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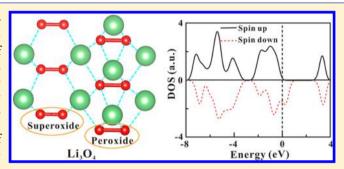


# A Stable, Magnetic, and Metallic Li<sub>3</sub>O<sub>4</sub> Compound as a Discharge Product in a Li-Air Battery

Guochun Yang, †,‡ Yanchao Wang,\*,† and Yanming Ma\*,†

Supporting Information

ABSTRACT: The Li-air battery with the specific energy exceeding that of a Li ion battery has been aimed as the nextgeneration battery. The improvement of the performance of the Li-air battery needs a full resolution of the actual discharge products. Li<sub>2</sub>O<sub>2</sub> has been long recognized as the main discharge product, with which, however, there are obvious failures on the understanding of various experimental observations (e.g., magnetism, oxygen K-edge spectrum, etc.) on discharge products. There is a possibility of the existence of other Li-O compounds unknown thus far. Here, a hitherto unknown Li<sub>3</sub>O<sub>4</sub> compound as a discharge product of the Li-



air battery was predicted through first-principles swarm structure searching calculations. The new compound has a unique structure featuring the mixture of superoxide  $O_2^{-1}$  and peroxide  $O_2^{-2}$ , the first such example in the Li–O system. The existence of superoxide O<sub>2</sub> creates magnetism and hole-doped metallicity. Findings of Li<sub>3</sub>O<sub>4</sub> gave rise to direct explanations of the unresolved experimental magnetism, triple peaks of oxygen K-edge spectra, and the Raman peak at 1125 cm<sup>-1</sup> of the discharge products. Our work enables an opportunity for the performance of capacity, charge overpotential, and round-trip efficiency of the Li-air battery.

SECTION: Molecular Structure, Quantum Chemistry, and General Theory

s a primary energy resource, oil in total amount is rather Alimited, and the ever-increasing energy demand will not be satisfied in the long run. It is more problematic that oil consumption generates CO2 and other pollutants. Clean and sustainable energy resources need to be developed in an effort to decrease our dependency on oil.<sup>1,2</sup> The rechargeable nonaqueous Li-air battery has received a great deal of interest in the past few years due to its extremely high specific energy, and it therefore has become one of the most promising candidates as an alternative energy resource for future electric vehicles.<sup>3–9</sup> However, the technology of nonaqueous Li–air batteries is still in its infancy, and key technical and scientific challenges remain to be overcome. 10-13

The performance of the nonaqueous Li-air battery correlates closely with the structure, composition, and physical proprieties of the actual discharge products,  $^{14-17}$  which are far from being well established.  $^{18}$  Previous studies suggested commonly that  $\text{Li}_2\text{O}_2$  is one of the discharge products.  $^{14,19-21}$ 

However, X-ray diffraction (XRD) patterns of the discharge products are considerably broadened compared to that of standard  $\text{Li}_2\text{O}_2$ , which may be attributed to poor crystallinity, structural defects, and/or composition non-stoichiometry of  $\text{Li}_2\text{O}_2$ . Moreover, oxygen (O) or Li Kedge spectra showed that local environments of O and Li in the discharged electrodes are different from those of Li<sub>2</sub>O<sub>2</sub>. For example, there is a strong peak that appeared at 535 eV in the O K-edge spectra that is not observable in standard Li<sub>2</sub>O<sub>2</sub>. The origin of this strong peak remains unknown although it was postulated to arise from structural defects and/or the nonstoichiometry composition (Li<sub>2-x</sub>O<sub>2</sub>) of Li<sub>2</sub>O<sub>2</sub>. More recently, an experimental study showed that the discharge products exhibit magnetism and a characteristic Raman peak at 1125 cm<sup>-1</sup>, both of which were not shown in the standard Li<sub>2</sub>O<sub>2</sub> sample. Subsequently, it was proposed that the magnetism and the unexplained Raman peak might be ascribed to the superoxide-like surface structure or nanocrystallization of Li<sub>2</sub>O<sub>2</sub>. Another study even found that the discharge products contain two components, an oxygen-rich component with superoxide-like character and a Li<sub>2</sub>O<sub>2</sub> component. The oxygenrich superoxide-like component has a much smaller charge potential than the Li<sub>2</sub>O<sub>2</sub> component. However, the actual structure and nature of the superoxide-like component are unknown.<sup>23</sup>

