$See \ discussions, stats, and author profiles for this publication \ at: \ https://www.researchgate.net/publication/242559603$ 

# (110) in Visible Light

ARTICLE	
CITATION	READS
1	280

# **5 AUTHORS**, INCLUDING:



Scott Walck

Army Research Laboratory

**53** PUBLICATIONS **1,589** CITATIONS

SEE PROFILE



Subscriber access provided by LIBRARY OF CHINESE ACAD SCI

# Article

# Photochemical Activity of Nitrogen-Doped Rutile TiO(110) in Visible Light

Oliver Diwald, Tracy L. Thompson, Tykhon Zubkov, , Scott D. Walck, and John T. Yates J. Phys. Chem. B, 2004, 108 (19), 6004-6008• DOI: 10.1021/jp031267y • Publication Date (Web): 17 April 2004 Downloaded from http://pubs.acs.org on March 27, 2009

#### **More About This Article**

Additional resources and features associated with this article are available within the HTML version:

- Supporting Information
- Links to the 28 articles that cite this article, as of the time of this article download
- Access to high resolution figures
- Links to articles and content related to this article
- Copyright permission to reproduce figures and/or text from this article

View the Full Text HTML



### Photochemical Activity of Nitrogen-Doped Rutile TiO<sub>2</sub>(110) in Visible Light

Oliver Diwald,† Tracy L. Thompson,† Tykhon Zubkov,† Ed. G. Goralski,‡ Scott D. Walck,‡ and John T. Yates, Jr.\*,†

Surface Science Center, Department of Chemistry, University of Pittsburgh, Pittsburgh, Pennsylvania 15260, and PPG Industries, Inc., Glass Technology Center, Pittsburgh, Pennsylvania 15238

Received: November 24, 2003; In Final Form: February 23, 2004

 $TiO_2(110)$  single crystals, doped with nitrogen via an NH<sub>3</sub> treatment at 870 K, have been found to exhibit photoactivity at photon energies down to 2.4 eV, which is 0.6 eV below the band-gap energy for rutile  $TiO_2$ . The active dopant state of the interstitial nitrogen that is responsible for this effect exhibits an N (1s) binding energy of 399.6 eV and is due to a form of nitrogen that is probably bound to hydrogen, which differs from the substitutional nitride state with an N (1s) binding energy of 396.7 eV. Optical absorption measurements also show enhanced absorption down to 2.4 eV for the NH<sub>3</sub>-treated  $TiO_2(110)$ . A co-doping effect between nitrogen and hydrogen is postulated to be responsible for the enhanced photoactivity of nitrogen-doped  $TiO_2$  materials in the range of visible light.

#### I. Introduction

The photocatalytic and photoelectrochemical applications of semiconducting oxide materials provide a strong scientific incentive for ongoing research.<sup>1–3</sup> Titanium dioxide (TiO<sub>2</sub>) belongs to one of the most extensively used and investigated systems and has the advantages of being inexpensive, chemically stable, and nontoxic. Other than the photocatalytic applications of TiO<sub>2</sub>, including the degradation of organic pollutants in air and aqueous solutions, it is used for photovoltaic applications<sup>4,5</sup> because it combines good electrical properties with excellent stability in many solvents over a wide pH range. However, because of the rather high intrinsic band gap of rutile TiO<sub>2</sub> (3.0 eV), only 4% of the incoming solar energy on the earth's surface can be utilized. Therefore, considerable efforts have been made to extend the photoresponse of TiO<sub>2</sub>-based systems further into the visible-light region, using dopants.

