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# Layer-by-Layer Self-Assembly: Toward Magnetic Films with Tunable Anisotropy

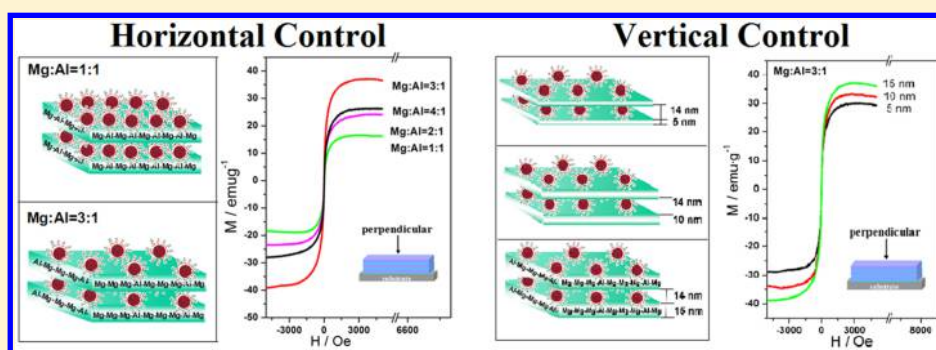
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## S Supporting Information



**ABSTRACT:** Highly anisotropic magnetic films are possible to fabricate by control of the coupling between individual magnetic particles. Selective control over coupling in the horizontal and vertical directions are of both fundamental and practical interest. Here we show such control in the multiple layer-by-layer (LBL) self-assembly of layered double hydroxide (LDH) nanosheets ( $x = 1, 2, 3$  and  $4$ ) with thicknesses of  $5\text{--}15\text{ nm}$ , co-assembled with 3-aminopropyl-trimethoxysilane (APTS) modified spherical  $\text{Fe}_3\text{O}_4$  nanoparticles (APTS- $\text{Fe}_3\text{O}_4$  NPs) on quartz substrates. The electrostatic charge density on the LDH sheets, controlled by the Mg/Al composition ratio, affects the NP packing in a single horizontal layer, while the thickness of the LDH sheets controls magnetic coupling between layers. The tunable magnetic properties (coercivity  $H_c$ , saturation magnetization  $M_s$ , anisotropy, and blocking temperatures) are measured as a function of these parameters. The maximum saturation magnetizations  $M_s$ ,  $36.3$  and  $25.1\text{ emu}\cdot\text{g}^{-1}$  in the perpendicular and parallel direction, respectively, are found for the sample of  $x = 3 = \text{Mg/Al}$  ratio in the LDH layer, and  $15\text{ nm}$  LDH layer thickness. This work provides a general method to adjust the anisotropy of magnetic films based on directional control of coupling of magnetic nanoparticles between and across, LDH nanosheets. We outline how higher anisotropy and even finer control could be achieved by pH and composition control over the electrostatic charge of the assembly components.

## 1. INTRODUCTION

The anisotropy of magnetic nanoparticle assemblies, resulting in magnetic hysteresis, has raised considerable attention, since it is required for the building blocks of nanoscale memory cells<sup>1,2</sup> and nontrivial quantum dynamics that are crucial for quantum information processing.<sup>3,4</sup> However, achieving 3D structures with control over magnetic anisotropy from the most common spherical nanoparticle still remains a great challenge because of the need to control spacing and orientation in several directions. Spherical magnetic nanoparticles can be densely packed into patterned templates (e.g., nanotubes, nanoholes, etc.), and by controlling interparticle distances their dipole–dipole interactions can have a great effect on anisotropy.<sup>5–8</sup> Most methods, however, are only effective in controlling the interparticle distance equally in all directions by tuning the diameter and spacing of the particles, or selectively only in 2D by the pattern template on which the particles adsorb. Another

common method employed for obtaining magnetic composite 3D structures is the dispersion of magnetic nanoparticles into a polymer matrix, but tunability of the magnetic anisotropy is still hard to achieve because magnetic nanoparticles are often randomly distributed throughout the polymer matrix and interparticle distance is difficult to control.<sup>9</sup> For obtaining spherical nanoparticle assemblies with high and tunable magnetic anisotropy it is highly desirable to have simultaneous, separate control of the interparticle distance in both the horizontal and vertical directions. Layer-by-layer (LBL) assembly methods have been used successfully to fabricate multilayer films of  $\text{Fe}_3\text{O}_4$  nanoparticles.<sup>10–14</sup> There is still a

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need, however to develop better separate control over the spacing in both the horizontal and vertical directions.

Recently, the intrinsic anisotropy of inorganic nanosheets of layered double hydroxides (LDHs) has attracted strong interest due to their ability to induce optical anisotropy.<sup>15,16</sup> In particular, optical films with high anisotropy fabricated by electrostatic self-assembly of LDH building blocks and optical molecules with highly symmetric structure<sup>17</sup> or encapsulated in micelles can also be obtained.<sup>18</sup> In addition, to their planar structural rigidity, it is easy to synthesize LDH sheets with precise composition and control over their surface electrostatic charge. Inspired by both factors, we hypothesized LBL assemblies of LDH nanosheets and magnetic nanoparticles may result in magnetic materials with highly tunable anisotropy. The rigid and positively charged LDH nanosheets can fix and disperse  $\text{Fe}_3\text{O}_4$  nanoparticles, avoiding aggregation and enhancing stability; while the controlled charge density and thickness of LDH nanosheets impart the ability for simultaneously tuning the interparticle distance of magnetic nanoparticles in both horizontal and vertical directions.

