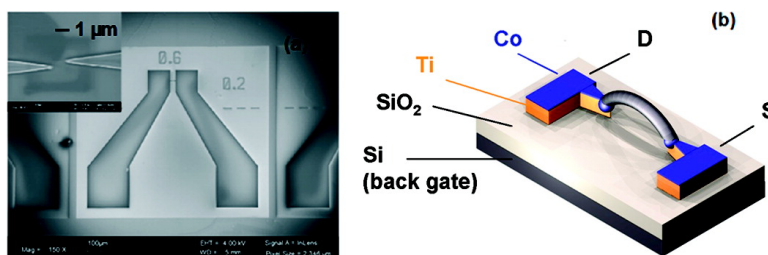


Negative Differential Resistance: Gate Controlled and Photoconductance Enhancement in Carbon Nanotube Intraconnects

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Negative Differential Resistance: Gate Controlled and Photoconductance Enhancement in Carbon Nanotube Intraconnects

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ABSTRACT

Intraconnects, as-grown single-walled carbon nanotubes bridging two metal electrodes, were investigated as gated structures. We show that even with a seemingly “ohmic” contact at zero gate voltage one observes negative differential resistance (NDR) at nonzero gate bias. Large differential photo conductance (DPC) was associated with the NDR effect raising hopes for the fabrication of novel high-speed optoelectronic devices.

Single-wall carbon nanotube (SWCNT) intraconnects (bridges spanning across two planar electrodes (Figure 1) have gained large attention from basic research and application^{1–11} points of view. Several investigations have been devoted to the theoretical^{12–14} and experimental^{15–20} aspects of the contact between the bridging tube(s) and metal. By comparison, fewer studies were devoted to a related negative differential resistance (NDR) effect, which is important for the construction of high-frequency oscillators. In the past, we have devised a way of fabricating SWCNT intraconnects at high yield; sharp electrode tips were used as seed catalysts thus, initiating the tube’s growth from tip to tip. We used this method to study the characteristics of gated structures with channels made of individual or small number of bundled tubes. Such electrode construction enabled us to interrogate intraconnects by Raman spectroscopy, current–voltage characteristics and their response to white light irradiation. As we shall see below, the intraconnects exhibited NDR, which was associated with a large differential photo con-

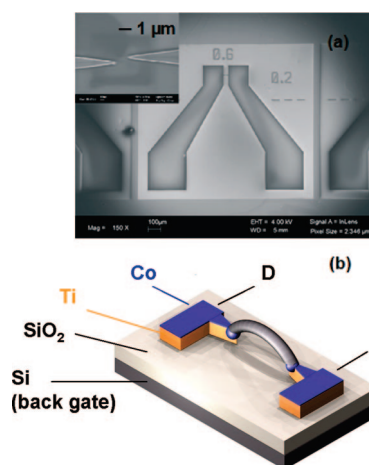


Figure 1. (a) SEM image of the metal electrodes. The distance between two sharp tips was 1 μm . Individually addressable CNT intraconnect was grown directly from tip to tip by CVD. (b) Schematic diagram of our CNT FET. Co was deposited as a catalyst on top of Ti electrode to initiate CNT intraconnect growth. Si substrate was used as a back gate.

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ductance for both ohmic and nonohmic contacts (defined at zero gate bias).

Experiments and Methods. SWCNT intraconnects were grown and self-aligned between prefabricated electrodes by use of chemical vapor deposition (CVD).^{21,22} For the present

study we replaced the CO precursor gas with methane/hydrogen mixture at elevated temperatures.²³ While some of the intraconnects could not be observed by scanning electron microscope (SEM), nonetheless they had a clear current–voltage (I – V) and Raman scattering characteristics. The electrode construction and labeling made it easy to assess the tubes even one year after fabrication. We made sure that nonmetalized surfaces, away from the electrode region, did not contribute to the electrical path. In addition, the electrode shape enabled focusing of the laser beam, used for Raman spectroscopy, right in-between the electrodes (see Supporting Information). All electrical and optical measurements were performed at room temperature.