Various transition-metal cation dopants have been extensively investigated.<sup>6–8</sup> However, the disadvantage of cationic dopants is that they can result in localized d-levels deep in the band gap of TiO<sub>2</sub>, which often serve as recombination centers for photogenerated charge carriers. Thus, anionic nonmetal dopants, such as carbon,<sup>9</sup> sulfur,<sup>10</sup> and nitrogen,<sup>11</sup> may be more appropriate for extension of photocatalytic activity into the visible-light region, because the related impurity states are supposed to be close to the valence band maximum. Furthermore, the position of the conduction band minimum, which must be kept at the level of the H<sub>2</sub>/H<sub>2</sub>O potential, when TiO<sub>2</sub> is used for the photoelectrolysis of water into hydrogen and oxygen, is not affected.<sup>12</sup>

With respect to nitrogen-related dopants, Sato reported for the first time that the calcining of mixtures of  $Ti(OH)_4$  and ammonium salts leads to  $TiO_2$ -based materials that can be activated with visible light.<sup>13</sup> He attributed this beneficial doping effect to  $NO_x$  impurities in the  $TiO_2$  lattice. On the basis of a theoretical analysis, Asahi et al.<sup>14</sup> suggested that if nitrogen ions

substitute for oxygen in the TiO<sub>2</sub> lattice, the corresponding N (2p) states are located above the valence band edge. By mixing N (2p) states with O (2p) states, the band gap of the nitrogendoped TiO<sub>2</sub> is reduced and the material should show photoactivity at energies below the intrinsic band gap edge (<3.0 eV). In fact, these authors have proven that films and powders of TiO<sub>2-x</sub>N<sub>x</sub> exhibit a dramatic improvement over pure titanium oxide in their optical absorption below 3.0 eV and of the level of photocatalytic activity for visible light. Irie et al. 15 measured higher quantum yields for the decomposition of gaseous 2-propanol on  $TiO_{2-x}N_x$  with UV light than with visible light and concluded that the nitrogen-induced effect allowing subband-gap excitation is due to an isolated N (2p) state rather than to narrowing of the band gap. Lindgren et al.16 used photoelectrochemical measurements to confirm that the nitrogencreated states are located close to the valence band maximum and that the conduction band edge remains unchanged by nitrogen doping.

Generally, polycrystalline nitrogen-doped samples have been produced by different preparative procedures: (i) reactive sputtering of TiO2 targets with nitrogen/argon gas mixtures for films, and (ii) hydrolysis of NH<sub>x</sub>/TiO<sub>2</sub> mixtures or thermal treatment of TiO<sub>2</sub> in a NH<sub>3</sub> atmosphere for powder samples. In a recent single-crystal study, we have shown that ion implantation of atomically clean TiO<sub>2</sub>(110) surfaces with mixtures of N<sub>2</sub><sup>+</sup> and Ar<sup>+</sup> ions, followed by subsequent annealing under ultrahigh-vacuum conditions, is an effective way to incorporate nitrogen into the TiO<sub>2</sub> lattice.<sup>17</sup> X-ray photoelectron spectroscopy (XPS) revealed exclusively a N (1s) feature at a binding energy of 396.6 eV that is attributed to substitutionally bound nitride nitrogen.<sup>11</sup> In contrast to expectations, these nitrogendoped crystals, which contained only nitride ions, exhibited a shift in the photothreshold energy of 0.2 eV to higher energy, compared to undoped TiO<sub>2</sub>(110), as measured by the photodesorption of molecular oxygen.<sup>16</sup>

In the present study, we have used another method to dope  $TiO_2(110)$  with nitrogen, and we have found that the photo-threshold energy can be shifted 0.6 eV in a *desirable* direction to *lower* energy. The new method of nitrogen doping involves

<sup>\*</sup> Author to whom correspondence should be addressed. E-mail address: jyates@pitt.edu.

<sup>†</sup> University of Pittsburgh.

<sup>‡</sup> PPG Industries.

high-temperature treatment of TiO<sub>2</sub>(110) with NH<sub>3</sub> gas and is observed to produce an additional subsurface N (1s) feature at 399.6 eV.