Here we demonstrate a simple method to fabricate nanocomposite films with high tunability of spacing and magnetic anisotropy. In order to realize the assembly of monodispersed  $\text{Fe}_3\text{O}_4$  nanoparticles and LDH nanosheets, 3-aminopropyltrimethoxysilane (APTS) was utilized to modify  $\text{Fe}_3\text{O}_4$  nanoparticles.<sup>19</sup> The APTS- $\text{Fe}_3\text{O}_4$ /LDH films were obtained by alternate LBL assembly of LDH nanosheets and APTS- $\text{Fe}_3\text{O}_4$  nanoparticles on a quartz substrate. SEM and AFM images show that the film surface is continuous and uniform, and UV-vis absorption spectra indicate a stepwise and regular film growth procedure. The magnetic anisotropy of the films can be tuned by controlling the interparticle distance of  $\text{Fe}_3\text{O}_4$  nanoparticles in the horizontal and vertical directions, through simply varying the charge density and thickness of the LDH layer. Though we only show assembly at pH = 7.5–8 and integer  $x = \text{Mg}/\text{Al}$  ratios, we show the method could be extended with potentially higher control at higher pH-s and non-integer ratios. The method could potentially be used in the fields of magnetic probes, separations and sensors for ions and biomolecules.

## 2. EXPERIMENTAL SECTION

**2.1. Materials.** Analytical grade chemicals including  $\text{Mg}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$ ,  $\text{Al}(\text{NO}_3)_3 \cdot 9\text{H}_2\text{O}$ ,  $\text{NH}_3 \cdot \text{H}_2\text{O}$ ,  $\text{H}_2\text{SO}_4$ ,  $\text{H}_2\text{O}_2$ ,  $\text{NaOH}$ ,  $\text{HCl}$ , formamide,  $\text{FeCl}_3 \cdot 6\text{H}_2\text{O}$ ,  $\text{FeCl}_2 \cdot 4\text{H}_2\text{O}$ , and 3-aminopropyltrimethoxysilane (APTS) were purchased from Sigma-Aldrich and used without further purification. Deionized and decarbonated water were used in all the experimental processes.

**2.2. Preparation of APTS- $\text{Fe}_3\text{O}_4$  Nanoparticles.**  $\text{Fe}_3\text{O}_4$  nanoparticles were prepared by coprecipitation of  $\text{Fe}^{2+}$  and  $\text{Fe}^{3+}$  in the presence of  $\text{NaOH}$ .<sup>20,21</sup> A mixture solution of 50 mL of 1 M  $\text{FeCl}_3 \cdot 6\text{H}_2\text{O}$ , 50 mL of 0.5 M  $\text{FeCl}_2 \cdot 4\text{H}_2\text{O}$ , and 50 mL of 0.4 M  $\text{HCl}$  was stirred vigorously under  $\text{N}_2$  at 80 °C, and then 250 mL of 0.5 M  $\text{NaOH}$  was rapidly added. The black color precipitates were obtained and stirring was kept for 1 h at 80 °C. After that, resulting  $\text{Fe}_3\text{O}_4$  nanoparticles were cooled to room temperature and then rinsing with water several times. The above nanoparticles (0.01 g) were then added into 100 mL of APTS (0.05 mol  $\text{L}^{-1}$ ) solution treated in an ultrasonic water bath for 20 min at room temperature.<sup>22</sup> The APTS- $\text{Fe}_3\text{O}_4$  nanoparticles were obtained by mechanical stirring for 5 h, followed by washing with methanol and water, after which 0.01

g of APTS- $\text{Fe}_3\text{O}_4$  nanoparticles was dispersed in 10 mL of water with pH = 7.5–8.0 using an ultrasonic water bath to form a stable dispersion of slightly negatively charged nanoparticles (Figure S1). Magnetic separation was employed when washing all of above samples in our work.

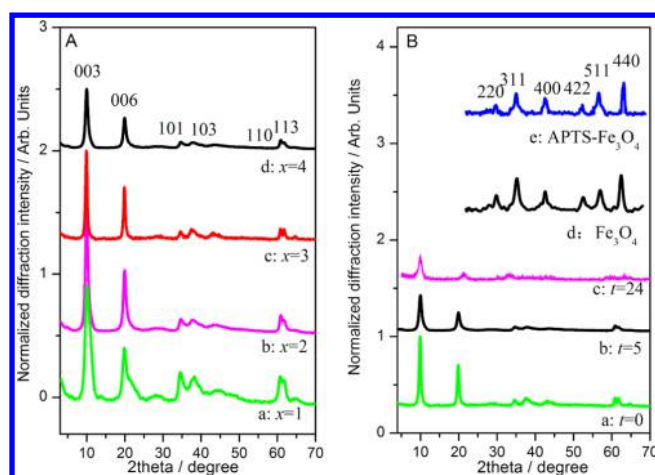
**2.3. Fabrication of the (APTS- $\text{Fe}_3\text{O}_4$ /LDH) $_n$  Films.** The  $\text{Mg}_x\text{Al-NO}_3$  LDH ( $x = 1, 2, 3$ , and 4) nanosheets were synthesized by the hydrothermal method reported previously.<sup>23</sup> The exfoliated  $\text{Mg}_3\text{Al-NO}_3$  LDH nanosheets of different thickness were obtained by shaking a 0.1 g of  $\text{Mg}_3\text{Al-NO}_3$  LDH in 100 mL of formamide solution for different times (0, 5, and 24 h). The cleaned quartz glass substrate was obtained by immersion in concentrated  $\text{NH}_3/30\% \text{H}_2\text{O}_2$  (7:3) and concentrated  $\text{H}_2\text{SO}_4$  for 30 min each. After each procedure, the quartz substrate was rinsed and washed thoroughly with water. The substrate was dipped in a colloidal suspension (0.1 g  $\text{mL}^{-1}$ ) of LDH nanosheets for 10 min followed by washing thoroughly, and then the substrate was treated with a 100 mL of APTS- $\text{Fe}_3\text{O}_4$  aqueous solution (0.01 wt %) for another 10 min followed by washing in water. Multilayer (APTS- $\text{Fe}_3\text{O}_4$ /LDH) $_n$  films were fabricated by alternate deposition of LDH nanosheets and APTS- $\text{Fe}_3\text{O}_4$  nanoparticle suspension for  $n$  cycles. All of the water used in the rinsing and washing process was with pH = 7.5–8.0.

