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Ru(0001) Model Catalyst under Oxidizing and Reducing Reaction Conditions: In-Situ High-Pressure Surface X-ray Diffraction Study

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With surface X-ray diffraction (SXRD) using a high-pressure reaction chamber we investigated in-situ the oxidation of the Ru(0001) model catalyst under various reaction conditions, starting from a strongly oxidizing environment to reaction conditions typical for CO oxidation. With a mixture of O₂ and CO (stoichiometry, 2:1) the partial pressure of oxygen has to be increased to 20 mbar to form the catalytically active RuO₂(110) oxide film, while in pure oxygen environment a pressure of 10⁻⁵ mbar is already sufficient to oxidize the Ru(0001) surface. For preparation temperatures in the range of 550–630 K a self-limiting RuO₂(110) film is produced with a thickness of 1.6 nm. The RuO₂(110) film grows self-acceleratedly after an induction period. The RuO₂ films on Ru(0001) can readily be reduced by H₂ and CO exposures at 415 K, without an induction period.

1. Introduction

In recent years, significant research has focused on the initial oxidation of metal surfaces by experimental as well as by theoretical techniques.^{1–9} This interest is not only spurred by its importance in practical applications such as heterogeneous catalysis, corrosion, protective oxide layers, or insulating layers in microelectronic devices but also because of a lack of information on the fundamental processes that determine the initial oxidation of metal surfaces on the atomic scale. In general, the chemisorption of oxygen precedes the actual oxidation of a metal surface.¹⁰ Only after a critical oxygen coverage has been reached, oxygen penetration including the development of a surface oxide may set in.^{6,11} A carefully studied exception from this generally accepted rule is the oxidation of Zr(0001).¹² On a Zr(0001) surface, oxygen dissociates and penetrates right away in the subsurface region.¹² Beyond a critical subsurface O coverage of 1 monolayer, additional oxygen is stabilized on the surface, forming an O–Zr–O trilayer. Similar processes are expected to be relevant for Ta, Ti, and V. In the meanwhile this O–Me–O trilayer motif (Me stands for a metal) has also been identified to play a key role in the initial growth of other metal oxides. Rhodium oxide, for example, forms O–Rh–O trilayers during the initial oxidation of Rh(111).⁷ For the oxidation of Ru(0001) it has been argued on the basis of density functional theory (DFT) calculations¹³ and supported by recent X-ray photoelectron spectroscopy (XPS) measurements^{14,15} that an O–Ru–O trilayer exists as a metastable structure, being a precursor to surface RuO₂.

The initial oxidation of Ru(0001) was proposed to proceed in an autocatalytic process.¹⁶ The saturation coverage of chemisorbed oxygen on Ru(0001) is reached with a Ru(0001)–(1 × 1)O surface at 1 monolayer,¹⁷ while the formation of an oxide has been observed only for coverages higher than 3

monolayers.¹⁸ The oxide formation requires large O₂ exposures since the sticking coefficient of oxygen on the Ru(0001)–(1 × 1)O surface is of the order of 10⁻⁶.¹⁹ However, as soon as RuO₂ nuclei are formed, the further oxidation of Ru(0001) takes place from these nuclei in a self-accelerated process^{16,20} since the dissociative sticking coefficient of O₂ on RuO₂ is as high as 0.8.¹⁹ In other words, the RuO₂ nuclei catalyze the oxidation of Ru(0001).

Here we report on “high-pressure” surface X-ray diffraction (SXRD) experiments which allow us to follow in-situ and on atomic scale the oxidation of the model catalyst Ru(0001) under a variety of reaction conditions, such as pure oxygen atmosphere and a mixture of CO and O₂ up to 30 mbar at various temperatures. Under these reaction conditions we observed after an induction period a self-limiting growth of RuO₂ films which consist of RuO₂ domains with a mean size of about 10 nm and a mean thickness of about 5 monolayers. We studied also in-situ the reduction of such prepared RuO₂ films on Ru(0001) by H₂ and CO exposures.