To determine the photoactivity in the threshold region of rutile (<3 eV), the photochemical reduction of Ag<sup>+</sup> ions in aqueous solutions was measured. 18-20 The capture of photogenerated electrons on the TiO<sub>2</sub> surface by Ag<sup>+</sup> ions causes the deposition of metallic silver and can, therefore, be directly used for quantification of the photoeffect on TiO<sub>2</sub>. This was done for the first time by Fleischauer et al., 18 who measured reaction rates and quantum efficiencies on ZnO and TiO2 single-crystal surfaces by monitoring changes in the optical reflectivity due to the formation of metallic silver. By means of atomic force microscopy (AFM), Farneth and co-workers used the photocatalytic reduction of Ag<sup>+</sup> ions on polycrystalline TiO<sub>2</sub> films to investigate changes in the photoactivity of different rutile crystal faces. 19,20 This method has been extended to assess the photoactivity of individual TiO<sub>2</sub> particles.<sup>21</sup>

#### II. Experimental Section

Polished TiO<sub>2</sub>(110) single-crystal substrates (Princeton Scientific Corporation, Princeton, NJ) were treated in a flow reactor system in an argon gas atmosphere (1 atm) at 870 K. Typical heating rates were 7 K/min. For doping, the argon flow was replaced by NH<sub>3</sub> for 5 min, after the target temperature had been reached. Subsequently, the crystals were kept in flowing argon for 1 h at 870 K and then cooled in flowing argon over a time period of 2-3 h. Optical absorption spectra were acquired with a Perkin-Elmer Lambda 9 spectrophotometer.

For the silver reduction assay, doped and undoped crystal substrates were cleaved into slices  $\sim$ 1 mm  $\times$  5 mm in size and then mounted pairwise on an optical microscope slide. Six to ten slices of the NH3-treated and the blank crystal were assembled in two columns and immersed in an aqueous 0.1 N AgNO<sub>3</sub> solution, which covered the sample surface to a depth of 3 mm. Illumination was performed with a 500-W highpressure mercury lamp (Oriel). For light exposure of selected energy, the light beam was first passed through a water filter, to eliminate effects from IR radiation. Interference filters of selected wavelength, with a bandwidth of 10 nm (Andover Corporation), were used. The radiation intensity was measured at the sample position over the spectral region of interest, using a thermopile detector (Oriel model 71751) as a primary standard. The incident UV light intensity at the position of the sample surface, depending on the photon energy, was in the range of 1-4 mW/cm<sup>2</sup>. Typical light exposure times were 1-5 min, which corresponded to photon fluxes of the order of 10<sup>17</sup> photons/cm<sup>2</sup>.

After sample exposure to light of defined energy and photon flux, the mounted crystals were rinsed, first in distilled water and then in ethanol. Remaining liquids were removed from the sample surface by blowing with dry air and subsequent drying in an oven for 40 min at 140 °C in air.

Imaging of the silver deposits on the TiO2 surface was performed using AFM (Nanoscope IIIa Dimension 3100, from Digital Instruments). The probe tips used in this investigation were silicon nitride and were obtained from Veeco (model DNP-20). Images were collected in contact mode and were used to directly obtain information relating to the height and overall shape of the deposited silver clusters. Typical lateral scan frequencies were 1.5 Hz.

For quantification of the deposited silver, a flooding algorithm of the WSxM 4.0 software package (Nanotec Electronica S.L.) was used. The volume covered by silver clusters, per square

# Optical Absorption Spectra of NH<sub>3</sub>-N-Doped and Undoped TiO<sub>2</sub> (110) Crystals

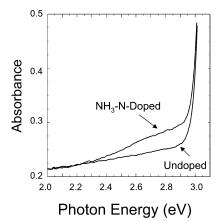


Figure 1. Absorption spectra of nitrogen-doped and undoped rutile TiO<sub>2</sub> single crystals measured at room temperature. The data have not been corrected for reflection loss.

centimeter, was taken as a measure of the quantity of deposited silver. Tip convolution effects, as well as the volume contribution from the detected silver clusters below a threshold of 2 nm, were neglected in this evaluation, and errors due to these two simplifications have a tendency to cancel each other. Values for the quantum yield were calculated as the number of Ag atoms produced per square centimeter, divided by the total number of photons per square centimeter impinging on the surface.

