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¹ Ettringite Strengthening at High Pressures Induced by the ² Densification of the Hydrogen Bond Network

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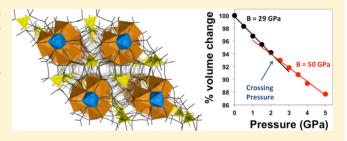
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ABSTRACT: Ettringite is a rare mineral with high-water content, more than half of its weight, and a relevant secondary product in Portland cement. Using density functional theory, we simulate the crystal structure and properties of ettringite under pressure. Our calculations predict a change in slope for all the lattice parameters versus pressure at about 2.5 GPa. Above such pressure, the elastic properties show a drastic increase of nearly 80% in the bulk modulus. This finding is explained in terms of a concurrent amorphization and densification of the hydrogen bond network. At low pressures,



ettringite can be compressed substantially without significant repulsion in the hydrogen bond network. At high pressures, the hydrogen bonds become stiff, and their contribution to the total repulsion is then important. These changes are also supported

by the evolution on the electronic density of ettringite with pressure.

6 INTRODUCTION

27 Ettringite is a mineral that forms near igneous contacts or in 28 xenoliths environments, associated with other species like 29 portlandite, gypsum or afwillite. Ettringite crystals are 30 hexagonal with a symmetry group P31c and a molecular 31 formula of $Ca_6Al_2(SO_4)_3(OH)_{12} \cdot nH_2O$ (n = 24–27). Its high 32 water content, more than one-half of its mass, confers to 33 ettringite a low density of approximately 1.745 g·cm⁻³. The unit 34 cell structure, shown in Figure 1, is based on calcium aluminate 35 columns aligned here parallel to the c-axis. Along the column, 36 the aluminum atoms are octahedrally coordinated to hydroxyl 37 groups and intercalated between planes with three calcium 38 atoms. The Ca atoms are coordinated to the hydroxyl groups of 39 aluminum and to four water molecules. The stoichiometry of 40 the columns leads to a charge excess, which is neutralized by 41 sulfate ions located in the channels between these columns. 42 Several experiments by X-ray diffraction reported that the space 43 between columns is filled with a variable number of water 44 molecules.²⁻⁷

Experimental works on ettringite have clearly pointed to the mportant role of water molecules in the cohesion between columns. Moore and Taylor, when first resolving the crystalline structure of ettringite, suggested that the water molecules transfer the charge from the calcium-aluminate to the negative

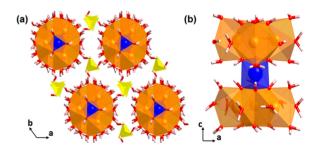


Figure 1. (a) Crystal structure of ettringite viewed along the c axis direction. Aluminate ions (AlO_3^{3-}) are shown as blue octahedra, sulfate ions (SO_4^{2-}) are shown as yellow tetrahedra, and calcium ions are inside orange polyhedra. Water and hydroxyl groups are denoted with red and white sticks for O and H atoms, respectively. (b) Detail of a calcium-aluminate column.

sulfate ions giving cohesion to the crystal.⁸ Ettringite can lose 50 about 27% of its structural water, including some of the 51 hydroxyl group when it is heated up to 50 °C.³ If the thermal 52 treatment continues and 30% of the water is lost, the crystalline 53

Received: February 23, 2012 Revised: June 28, 2012 54 order disappears,³ a process that suggests a strong role of water 55 in its structural integrity ettringite. The high water content of 56 ettringite also affects its elastic properties. Using Brioullin 57 spectroscopy,² the elastic moduli along the *c* direction was 58 found to be larger than in the *ab* plane by nearly a factor of 2. 59 Such anisotropy is attributed to the stiffer iono-covalent bonds 60 in the calcium-aluminate columns compared to the electrostatic 61 bonding mediated by water in the ab plane.² These 62 experimental studies revised so far indicate the basic role of 63 water, which must be carefully studied when looking at the 64 bonding and the elastic properties of ettringite.

