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# ***Slow Diffusion by Singlet State NMR Spectroscopy***

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## **ABSTRACT**

Small diffusion coefficients can be measured by using populations of singlet states that have a relaxation time constant  $T_s$  that can be much longer than the longitudinal relaxation time  $T_1$ . Spatial information can be encoded with pulsed field gradients in the manner of stimulated echo sequences. Singlet states can be excited via double-quantum coherences to enhance the efficiency of phase encoding and decoding.

# Diffusion Ordered NMR Spectroscopy

## - DOSY -

Nuclear magnetic resonance (NMR) has long been known as a flexible tool to study transport phenomena such as diffusion, flow, convection, or electrophoretic mobility [1, 2]. Information about the localization of molecules can be encoded and decoded by pulsed field gradients (PFG) before and after a delay where translational motion can occur. A pulsed field gradient can be characterized by a product  $\kappa = \gamma p s G_{\max} \delta$ , where  $\gamma$  is the gyromagnetic ratio,  $p$  the coherence order,  $G_{\max}$  the peak intensity of the gradient, and  $\delta$  its duration.  $s$  is a dimensionless shape factor ( $0 < s \leq 1$ ). When the signal  $S$  is observed as a function of the gradient strength and compared to a signal  $S_0$  obtained with very weak gradients, all other parameters remaining the same, the decay of the ratio obeys a Gaussian function:

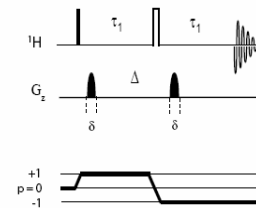
$$S/S_0 = \exp(-D\kappa^2\Delta')$$

where  $D$  is the diffusion coefficient and  $\Delta'$  the effective interval between encoding and decoding by the pulsed field gradients. The analysis of such Gaussian decays allows one to determine  $D$ . The resulting 2D representation with the isotropic chemical shift along the x axis and the diffusion coefficient along the y axis is referred as “diffusion ordered NMR spectroscopy” [2].

# Measurements of slow diffusion are limited by the relaxation time constants ( $T_1$ or $T_2$ )

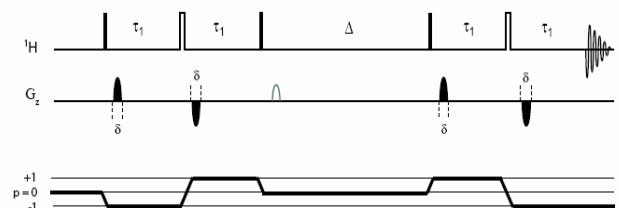
Depending on the details of the pulse sequences, one must take into account attenuation factors that depend on various delays where transverse and longitudinal magnetization components suffer from  $T_2$  or  $T_1$  relaxation.

In the most elementary spin-echo (SE) experiments,  $T_2$  relaxation limits the interval where diffusion can be observed to  $\Delta \approx T_2$ .



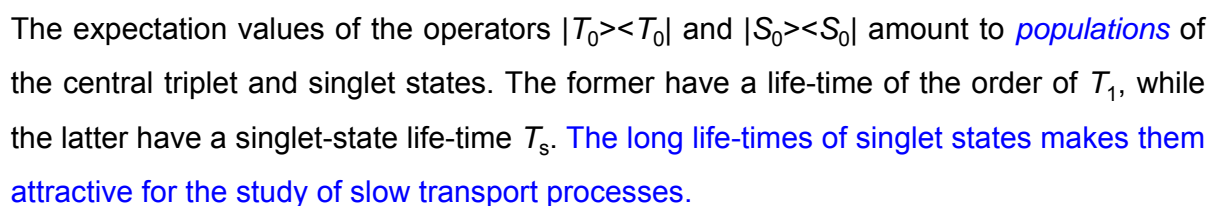
**Figure 1.** The simple spin echo (SE) sequence.

In stimulated echo experiments (STE), where the information is stored in the form of longitudinal magnetization, the window can be extended to  $\Delta \approx T_1$ , which allows one to probe slower transport processes since  $T_1 \geq T_2$ .