XPS measurements were performed at the PPG Glass Technology Center. The XPS spectra were acquired with a spectrometer (Vacuum Generators, model ESCALAB Mk II), using unmonochromatized Mg K\alpha X-rays (240 W). Measurements via XPS and Ar<sup>+</sup> sputtering steps (4 keV,  $I = 4 \mu A/$ cm<sup>2</sup>) were performed in a high-vacuum chamber with a base pressure of  $10^{-8}$  Torr. The binding energy of the N (1s) peaks were calibrated with respect to the O (1s) peak from TiO2 at 530.2 eV. The nitrogen concentration was estimated from the relative area intensities of the N (1s), Ti (2p), and O (1s) peaks. They were normalized using relative sensitivity factors in the Vacuum Generators software package.

#### III. Results and Discussion

A. Photoabsorption Spectrum. The absorption spectrum in Figure 1 reveals the fundamental absorption edge of rutile at 3.0 eV on both a nitrogen-doped rutile and an undoped rutile TiO<sub>2</sub>(110) single crystal. However, the nitrogen-doped crystal, which was produced via thermal treatment in an ammonia atmosphere, also has an increased optical absorption in the range of 2.4-3.0 eV. These crystals, after they have been doped, are bluish-green in color, whereas the undoped crystals are slightly yellowish in color, yet clear. The color of the nitrogen-doped crystals is partially due to the bulk reduction of the crystal, because the thermal decomposition of NH<sub>3</sub> on the TiO<sub>2</sub> surface results in the evolution of molecular hydrogen, which reduces the crystal electronically.<sup>22,23</sup> In addition, the presence of the absorption band below the edge at 3 eV indicates that nitrogen penetrates into the single crystal and effectively changes the electronic structure of the solid.

B. Silver Deposition on the Surface of Doped and Undoped TiO<sub>2</sub> Crystals. For investigation of the photocatalytic activity in the near-threshold region (2.3 eV  $\leq h\nu \leq$  3.0 eV), the silver deposition method was applied, using defined numbers of

TABLE 1: Energy Dependence of Silver Deposition on Undoped and Nitrogen-Doped Rutile TiO<sub>2</sub> Single Crystals

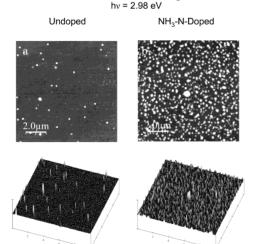
energy (eV)	number of photons $(\times 10^{17} \text{ cm}^{-2})$	Undoped TiO <sub>2</sub> Crystal		Nitrogen-Doped TiO <sub>2</sub> Crystal			
		number of images	silver concentration ( $\times 10^{-9} \text{ mol/cm}^2$ )	standard deviation	number of images	silver concentration ( $\times 10^{-9} \text{ mol/cm}^2$ )	standard deviation
dark	0	8	0.2	0.3	6	0.6	0.7
3.16	4.2	12	241.4	81.8	12	136.3	124.6
2.98	5.6	12	5.6	4.5	11	21.7	4.6
2.85	6.5	18	4.3	2.0	10	14.7	7.3
2.70	2.3	10	1.1	0.7	8	3.6	1.5
2.47	8.6	10	0.3	0.2	10	1.6	0.8
2.27	8.2	6	0.3	0.3	10	0.3	1.1

Ag Photoreduction on  $TiO_2(110)$  hv = 2.47 eVUndoped NH<sub>3</sub>-N-Doped

a b500mm

**Figure 2.** Contact-mode AFM images of silver deposits on the surface of (a and c) undoped and (b and d) NH<sub>3</sub>-nitrogen-doped TiO<sub>2</sub> single crystals after photoreduction with  $h\nu=2.47$  eV. Black-to-white vertical contrast in panels a and b (two-dimensional (2D) images) is 5 nm, whereas the three-dimensional (3D) representations in panels c and d have a vertical scale of 20 nm.