By all appearance, there is a necessity to re-examine the phase diagram of Li-O systems in an effort to explore other possible Li-O compounds, which are potential discharge products. In fact, Li<sub>2</sub>O and LiO<sub>2</sub> have previously been proposed, but they were subsequently ruled out. 12,25 The

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XRD pattern of Li<sub>2</sub>O has a significant deviation from the experimental data of the discharge product.<sup>26</sup> LiO<sub>2</sub> is unstable at room temperature,<sup>27</sup> the operating temperature of a Li—air battery. Other stoichiometries in the Li—O system remain unknown thus far, but in view of the various puzzles related to the discharge product, they might be potentially important to the understanding of discharge products.

Here, we explore the entire phase diagram of the Li–O system by using the first-principles swarm-intelligence structural prediction calculations  $^{28,29}$  unbiased by any prior known structures. A hitherto unknown  $\rm Li_3O_4$  compound was uncovered for the first time as a stable discharge product and found to exhibit intriguing half-metallic magnetism. The finding of  $\rm Li_3O_4$  enables the natural explanation of earlier unsolvable experimental observations. Our work represents a significant step forward toward the understanding of discharge products in a Li–air battery.

The phase stabilities of various  $\text{Li}_x \text{O}_y$  (x = 1-4, y = 1-4) compounds were investigated by calculating their formation energies at 0 K relative to the products of dissociation into constituent elements, as summarized in Figure 1. In general, a

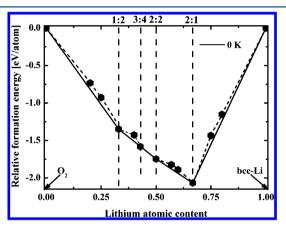
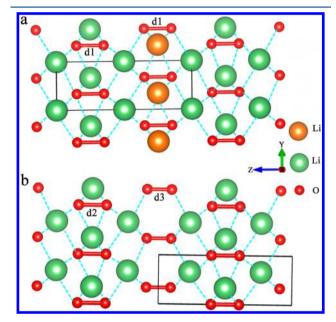


Figure 1. Relative formation energy with respect to solid Li and oxygen gas  $(O_2)$  for various Li-O stoichiometries at 0 K. The bodycentered cubic (bcc) structure of elemental Li was used. Because Li has a light atomic mass, the zero-point energy was included in the formation energy calculations. The convex hulls are shown by solid lines. Dotted lines that directly connect data points are guides for the eyes.

structure whose formation energy lies on the convex hull (i.e., the phase stability line) is deemed stable with respect to decomposition into other Li-O compounds or elemental solids and thus can be experimentally synthesized.<sup>30</sup> As shown in Figure 1 at 0 K, the formation energy data of LiO2, Li3O4, Li<sub>2</sub>O<sub>2</sub>, and Li<sub>2</sub>O sit right on the convex hull, while that of Li<sub>2</sub>O<sub>3</sub> is slightly above the convex hull. This indicates that under controllable experimental conditions, LiO2, Li3O4, Li2O2, and Li<sub>2</sub>O are synthesizable at low temperature. Note that the formation energy data were obtained by using the structures derived from our first-principles structure searching simulations, and they will be described in detail below. With the only input of chemical compositions in our CALYPSO structure searching calculations, the experimental antifluorite structure of Li<sub>2</sub>O (space group Fm3m, four formula units per cell) and the hexagonal structure of Li<sub>2</sub>O<sub>2</sub> (space group P6<sub>3</sub>/ mmc, two formula units per cell) were successfully reproduced, validating our structure searching methodology in application to the Li-O system.

The bulk crystalline structure of LiO<sub>2</sub> has not yet been fully resolved.<sup>27</sup> Most of the studies assumed the orthorhombic *Pnnm* structure as obtained by replacing Na with Li for the known NaO<sub>2</sub> structure.<sup>31–33</sup> Our structure searching calculations unbiased by any prior known structure confirmed that the *Pnnm* structure is indeed the most stable structure. The lattice parameters of the *Pnnm* structure are summarized (Table S2, Supporting Information).

