) on Au₃₈. These interactions allow OH to adsorb on the original unstable hcp site, but they are also why the reaction of water dissociation on Co₆@Au₃₂ is more endothermic than that on Au₃₈, because the dissociated H atom is adsorbed on a B2 site, which is less stable than the other sites, such as B_1 and T_1 sites (Table 2).

In TS2, OH(ads) and CO(ads) form a partial bond with each other, and OH(ads) moves upward and away from the

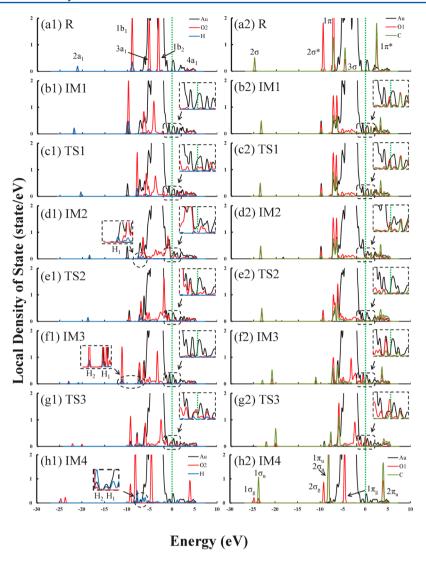


Figure 6. Calculated electronic local densities of states of the system projected on the orbitals for the adsorbed constructs of H_2O (left panel) and CO (right panel) species on Au_{38} , as well as the d-projected electron density of the bound Au (black) atoms. The blue, red, olive green, and black lines represent the LDOS of hydrogen, oxygen, carbon, and Au(d band), respectively. The green dashed line represents the Fermi level.

nanoparticle to make the O2-related orbitals become sharp, similar to those found in the gas phase (Figures 5e1 and 6e1). The lower energetic barrier of this associated reaction on Co6@ Au_{32} compared to Au_{38} is due to the greater π back-donation of Au d electrons on Co₆@Au₃₂, such as that found in IM1. After IM3 is formed, the LDOS patterns of the two panels for the oxygen orbitals become similar, except that O2 still bonds to H2 and O1 has the characteristics of metal back-donation found in both nanoparticles. The main reason for the lower exothermicity in the carboxyl formation reaction on Co₆@Au₃₂ compared to Au₃₈ is mainly because the adsorption energy of the H atom on a B₂ site of Co₆@Au₃₂ is less than that on Au₃₈ (Table 2). Because the dissociated H atoms act as spectators after their dissociation except for the formation reaction of hydrogen molecules (reaction vii), the migration of H1 to other sites can be considered as the isomerization reaction of IM3. Because reaction heat is a thermodynamic property, the system should be allowed to search the conformation space to reach thermal equilibrium. If H1 moves to the site with a higher adsorption energy, such as B₁ or T₁, the reaction of carboxyl formation on Co6@Au32 would become approximately 12 kcal/ mol more exothermic than that with H1 adsorbed on a B2 site.

Thus, carboxyl formation on $Co_6@Au_{32}$ would be more exothermic than that on Au_{38} .

In the final hydrogen dissociation reaction, the LDOS of TS3 on $\text{Co}_6 @ \text{Au}_{32}$ resembles that of IM3. Compared with Au_{38} , there are some differences between IM3 and TS3 in terms of LDOS, especially around the Fermi energy. Because of the higher degree of similarity in electronic structures, the energetic barrier of the final hydrogen dissociation on $\text{Co}_6 @ \text{Au}_{32}$ is lower than that on Au_{38} . The greater endothermicity of this hydrogen dissociation on $\text{Co}_6 @ \text{Au}_{32}$ compared to Au_{38} is due to the adsorption site of the dissociating H atoms, similar to that found in the carboxyl formation reaction.

CONCLUSIONS

We have compared the reaction mechanism of the water—gas shift reaction catalyzed by a $Co_6@Au_{32}$ nanoparticle against that for a Au_{38} nanoparticle, as well as the related adsorption configurations and their electronic properties at the DFT level in conjunction with the PAW approach. Our results show that the adsorptions and reactivity of $Co_6@Au_{32}$ are more similar to those of Au_{32} than to those of Co_{32} . However, the related energies of all of the WGSR chemical species on a $Co_6@Au_{32}$

nanoparticle are lower than those of the gas-phase reactants with the clean nanoparticle, and therefore, the WGSR on a $\text{Co}_6@\text{Au}_{32}$ nanoparticle should be able to proceed under ambient pressure, without the high pressure required for that taking place on a Au_{38} nanoparticle. The calculated reaction rate constant for the rate-limiting step $[\text{H}_2\text{O}(\text{ads}) \to \text{H}(\text{ads}) + \text{OH}(\text{ads})]$ predicts that the WGSR can take place on a $\text{Co}_6@\text{Au}_{32}$ nanoparticle with a reasonable reaction rate at low temperature. These advantages are mainly contributed from the core Co atoms, which control not only the geometry of the surface Au atoms but also the electronic structure near the Fermi energy. Accordingly, we predict that the bimetallic $\text{Co}_6@\text{Au}_{32}$ nanoparticle has better catalytic activity for the WGSR, as it combines the advantages of the catalytic behavior of Au metal and the better adsorption ability of Co metal.

ASSOCIATED CONTENT

S Supporting Information

Adsorption sites of WGSR-related chemical species on the nanoparticles (Table S1). Isomers of adsorbed CO, H₂O, OH, and H on the Co₆@Au₃₂ nanoparticle and their optimized geometric parameters (Figures S1–S4, respectively). Geometry simulations of intermediates, transition states, and products for the WGSR (Figure S5). The two major reactions related to the redox mechanism for the WGSR (Figure S6). This material is available free of charge via the Internet at http://pubs.acs.org.

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Notes

The authors declare no competing financial interest.

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