2. Experimental Details

The SXRD experiments were performed at beam line ID03 at the ESRF, using a specially designed ultrahigh-vacuum/high-pressure chamber mounted on a high-precision diffractometer.²¹ A photon energy of 17 keV was selected with a Si(111) double monochromator corresponding to a wavelength of 0.73 Å. This experimental setup is truly unique in the in-situ study of structural changes on model catalyst's surfaces under realistic reaction conditions such as those encountered with the initial oxidation of metal surfaces. The cleanliness of the surface is maintained as good as under ultrahigh-vacuum conditions. So far very few in-situ experiments on the initial oxidation of metal surfaces have been reported in the literature^{9,22–24} which aim at unraveling microscopic reaction steps under realistic reaction conditions.

The Ru(0001) sample was prepared by Ar-ion etching ($U = 2$ kV, $I = 15$ mA, $p(\text{Ar}) = 10^{-5}$ mbar, for 20 min at 760 K) and oxygen treatment cycles at elevated temperatures. The initial

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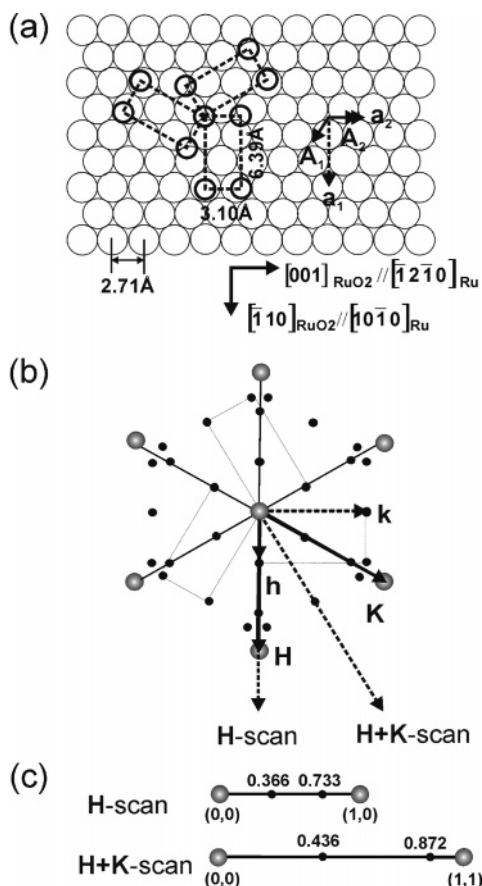


Figure 1. (a) Top view of RuO₂(110) grown on Ru(0001). The in-plane lattice vectors **A**₁ and **A**₂ and the out-of-plane lattice vector **A**₃ define the crystal lattices of Ru(0001): **A**₁ = **A**₂ = 2.71 Å; angle = 120°; **A**₃ = 4.28 Å. The RuO₂(110) surface is oriented along the **A**₃ axis of Ru(0001). The in-plane lattice vectors of RuO₂(110) are depicted in a, forming a rectangular surface unit cell, defined by the lattice vectors **a**₁ and **a**₂. The in-plane epitaxial growth relationship between the RuO₂(110) film and the Ru(0001) substrate is indicated, i.e., RuO₂[001]//Ru[1210] and RuO₂[110]//Ru[1010]. In reciprocal space (b), **H** and **K** are the in-plane lattice vectors and **L** is parallel to **A**₃. The reciprocal lattice vectors of RuO₂(110) are **h**, **k**, **l**, where **h** is parallel to **H** and **k** is aligned along an equivalent direction of **H** + **K**, while **l** and **L** are parallel. (c) Illustration of the **H** and **H** + **K** scans, indicating the positions of expected diffraction maxima of RuO₂(110) with respect to Ru(0001).

oxidation of the Ru(0001) was monitored in situ by SXRD while exposing the sample to various O₂ and CO partial pressures (ranging from 10⁻⁵ mbar to some 10 mbar) at sample temperature ranging from 550 to 750 K. Reduction of RuO₂ films on Ru(0001) was carried out by H₂ or CO exposure at 415 K.

3. Results and Discussion

In Figure 1 the lattice vectors of the Ru(0001) surface are shown which define the crystal lattice of the hexagonal close packed (hcp) crystal. In reciprocal space, **H** and **K** are the in-plane lattice vectors corresponding to **A**₁ and **A**₂ in Figure 1a, and **L** is the out-of-plane vector. The reciprocal lattice vectors of the tetragonal RuO₂(110) on Ru(0001) are given in fractions of **H**, **K**, and **L**, and are annotated as **h**, **k**, **l**, where **h** is parallel to **H** and **k** is equivalent to the **H** + **K** (by mirror plane symmetry) direction (cf. Figure 1), while **l** is aligned along the **L** direction.