Our structure searching calculations identified a new stable  $\text{Li}_3\text{O}_4$  compound with a hexagonal structure (space group *P-6m2*, one formula unit per cell; see Figure 2b). The structure



**Figure 2.** Crystal structures of  $\text{Li}_2\text{O}_2$  (a) and  $\text{Li}_3\text{O}_4$  (b). The solid rectangle represents one unit cell. The O–O distance (d1) in  $\text{Li}_2\text{O}_2$  is 1.55 Å, while O–O distances (d2 and d3) in  $\text{Li}_3\text{O}_4$  are 1.54 and 1.35 Å, respectively.

contains two inequivalent Li's occupying 2h (0.3333, 0.6667, 0.7671) and 1f (0.6667, 0.3333, 0.5000) positions and two inequivalent O's sitting at 2g (0.0000, 0.0000, 0.3956) and 2i (0.6667, 0.3333, 0.0914) sites. Each Li atom forms a six-fold coordination with O atoms, having three different Li-O distances of 1.97, 2.17, and 2.09 Å. All O atoms exist in quasimolecular O2 forms. There are two kinds of O2 molecules with different intramolecular O-O bond lengths (1.35 and 1.54 Å), which are apparently longer than the bond length in gas-phase O<sub>2</sub> (1.21 Å).<sup>34</sup> These enlarged intramolecular bonds are easily understood because O2 molecules in Li3O4 are negatively charged (Table S3, Supporting Information). The transferred charges from Li to  $O_2$  occupy the anithonding  $\pi^*$  orbitals of  $O_2$ molecules to weaken and lengthen the intramolecular bond. Li<sub>3</sub>O<sub>4</sub> is structurally very similar to Li<sub>2</sub>O<sub>2</sub> (Figure 2a). The removal of brown Li atoms in Li<sub>2</sub>O<sub>2</sub> leads to the formation of Li<sub>3</sub>O<sub>4</sub> and thus shorter O-O bonds from d1 to d3. The shorter d3 bond length of 1.35 Å is within the length range (1.3–1.4 Å) of the superoxide  $O_2^{-35}$  All of the intramolecular O–O bond lengths in Li<sub>2</sub>O<sub>2</sub> are equal to 1.55 Å, a characteristic bond length for a peroxide  $O_2^{\frac{1}{2}}$  seen also as the longer d2 bond in Li<sub>3</sub>O<sub>4</sub>. In this regard, the structure of Li<sub>3</sub>O<sub>4</sub> can be well recognized as a mixture of superoxide  ${\rm O_2}^-$  and peroxide  ${\rm O_2}^{2-}$ groups. Note that our structural model of Li<sub>3</sub>O<sub>4</sub> is

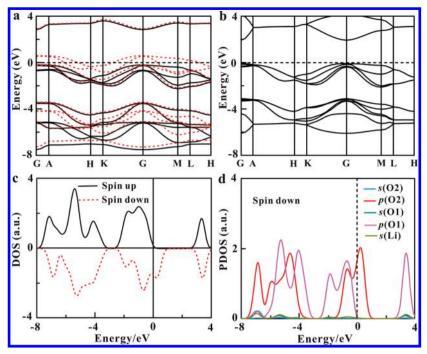


Figure 3. (a,b) Spin and electron band structures for  $Li_3O_4$  and  $Li_2O_2$ , respectively. Spin-up and spin-down bands are shown in solid and dotted lines, respectively. (c) Spin DOS for  $Li_3O_4$ . (d) Projected spin-down DOS for  $Li_3O_4$ . O1 and O2 are within peroxide and superoxide groups, respectively. The dashed lines represent the Fermi levels.

fundamentally different from the earlier Li-vacancy or O-rich surface counterparts of  ${\rm Li_2O_2}$ .  $^{18,36-38}$ 