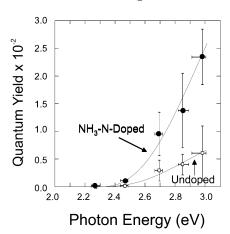
photons passing through 10-nm windows in the optical filters. Although unilluminated crystal surfaces exhibit only small noticeable changes before and after contact with the AgNO<sub>3</sub> solution (for quantification, see Table 1), the deposition of silver clusters on the doped crystal surfaces was significant after exposure to visible light below  $h\nu = 3.0$  eV. Figures 2a and b show typical topographic AFM images that have been obtained on samples after light excitation with  $h\nu = 2.47$  eV. The contrast results from vertical height differences on the surface, where black is low and white is high. Three-dimensional (3D) topographic reconstructions from the same data sets are displayed in Figures 2c and d. The doped and irradiated surface clearly is homogeneously covered by silver clusters with heights of 3-10 nm, whereas the surface of the undoped crystal contains only a limited number of clusters, which are inhomogeneously distributed. When the same experiment was performed with  $h\nu$ = 2.98 eV (Figures 3a and b), the silver deposition rate was dramatically enhanced and silver clusters with average heights of 10 nm are imaged via AFM. For quantification of the silver concentration, the average volume of these silver clusters, per square centimeter, was derived from the AFM data. Based on 6-18 scans of different regions of every sample surface, the average concentration is shown in Table 1 as a function of photon energy. Assuming that each Ag atom originates from a photogenerated electron, the quantum yield<sup>22</sup> was calculated for the near-threshold region (<3.0 eV) and plotted as a function of photon energy, as shown in Figure 4.24 In the near-threshold region (<3.0 eV), the concentration of deposited silver on the nitrogen-doped crystal was greater than that for the undoped



Ag Photoreduction on TiO<sub>2</sub>(110)

**Figure 3.** Contact-mode AFM images of silver deposits on the surface of (a and c) undoped and (b and d) NH<sub>3</sub>-nitrogen-doped TiO<sub>2</sub> single crystals after photoreduction with  $h\nu=2.98$  eV. Black-to-white vertical contrast in panels a and b (2D images) is 10 nm, whereas the 3D representations in panels c and d have a vertical scale of 50 nm.

# Action Curve for Photoinduced Ag-Deposition on TiO<sub>2</sub>(110)



**Figure 4.** Action curve, showing the quantum yield versus excitation energy for the photoreduction of silver on the surface of undoped and NH<sub>3</sub>-nitrogen-doped  ${\rm TiO_2}$  single crystals. Error bars in the quantum yield are based on the standard deviation from measurements on 8–18 images from different regions of the crystal (see Table 1). Error bars in photon energy correspond to the bandwidth of the interference filters used, where the location of the point is the band maximum of the respective filter transmission spectrum.

crystal (by a factor of 2 (hv = 2.47 eV) to 4 (hv = 2.98 eV)), which proves the substantial photocatalytic activity of the NH<sub>3</sub>-treated crystal in the visible-light spectral range. The threshold

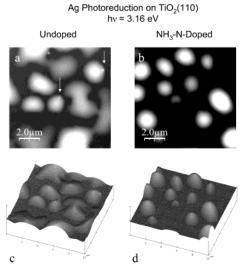


Figure 5. Contact-mode AFM images of silver deposits on the surface of (a and c) undoped and (b and d) NH3-nitrogen-doped TiO2 single crystals after photoreduction with  $h\nu = 3.16$  eV. Black-to-white vertical contrast in panels a and b (2D images) is 100 nm, whereas the 3D representations in panels c and d have a vertical scale of 200 nm.

photon energy for silver deposition on the nitrogen-doped TiO<sub>2</sub>-(110)—which is  $\sim$ 2.4 eV—is in good agreement with the optical absorption spectrum in Figure 1.

