See discussions, stats, and author profiles for this publication at: https://www.researchgate.net/publication/244964717

# Fusion and Fragmentation Dynamics at Equilibrium in Triblock Copolymer Micelles

ARTICI F	in	MACROMO	I FCI II FS	<ul> <li>DECEMBER 201</li> </ul>	12

Impact Factor: 5.8 · DOI: 10.1021/ma3018298

CITATIONS	READS
5	29

### 1 AUTHOR:



Yahya Rharbi French National Centre for Scientific Research

**60** PUBLICATIONS **1,094** CITATIONS

SEE PROFILE

pubs.acs.org/Macromolecules

## 1 Fusion and Fragmentation Dynamics at Equilibrium in Triblock **2 Copolymer Micelles**

3 Y. Rharbi\*

5

10

11

12

13

14

15

16

17

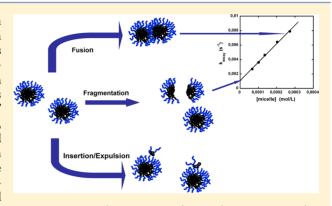
18

19

20

4 Laboratoire de Rhéologie et procédés, CNRS/UJF/INPG, UMR 5520, B.P.53, F-38041 Grenoble Cedex 9, France

ABSTRACT: Amphiphilic block copolymers autoassemble in water to form either dynamically active micelles or frozen particles at high surface tension. The dynamics of these systems is dominated by an individual process, which involves insertionexpulsion of copolymer chains, and a collective one, which involves fusion and fragmentation of proper micelles. The details of these mechanisms can drastically affect the micelles' morphology and some of their applications (drug delivery, template for mesoscopic structures, etc.). While fusion and fragmentation were found to be important in out-of-equilibrium kinetics such as sphere-to-rod transition, they were reported to be irrelevant at equilibrium by both theories and chain randomization experiments. We show, for the first time, that fusion and



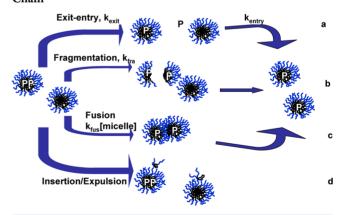
fragmentation do in fact take place at equilibrium in triblock copolymer micelles poly(ethylene oxide)—poly(propylene oxide) poly(ethylene oxide). This was achieved using a fluorescent technique, which probes the randomization of hydrophobic pyrene derivatives between micelles.

hen amphiphilic block copolymers are dissolved in aqueous solution, they form micelle aggregates with the 23 hydrophobic block in the core and the hydrophilic part in the 24 corona. Unlike small surfactants, which are dynamically active, 25 block copolymers exhibit slow to frozen kinetics.<sup>2</sup> Particularly 26 when the surface tension between blocks is high, they can be 27 trapped in metastable states without reaching their thermody-28 namic equilibrium.<sup>2</sup> The kinetic in surfactant micelles is 29 dominated by two mechanisms.<sup>3,4</sup> The first, described by 30 Anniasson and Wall (A-W), involves unimer/micelle inter-31 actions via insertion-expulsion of unimers.<sup>3</sup> The second 32 involves micelle-micelle interactions via fusion and fragmenta-33 tion. 5,6 Yet, the dynamics in block copolymer micelles differs 34 from the surfactant kinetics due to the chain correlation in the 35 core and the strong steric repulsion of the corona.<sup>2b</sup>

Halperin and Alexander predicted insertion-expulsion to be 37 the main dynamic process in diblock copolymers (Chart 1, path 38 d), whereas Dormidonto argued that fusion and fragmentation 39 (Chart 1, paths b and c) are favorable in the early stage of 40 micellization while unimer insertion-expulsion is the main 41 process at equilibrium.<sup>8</sup> Kinetic experiments at equilibrium 42 involving randomization of diblock copolymers between 43 micelles shows that insertion-expulsion is the main path for 44 the chain exchange. 2b Fusion and fragmentation of block 45 copolymer micelles were observed in experiments involving 46 morphological transition such as spheres-to-rod-like mi-47 celles. 9-11 Yet, to our knowledge there are no reports on the 48 fusion and fragmentation at equilibrium in block copolymers. In this work, we monitor the fusion and fragmentation 50 dynamic at equilibrium in the triblock copolymer poly(ethylene

51 oxide)-poly(propylene oxide)-poly(ethylene oxide) with

Chart 1. Various Processes for Exchanging the Copolymer Chains and Probes (P) between Micelles: (a) Exit-Entry of the Probe, (b) Fragmentation-Growth, (c) Fusion-Fragmentation, and (d) Insertion-Expulsion of Individual Chain



relatively large molecular weights (PEO<sub>17</sub>PPO<sub>60</sub>PEO<sub>17</sub>). We 52 use a fluorescent technique that was discovered few years ago 53 by some of the authors of this paper, which exploits the 54 randomization of hydrophobic pyrene derivatives (PyC<sub>18</sub>) 55 between micelles as a tool to probe the fusion and 56 fragmentation.<sup>6</sup> We show that fusion—fragmentation and 57

Received: September 5, 2012 Revised: October 9, 2012

Macromolecules Article

58 fragmentation—growth take place in the triblock copolymer 59 P103 micelle with very slow rates.