The technological interest of ettringite arises from its 66 presence in the hydration of Portland cement. Ettringite 67 forms as a hydration product of calcium aluminate phases and 68 gypsum, and it establishes a complex equilibrium with other 69 crystalline phases in cement, such as monosulphoaluminate. 70 During the early hydration, ettringite crystals played an 71 important role in the cement setting and workability of 72 cements, and their presence must be carefully controlled. 9,10 At 73 later ages, ettringite might form in concrete under humid 74 environments owing to sulfate attack. As a result of the 75 crystallization pressure and water adsorption, this secondary 76 formation of ettringite often causes microcracks in cement 77 pastes, which reduce their performance. 11 Ettringite is also the 78 main binding phase of calcium sulphoaluminate cements, a 79 more environmentally friendly alternative to ordinary cement, 80 and a current focus in scientific and technological research. 12 81 Another potential application of ettringite is waste manage-82 ment. 13 Oxyanions such as B(OH)4-, SeO₄2-, CrO₄2-, and $83 \text{ MoO}_4^{\ 2-}$ can be stabilized in the channels between columns, 84 occupying the position of the sulfate groups. 14 Divalent and 85 trivalent metals can compete with aluminum for the octahedral 86 positions within the columns, and ettringite with Cu, Cr, Cd, 87 Fe, Pb, and Zn substitutions have been synthesized.¹

The practical manufacturing of cement-based materials and the efficient stabilization of heavy ions require that we are able to understand and control the properties of ettringite under a wide range of thermodynamic conditions, such as high pressures. Despite the fact that the effect of temperature on the stability of ettringite has been studied in depth, just one article reported the structural changes as the pressure increases. In that work, the authors determined that at a pressure of 3 GPa, ettringite lost the long-range crystalline order because its X-ray diffraction pattern showed a broad reflection instead of defined peaks. The recovery of crystalline order when pressure decreases suggests that water does not loo leach in the process. Measured by IR spectroscopy, the O–H and S–O vibrational bands suggest an amorphization of the water network, preserving the basic structure of ettringite.

In this work we aim to further understand the structure and properties of ettringite with or without pressure using density functional theory (DFT) simulations. We analyzed the atomic structure and electronic properties up to 5 GPa. Special attention was paid to the changes on the water network, which we will show is indeed the critical factor that drives the singular behavior of ettringite crystals under pressure. Due to the highwater content, ettringite serves as a model to study the impact of water in the structure and properties of inorganic crystals.

COMPUTATIONAL DETAILS

113 The simulations were carried out in the framework of the DFT, 114 as implemented in the SIESTA code. 16 The exchange-115 correlation functional was chosen from the generalized gradient

approximation (GGA)¹⁷ following the Perdew-Burke-Ernzer- 116 hof (PBE) scheme. 18 SIESTA is based on a combination of the 117 localized atomic orbitals basis set and pseudopotentials to 118 perform calculations of the electronic structure. The 119 pseudopotentials and double zeta polarized (DZP) basis sets 120 for Al-Ca-O-H were tested previously, which showed 121 an excellent agreement with the experimental crystalline 122 structures and elastic properties of calcium aluminates. We 123 considered a 3s² 3p⁴ atomic configuration for sulfur atoms using 124 a DZP basis set. The converged simulations were done using a 125 mesh cutoff of 350 Ry and several k-points given by the grid 126 cutoff of 10.78 Å.²² The unit cell is relaxed by the conjugate 127 gradient method simultaneously minimizing the atomic 128 positions and the lattice parameters. The convergence criterion 129 for the atomic forces is $0.02 \text{ eV}\cdot\text{Å}^{-1}$. The pressure was applied 130 isotropically to the unit crystalline cell. The elastic tensor 131 coefficients were finally calculated from the strain-stress 132 relationship. The strains were applied in steps of 0.5% from 133 -3% to +3% to calculate the corresponding stresses, so that 134 after a linear fit, the elastic coefficients are obtained. 135

■ RESULTS AND DISCUSSION

We first looked at the effect of different water amounts on the 137 structure of ettringite, since they partially occupy the empty 138 spaces in the intercolumn channels. He have analyzed the 139 crystalline structure of ettringite for the whole range of possible 140 water occupancies in the unit cell, from 24 to 27. Table 1 141 t1