Before presenting the experimental data, we summarize briefly those structural properties of RuO₂,²⁵ which are important for the discussion of our SXRD experiments. RuO₂ crystallizes

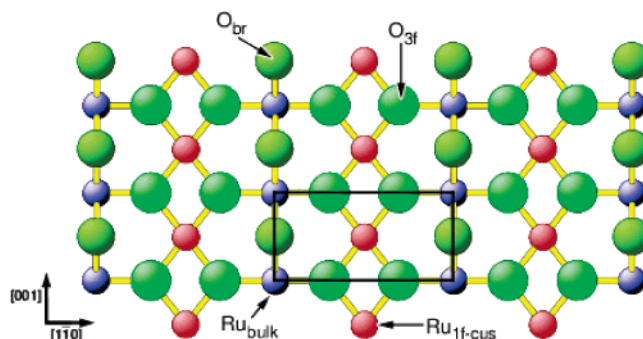


Figure 2. Stick and ball model of the stoichiometric RuO₂(110) surface. The small balls (red and blue) are Ru atoms; the large balls represent oxygen atoms. The catalytic activity of the RuO₂(110) surface is determined by the under-coordinated Ru (Ru_{1f-cus}, red) and bridging oxygen atoms (O_{br}). Bulk Ru atoms are small blue balls. The Ru atoms form in-plane a centered unit cell indicated by the rectangle.

in the rutile structure. On the hexagonal Ru(0001) surface an unstrained RuO₂ film (tetragonal structure) grows epitaxially and incommensurately. The flat oxide surface is oriented along the (110) direction, while in-plane three rotational domains (rotated by 120° with each other) coexist for symmetry reasons (*P2mm* of RuO₂ versus *P3m1* of Ru(0001); cf. Figure 1a,b). The Ru sublattice of RuO₂(110) forms a centered unit cell (see Figure 2). Because Ru is a much stronger X-ray scatterer than O, we expect for a RuO₂(110) film strong intensities only for even diffraction order of **h** + **k**.

The first experiment was designed to study in-situ the oxidation of Ru(0001) in a pure molecular oxygen atmosphere (the background pressure in the high-pressure chamber was 2 × 10⁻⁹ mbar). With SXRD we measured **H** and **H** + **K** scans (defined in Figure 1c) at *L* = 0.2 as well as the **L** scan at (*H*, *K*) = (0.733, 0), while the Ru(0001) surface was exposed to 10⁻⁵ mbar of molecular oxygen, keeping the sample temperature at 630 and 580 K, respectively (see Figure 3). After an induction period of about 20 min a distinct diffraction peak occurs at 0.733 in the **H** scan (cf. Figure 3a), which corresponds to the second-order diffraction; the peak at 0.3665 is too small to be detectable. After the induction period the X-ray diffraction intensities at *H* = 0.733 increase steeply with exposure time, finally saturating after 60 min. The found induction period together with the steep raise in intensity (see inset of Figure 3a) is reconciled with an autocatalytic oxidation process of Ru(0001). In the **H** + **K** scan additional peaks appear at 0.436 and 0.872 (cf. Figure 3b). The intensity of the 0.436 peak is much smaller than that of the 0.872 peak consistent with the centered surface unit cell of the Ru sublattice of RuO₂(110).

From the in-plane lattice parameters *H* = 0.733/2 and *H* = *K* = 0.436 in reciprocal space we derive the (real space) in-plane lattice parameters of the RuO₂ film to be 3.10 Å × 6.39 Å. These values correspond nicely to the bulk-truncated unit cell of RuO₂(110), namely, 3.11 Å × 6.38 Å as expected for an unstrained RuO₂(110) film on the Ru(0001) surface.