**Figure 2.** Standard stimulated echo sequence with bipolar pulse pairs (BPPSTE), often used for “diffusion-ordered spectroscopy” (DOSY).

Temperature (K)	$T_s$ (s)	$T_1$ (s)
220	8	0.5
230	9	1.0
245	14	0.5
255	20	1.5
265	25	2.0
275	33	3.0
285	37	4.0
295	43	5.0

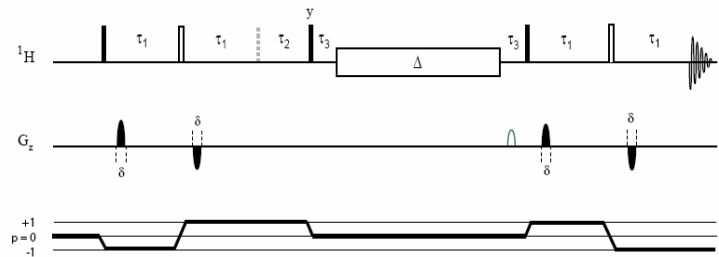


# Singlet-state-single-quantum-DOSY

## - SS-SQ-DOSY -

The key idea is shown in Figure 4, which results from a simple combination of Figures 2 and 3. A singlet state is populated at the beginning of the diffusion interval  $\Delta$  in Fig. 4. The spatial information is encoded and later decoded by a bipolar pair of pulsed field gradients. The signals are attenuated in proportion to  $\exp\{-D4k^2\Delta\} \exp\{-\Delta/T_s\}$ . Thus the main difference lies in the fact that  $T_1$  is replaced by the much longer singlet-state relaxation time  $T_s$ .

**Figure 4.** "Singlet-state-single-quantum-DOSY" - The information about spatial localization is stored in the form of singlet state populations with a relaxation time  $T_s$  in the interval  $\Delta$ . The intervals  $\tau_1$ ,  $\tau_2$  and  $\tau_3$  are defined as in sequence (a).