Figure 1 shows an SEM image of the electrodes. Here, the distance between the two electrode tips was 1 μm , though the electrode layout had coaligned and laterally shifted patterns (Figure 1a inset). The schematic of a gated structure, made of CNT intraconnect, is shown in Figure 1b. The two metal electrodes were used as source and drain electrodes, respectively. A silicon substrate was used as a back gate electrode. The silicon surface was oxidized to a thickness of 20 nm prior to the electrode deposition.

Raman spectra were measured using the 514.5 nm line of an ion Ar^+ laser, a single grating spectrometer, which was equipped with two notch filters for laser line rejection and a N_2 cooled CCD detector. The system was equipped with a $\times 50$ microscope to aid focusing of the laser light between the sharp tips of the electrodes. The Raman spectra of SiO_2 and Si away from the bridge were used as reference and were subtracted from the spectra of the CNT intraconnect. Broadband light illumination was made with a white-light tungsten source, which had a high energy cutoff at 3 eV. The white-light source intensity was 0.25 mW/cm^2 , and it was illuminating the entire structure.

Results. Figure 2 shows two examples, one exhibiting nonlinear and the other linear electrical characteristics. Figure 2a–c shows SEM, $I_{\text{ds}}-V_{\text{ds}}$, and Raman spectra of Sample 1. The nonlinear I – V curve at zero gate voltage indicates a presence of potential barrier between the tube and metal contact(s). The asymmetry in the curve points to one dominant barrier (otherwise, the curve would have been symmetric). Its Raman spectra are shown in Figure 2c. The low frequency spectra exhibit a narrow single peak (5 cm^{-1} wide and limited only by the system resolution) for the radial breathing mode (RBM) at 191.9 cm^{-1} . The peak may indicate a metallic tube (12,6) with an average diameter of 1.252 nm.²⁴ It exhibits a very large graphitic line at 1355 cm^{-1} , which may indicate defects or, large stress. The high-frequency region exhibits two peaks at 1594 and 1567 cm^{-1} , respectively. Their frequency difference 27 cm^{-1} corresponds to a tube diameter $d = 1.716$ nm if we follow²⁵ and assume a metallic type tube. It corresponds to a tube diameter $d = 1.329$ nm assuming semiconductor type. Sample 2 exhibits linear I – V characteristics at zero gate bias (Figure 2e). Its low frequency Raman spectrum consists of a broader single RBM peak (30 cm^{-1} wide) centered at 176.2 cm^{-1} , which implies semiconductor tube(s) (11,9) with an average diameter of 1.365 nm. It exhibits a large defect line at 1352 cm^{-1} .