The evolution of the **H** scan does not indicate any variation in the full-width at half-maximum (fwhm). The fwhm of the peak at 0.733 in the **H** scan keeping *L* = 0.2 and 1.3 (not shown) is in both cases 0.018. Therefore the averaged lateral dimension of the RuO₂ domains along the [110] direction is approximately 150 Å. The fwhm of the peak at 0.872 in the **H** + **K** scan at *L* = 1.3 is 0.016 which corresponds to a domain size of approximately 80 Å along the [001] direction. Both estimated dimensions are in line with recent scanning tunneling microscopy (STM) images of RuO₂(110) films on Ru(0001).^{26,27} In a previous STM study²⁶ we observed that RuO₂(110) domains

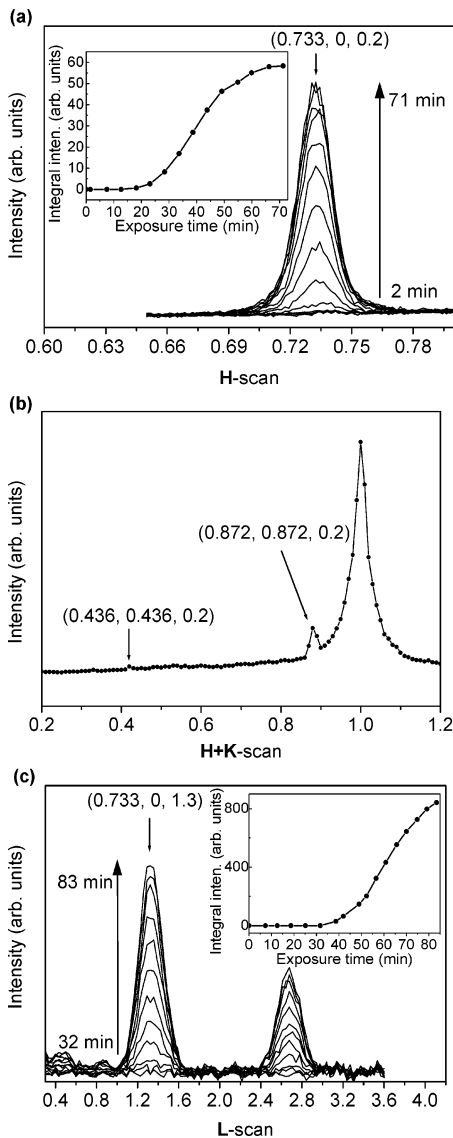


Figure 3. Monitor the oxidation of the Ru(0001) sample by in-situ SXR D. (a) Evolution of the **H** scan with exposure time when exposing the Ru(0001) sample to $p(\text{O}_2) = 10^{-5}$ mbar at a sample temperature of 630 K. In the inset the integral intensity is shown as a function of oxygen exposure time. The peak at 0.733 is the second-order reflection of $\text{RuO}_2(110)$ along the **H** direction. (b) Corresponding **H** + **K** scan of the final RuO_2 film. (c) Evolution of **L** scans with exposure time when exposing the Ru(0001) sample to $p(\text{O}_2) = 10^{-5}$ mbar at a sample temperature of 580 K. In the inset the integral intensity is displayed as a function of oxygen exposure time.

coexist with $(1 \times 1)\text{O}$, covering comparable fractions of the surface. Since $\text{RuO}_2(110)$ grows incommensurately on Ru(0001), the intensity of RuO_2 -related diffraction spots are only affected by the oxide layer but not by the $(1 \times 1)\text{O}$ phase. In contrast the crystal truncation rods depend strongly on both the $(1 \times 1)\text{O}$ patches and the $\text{RuO}_2(110)/\text{Ru}(0001)$ interface therefore being difficult to analyze. We thus restrict our current discussion only to RuO_2 -related diffraction spots.

We also measured the evolution of the **L** scans for $(H, K) = (0.733, 0)$ during the oxidation of Ru(0001) (cf. Figure 3c). **L** scans probe the out-of-plane periodicity of the $\text{RuO}_2(110)$ film. The $(H, K) = (0.733, 0)$ diffraction beam is predominantly sensitive to the Ru sublattice of the $\text{RuO}_2(110)$ film. After an induction period of about 30 min (cf. Figure 3c) two distinct maxima evolve in the **L** scan, namely, at $L = 1.3$ and at $L = 2.6$. The different induction periods of the **H** (cf. Figure 3a)

TABLE 1: Oxidation of Ru(0001) in Pure Oxygen Environment for Various Preparation Temperatures T