Higher rates for silver deposition were observed after irradiation experiments with a supra-band-gap photon energy (e.g.,  $h\nu = 3.16$  eV; see Figures 5a and c), leading to a much more complex topography of silver deposits. Different from all subband-gap experiments, both doped and undoped TiO<sub>2</sub> surfaces produce large silver islands with heights of 50-170 nm. The existence of much-smaller clusters, as indicated by arrows in Figure 5, suggests that the large features emerge from the coalescence of many primary particles during photodeposition or during the drying procedure at 310 K. The average silver concentration was determined to be  $(2.41 \pm 0.8) \times 10^{-7}$  mol/ cm<sup>2</sup> for an undoped crystal and (1.36  $\pm$  1.25)  $10^{-7}$  mol/cm<sup>2</sup> for the nitrogen-doped crystal. The large values for the standard deviations reflect the fact that, on both sample surfaces, some areas are uniformly coated, whereas others exhibit an inhomogeneous distribution of elevated features. The quantum yields at  $h\nu = 3.16$  eV for the undoped and the doped samples are 0.35 and 0.2, respectively. These values are too high, considering that the (110) orientation of rutile TiO<sub>2</sub> is an order of magnitude less photochemically reactive than the (101) plane, for which quantum yields between 0.6 and 0.4 were reported. 19,20 Together with the inhomogeneous distribution of silver, this suggests the presence of autocatalytic reactions during silver reduction with light that has supra-band-gap energy.

C. X-ray Photoelectron Spectroscopy. For chemical identification of the valence state of the dopant nitrogen, XPS was used. Surface analysis on doped and undoped TiO<sub>2</sub>(110) showed an impurity XPS feature that was centered at 399.6 eV. UV treatment in air and/or removal of the first 5 Å of the surface by sputtering with Ar<sup>+</sup> ions lead to the extensive depletion of these signals (Figure 6a). These traces of nitrogen must be attributed to surface contamination, which probably originates from spurious NH<sub>3</sub> in the flow reactor chamber. Their presence only on the surface excludes their role in the enhanced optical activity in the range of visible light. However, after the removal of 60 Å of TiO<sub>2</sub> by Ar<sup>+</sup> sputtering, two N (1s) features at 399.6 and 396.7 eV are observed, which indicates the presence of bulk nitrogen in two different chemical states. The total nitrogen



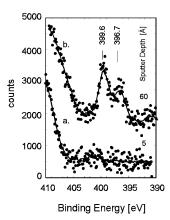


Figure 6. Nitrogen (N 1s) spectrum from a doped TiO<sub>2</sub>(110) crystal at two different Ar<sup>+</sup> sputter depths: (a) 5 Å and (b) 60 Å.

concentration at 60 Å was determined to be  $\sim$ 0.6 at. %. The feature at 396.7 eV is generally known as the "nitride" peak and has been assigned to N2- anions, which substitute for oxygen in the TiO2 lattice.25 As indicated by the smaller chemical shift, the nitrogen species that is associated with the feature at 399.6 eV carries less negative charge than a nitride ion. The value for the binding energy agrees well with XPS data from NH<sub>x</sub>-containing compounds<sup>26</sup> and suggests that the 399.6 eV N (1s) species is chemically bound to hydrogen.

Our previous studies of nitrogen implantation into TiO<sub>2</sub>(110) have shown that the photochemical absorption edge is shifted 0.2 eV to higher energy and that the substitutional nitride ion with an N (1s) binding energy of 396.7 eV is the only subsurface dopant species present.<sup>17</sup> This nitride species is probably also produced by the Ar+ ion sputtering method used to expose the subsurface region for XPS analysis in these experiments. This causes partial conversion into the nitride of nitrogen introduced by NH<sub>3</sub> doping at 870 K. We thus assign the nitrogen state (N (1s) = 399.6 eV) introduced by NH<sub>3</sub> as being responsible for the shift of the photochemical threshold, down to  $\sim$ 2.4 eV. This form of nitrogen is most likely located in an interstitial site, and this form of nitrogen dopant is probably bound to hydrogen. These results disagree with the conclusions of Asahi et al.,14 who reported that nitride ions that substitute for O2- ions in the TiO<sub>2</sub> lattice are the necessary dopant species for TiO<sub>2</sub> photocatalysis in the visible-light region.