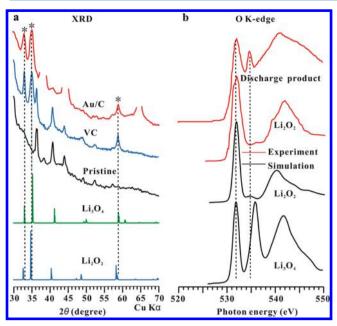
To assess the phase stability of Li<sub>3</sub>O<sub>4</sub> at room temperature, the quasi-harmonic model<sup>39</sup> was adopted to calculate the Gibbs free energy at 300 K through calculations of phonon spectra using the supercell method. The phonon results (Figure S1, Supporting Information) confirmed that the predicted Li<sub>3</sub>O<sub>4</sub> structure is dynamically stable without showing any imaginary phonon modes. Because the considered temperature of 300 K is far below the melting temperatures (e.g., melting temperature at 1711 K for  $\text{Li}_2\text{O}$ ), <sup>40</sup> the quasi-harmonic model should be rather valid. As a further test, the entropic curve of Li<sub>2</sub>O was calculated to compare with the experimental data.<sup>25</sup> It is evident that theory matches nicely with experiment (Figure S2, Supporting Information). It is seen (Figure S3, Supporting Information) that Li<sub>3</sub>O<sub>4</sub> is a thermodynamically stable phase at 300 K with respect to various dissociation routes into Li<sub>2</sub>O<sub>2</sub> +  $O_2$ ,  $Li_2O + O_2$ , and  $Li_2O_2 + LiO_2$ . From our calculations,  $Li_2O$ is the energetically most favorable stoichiometry, but it was seldom seen in the experiments on the discharge products. 12,25 This stemmed from the fact that the Li-air battery is working in O-rich conditions, where Li<sub>2</sub>O<sub>4</sub> and Li<sub>2</sub>O<sub>2</sub> are more stable with respect to  $Li_2O + O_2$  (Figure S3, Supporting Information). Below, we focus on the discussion of the stable Li<sub>3</sub>O<sub>4</sub> compound and its correlation with the discharge product.

The spin band structures for  $\mathrm{Li_3O_4}$  are depicted in Figure 3a. It can be clearly seen that  $\mathrm{Li_3O_4}$  exhibits a peculiar hole-doped metallic behavior. More intriguingly, a spin density of states (DOS) plot (Figure 3c) reveals novel half-metallic magnetism. The spin-up states are insulating, while the spin-down states are conducting. Analysis of the projected spin DOS suggested that the magnetism and conducting states arise from the superoxide  $\mathrm{O_2}^-$  groups (Figure 3d). The same mechanism on the superoxide-induced magnetism and metallicity was also reported in the Li-vacancy or O-rich surfaces of  $\mathrm{Li_2O_2}^{.18,36}$  In

an effort to better understand the hole-doped metallic behavior in Li $_3O_4$ , the insulating band structures of Li $_2O_2$  are plotted in Figure 3b. Line shapes of Li $_2O_2$  resemble those of the spin-up energy band of Li $_3O_4$ . Li acts as an electron donor, and the half depletion of Li out of Li $_2O_2$  leads to the substantial electron loss and thus lifts the spin-down states above the Fermi level. Magnetism and metallic holes are inevitably created. The appearance of magnetism in Li $_3O_4$  gives rise to the direct understanding of the experimental magnetism of the discharge product. <sup>24</sup>

The primitive unit cell of Li<sub>3</sub>O<sub>4</sub> contains seven atoms. As a result, 18 optical phonon modes are expected at the zone center. Group theory analysis (Table S4, Supporting Information) showed that there are 10 Raman-active modes  $(1E'' + 2E'' + 3A'_1 + 4E')$ , 8 infrared-active modes (4E' + $4A''_{2}$ ), and three silent modes  $(3A''_{1})$ . Of particular interest is the existence of a high-frequency Raman  $A'_1$  vibration at 1101 cm<sup>-1</sup>, which is not seen in standard Li<sub>2</sub>O<sup>41</sup> and Li<sub>2</sub>O<sub>2</sub><sup>42</sup> (Tables S5 and S6, Supporting Information). This high-frequency mode corresponds to the intramolecular O-O stretching vibration of the superoxide group (Figure S1, Supporting Information). It is interesting to find that the two Raman-active vibrations at 788 and 256 cm<sup>-1</sup> of Li<sub>2</sub>O<sub>2</sub> appeared also in Li<sub>3</sub>O<sub>4</sub> (739 and 224 cm<sup>-1</sup>). They are intimately related to the O-O vibrations of the peroxide group. The appearance of the characteristic Raman vibration at 1101 cm<sup>-1</sup> in Li<sub>3</sub>O<sub>4</sub> correlates nicely with the experimental Raman peak at 1125 cm<sup>-1</sup> measured from the discharge product.<sup>24</sup>

The XRD spectra of various Li–O compounds (LiO<sub>2</sub>, Li<sub>2</sub>O, Li<sub>3</sub>O<sub>4</sub>, and Li<sub>2</sub>O<sub>2</sub>) were simulated (Figures 4a and S4, Supporting Information) to compare with the experimental data. Except for slight peak shifts, the major XRD features of Li<sub>3</sub>O<sub>4</sub> are nearly identical to those of Li<sub>2</sub>O<sub>2</sub>, and fundamentally, they both can explain the experimental data. In contrast, LiO<sub>2</sub> and Li<sub>2</sub>O have XRD patterns in substantial deviation with the