T/K	$p(\text{O}_2)/\text{mbar}$	fwhm(L)	no. of RuO_2 layers
580	$1.3 \times 10^{-5}, 1.1 \times 10^{-3}, 1.2 \times 10^{-2}$	0.29	5.0
600	$1.0 \times 10^{-4}, 2 \times 10^{-3}, 2 \times 10^{-2}, 1.0 \times 10^{-1}$	0.29	5.0
630	5.0×10^{-5}	0.26	5.6
670	$5.0 \times 10^{-5}, 1 \times 10^{-1}$	0.19	7.7
720	5.0×10^{-5}	0.13	11.2

The fwhm(**L**) are taken for the maximum at $L = 1.3$ in the **L** scan for $H = 0.733$ and $K = 0.0$. From these values we estimate the number of RuO_2 layers in the oxide by $1.33/(\text{fwhm}(\text{L}) \times 0.91)$. The correction factor 0.91 was taken from the literature.³⁶ The separation between the two maxima in the **L** scans is 1.33.

and **L** scans (cf. Figure 3c) are related to the different sample temperatures (630 K versus 580 K). An increase of the background intensity in the **L** scans upon oxidation is not observed so that the presence of a single well-ordered O—Ru—O trilayer or a single crystalline RuO_2 wetting layer cannot be identified with SXR D. From the fwhm of the peaks in the **L** scan we estimate the thickness of the $\text{RuO}_2(110)$ to be 1.6 nm; i.e., the RuO_2 film consists of five RuO_2 layers. The fwhm of the maxima in the **L** scans does not change during the oxidation of Ru(0001). The vertical growth of the $\text{RuO}_2(110)$ oxide film is self-limited. From the separation of the maxima in the **L** scan we infer that the Ru subplanes in the $\text{RuO}_2(110)$ domain are separated by 3.2 Å, in perfect agreement with the bulk layer spacing of $\text{RuO}_2(110)$ structure (3.23 Å^{25}).

Similar SXR D results were obtained for the oxidation at various preparation temperatures in the range of 580–650 K. The thicknesses and domain sizes of these RuO_2 films are compiled in Table 1. Obviously, the thickness of the RuO_2 domains is practically constant (5–6 RuO_2 layers) in the temperature range of 580–630 K. We varied also the partial pressure of oxygen in the range of 10^{-5} – 10^{-1} mbar. The thickness of the RuO_2 film was not affected by the variation of O_2 pressure. For sample temperatures of 670 and 720 K the thickness of the RuO_2 film increases significantly. However, also a faceting of the RuO_2 film took place. The lateral RuO_2 domain size does not change when the preparation temperature is varied from 580 to 720 K.

If the preparation temperature was chosen to be lower than 540 K (even for an O_2 pressure of 10^{-1} mbar), the Ru(0001) surface did not form a crystalline oxide film. This result is consistent with recent scanning photoelectron microscopy (SPEM) investigations.¹⁵ Here it is important to recall that SXR D is only sensitive to those parts of the surface which are well-ordered. This means that SXR D is “blind” to amorphous or weakly ordered structures such as Ru_xO_y which was recently proposed by a XPS study.¹⁴

The most intriguing feature of these SXR D measurements is that the fwhm of RuO_2 -related diffraction peaks does not vary during the oxidation process, neither in-plane (**H**, **H** + **K** scan) nor out-of-plane (**L** scan). We have to recall that the transfer width of the incident X-ray beam is several micrometers. Rocking scans around the surface normal from the Ru(0001) crystal surface at the minimum of a crystal truncation rod indicates an averaged terrace width of Ru(0001) of 500 nm. Therefore, the observed fwhm values are solely determined by the RuO_2 domain size which is around 10 nm across and 1.6 nm in thickness. The observed increase of the X-ray diffraction intensities of RuO_2 -related diffraction spots in the **H**, **H** + **K** and **L** scans upon O_2 exposure is, therefore, associated with a growing number of such RuO_2 domains whose diffraction

intensity contributes incoherently to the measured diffraction intensity. To ensure that the diffraction intensity of each RuO₂ domain sums up incoherently, the RuO₂ domain must not be in registry with respect to the in-plane RuO₂(110) lattice; otherwise the fwhm of the maxima in the **H** and **H** + **K** scans would decrease with an increase of the O₂ exposure time. Since the RuO₂(110) domains grow incommensurately on the substrate Ru(0001) surface, they are incoherent to each other, thus contributing additively to diffraction intensity. We note that the growth of a single RuO₂ domain cannot be time-resolved by the present SXRD experiment with its time resolution of about 100 s.