Table 1. Lattice Parameters of Ettringite from Our Calculations and Experimental Values from the Literature

structure	а-b (Å)	c (Å)	volume (Å)	density (g/ cm³)	water
this work	11.217	21.575	2357.90	1.717	24
	11.241	21.604	2364.15	1.737	25
	11.289	21.617	2385.57	1.747	26
	11.338	21.6437	2411.86	1.753	27
ICDD 41-1451	11.224	21.408	2335.62	1.79	26
Speziale ²	11.24	21.468	2348.80	1.78	26
Hartman ³ *	11.116	21.355	2285.22	1.81	25
Berliner ⁷ *	1.117	21.360	2306.81	1.90	26
Taylor ⁴	11.260	21.48	2358.53	1.76	25.7
Neubauer ⁵ *	11.229	21.478	2345.46	1.77	25.7
Clark ⁶	11.240	21.468	2348.80	1.78	26

collects our computed lattice parameters. The increase from 24 142 to 27 water molecules in the ettringite unit cell has a small 143 effect on its lattice parameters. The lattice constants a and b 144 increase by 1%, and the change in the c dimension remains 145 small within 0.25%. This results in a maximum cell volume 146 increase of 2.3%. Although our lattice constants are slightly 147 larger than the experimental data, also given in Table 1, they 148 still agree well.

Figure 2 shows a simulated X-ray powder diffraction pattern 150 £2 of ettringite structure with 26 water molecules in the unit cell, 151 together with the experimental spectrum generated after the 152 recent data reported by Hartman and Berliner. Note the 153 excellent match in the position and intensity of the peaks, just 154 slightly shifted to larger d-spacing due to the increase of the 155 lattice dimensions. This match ensures that the atomic 156 arrangement and structural features are fully conserved after 157 relaxation. Also of interest is the presence of a small peak at 7.3 158 Å. It has been suggested that such peak might arise from 159

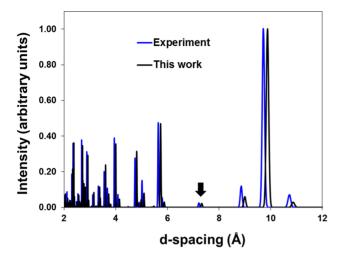


Figure 2. X-ray powder diffraction patterns of ettringite obtained from the computed structure using the Mercury code²⁴ (black) and from experiments²³ (blue).

160 impurities in the sample.⁶ However, its presence in our 161 simulated X-ray diffraction pattern shows that the peak is 162 intrinsic to the ettringite structure and independent of 163 impurities.

Our simulations accurately described the unit cell and atomic arrangement of ettringite, and we have determined that the variation in stoichiometry due to the partial occupation of water molecules expands or contracts the unit cell by ~2%. Because ettringite phases are important for cement materials and contribute to the overall elastic behavior, we next focused on their elastic properties. For ettringite with 24 water molecules in the stoichiometric formula, we computed the full elastic tensor from which the moduli were derived through standard definitions. Table 2 presents our elastic moduli, together with

Table 2. Theoretical Elastic Properties of Ettringite (in GPa)^a

property (GPa)	DFT calculations	experiment 2
K	29.96	27.30
G	12.57	9.45
E	33.08	25.95
η	0.32	0.34
$E_{100-010}$	27.95	20.1
E_{001}	55.62	41
anisotropy index	0.505	0.529

^aThe uniaxial Young moduli (E) are given in the [100], [010], and [001] directions corresponding to the a, b, and c lattice parameters, respectively. The experimental data¹² are also included.

174 the experimental data obtained by Brioullin spectroscopy. The 175 Hill average is presented here. The values obtained for the bulk 176 and shear moduli, as well as the Poisson ratio, agree very well 177 with experiments. The Young's moduli are slightly higher than 178 the experimental values. However, they change systematically in 179 all the directions, preserving the anisotropy of the elastic 180 moduli. Such anisotropy was quantified by the anisotropy 181 index, 26 which represents the difference in elasticity between 182 the crystallographic axes. In ettringite, the elasticity is twice 183 higher in the direction of the columns c than in the 184 perpendicular ones (the a and b plane). Such anisotropy is

intimately linked to the water content in ettringite, as we will 185 discuss below.