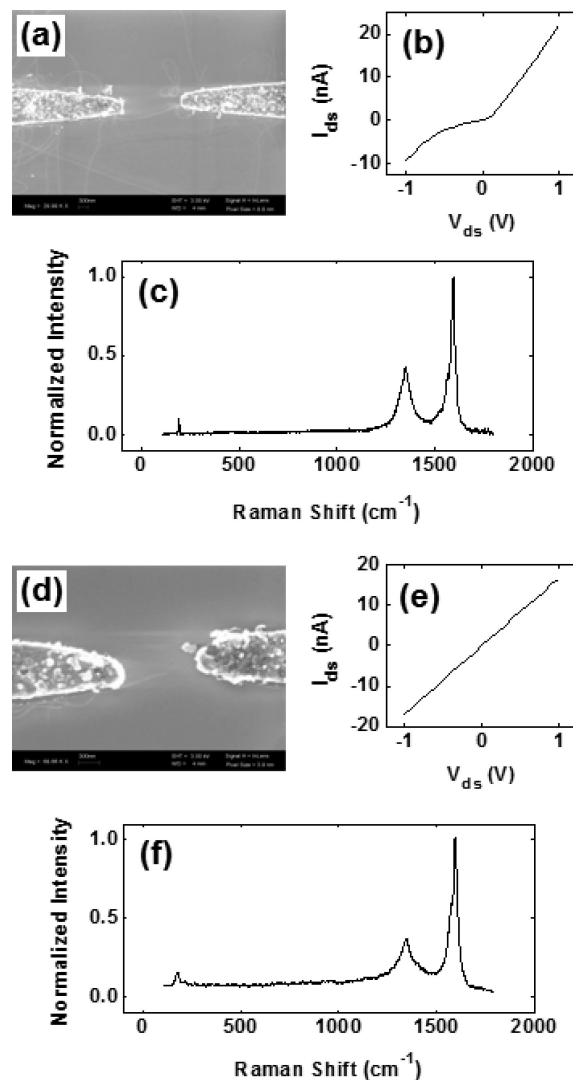


Figure 2. Sample 1: (a) SEM image of CNT intraconnect; (b) nonlinear $I_{\text{ds}}-V_{\text{ds}}$ characteristic; (c) Raman spectra with RBM at 191.9 cm^{-1} . Sample 2: (d) SEM image of CNT intraconnect; (e) linear $I_{\text{ds}}-V_{\text{ds}}$ characteristic; (f) Raman spectra with RBM at 176.2 cm^{-1} .

The high-frequency frequency exhibits two peaks at 1596 and 1578 cm^{-1} , respectively. Their difference 18 cm^{-1} implying tube diameters $d = 2.101$ nm and $d = 1.628$ nm for metallic or semiconductor type tubes, respectively.

Electrical properties in the dark and under white light illumination were measured for Sample 1 (Figure 3) and Sample 2 (Figure 4). The channel formed by the intraconnects exhibited a natural p -type characteristics owing to the presence of oxygen in the tubes. Figure 3a shows a map of V_{gs} vs V_{ds} for the nonilluminated intraconnect. Two negative differential resistance (NDR) regions may be observed in the $I_{\text{ds}}-V_{\text{gs}}$ plot: one, in the range between -3 and -6 V and the other, around -9 V (Figure 3b–c). The NDR peak for the first region was shifting to the negative side as the V_{ds} value increased from 0.25 to 0.75 V (Figure 3c). The second NDR region was stable around -9 V. Characteristics of the same intraconnect under illumination are shown in Figure 3d–f. Surprisingly, the light did not mask the NDR effect (Figure 3f) but locked it to the same V_{gs} range for all

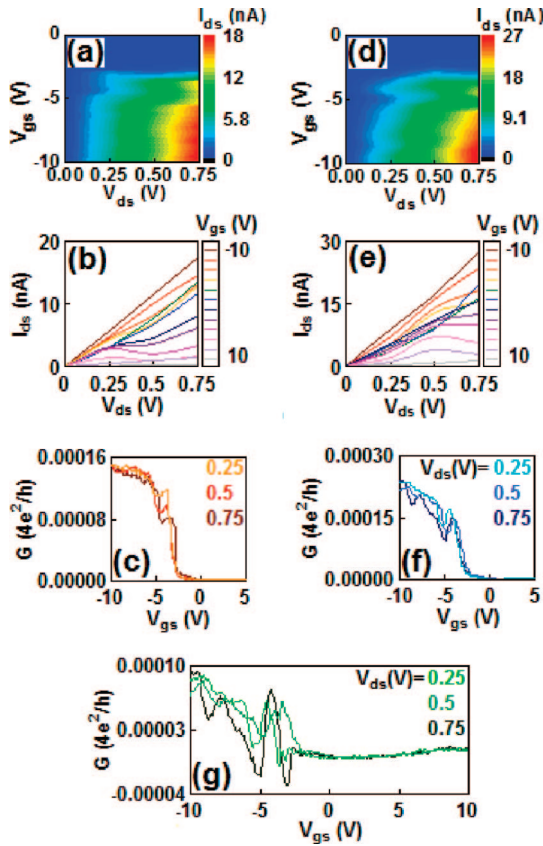


Figure 3. Electrical properties of Sample 1 before and after white light illumination. (a) I_{ds} vs V_{gs} and V_{ds} characteristics. Negative differential resistance was observed for V_{gs} between -2 and -6 . (b) I_{ds} – V_{ds} characteristic for various V_{gs} values from $+10$ to -10 . I_{ds} – V_{ds} curves become nonlinear when the gate voltages for a negative differential region are applied. (c) G – V_{gs} characteristics. Panel a is converted into conductance, G . (d) I_{ds} vs V_{gs} and V_{ds} characteristics after irradiation of white light. (e) I_{ds} – V_{ds} characteristic for various V_{gs} values from $+10$ to -10 . I_{ds} – V_{ds} curves become nonlinear when the gate voltages for a negative differential region are applied. (f) G – V_{gs} characteristics under illumination. (g) Photo differential conductance.