The time evolution of RuO₂-related diffraction intensities during oxygen exposure discloses an induction period (cf. Figure 3) which can be explained by the time required to form the first RuO₂ domains on the Ru(0001) surface. The further lateral growth of the RuO₂(110) film is then promoted by these already existing RuO₂(110) domains in that new RuO₂ domains preferentially form in their vicinity. Consequently, the oxide growth becomes self-accelerated. Accordingly, the RuO₂(110) film consists of RuO₂(110) domains which bunch together consistent with recent STM images, indicating a flakelike pattern of the RuO₂(110) film.²⁶

For the reduction experiments we produced first a 1.6 nm thick RuO₂(110) film on Ru(0001), consisting of RuO₂ domains with dimensions of about 10 nm across. This oxide film was subsequently reduced by hydrogen exposure of $p(\text{H}_2) = 10^{-5}$ mbar, keeping the sample temperature at 417 K. The progressing reduction of RuO₂ by H₂ was monitored in-situ by both the evolution of the **L** scan and that of the **H** scan (cf. Figure 4). The integral X-ray intensities of the $H = K = 0.872$ and $L = 1.3$ peak are shown in the insets of Figure 4a,b as a function of the H₂ exposure time. No induction period occurred within the experimental time resolution of 100 s. The integral intensity of this peak decreases linearly with H₂ exposure time. The fwhm of the peaks in the **L** and **H** + **K** scans does not change with H₂ exposure time, indicating that neither the averaged size nor the thickness of the RuO₂ domains varies during the reduction process by H₂ exposure. The reduction of the RuO₂(110) film thus appears to proceed via the annihilation of complete RuO₂ domains. Since no induction period is observed, the H₂ reduction of RuO₂(110) film differs from the reduction of Rh₂O₃ for which an autocatalytic process was identified.²⁸ For lower temperatures we have not observed any reduction of the RuO₂ film by hydrogen exposure. The threshold temperature, where the reduction by hydrogen exposure sets in, is around 400 K, consistent with recent H₂ oxidation experiments.²⁹

Similar reduction experiments were performed by dosing CO. The thickness of the starting RuO₂ film was 1.6 nm; the lateral domain size was 10 nm. Upon CO exposure ($p(\text{CO}) = 10^{-5}$ mbar) at a sample temperature of 415 K, the integral intensity of the $H = 0.733$ peak at $K = 0$ and $L = 0.5$ decreased linearly with time (cf. Figure 4c). Again neither the fwhm of the **H** scans nor that of the **L** scans changed, indicating a reduction process similar to that for hydrogen.

In a subsequent experiment we oxidized the Ru(0001) surface by a mixture of CO and O₂ (CO:O₂ stoichiometry, 1:2). The scientific background to this experiment is that we had not been able to produce an oxide under ultrahigh-vacuum condition as soon as we used a gas mixture of CO and O₂.³⁰ However, supported and unsupported Ru catalysts oxidize readily under oxidizing reaction conditions, where CO and O₂ are both present in the gas feed.^{31,32} Another motivation for this experiment is to understand the CO oxidation experiments by Peden and