**2.4. Characterizations.** UV-vis absorption spectra were recorded (PerkinElmer Lambda 750) in the range from 200 to 800 nm. X-ray powder diffraction (XRD) patterns were measured using a Philips PW1820 diffractometer under the conditions 40 kV, 50 mA, Cu  $K\alpha$  radiation ( $\lambda = 0.154 \text{ nm}$ ) step-scanned with a scanning rate of  $0.5^\circ/\text{min}$  and a  $2\theta$  angle ranging from  $2^\circ$  to  $70^\circ$ . The morphology was investigated by using a scanning electron microscope (SEM, ZEISS), transmission electron microscope (TEM, JEOL JEM-2100), and atomic force microscope (AFM, Digital Instruments, version 6.12). The starting materials and the assemblies were analyzed by X-ray photoelectron spectroscopy (XPS) with monochromatized Al  $K\alpha$  exciting X-radiation (PHI Quantera SXM) and Raman spectra with a confocal Raman microspectrometer, 514.5 nm of excitation and 0.50 mW laser power (Renishaw Instruments Co. Ltd., RM2000). The field-dependence magnetization measurements were carried out by a vibrating sample magnetometer (VSM, MicroMag 2900/3900) at room temperature, and the calculation of the magnetization value was based on the film weight (subtracting the weight of the substrate). The curves of magnetization vs temperature were recorded with a superconducting quantum interference device (SQUID) magnetometer (Quantum Design, MPMS MultiVu Application, Revision 1.60). For ZFC measurements, the film samples were first cooled from 300 to 10 K in zero field. Then, under the magnetic field of 10 Oe, the magnetization was recorded during heating the sample up to 300 K. On the other hand, for FC measurements, the film samples were cooled from 300 to 10 K under the magnetic field of 10 Oe, and magnetization was measured in the heating regime under the same magnetic field.

## 3. RESULTS AND DISCUSSION

### 3.1. Structural and Morphological Characterizations.

Figure 1 shows the powder XRD patterns of  $\text{Mg}_x\text{Al-LDHs}$  ( $x = 1, 2, 3$  and 4),  $\text{Fe}_3\text{O}_4$ , APTS- $\text{Fe}_3\text{O}_4$ , and exfoliated  $\text{Mg}_3\text{Al-LDHs}$  ( $t = 0, 5$ , and 24 h). The patterns of  $\text{Mg}_x\text{Al-LDHs}$  ( $x = 1, 2, 3$ , and 4) nanosheets can be indexed to a hexagonal lattice (Figure 1A, curves a–d). The two sharp and intense characteristic



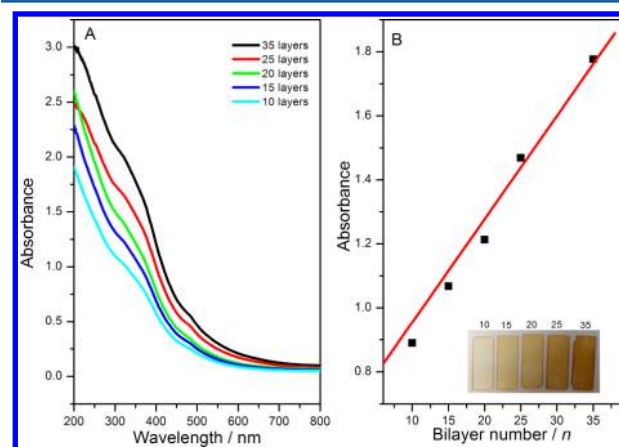
**Figure 1.** XRD patterns of (A) the Mg<sub>x</sub>Al-LDHs ( $x = 1, 2, 3$ , and  $4$ ) nanosheets and (B) the exfoliated Mg<sub>3</sub>Al-LDHs nanosheets ( $t = 0, 5$ , and  $24$  h), Fe<sub>3</sub>O<sub>4</sub> nanoparticles, and APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles. All diffraction intensities were normalized with respect to their corresponding strongest peak.

diffraction peaks appearing as symmetric lines at  $2\theta$  10° and 20° are ascribed to the (003) and (006) plane, consistent with the presence of nitrate as charge compensating anions in the interlayer space. For the exfoliated Mg<sub>3</sub>Al-LDHs ( $t = 0, 5$ , and  $24$  h, Figure 1B) nanosheets, the decreased intensity and increased half-peak width (from 0.76° to 0.92°) of the (003) reflection upon prolonging the exfoliated time indicates decreased crystallinity. The diffraction peaks of Fe<sub>3</sub>O<sub>4</sub> are indexed a face centered cubic structure (JCPDS No. 65-3107, Figure 1B, curve d).<sup>24</sup> Furthermore, the coated APTS molecule shows no influence on the Fe<sub>3</sub>O<sub>4</sub> crystallinity (Figure 1B, curve e).

In order to avoid the effect of LDH nanosheet size on the dispersion degree of APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles, and directly compare the effects of composition, LDH nanosheets with nearly equal sizes were prepared. TEM images (Figure 2A–D) show the diameter of the Mg<sub>x</sub>Al-LDHs ( $x = 1, 2, 3$ , and  $4$ ) nanosheets is  $\sim 100$  nm. HRTEM images show the Fe<sub>3</sub>O<sub>4</sub> nanoparticles have an average size of  $\sim 13$  nm, and the particle size distributions are in good agreement with the normal distribution (Figures 2E and 2G, respectively). The HRTEM image shows the clear and continuous lattice fringes (Figure 2F). The lattice fringe pitch of 0.25 nm is magnified and