V_{ds} values used. The differential photo conductance (DPC, the difference between conductance in the dark and under light illumination) was fluctuating rapidly within the first NDR range (Figure 3g).

Figure 4a shows a map of V_{gs} vs V_{ds} for nonilluminated intraconnect. Two negative differential resistance (NDR) regions are observed in the I_{ds} – V_{gs} curve as well: one in the range between -3 and -6 V and another around -9 V (Figure 4b,c). The NDR peak for the first region was shifting albeit in both negative and positive directions as V_{ds} increased from 0.25 to 0.75 V. The second NDR region was stable around -9 V. The same intraconnect under white light illumination is shown in Figure 4d–f. Here, the light somewhat masked the NDR effect (Figure 3f) for the first V_{gs} region but left a small but noticeable presence around $V_{gs} \sim -9$ V. The DPC was stable and exhibited a marked increase for increasing V_{ds} (Figure 4g).

Discussion. NDR effect is the basis for many microwave oscillators.²⁷ The effect may be attributed to two competing transport venues, for example, nonlinear drift and thermionic venues or tunneling and thermionic excitation processes. In

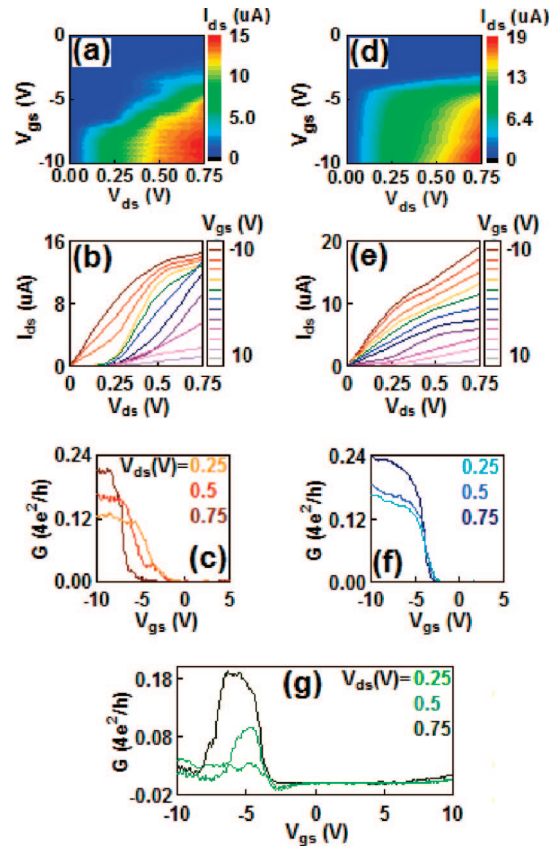


Figure 4. Electrical properties of Sample 2 before and after white light illumination. (a) I_{ds} vs V_{gs} and V_{ds} characteristics. (b) I_{ds} – V_{ds} characteristic for various V_{gs} values from $+10$ to -10 . (c) G – V_{gs} characteristics. Panel a is converted into conductance, G . (d) I_{ds} vs V_{gs} and V_{ds} characteristics under white light illumination. Overall current was enhanced and the device became stable with less noise. (e) I_{ds} – V_{ds} characteristic for various V_{gs} values from $+10$ to -10 . I_{ds} – V_{ds} curves become nonlinear when the gate voltages for a negative differential region are applied. (f) G – V_{gs} characteristics under illumination. (g) Photo differential conductance.