D. Comparison of Nitrogen-Doping Effects in TiO2 with Those in ZnO. Although the effect of nitrogen doping in TiO<sub>2</sub>-(110) is clearly shown in this work, the explanation for this effect is still somewhat under debate. The effect of doping in other wide-band-gap semiconductors, such as zinc oxide, for example, has recently also been studied. Zinc oxide (ZnO) is another n-type semiconductor, similar to  $TiO_2$ , with an intrinsic band gap of  $E_g = 3.3$  eV. In ZnO, nitrogen as a dopant should produce shallow acceptor levels within the band gap. However, at this point in time, no one has successfully obtained p-type ZnO using a pure nitrogen source.<sup>27</sup> On the other hand, the codoping of nitrogen, in conjunction with hydrogen or gallium, has been shown to be successful.<sup>28,29</sup> The presence of donor dopants, as suggested by Yamamoto,<sup>30</sup> has two consequences that may also apply to N:H-doped TiO2. The first is an enhanced incorporation of acceptors, because of their strong electrostatic attractive interaction, which leads to higher acceptor (nitrogen) incorporation in TiO<sub>2</sub>. The second consequence is the lowering of the acceptor energy level, bringing the level closer to the valence band. A hybridized state (N p-orbitals and H s-orbitals) is postulated to be formed in  $TiO_2$ , by analogy to ZnO. Both effects are beneficial for the enhanced optical absorption of  $ZnO^{28}$  and might also explain the photoactivity of N:H-doped  $TiO_2$  in the visible spectral range.

#### **IV.** Conclusion

These findings, in conjunction with a recent investigation of nitrogen-implanted  $TiO_2(110)$ ,  $^{16}$  indicate that nitrogen with an N (1s) binding energy of 399.6 eV is effective in reducing the threshold photon energy for photochemistry from 3.0 eV to 2.4 eV, which is a shift of 0.6 eV into the visible spectral region. This form of nitrogen can be introduced into the  $TiO_2$  bulk via heat treatment in NH<sub>3</sub> at 870 K. This active nitrogen is likely to be interstitial and chemically bound to hydrogen in the first few hundred Ångstroms below the  $TiO_2$  surface. In contrast, substitutional nitrogen, in the form of implanted nitride species, is inactive for lowering the photochemical threshold energy below the band-gap energy of rutile  $TiO_2$  ( $E_g = 3.0$  eV). These findings are in contrast to current scientific opinion.

**Acknowledgment.** This work was supported by the DoD Multidisciplinary University Research Initiative (MURI) program, administered by the Army Research Office, under Grant No. DAAD-19-01-0-0619. We would like to thank Yujie Sun for helpful advice and direction with the AFM measurements made at the University of Pittsburgh. Special thanks to Pittsburgh Plate Glass (PPG) for scientific collaboration and analytical services.

#### References and Notes

- (1) Linsebigler, A.; Lu, G.; Yates, J. T., Jr. Chem. Rev. 1995, 95, 735–758.
- (2) Hoffmann, M. R.; Martin, S. T.; Choi, W.; Bahnemann, D. W. *Chem. Rev.* **1995**, *95*, 69–96.
- (3) Serpone, N., Pilezzetti, E. Eds. *Photocatalysis: Fundamentals and Applications*; Wiley-Interscience: New York, 1989.
  - (4) Fujishima, A.; Honda, K. Nature 1972, 238, 37-38.
  - (5) O'Regan, B.; Gratzel, M. Nature 1991, 353, 737-739.
- (6) Choi, W.; Termin, A.; Hoffmann, M. R. J. Phys. Chem. B 1994, 13669-13679.