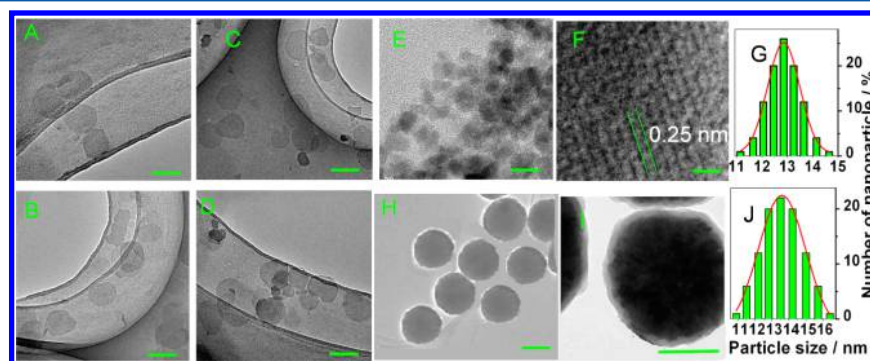
corresponds well with the  $d$ -spacing of the (311) reflections for Fe<sub>3</sub>O<sub>4</sub> in agreement with the values of 2.53 Å obtained from the JCPDS database.<sup>25</sup> After coating with APTS, the dispersion degree displays an obvious increase because the APTS coating inhibits aggregation of Fe<sub>3</sub>O<sub>4</sub> nanoparticles (Figure 2H). The core–shell structure can be observed in the HRTEM image with an average size of 14 nm (Figure 2I), and the coating layer is uniform with a thickness about 1 nm. The particle size distributions are in good agreement with the normal distribution (Figure 2J).

In order to investigate the induced magnetic anisotropy behavior of spherical Fe<sub>3</sub>O<sub>4</sub> nanoparticles, the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>*n*</sub> films were fabricated through alternate deposition of the LDHs nanosheets and APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles on quartz substrates. Figure 3A shows the UV–vis absorption spectra of



**Figure 3.** (A) UV–vis spectra of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>*n*</sub> films ( $x = 3$  and  $t = 0$ ) with various bilayer numbers. (B) Absorption intensity as a function of bilayer number (inset: photographs).

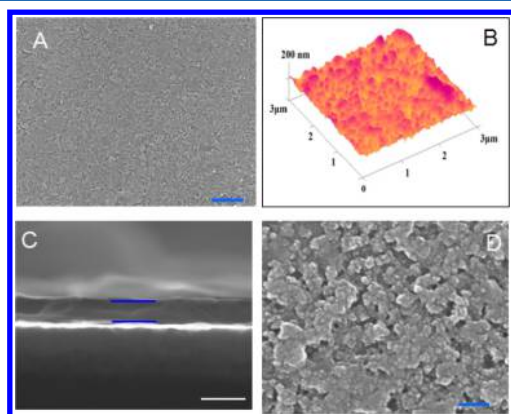
the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>*n*</sub> films ( $x = 3$  and  $t = 0$ ,  $t$  stands for exfoliating time of LDH nanosheets) with  $n$  indicating different number of bilayers. Although no obvious absorption peaks of Fe<sub>3</sub>O<sub>4</sub> were observed, the absorption intensity of the films across the whole spectrum range gradually increases with the increase of the number of bilayers. The absorption intensity at  $\sim 350$  nm correlates linearly with  $n$  (Figure 3B), indicating a stepwise and regular film growth procedure. In the cases where  $x \neq 3$ , similar phenomena are also observed (not shown here). This was further confirmed by the gradual color enhancement



**Figure 2.** (A–D) TEM images of the Mg<sub>x</sub>Al-LDHs ( $x = 1, 2, 3$ , and  $4$ , scale bar = 100 nm). (E) HRTEM image of Fe<sub>3</sub>O<sub>4</sub> nanoparticles at low magnification (scale bar = 10 nm), (F) same high-magnification (scale bar = 0.5 nm). (G) Particle size distribution of the Fe<sub>3</sub>O<sub>4</sub> nanoparticles. (H) HRTEM image of APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles low magnification (scale bar = 10 nm), (I) same at high magnification (scale bar = 0.5 nm). (J) Particle size distributions of APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles.



with the increase of bilayer number (inset in Figure 3B). The surface morphology of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>10</sub> ( $x = 3$  and  $t = 0$  h) film was characterized by SEM and AFM images. SEM image (Figure 4A) of the film from a top view shows that the



**Figure 4.** (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>10</sub> film ( $x = 3$  and  $t = 0$  h) (A) Top view SEM image of the film (scale bar = 2  $\mu$ m). (B) AFM image. (C) side view SEM image (scale bar = 500 nm). Blue markers indicate the bottom and top of the film. (D) high-magnification SEM image (scale bar = 200 nm).

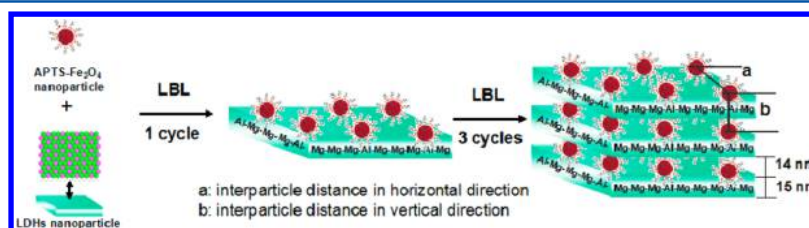
film surface is continuous and uniform, with a root-mean-square (RMS) roughness of  $\sim 15.8$  nm (Figure 4B). The film thickness is around 300 nm (Figure 4C), from which the thickness of one bilayer (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>1</sub> ( $x = 3$  and  $t = 0$  h) was estimated to be  $\sim 30$  nm. According to the thickness of  $\sim 15$  nm for an LDH nanosheet and the diameter of  $\sim 14$  nm for an APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticle, APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH supramolecular structure demonstrates the ideal single-layered arrangement model (Figure 5). The film thickness approximately linear increases as a function of bilayer number, confirming a uniform and periodic layered structure (Figure S2). A high-magnification SEM image (APTS-Fe<sub>3</sub>O<sub>4</sub> as the terminal layer, Figure 4D) demonstrates the APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticle is well-dispersed on the layer of LDH.

**3.2. Magnetic Properties of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>n</sub> Films.** In order to obtain films with magnetic anisotropy we control the Fe<sub>3</sub>O<sub>4</sub> interparticle spacing, in the horizontal and vertical directions and measure the magnetic properties (coercivity  $H_c$  and saturation magnetization  $M_s$ ) of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>n</sub> films. The main control tool, as discussed, are the charge density and layer thickness of LDH nanosheets.