**Figure 4.** (a) The simulated XRD spectra of  $\text{Li}_2\text{O}_2$  and  $\text{Li}_3\text{O}_4$  to compare with the experimental data on discharge and pristine electrodes supported on a Celgard 480 separator (100 mA  $g_{\text{carbon}}^{-1}$ ) for Vulcan carbon (VC) and Au-catalyzed Vulcan carbon (Au/C) (ref 22). The star symbols are indications of the characteristic peaks of the discharge products. (b) The simulated O K-edge spectra of  $\text{Li}_2\text{O}_2$  and  $\text{Li}_3\text{O}_4$  to compare with experimental data of discharged Au/C electrodes at 100 mA  $g_{\text{carbon}}^{-1}$  (ref 22).

experiments, and they thus should be ruled out as the discharge product. Further comparison (Figure S5, Supporting Information) was also made by using another experimental XRD measurement.<sup>21</sup>

The calculated O K-edge spectra for  $\text{Li}_2\text{O}_2$  and  $\text{Li}_3\text{O}_4$  together with the available experimental data<sup>22</sup> are shown in Figure 4b. The experimental spectrum of the pure  $\text{Li}_2\text{O}_2$  sample was well reproduced by our calculations. The strong experimental peak at about 535 eV on the discharge product is apparently missing in the  $\text{Li}_2\text{O}_2$  sample.<sup>43,44</sup> In contrast,  $\text{Li}_3\text{O}_4$  accounts for this strong peak. Further analysis revealed that such a peak in  $\text{Li}_3\text{O}_4$  originates from the intramolecular O–O bonding of the superoxide group (Figure S6, Supporting Information).

It is known that the decomposition of the electrolyte and the oxidation of the carbon matrix in the Li–air cell will form side products such as Li<sub>2</sub>CO<sub>3</sub>, HCO<sub>2</sub>Li, and CH<sub>3</sub>CO<sub>2</sub>Li. 45,46 It is apparent that our finding of the Li<sub>3</sub>O<sub>4</sub> compound is not the side products from the standpoint of IR, Raman, and XRD spectra (Figures S7 and S8, Supporting Information). Moreover, the Li<sub>3</sub>O<sub>4</sub> compound with the half-metallic magnetism is greatly different from these side products, which are insulators.

By all above evidence, it is suggested that in addition to  $\text{Li}_2\text{O}_2$ ,  $\text{Li}_3\text{O}_4$  is deemed to exist as one of the discharge products. Unfortunately, we cannot determine the ratio between  $\text{Li}_2\text{O}_2$  and  $\text{Li}_3\text{O}_4$  based on our calculation, which needs further experimental confirmation. Moreover, our finding of  $\text{Li}_3\text{O}_4$  is not against the fact that  $\text{Li}_2\text{O}_2$  is the main discharge product because the coexistence of  $\text{Li}_2\text{O}_2$  with  $\text{Li}_3\text{O}_4$  does not contradict the experimental observation. A viable competition between  $\text{Li}_3\text{O}_4$  and  $\text{Li}_2\text{O}_2$  in the discharge product might be expected. First, energetically,  $\text{Li}_3\text{O}_4$  and  $\text{Li}_2\text{O}_2$  are competitive. The formation free energy of  $\text{Li}_3\text{O}_4$  is -3.09 eV/Li, which is

only slightly favorable over -3.01 eV/Li in Li<sub>2</sub>O<sub>2</sub>. Second, although both Li<sub>2</sub>O<sub>2</sub> and Li<sub>3</sub>O<sub>4</sub> can explain the experimental XRD data, the understanding of considerably broadened XRD spectra is still missing. If the contribution from both Li<sub>2</sub>O<sub>2</sub> and Li<sub>3</sub>O<sub>4</sub> is assumed, the broadened XRD feature is explainable (Figure S5, Supporting Information).<sup>21</sup>