addition, optimal construction typically breaks the symmetry between the contact to and from the channel. For example, a CNT channel extending between two different metals,¹³ which takes advantage of the metals' different work functions and the distribution of the tube's electronic density of states. In contrast, both of our electrodes were made of the same metal and hence we resorted to a gated configuration. The gate bias affects the distribution of carriers along the tube. If we postulate the existence of a secondary interface between the CNT channel and electrode, possibly in a form of a point Schottky contact or a quantum dot (QD), gate biasing affects the distribution of electronic states across such barrier, as well. We point to the very small, 1 nm thick catalytic seed film in our structures, which may result in a point contact during the tube's growth process. Tunneling may now occur between the tube and discrete states in this metal dot. The secondary NDR region for both samples may imply tunneling through a second QD band. Tunneling and transport must depend on the relative energy states' distribution between dot and channel, as well as, on the gate bias. NDR was also noted in our previous experiments while using another

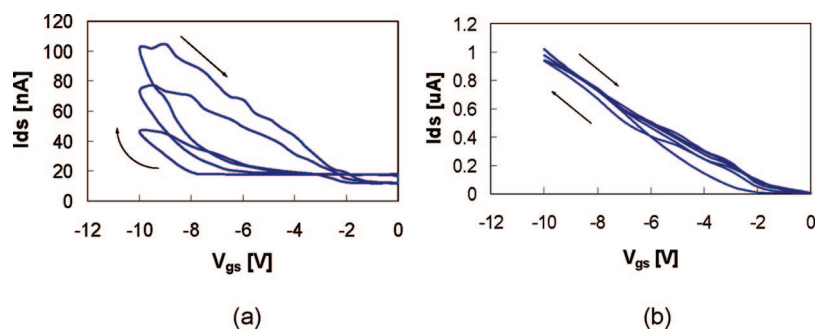


Figure 5. Hysteresis curves in the dark: (a) Nonlinear case, sample 1. (b) Linear case, sample 2. The arrows point the direction of the scan. In panel a, the overall current increased for increasing numbers of scans until stabilizing.

growth technique [ref 21, Figure 3] albeit at larger gate bias.

Sample 1 is probably made of individual or of a few tubes as judged by the narrow Raman RBM spectra; its current only weakly dependent on gate voltage and one may attribute its behavior to mostly the Schottky point interface. Its sensitivity to the contact with electrodes led to the asymmetry in the $I_{ds}-V_{ds}$ curve. In contrast, sample 2 is probably made of a small bundle thus is more likely to portray ohmic contact at zero gate bias. Sample 2 is more responsive to the gate bias as expected of a semiconductor channel.

Raman classification of the tubes may be distorted by stress and contact to the substrate.²⁶ In addition and despite careful alignment (see Supporting Information), the laser spot may cover tubes, which do not contribute to electrical conduction yet, may be at resonance with the laser line thus, affecting the Raman spectra. This could imply that both samples were of semiconductor type if we mainly base our conclusion on a self-consistent value for the tube's diameter. In that case, the samples differ only by the number of tubes in their respective bundles. Nevertheless, Raman spectroscopy is an important characterization tool, which corroborates the existence of nanotubes between electrodes and helps determine their type, thus complementing conductance measurements.

The characteristics of the channel under broadband optical excitation corroborated these findings, as well. Optical induced carriers in the channel overcame the tunneling process, hence the marked conductivity change for the first NDR region (around -5 V). The optical effect is smaller for the “nonlinear” sample 1 because of the interfering effect from the contact barrier. Moreover, since photoexcited carriers are composed of both types (holes and electrons), n -type characteristics at large positive gate voltages ought to be exhibited. Indications to such behavior are provided in Figure 4g. Furthermore, when the graphs are plotted as log plots (not shown) such characteristics are greatly accentuated. Both panels c and f of Figure 3 indicate a relatively sharp “ON” transition, which implies large channel mobility with back-to-back leaky Schottky diodes.²⁸ The ON transition in Figure 4c,f tends to have quadratic behavior as expected from an ordinary semiconducting channel.