We turn now to investigate the behavior of ettringite under 187 pressure. Figure 3 shows the lattice parameters and volume 188 f3 against pressure from our calculations, together with the 189 experimental values after X-ray diffraction measurements. The 190 change of the lattice parameters with pressure matches very 191 well with the experiments, with the c axis having a systematic 192 small shift. While in our simulations we reached pressures up to 193 5 GPa, the experimental data stop at pressures about 1.3 GPa 194 because a quality fit of the lattice parameters was obtainable.⁶ 195 Interestingly, a change in slope is found between 2.0 and 3.0 196 GPa in all the lattice parameters and the volume. The nonlinear 197 evolution for the lattice vectors of ettringite with pressure is 198 coupled with different elastic properties. For instance, if we 199 computed the bulk modulus from the volume versus pressure 200 data in the first slope range for ettringite with 24 water 201 molecules per unit cell, we find a value of 29.7 GPa, in excellent 202 agreement with the 27.8 GPa obtained from the experiments, 203 and the 29.96 GPa from the computed elastic tensor. However, 204 the secondary slope gives an increase of nearly 80% in the bulk 205 modulus, up to 50 GPa. Similar increase takes place for all the 206 studied water contents. Therefore, our simulations reveal a 207 nonlinear response of ettringite crystals to isostatic pressure, 208 independent of the water content, which implies a strengthen- 209 ing of the crystal for pressures larger than 2 GPa.

Experimental studies found a change in the S-O and O-H 211 bond properties at around 2.5 GPa,6 which might be linked to 212 our discovery of a lattice parameter crossover. We thus analyzed 213 the bonds in more detail by calculating their length distribution 214 d_{O-X} (where X = Ca, Al, S, H) at different pressures. Although 215 all the bond lengths decrease slightly as the pressure increased 216 (up to 0.05 Å), the structure of Ca-O, Al-O, S-O, and H-O 217 bonds was kept under pressure. More interesting is the length 218 distributions of hydrogen bonds $d_{O\cdots H}$. The network of 219 hydrogen bonds for ettringite includes water-water, water- 220 hydroxyl, and water-sulfate bonds, as schematically shown in 221 Figure 4a. At 0.0 GPa, all the hydrogen bonds clearly lie in a 222 f4 range from 1.7 to 2.3 Å with a bimodal distribution, as seen in 223 Figure 4b. When increasing pressure, the first peak maximum in 224 $d_{\text{O...H}}$ distances decreased substantially, about 0.2 Å. The second 225 maximum peak, at 2.1 Å, decreases even more noticeably with 226 pressure. In fact, the distances of both peaks became equal at 227 5.0 GPa. The strength of hydrogen bonds is usually related to 228 their length, 27 as well as the decrease in the accompanying X- 229 H covalent bond strength. 28 Therefore, their interaction 230 becomes stronger under pressure, with a remarkable change 231 after 3 GPa when the water-hydroxyl distance $d_{\mathrm{O}\cdots\mathrm{H}}$ shrinks 232 even further. This mechanism also explains the increase of the 233 elastic properties. At low pressures, the hydrogen bonds absorb 234 most of the stress by decreasing mainly their bond distances. At 235 higher pressures, the hydrogen bonds interact more closely, 236 with the corresponding increase of the elastic properties of 237 ettringite.

A second parameter that characterizes the network of 239 hydrogen bonds is the OH-O angle. In ettringite, this angle 240 distribution changes with the pressure from a narrow to a wider 241 distribution, indicating that not only do the hydrogen bond 242 distances decrease, but also a reorientation of the water 243 molecules occurs. As shown in Figure 4c, the bond angles at 0.0 244 GPa have three discrete values: $\sim 4.5^{\circ}$, 9° , and 17° . At high 245 pressures the angles become wider and spread between 5° and 246 25° . This finding indicates a more disordered hydrogen bond 247