Besides the thermodynamical stability of  $\text{Li}_3\text{O}_4$ , it is necessary to assess its electrochemical behavior because the electrochemical kinetics, that is, the overpotentials of formation, are also important to govern the actual formation of  $\text{Li}_3\text{O}_4$  during the discharge process. We have explicitly calculated the electrochemical behavior of  $\text{Li}_3\text{O}_4$  and  $\text{Li}_2\text{O}_2$  through calculation of their reduction potentials (calculation details can be found in the Supporting Information). The obtained reduction potentials of  $\text{Li}_2\text{O}_2$  (0.45 V) and  $\text{Li}_3\text{O}_4$  (0.38 V) are rather similar. Therefore, the electrochemical behavior of  $\text{Li}_3\text{O}_4$  has a great similarity to that of  $\text{Li}_2\text{O}_2$ . Because  $\text{Li}_2\text{O}_2$  is reversible during the discharge and charge processes, we expect that  $\text{Li}_3\text{O}_4$  might be also reversible.

It is known that the high resistance of  $\text{Li}_2\text{O}_2$  limits greatly the rate capacity, results in high charge overpotential, and hampers the rechargeability of the Li–air battery. \(^{14,47,48}\) Earlier studies suggested that electron transport can be enhanced by creation of Li-vacancy or O-rich surfaces of  $\text{Li}_2\text{O}_2$ . \(^{18,36}\) However, the generation of a substantial amount of Li vacancies is unfavorable, and O-rich surfaces cannot offer a transportation path across the film. \(^{49}\) The finding of metallic  $\text{Li}_3\text{O}_4$  opens up the possibility for optimizing electron conduction to improve the performance of the Li–air battery. A proper control on the experimental condition to generate an optimal amount of  $\text{Li}_3\text{O}_4$  in discharge products is greatly demanded.

In summary, the use of the advanced swarm structure searching technique allowed us to identify a hitherto unknown  ${\rm Li_3O_4}$  compound, which is energetically stable in O-rich conditions with respect to the dissociation into  ${\rm Li_2O_2} + {\rm O_2}$  and  ${\rm Li_2O} + {\rm O_2}$ .  ${\rm Li_3O_4}$  has a peculiar structural feature consisting of both superoxide  ${\rm O_2}^-$  and peroxide  ${\rm O_2}^{2-}$  groups, the only known mixed structure example so far in the Li–O system. The superoxide  ${\rm O_2}^-$  groups induced the magnetism and the hole-doped metallicity of the system. The unresolved experimental observations on magnetism, the O K-edge spectrum, and the Raman peak at 1125 cm<sup>-1</sup> were fundamentally understood. Our work represents a significant step forward toward the understanding of a Li–O phase diagram and possibly provides a pathway for optimizing the performance of the Li–air battery.

# **■ COMPUTATIONAL METHODS**

Our structure searching simulations are performed through the CALYPSO structure prediction method <sup>28,29</sup> based on a particle swarm optimization algorithm via a global minimization of free-energy surfaces merging ab initio total energy calculations as implemented in the CALYPSO code, <sup>28</sup> which is unbiased by any prior known structural information. Our method has been benchmarked on various known systems, ranging from elements to binary and ternary compounds. <sup>28,50-53</sup> Detailed description of the method and predictions can be found in the Supporting Information.

Total energy calculations were performed in the framework of density functional theory within the generalized gradient approximation<sup>54</sup> as implemented in the VASP code.<sup>55</sup> The electron—ion interaction was described by means of a projector augmented wave with s<sup>1</sup>p<sup>0</sup> and s<sup>2</sup>p<sup>4</sup> electrons as valence for Li and O atoms, respectively. The cutoff energy of 700 eV and

appropriate Monkhorst–Pack k-meshes were chosen to ensure that the total energy calculations were well converged to less than 1 meV/atom. The relative stability of different Li–O compounds at room temperature was calculated within a thermodynamic model<sup>56</sup> defined as  $G(T) = G(T,N_{\rm Li},N_{\rm O}) - N_{\rm Li}\mu_{\rm Li}(T) - N_{\rm O}\mu_{\rm O}(T)$ , where  $G(T,N_{\rm Li},N_{\rm O})$  are the Gibbs free energies of Li–O compounds and  $\mu_{\rm Li}(T)$  and  $\mu_{\rm O}(T)$  are the chemical potentials of Li and O atoms, respectively.

### ASSOCIATED CONTENT

## Supporting Information

The main structural parameters, the calculated Mulliken atomic charge, phonon entropy, phonon dispersion curve, vibrational frequency, simulated XRD, and O K-edge spectra. This material is available free of charge via the Internet at http://pubs.acs.org.

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#### Notes

The authors declare no competing financial interest.

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