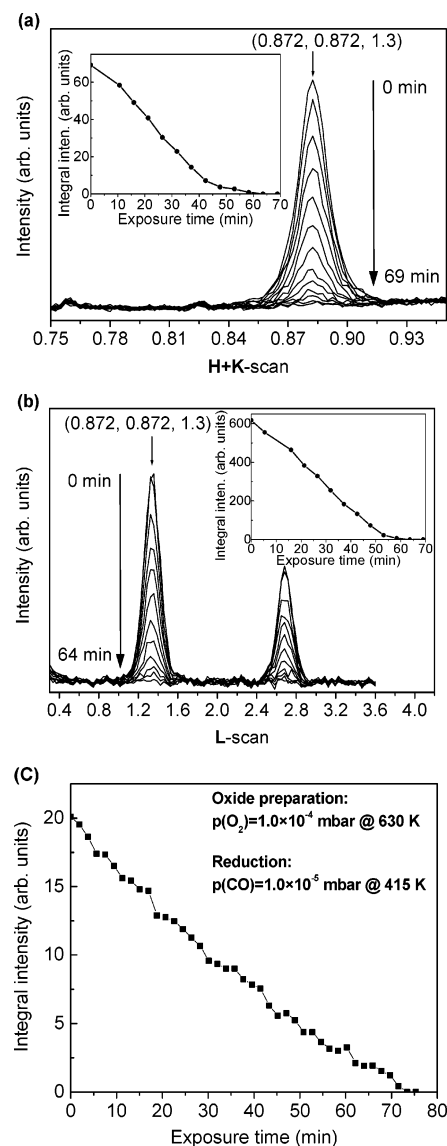


Figure 4. Reduction of a 5 monolayer thick RuO₂(110) film grown on Ru(0001) by exposure to 10^{-5} mbar of H₂ at a sample temperature of 417 K. (a) Evolution of **H** + **K** scans depending on the exposure time. In the inset the integral intensities are shown as a function of H₂ exposure time. (b) Evolution of **L** scan depending on the exposure time. In the inset the integral intensities are indicated as a function of H₂ exposure time. (c) Reduction of a 5 monolayer thick RuO₂(110) film on Ru(0001) by exposure to 10^{-5} mbar of CO at a sample temperature of 415 K. The integral intensities of the peak at 0.733 along the **H** direction are depicted as a function of exposure time.

Goodman.³³ Peden and Goodman showed that under ultrahigh-vacuum conditions Ru(0001) is inactive in the oxidation of CO. However, under higher pressures in the millibar range, Ru(0001) turned into a very active oxidation catalyst, whose active phase has not been identified unequivocally.

With a mixture of CO and O₂ of up to 10 mbar in total pressure no oxidation of the Ru(0001) was observed with in-situ SXRD and within 2000 s, keeping the sample temperature at 630 K. However, if we increased the total CO + O₂ pressure to 28 mbar, then after an induction period of 700 s the oxidation of Ru(0001) set in spontaneously and was completed within one single **H** scan (i.e. within 100 s; cf. Figure 5a,b). In comparison to the oxidation at 10^{-5} mbar of pure oxygen, the complete surface oxidation of Ru(0001) by 28 mbar of CO + O₂ proceeds much faster, indicating that kinetic hindrance was the main reason for the retarded oxidation of Ru(0001) by a

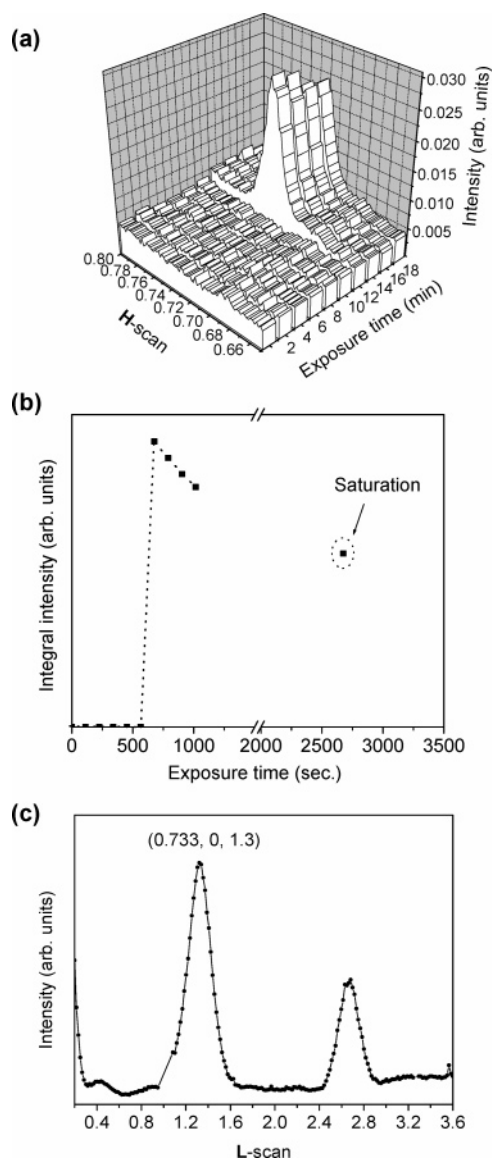


Figure 5. Oxidation of the Ru(0001) sample by exposing it to a gas mixture of O_2 and CO: $p(O_2) = 18$ mbar and $p(CO) = 10$ mbar at a sample temperature 630 K. (a) Evolution of the **H** scan with exposure time. (b) Integral intensity as a function of exposure time. (c) **L** scan for the final $RuO_2(110)$ film. From the separation between the two peaks at $L = 1.3$ and 2.6 a layer distance of 3.2 \AA is inferred, while the fwhm of these peaks indicates a RuO_2 film thickness of 1.6 nm .