Surface states play an important role in such devices. We observed hysteresis upon ramping the gate voltage in the negative and positive directions (Figure 5). The lower curves

are always in the negative direction while the top curves are scanned from negative to positive gate bias values. The hysteresis did not change the presence of NDR which means that the latter is inherent to the tube/electrode contact but not to its surface. After a few scans, the hysteresis became narrower. This is to be expected of surface states. In addition, clockwise hysteresis is indicative of movable ions or, adsorbed water molecules (in contrast to charge trapped in the underlined oxide layer – the latter results in counter-clockwise hysteresis curve). Yet, sample 1 exhibited a large hysteresis and sample 2 exhibited almost a replica of the curve in the negative and positive gate bias ramp. That may be understood if the channel of sample 1 is composed of only few or even individual carbon nanotube. The effect of surface states must have been larger here than their effect on sample 2 with a larger bundle of tubes.

In summary, we demonstrated NDR in carbon nanotube gated structures under bias regardless of the linearity (or nonlinearity) of the contacts. The effect was accompanied by a large DPC. Since NDR enables the realization of microwave oscillators, one may envision that such DPC will be instrumental in designing new high-speed opto-electronic oscillators.

Supporting Information Available: This material is available free of charge via the Internet at <http://pubs.acs.org>.

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Negative Differential Resistance: Gate Controlled and Photoconductance Enhancement in Carbon Nanotube Intra-connects

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Supporting Information. The metal electrodes were composed of 1 nm of Co on top of 30 nm of Ti layer. The silicon substrate was oxidized (20 nm) prior to the electrode deposition. The electrodes were annealed at 550 °C for 10 min prior to the tube growth in order to break the electrode tip into smaller islands. Hydrogen gas was introduced into a quartz tube for 5 min in order to reduce the oxidized Co. Methane and hydrogen gas mixture were used as gas precursors for an hour at 900 °C at a flow rate of 1900 sccm for methane and 20 sccm for hydrogen. Samples were cooled down to room temperature under argon gas flow. A typical FE-SEM system has a resolution of 5 nm and 1-5 nm tubes may be resolved by their shadow or halo on the substrate. That in turn, depends on the substrate's surface states and their ability to charge. However, we later found that in fact, good quality bridges are often barely seen in these SEM pictures, hence the use of Raman spectroscopy and conductivity measurements. Our present batch consisted of 14 samples (9 samples conducting: 4 – linear; 5 – nonlinear), is taking into

account cases, which do not show up in SEM pictures yet, exhibited conductance and Raman signature. Temperature is critical to a successful SWCNT growth: lower yield (2 conductive samples out of 14) was exhibited at lower temperatures deviating only by 50 °C.

Current-voltage (I-V) curves were measured using Keithley 236 with current noise level <1 pA using sharp probes (~5 µm, Cascade Microtech). The probes were placed under a microscope interfaced with a CCD camera for better inspection. Gate-current measurements were made with the Si wafer acting as a back electrode. The conduction through a CNT bridge at $V_{ds}=1$ V was at least three orders of magnitudes larger than the noise and/or the current through or, along the silicon and/or electrode/silicon path. The signal, carried by triax cables, was averaged during a period larger than the RC constant ($\tau < 0.03$ sec) of the circuit.

Two more examples of ‘linear’ and ‘non-linear’ contacts are shown in Figure S1 and Figure S2, respectively. The SEM images of the bridge under test are provided, as well. Clearly seen is the breakdown of the upper Co layer. The resistance of the electrode (~10 ohms/ 100 microns) did not change much before and after annealing. This is to be compared to the bridge resistance of several tens of megaohms. Both samples exhibited NDR around $V_{gs}=-5$ V and -10 V, respectively. The clearer picture is again for the ‘linear case’. We note that previous measurements [21-22] also exhibited NDR, which at the time, were not fully pursued. These bridges were fabricated using CVD with another precursor (CO) and similar layout of electrodes (with the exception of a thicker Co layer of 30 nm). Such precursor requires lower processing temperatures (~750 °C). Thicker Co layer resulted in occasional growth of multi-walled carbon nanotubes. Due to the sensitivity of the process to temperature fluctuations, a critical factor is therefore, the actual temperature at the substrate surface. This is affected by several factors such

as, the gas flow, the location of the sample in the CVD chamber, the heat capacity of the sample holder, etc.