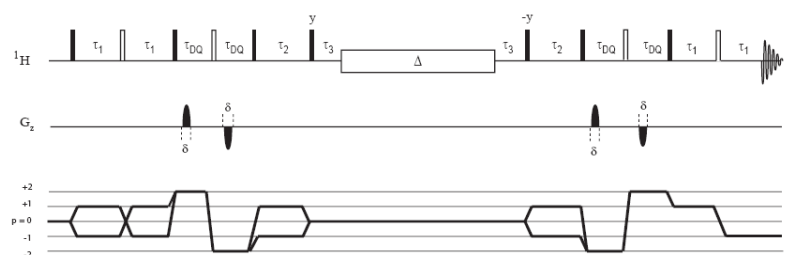


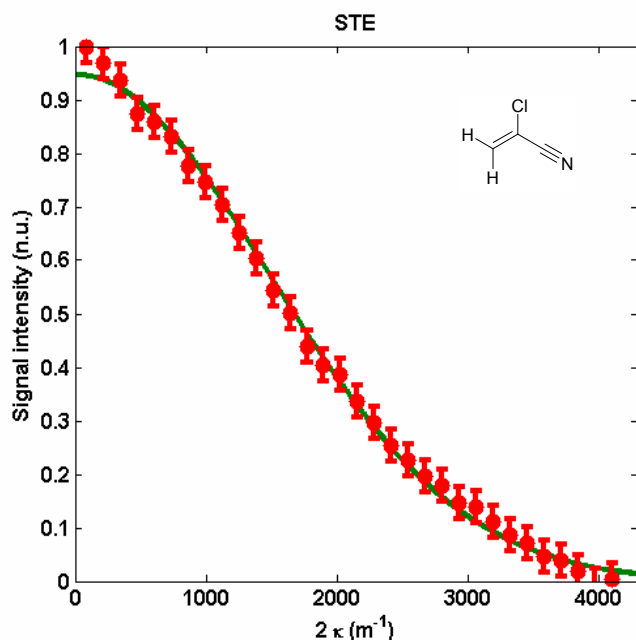
# Singlet-state-double-quantum-DOSY

## - SS-DQ-DOSY -

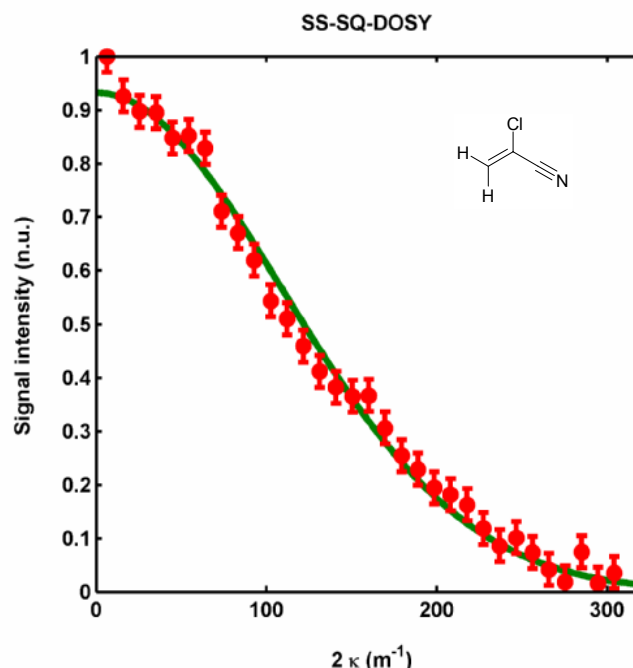
Single-quantum coherences generated at the end of the second  $\tau_1$  interval,  $\sigma = 2I_xS_z + 2I_zS_x$ , are temporarily converted into double-quantum coherences  $\sigma = -(2I_xS_y + 2I_yS_x) = i(I_+S_+ - I_-S_-) = -2DQC_y$ . The bipolar pulsed field gradients are applied to double-quantum coherences. Thus, compared to Figure 4, it is possible to determine diffusion constants that are four times smaller, or to obtain the same effect with field gradients that are half as strong or half as long.

**Figure 5.** "Singlet-state-double-quantum-DOSY" - The intervals  $\tau_1$ ,  $\tau_2$  and  $\tau_3$  must be like in sequence (a)





**Figure 6.** Gaussian decays of proton signal intensities of 10 mM 2-Chloroacrylonitrile in DMSO-D<sub>6</sub>/D<sub>2</sub>O = 1:3, measured at 300 MHz and  $T = -18.8\text{ }^{\circ}\text{C}$  (254 K) obtained with the **bipolar pulse pairs stimulated echo** of Fig. 2 with  $\Delta' = 0.21\text{ s}$ ,  $\delta = 2\text{ ms}$ . The parameter  $\kappa$  is defined by  $\kappa = \gamma ps G_{\max} \delta$ . Note the scales of the horizontal axes



**Figure 7.** Gaussian decays of proton signal intensities of 10 mM 2-Chloroacrylonitrile in DMSO-D<sub>6</sub>/D<sub>2</sub>O = 1:3, measured at 300 MHz and  $T = -18.8\text{ }^{\circ}\text{C}$  (254 K) obtained with the **singlet-state sequence** of Fig. 4 with  $\Delta' = 19.5\text{ s}$ ,  $\delta = 155\text{ }\mu\text{s}$ . The parameter  $\kappa$  is defined by  $\kappa = \gamma ps G_{\max} \delta$ . Note the scales of the horizontal axes

## CONCLUSIONS

It has been shown that slow transport phenomena can be characterized by combining methods for the measurement of diffusion coefficients with the excitation and detection of singlet states. The transient excitation of double-quantum coherences allows one to use weaker gradients. In 2-Chloroacrylonitrile, the life-times  $T_s$  of the singlet states were found to be about an order of magnitude longer than the longitudinal relaxation times  $T_1$  over a wide range of temperatures.

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