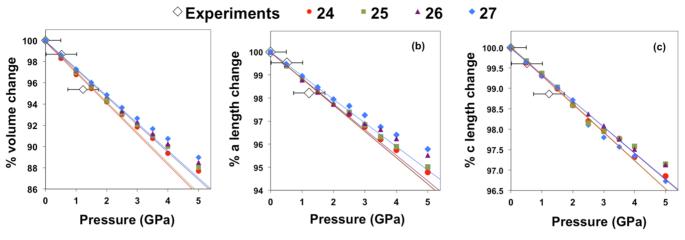


Figure 3. Pressure dependence of the (a) volume and the (b,c) lattice parameters of ettringite with different water occupancies in the unit cell (from 24 to 27). The empty rhombi denote the experimental data from ref 6. Note that the trends are divided into two regions according to the change of slope with pressure. The lines correspond to linear fits of our data at up to 2.0 GPa, and are extended for a better view of the slope change.

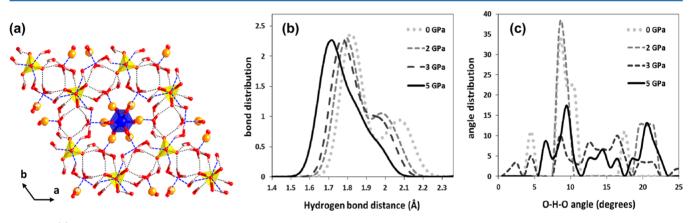


Figure 4. (a) Hydrogen bonded-network structure in ettringite at 0 GPa. Black-dotted and blue-dashed lines represent the hydrogen bonds shorter and larger than 2 Å, respectively. (b) Distance and (c) angle distributions for the hydrogen bonds in ettringite at several pressures.

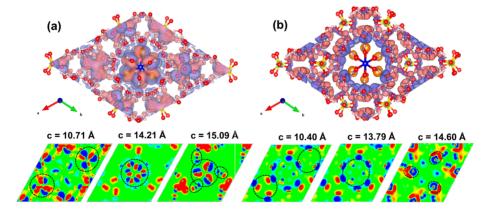


Figure 5. Difference in electronic densities between ettringite with and without water molecules at (a) zero and (b) high (5 GPa) pressures. The unit cell is shifted with respect to Figures 1 and 3 for a better view. Blue and red areas represent excess and defect in electronic density, respectively. Bottom pannels in (a) and (b) are cuts for several *c* values showing the charge polarization of (left) sulfate groups, (middle) hydroxyl groups, and (right) calcium ions and water molecules. For clarity, the mentioned groups are marked by dashed circles.

248 network, which might be related to the amorphization observed 249 at high pressures in X-ray experiments.⁶

We next analyzed the charge density of ettringite to 251 understand its chemical bonding, looking at the role of water 252 on the cohesion and elastic properties. We have computed the 253 difference in electronic densities between ettringite with and 254 without water. We represented this difference taking $\Delta \rho_{(r)}$ =

 $\rho_{\rm (r)}({\rm ett}) - \rho_{\rm (r)}({\rm dry}) - \rho_{\rm (r)}({\rm wat})$, where $\rho_{\rm (r)}$ is the electronic 255 density for the unit cell of ettringite, the "dry" means without 256 water, and the "wat" with only water. The difference $\Delta\rho_{\rm (r)}$ is 257 plotted in Figure 5a with the excess and defect electronic 258 f5 density represented with the blue and red areas, respectively. 259 The difference $\Delta\rho_{\rm (r)}$ changes through the entire unit cell, 260 except around the Al atoms that are without polarization. There

262 is a clear charge polarization in the sulfate groups, hydroxyl 263 groups, calcium ions, and water molecules, as seen in the lower 264 panels of Figure 5a. The cohesion between columns in 265 ettringite was previously attributed to a positive charge transfer 266 from the columns to the hydrogen in the water molecules, ²⁶⁷ connected to the negative $SO_4^{\ 2-}$ groups.⁸ Rather than charge 268 transfer, the cohesion involves the charge polarizations of 269 columns, water molecules, and sulfate ions.