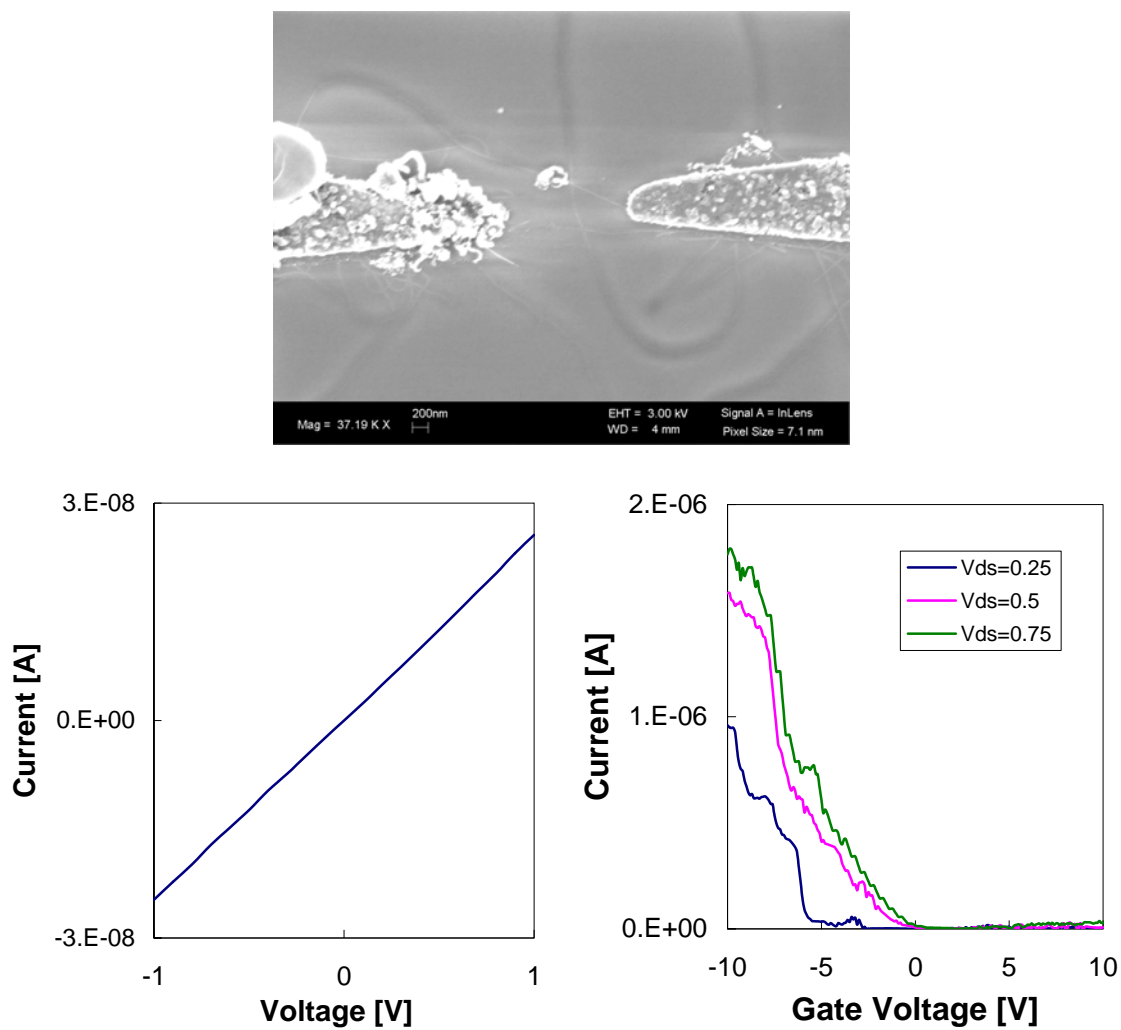


Figure S1. A “linear” intra-connect. Bottom Left: I- V_{ds} . Bottom Right: I- V_{gs} .