1:2 CO + O_2 gas mixture. From the fwhm of the **H** scan and **H** + **K** scan we estimate the averaged RuO_2 domain size again to be 10 nm across. From the fwhm of the **L** scan the averaged thickness of the $RuO_2(110)$ domains is derived to be 1.6 nm . The $RuO_2(110)$ film is very flat as indicated by the small modulations between the Bragg reflection in the **L** scans (cf. Figure 5c).

The long induction period together with the high CO + O_2 pressure of 28 mbar required to oxidize the Ru(0001) sample suggests that the oxidation of the first RuO_2 domain imposes the rate-determining step. Several scenarios are conceivable. Either we need a large critical size of a RuO_2 nucleus to initiate the oxidation of the first RuO_2 domain. The formation of such a RuO_2 nucleus is subject to reduction by CO so that the critical size of the nucleus is only stable when the pressure of oxygen is high enough. Or an alternative scenario, which we do not favor, is that the nucleation of RuO_2 domains needs a critical mass of incorporated oxygen. Once it is reached the growth of

$RuO_2(110)$ proceeds rather fast for thermodynamical reasons.³⁴ In any case the oxide growth is self-accelerated and spreads across the Ru(0001) surface. From Figure 5b it is evident that when the complete RuO_2 film has formed, further CO + O_2 exposure leads even to a partial decrease of the integral intensity of the $H = 0.733$ peak. We may speculate that this decrease is related to a partial roughening of the RuO_2 film. Finally, for exposure times longer than 3000 s the RuO_2 film is stable; i.e., the model catalyst has reached steady-state conditions.

These CO + O_2 experiments clearly show that $RuO_2(110)$ is the active phase in the high-pressure CO oxidation experiments over Ru(0001) of Peden and Goodman.³³ The CO + O_2 experiment confirms also nicely the *core shell model* of Assmann et al., which was proposed for supported³¹ and powder^{32,35} Ru catalysts. For these practical Ru-based catalysts an ultrathin (self-limiting) RuO_2 film of about $1\text{--}2 \text{ nm}$ was postulated to cover the Ru metal core of the Ru particles in order to explain the in-situ reaction experiments and the in-situ Fourier Transform Infrared (FTIR) experiments in combination with the hydrogen reduction experiments.^{31,32,35} Here with the Ru(0001) surface, serving as the corresponding model catalyst, we identify also a kind of core shell model in which the Ru(0001) surface is covered by a self-limiting grown $RuO_2(110)$ film whose thickness is 1.6 nm .

4. Conclusion

With surface X-ray diffraction experiments using a dedicated high-pressure reaction chamber we explored the oxidation of the Ru(0001) model catalyst under reaction conditions which are far beyond the ultrahigh-vacuum regime. In particular, we investigated in-situ the oxidation of the Ru(0001) model catalyst surface under various reaction conditions starting from a strongly oxidizing environment to reaction conditions typical for CO oxidation. With a mixture of O_2 and CO (stoichiometry CO: $O_2 = 1:2$) the partial pressure of oxygen has to be increased to 20 mbar in order to form the catalytically active $RuO_2(110)$ oxide film, while with pure oxygen exposure a pressure of 10^{-5} mbar is sufficient to oxidize the Ru(0001) surface. In all oxidation experiments a self-limiting surface oxide is produced after an induction period. The RuO_2 film consists of $RuO_2(110)$ domains with mean lateral dimensions of 10 nm across. For preparation temperatures in the range of $550\text{--}630 \text{ K}$ the $RuO_2(110)$ film reveals a thickness of 1.6 nm . The chemical reduction of such ultrathin and flat RuO_2 films on Ru(0001) by hydrogen or CO exposure (sample temperature, 415 K) proceeds without an induction period.

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