- (7) Hermann, J. M.; Disdier, J.; Pichat, P. Chem. Phys. Lett. 1984, 108. 618–622.
  - (8) Anpo, M.; Takeuchi, M. J. Catal. 2003, 216, 505-516.
- (9) Khan, S. U. M.; Al-Shahry, M.; Ingler, W. B., Jr. Science 2002, 297, 2243-2245.
- (10) Umebayashi, T.; Yamaki, T.; Itoh, H.; Asai, K. Appl. Phys. Lett. **2002**, 81, 454–456.
- (11) Morikawa, T.; Asahi, R.; Ohwaki, T.; Aoki, K.; Taga, Y. *Jpn. J. Appl. Phys.* **2001**, *40*, L561–L563.
  - (12) Hagfeldt, A.; Grätzel, M. Chem. Rev. 1995, 95, 49-68.
  - (13) Sato, C. S. Chem. Phys. Lett. 1986, 123, 126-128.
- (14) Asahi, R.; Morikawa, T.; Ohwaki, T.; Aoki, K.; Taga, Y. Science **2001**, 293, 269–271.
- (15) Irie, H.; Watanabe, Y.; Hashimoto, K. J. Phys. Chem. B **2003**, 107, 5483–5486.
- (16) Lindgren, T.; Mwabora, J. M.; Avendaño, E.; Jonsson, J.; Hoel, A.; Granqvist, C.-G.; Lindquist, S.-E. J. Phys. Chem. B 2003, 107, 5709—5716
- (17) Diwald, O.; Thompson, T. L.; Goralski, E. G.; Walck, S. D.; Yates, J. T., Jr. J. Phys. Chem. B **2004**, 108, 52–57.
- (18) Fleischauer, P. D.; Kann, H. K. A.; Shepherd, J. R. *J. Am. Chem. Soc.* **1972**, *94*, 283–285.
- (19) Hotsenpiller, P. A. M.; Bolt, J. D.; Farneth, W. E.; Lowekamp, J. B.; Rohrer, G. S. J. Phys. Chem. B 1998, 102, 3216–3226.
- (20) Lowenkamp, J. B.; Rohrer, G. S.; Hotsenpiller, P. A.; Bolt, J. D.; Farneth, W. E. *J. Phys. Chem. B* **1998**, *102*, 7323–7327.
- (21) Farneth, W. E.; McLean, R. S.; Bolt, J. D.; Dokou, E.; Barteau, M. A. *Langmuir* **1999**, *15*, 8569–8573.
- (22) Rekoske, J. E.; Barteau, M. A. J. Phys. Chem. B 1997, 101, 1113–1124.
- (23) This produces a broad optical absorption (from 0.5 eV to 2 eV) that is associated with electronically reduced defects in the bulk (not shown in Figure 1).
- (24) The accuracy of the quantum yield in Table 1 is dependent on the validity of the assumptions made, i.e., neglect of tip convolution effects and neglect of features below the 2-nm threshold. The AFM probe has a finite radius; therefore, the lateral dimensions of small protruding features from the surface are systematically overestimated. Thus, the values in Table 1 are an upper limit for the actual quantum yields.
- (25) Saha, N. C.; Tompkins, H. G. J. Appl. Phys. **1992**, 72, 3072–3079.
- (26) Souto, S.; Alvarez, F. Appl. Phys. Lett. 1997, 70, 1539-1541 and references therein.
- (27) Lee, E.-C.; Kim, Y.-S.; Jin, Y.-G.; Chang, K. J. *Phys. Rev. B* **2001**, 64, 085120 (1–5).
- (28) Minegishi, K.; Koiwai, Y.; Kikuchi, Y.; Yano, K.; Kasuga, M.; Shimizu, A. *Jpn. J. Appl. Phys.* **1997**, *36*, L1453–L1455.
- (29) Joseph, M.; Tabata, H.; Kawai, T. Jpn. J. Appl. Phys. **1999**, 38, L1205-L1207.
  - (30) Yamamoto, T. Thin Solid Films 2002, 421, 100-106.