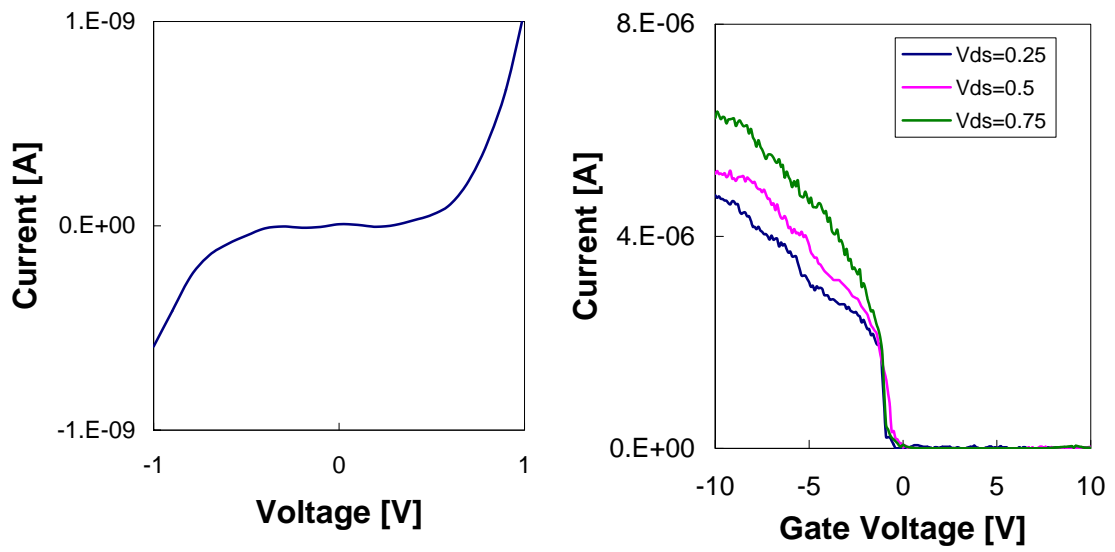
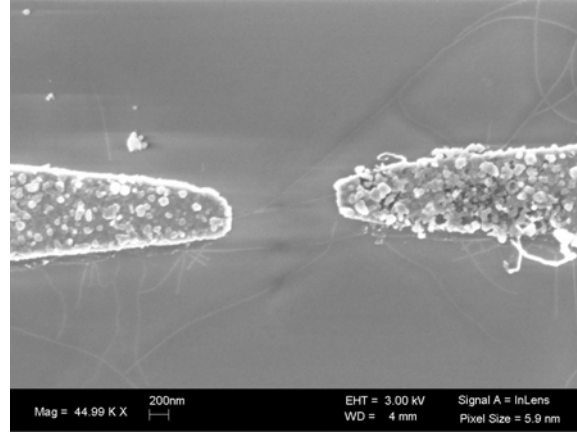


Figure S2. A “non-linear” intra-connect. Bottom Left: I-V_{ds}. Bottom Right: I-V_{gs}.

The electrode shape made it easy to place the laser spot right between them. The Raman system was interfaced with a microscope which enabled imaging the electrodes and laser spot alike. In Figure S3 we show an example of such image. The laser spot is slightly larger than the spacing between electrodes and therefore, the resultant Raman spectra may be affected by tubes, which are not forming a conducting bridge. By scanning the spot in the vicinity of the bridge we concluded that Raman signal for carbon nanotubes is detected only when the spot is placed directly between the electrodes. Otherwise, the detected signal consisted of only silicon and silicon oxide. The Raman spectra of SiO_2 and Si away from the bridge were used as reference and were subtracted from the spectra of the SWCNT intra-connect.



Figure S3. An example of a laser beam partially focused between the electrode tips for Raman spectroscopy measurements. Slight defocusing was necessary for picture taking (otherwise, the intensity of the laser beam would have overwhelmed the white light background intensity).