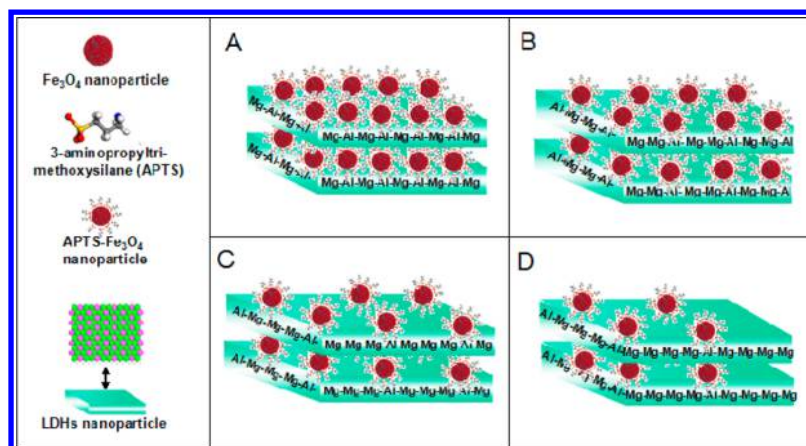
**3.2.1. Tuning of the Particle Density in the Horizontal Direction by Changing the Charge Density of the LDH Nanosheets.** The tunability of positive charge density on the surface of the LDH nanosheets originates from their tunable metal ratio of bivalent and trivalent. This provides the chance of

controlling dispersion degree and interparticle distance for APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles, which is absolutely critical for the purpose of tuning the magnetic property for collective films. The interaction of adjacent particles can be neglected when the interparticle distance is more than a critical value.<sup>26</sup> Therefore, in order to eliminate the interparticle interaction in vertical direction as far as possible, the unexfoliated LDH nanosheets with a maximum thickness of  $\sim 15$  nm were used to assemble with spherical APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles and are shown in Figure 6.

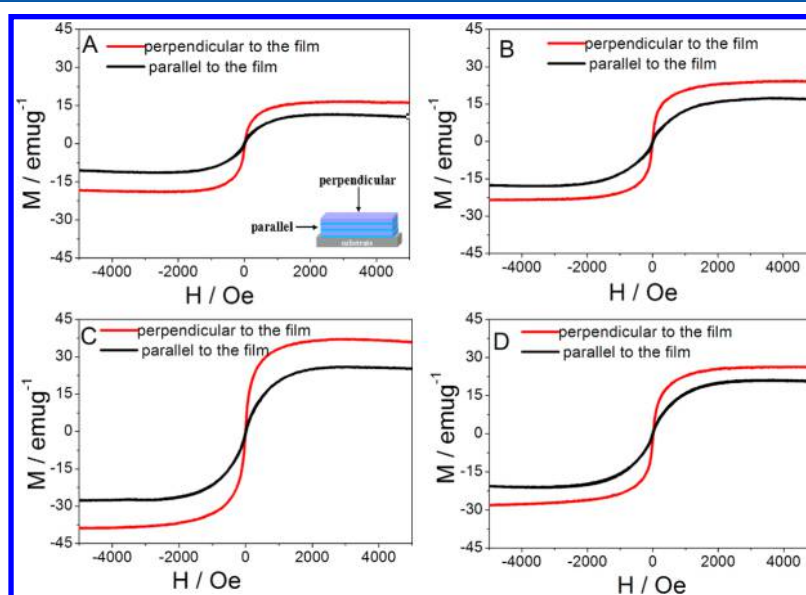
The magnetic hysteresis loops for the Fe<sub>3</sub>O<sub>4</sub>, APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticle (Figure S3), and the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 1, 2, 3$ , and 4, Figure 7) measured at room temperature are recorded. All of the samples exhibit the superparamagnetic behavior with negligible  $H_c$  (less than 14.0 Oe) and  $M_r$  (less than 0.90 emu g<sup>-1</sup>).<sup>27,28</sup> The (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 1, 2, 3$ , and 4) display noncoincident hysteresis loops in the parallel and perpendicular directions of magnetic field, compared to the coincident curves of the Fe<sub>3</sub>O<sub>4</sub> and APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles, indicating magnetic anisotropy of films. The  $H_c$  value in the parallel direction is more than 3-folds higher than that in the perpendicular direction, but the  $M_s$  in the parallel direction is only  $\sim 80\%$  of the perpendicular one (Table S1). This demonstrates that the rigid and oriented LDH nanosheets induce ordered alignment of the spherical Fe<sub>3</sub>O<sub>4</sub> nanoparticles, resulting in magnetic anisotropy property of films. When magnetic fields are applied both parallel and perpendicular to the film,  $M_s$  increases at first and then decreases with further increase of  $x$ . The maximum values of 36.3 and 25.1 emu g<sup>-1</sup> in perpendicular and parallel directions, respectively, are reached in the case of  $x = 3$  due to excellent degree of dispersion and interparticle distance of APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles on LDH layer. This result indicates that the magnetic properties ( $H_c$ ,  $M_s$ , and anisotropy) of films can be effectively controlled by changing the degree of dispersion and interparticle distance of APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles through simple tuning of the charge density of LDH layer. This can be explained by the “charge density-matching” assembly between the APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles and LDH nanosheets. In the case of  $x \neq 3$ , the charge density-mismatch may lead to the formation of aggregates ( $x = 1$  and 2) and reduced stack density ( $x = 4$ ) of APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles, consequently resulting in the decrease of  $M_s$ . Based on results above, particle shape is not the cause for the anisotropy of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 1, 2, 3$ , and 4), since all the particles are spherical. Thus, the changes in anisotropy originate from the different dipolar interactions between NPs in the same layer. Moreover, it is noted that the  $M_s$  values for all of films in our work are lower than previously reported values (80–90 emu g<sup>-1</sup>),<sup>29,30</sup> which is caused by the lower content of pure Fe<sub>3</sub>O<sub>4</sub> nanoparticle on the film. Therefore, the above results demonstrate that the magnetic properties of the assembled



**Figure 5.** A schematic representation for the assembly and structure of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>3</sub> film ( $x = 3$  and  $t = 0$  h).