To understand the changes with pressure, we have also 271 analyzed the difference in electronic density with pressure at 5 272 GPa, plotted in Figure 5b. It is significant that the maximum value of $\Delta \rho_{\rm (r)}$ decreases with respect to the one at 0 GPa pressure. Only small charge excess is present in sulfate oxygen atoms accompanied with the charge polarization of surrounding water molecules. Compared with panel a, the charge is no 277 longer polarized in the calcium and hydroxyl ions. These changes in $\Delta \rho_{\rm (r)}$ are consistent with a drastic decrease of polarization under pressure. At high pressures, the atoms pack 280 in such a way that their polarizations are negligible. Not only 281 the analysis of the $d_{\text{O} ext{--} \text{H}}$ distances, but also the electronic 282 density distribution confirms the densification of the hydrogen 283 bond network at high pressures.

The crossover of stiffness with pressure is explained by an 285 atomic packing effect. At a zero pressure, the atoms are at their 286 equilibrium distances, and the charge polarization is the 287 dominant interaction. As the pressure increases, the bond 288 distances decrease, entering into the interatomic repulsive 289 regime, a common feature in most crystals. With respect to 290 more homogeneous systems, ettringite includes a diversity of chemical bond types. At low pressures, the hydrogen bonds 292 became shorter without a significant contribution to the 293 repulsive forces, because their potential energy surface is rather 294 flat. 29,30 The crystal then compresses without important 295 repulsive forces in the stiffer iono-covalent bonds. At high 296 pressures, the hydrogen network becomes so packed that their 297 repulsion is comparable to that of the iono-covalent bonds. The whole crystal is now a repulsive system, hence the nonlinear evolution of the volume versus pressure. Because this crossover 300 is related to heterogeneity in the chemical bonds, we expect 301 that this feature is general to crystals with diverse bonding 302 schemes such as layered materials or molecular crystals. 31-33

There are noteworthy comments on zero point effects. It is 304 well-known that the isotopic substitutions modify the internal 305 vibrations of solids and change the zero-point energy. Ice has 306 shown an anomalous isotope effect with a higher volume for deuterated hexagonal ice with respect to the light one.^{34,35} The same effect on zero point vibrations can be reached by 309 increasing either the mass of the isotopes or the pressure. This 310 effect can indeed play a role in the volume changes we found 311 for ettringite at high pressures, and therefore in the increase of 312 the bulk modulus. However, we still lack experimental 313 information about elastic properties or the pressure behavior 314 of deuterated ettringite. Experimental determination of the 315 isotope effect together with a theoretical vibrational analysis will 316 definitely be an interesting topic, although it is beyond the 317 scope of this work.

318 CONCLUSIONS

319 Ettringite is a natural mineral with more than half water 320 content. It forms a hydration product in Portland cements and 321 plays a key role in their physical and chemical stability. In this 322 work we have studied its structure, elastic properties, and electronic properties under pressure using simulations within 323

We calculated the structure and elastic properties of ettringite 325 with pressure. At a zero pressure, our values agree well with 326 experiments. The elastic moduli are highly anisotropic along 327 different crystallographic directions, such anisotropy is directly 328 related to the nature of the chemical bond in those directions: 329 iono-covalent along the c axis, and hydrogen bonding in the a 330 and b plane. Our calculations indeed show that the charge 331 polarization changes in the ions and molecules between 332 columns, changes that are held responsible for the cohesion 333 in ettringite. At high pressures, above 2.5 GPa, the lattice 334 constants and the crystal elastic properties enter into a different 335 region, where the bulk modulus increases by about 80%. Such 336 drastic change in the elastic properties is related to the 337 densification of the hydrogen bond network. The weak 338 hydrogen bonds compress without inducing big repulsive 339 interatomic forces up to about 2.5 GPa (the crossover 340 pressure), when they become closer, and the oxygen-hydrogen 341 repulsive energy is comparable with the rest of the iono- 342 covalent bonds.

Because the modification of elastic properties with pressure is 344 driven by the heterogeneity of interatomic forces, our results 345 might be generalized to crystals with dissimilar types of 346 chemical bonds, such as molecular crystals, layered materials, 347 etc. These results also point to the importance of hydrogen 348 bonding in the elasticity of crystalline materials, a factor that, to 349 date, has been commonly explored in soft matter such as 350 proteins or polymers.36,37

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