**Figure 6.** A schematic representation for the assembly and structure of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>2</sub> films and the effect of changing the composition (and charge) of the Mg<sub>x</sub>Al<sub>1-x</sub>-LDH material: (A)  $x = 1$  and  $t = 0$  h, (B)  $x = 2$  and  $t = 0$  h, (C)  $x = 3$  and  $t = 0$  h, and (D)  $x = 4$  and  $t = 0$  h.



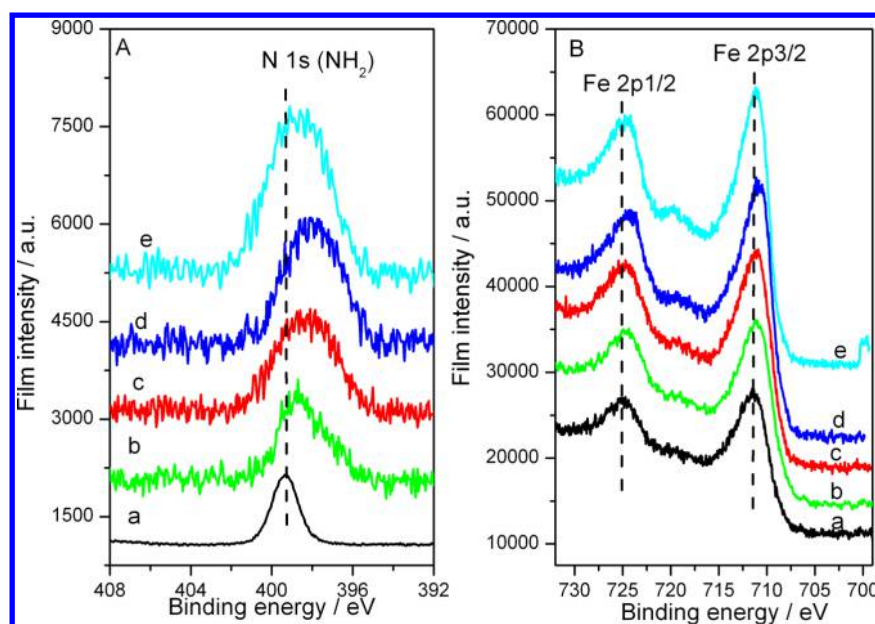
**Figure 7.** (A–D) Magnetic hysteresis loops of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 1, 2, 3$ , and  $4$ , respectively), at room temperature.

films can be tuned by simply varying the charge density of the LDH layer. As Mg<sup>2+</sup> and Al<sup>3+</sup> can be mutually substituted in the LDH lattice, using non-integer  $x = \text{Mg}/\text{Al}$  ratios would allow precise tuning of the electrostatic charge on the surface.

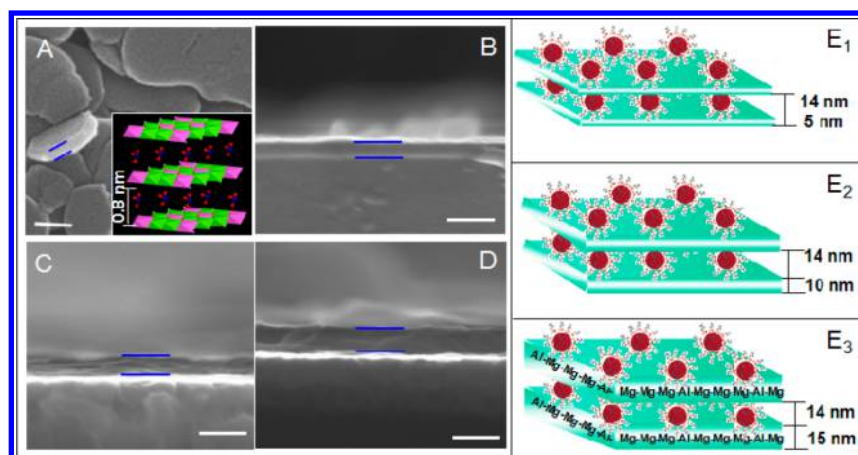
The XPS measurement is used to explain the effect of charge density of LDH nanosheets on the magnetic property of (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 1, 2, 3$ , and  $4$ ) and is shown in Figure 8. Compared to the original APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles (NH<sub>2</sub>: 399.2 eV, Figure 8A, curve a), the shift in the XPS peak of films for N 1s (NH<sub>2</sub>: 398.7–398.1 eV, Figure 8A, curves b–e) to lower binding energy indicates APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles interact strongly with the LDH layers. The shift is monotonic, reaching a maximum in the film with  $x = 3$  which we attribute to the best electrostatic charge match of LDH nanosheets to that of the APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles, and the interaction decreases again for  $x = 4$ . Similar shift is observed for the XPS peak of Fe 2p<sub>3/2</sub> (from 711.5 eV of nanoparticle to 711.1–710.7 eV of films) and Fe 2p<sub>1/2</sub> (from 725.1 eV of nanoparticle to 724.6–724.2 eV of films, Figure 8B), displaying the charge distribution of encapsulated Fe<sub>3</sub>O<sub>4</sub> nanoparticles can also be affected by charge density out of APTS shell. Therefore, the tunable electrostatic interaction between LDH and APTS-

Fe<sub>3</sub>O<sub>4</sub> nanoparticles can be obtained by tuning  $x$ , which provides the tunability for the interparticle distance of Fe<sub>3</sub>O<sub>4</sub> nanoparticles, leading to tunable magnetic properties. Though here we investigate only integer values of  $x$ , non-integer values of  $x$  are also possible (for  $x < x_{\text{max}} = 4$ ), providing opportunities for fine control over the LDH surface charge.

**3.2.2. Tuning of the Thickness of LDH Nanosheets for Particle Spacing in the Vertical Direction.** Because of the multilayer feature of LDH, the originally synthesized LDH nanosheets can be further exfoliated to nanosheets with smaller thicknesses, which makes the tuning of the interparticle distance of Fe<sub>3</sub>O<sub>4</sub> nanoparticles in vertical direction possible. Figure 9A shows the thickness of original LDH nanosheet is  $\sim 15$  nm, which includes about 20 LDH single layers calculated by the ideal model (inset in Figure 9A). The Mg<sub>3</sub>Al-LDHs sheets with a different thickness ( $\sim 5, 10$ , and  $15$  nm) were obtained by exfoliating in formamide for different durations ( $t = 24, 5$ , and  $0$  h) before LBL assembly with APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles. The thickness of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>10</sub> films ( $t = 24, 5$ , and  $0$  h) is 200, 250, and 300 nm, respectively (Figure 9B–D). This is in agreement with the ideal single-layered arrangement model of the APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH supra-



**Figure 8.** (A) N 1s and (B) Fe 2p XPS spectra of (a) APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles and (b–e) the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 1, 2, 3$ , and  $4$ , respectively).



**Figure 9.** (A) Top view of SEM for unexfoliated Mg<sub>3</sub>Al-LDH nanosheets (inset: ideal model of LDH nanosheet, scale bar = 30 nm). (B–D) Side view of SEM for the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>10</sub> films ( $x = 3$ ,  $t = 24, 5$ , and  $0$  h, scale bar = 500 nm), respectively. (E) a representation for the structure of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>2</sub> films: (E<sub>1</sub>)  $x = 3$  and  $t = 24$  h, (E<sub>2</sub>)  $x = 3$  and  $t = 5$  h, and (E<sub>3</sub>)  $x = 3$  and  $t = 0$  h.

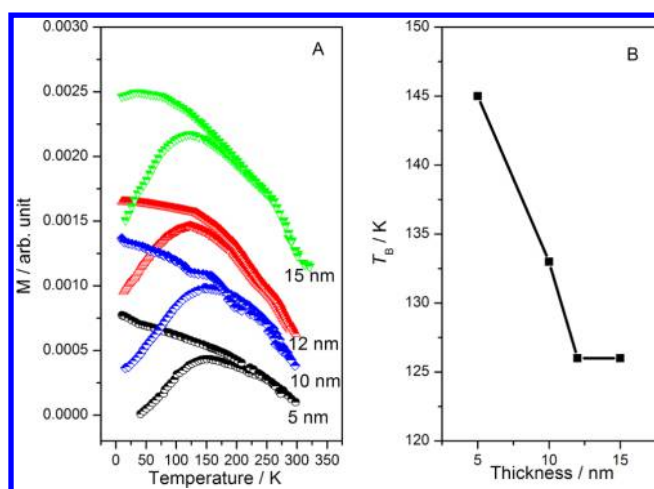
molecular structure with the thicknesses of  $\sim 5$ , 10, and 15 nm for different layers of LDH and  $\sim 14$  nm for the spherical APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles (Figure 9E). Therefore, the interparticle distance of adjacent APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles in the vertical direction can be tuned by assembling with a different thickness of LDH nanosheets.

Generally, the interparticle spacing will have a great influence on the collective magnetic behavior, which can be observed from temperature-dependent ZFC/FC magnetization.<sup>26</sup> Thus, the ZFC/FC magnetization of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>10</sub> films ( $x = 3$ , thickness of LDH layer is 0, 5, 10, 12, and 15 nm, respectively) has been tested at magnetic fields 10 Oe and shown in Figure 10. With increasing thickness (interparticle spacing in the vertical direction) of LDH layer, the maxima of the ZFC curves assigned to the mean blocking temperatures ( $T_B$ ), as well as the splitting points ( $T_{irr}$ ) between the ZFC and FC curves, shift toward lower temperatures.  $T_B$  decreased from 145 to 126 K as the interparticle spacing increased to 24 nm. A similar phenomenon was previously observed by García-Otero

in dilute dispersed magnetic NPs.<sup>31</sup> Along further increasing interparticle spacing from 24 to 29 nm, no change of  $T_B$  is observed, displaying the disappearance of interparticles interaction. It is noted that this value (24 nm) is lower but similar to a previously reported value (31.5 nm),<sup>26</sup> perhaps reflecting the fact that for loosely-packed particles between LDH sheets, the interparticle spacing is larger than just accounting for the LDH layer thickness.

The magnetic properties of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 3$ ,  $t = 24, 5$ , and  $0$  h) were studied by measuring the magnetic hysteresis curves at room temperature (Figure 11A–C). The (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 3$ ,  $t = 24, 5$ , and  $0$  h) display apparent magnetic anisotropy. The  $H_c$  in the parallel direction is more than 4-folds than that in the perpendicular direction, but the  $M_s$  in parallel is less than 75% of perpendicular (Table S2). The maximum  $M_s$  is present in the assembly with the thickest LDH ( $t = 0$  h) with 29 nm particle distance in the vertical direction, which implies the absence of (or negligible) interparticle interactions for nanoparticles





**Figure 10.** (A) FC (top) and ZFC (bottom) magnetization of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>10</sub> films ( $x = 3$ , thickness of LDH layer = 5, 10, 12, and 15 nm, respectively). (B)  $T_B$  as a function of thickness of LDH layer.

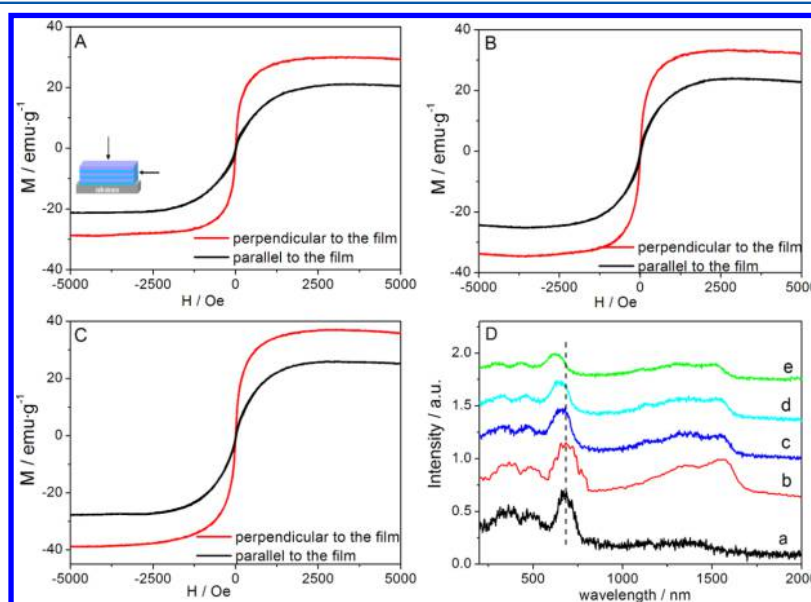
spacing beyond 24 nm. In the case of smaller particle spacing of 14–24 nm, the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 3$ ,  $t = 24$  and 5) can be referred to as a coupled particles system. This result shows interparticle spacing of Fe<sub>3</sub>O<sub>4</sub> has an effect on the collective magnetic behavior of the films, demonstrating the tunable magnetic anisotropy can be achieved by simply controlling the thickness of LDH nanosheets.

Raman spectra were measured to illustrate the effect of thickness change of LDH nanosheets on the magnetic property of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 3$ ,  $t = 24$ , 5, and 0 h). Figure 11D displays the Raman spectra of the Fe<sub>3</sub>O<sub>4</sub>, APTS-Fe<sub>3</sub>O<sub>4</sub>, and (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 3$ ,  $t = 24$ , 5, and 0 h) in the range 200–2000 cm<sup>-1</sup>. The peak located at 670 cm<sup>-1</sup> can be assigned to the A<sub>1g</sub> mode of Fe<sub>3</sub>O<sub>4</sub> (curve a).<sup>32–35</sup> After modifying by APTS molecules, the APTS-Fe<sub>3</sub>O<sub>4</sub> spectra show two new peaks at 1336 cm<sup>-1</sup> ( $\tau$  CH<sub>2</sub> and  $\omega$  CH<sub>2</sub>) and 1564

cm<sup>-1</sup> ( $\delta$  NH<sub>2</sub>) correspond to the APTS molecule, and indicating successful coating of APTS on the Fe<sub>3</sub>O<sub>4</sub> (curve b). For the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 3$ ,  $t = 24$ , 5, and 0 h), the blue-shift (from 670 to 645 cm<sup>-1</sup>) and reduced peak intensity compared to APTS-Fe<sub>3</sub>O<sub>4</sub> indicate that the intramolecular mobility (collisions, vibrations, and rotations, etc.) of Fe<sub>3</sub>O<sub>4</sub> is decreased. This result is caused by the host–guest-like interactions and the influence of changed charge density of LDH layer. Moreover, the gradually increased blue-shift (from 659 to 622 cm<sup>-1</sup>) and decreased intensity with increase of the thickness of LDH nanosheets (from 5 to 15 nm) demonstrate the increased immobilization and confining force of LDH nanosheet for Fe<sub>3</sub>O<sub>4</sub> nanoparticles. Therefore, LDH layer induces the ordered and oriented stack of Fe<sub>3</sub>O<sub>4</sub> nanoparticles in the interlayer of the LDH gallery, which imparts the anisotropic magnetic behavior for the films.

#### 4. CONCLUSION

In summary, the APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH films with high anisotropy and tunable magnetic properties were obtained by alternate assembly of LDH nanosheets and spherical APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles through the layer-by-layer (LBL) method on a quartz substrate. The relationship between particle spacing (in the horizontal and vertical directions) of the Fe<sub>3</sub>O<sub>4</sub> nanoparticles and magnetic properties of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>n</sub> films is discussed from the viewpoint of charge density and layer thickness of the LDH nanosheets, respectively. The maximum saturation magnetization value is present in the sample of  $x=3$  and  $t = 0$  ( $M_s$  is 36.3 and 25.1 emu g<sup>-1</sup> in perpendicular and parallel directions, respectively), which results from the excellent dispersion degree of APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles and charge density-matching assembly. XPS and Raman confirm that the electrostatic, immobilization, and confining forces between LDH and the APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles can be tuned. Both the interparticle distance of Fe<sub>3</sub>O<sub>4</sub> nanoparticles and the magnetic anisotropy of the films can be controlled. The work also opens up possibilities for further



**Figure 11.** (A–C) Magnetic hysteresis loops of the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 3$ ,  $t = 24$ , 5, and 0 h, respectively) at room temperature. (D) Raman spectra of (a) Fe<sub>3</sub>O<sub>4</sub> nanoparticles, (b) APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles, and (c–e) the (APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH)<sub>35</sub> films ( $x = 3$ ,  $t = 24$ , 5, and 0 h, respectively).



control. By using washing/assembly solutions of higher pH, one could adjust the charge of the APTS-Fe<sub>3</sub>O<sub>4</sub> NPs to more negative values. Similarly, ability to use non-integer ratios  $x = \text{Mg}/\text{Al}$  ratios in the LDH would allow finer control of the surface charge, and better optimization of both the spacing and magnetic properties. Therefore, this work not only provides a method to fabricate magnetic film with high anisotropy and tunability based on LDH nanosheets and APTS-Fe<sub>3</sub>O<sub>4</sub> nanoparticles but also systematically deepens the understanding of the correlation between structure and magnetic anisotropic property of the APTS-Fe<sub>3</sub>O<sub>4</sub>/LDH films. According to this strategy, other anisotropy materials and devices can be obtained by virtue of inherent anisotropy of LDH nanosheets and the variety of functional molecules and nanoparticles.

## ■ ASSOCIATED CONTENT

### Supporting Information

Magnetic parameters, SEM, and magnetic hysteresis loops. The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpcc.5b01065.

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### Notes

The authors declare no competing financial interest.

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