Triblock copolymer Pluronic P103 (PEO $_{17}$ PPO $_{60}$ PEO $_{17}$ ) 61 with  $M_{\rm w}=4.95$  kg/mol (BASF Corp.) was used as received. 62 Doubly deionized water was used in the preparation of the P103 solutions. The probe 1-pyrenyloctadecanone,  $C_{34}H_{44}O$  64 (PyC $_{18}$ ), was prepared via a Friedel—Crafts acylation of pyrene 65 with stearoyl chloride in dichloroethane and in the presence of 66 aluminum chloride (AlCl $_3$ ). The PyC $_{18}$  was solubilized in 67 P103 micelles by mixing P103 (20 g/L) with traces of PyC $_{18}$  at 68 85 °C. The solution was strongly agitated for 10 min with a 69 Vortex genie 2 model G 650 mechanical shaker at its maximum 70 frequency (>10 Hz). The solution was centrifuged at 5000 rpm 71 for 15 min at 25 °C to remove the nondissolved probe. 72 Fluorescence measurements were carried out with a Fluorolog 73 III (2-2) of Jobin Yvon spectrometer in the S/R mode.

Kinetic experiments were carried out by mixing a P103 (20 75 g/L) containing PyC<sub>18</sub> with a probe-free P103 solution (at different concentrations) in a 2 mm thick cell. The ratio of P103 containing PyC<sub>18</sub> to the probe-free P103 was 1/20. All 78 the measurements were carried out at 25 °C. The excitation 79 wavelength was 344 nm, and the emission was monitored every 80 30 s at  $\lambda_{\rm em}=480$  nm for the excimer and  $\lambda_{\rm em}=375.5$  nm for 81 the monomer.

At 25 °C, the P103 form micelles above the critical micelle so concentration (cmc = 0.7 g/L)<sup>13</sup> with a hydrodynamic radius  $R_h = 8 \text{ nm}^{14}$  and an aggregation number  $N_{\text{agg}}$  around 59. When the P103 solution is heated above the cloud point and cooled down to room temperature, the micelles dissolve randomly the hydrophobic probe  $PyC_{18}$ . In Figure 1 we present

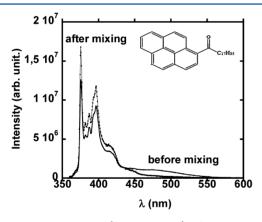


Figure 1. Emission spectra ( $\lambda_{\rm ex}=344~{\rm nm}$ ) of PyC $_{18}$  in aqueous solution of P103 micelles. The spectrum labeled "before mixing" refers to a solution of P103 20g/L containing PyC18. The spectrum "after mixing" refers to the solution obtained by mixing 0.05 mL of the full micelles with 1 mL of empty P103 micelles (20 g/L). Inset: molecular structure of the probe PyC $_{18}$ .

88 the fluorescence spectrum of P103 (20g/L) micelles containing PyC<sub>18</sub>. The spectrum has a broad excimer emission with a peak at 480 nm and monomer fluorescence at 375.5–400 nm. The existence of the excimer emission at 480 nm infers the presence of micelles bearing two or more PyC<sub>18</sub> molecules. When the P103 20 g/L with PyC<sub>18</sub> is mixed with an excess PyC<sub>18</sub>-free P103 solution, the spectrum evolves and shows a higher monomer emission and no discernible excimer band, which infers the distribution of PyC<sub>18</sub> among all the micelles yielding only micelles with one probe. The ratio of excimer-to-monomer intensities  $(I_{\rm F}/I_{\rm M})$  increases linearly with increasing the average

number of probes per micelle  $\langle n \rangle$  (not shown here):  $\langle n \rangle = 99$  [PyC<sub>18</sub>]/([P103] – cmc). This infers that PyC<sub>18</sub> undergo a 100 random Poisson distribution among the P103 micelles up to 101  $\langle n \rangle \approx 3$ . 102

In Figure 2, we show the result of a time-scan experiment in 103 f2 which we monitor the decrease in the excimer intensity ( $\lambda_{\rm em}$  = 104

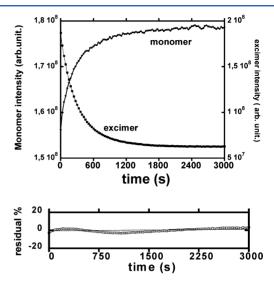
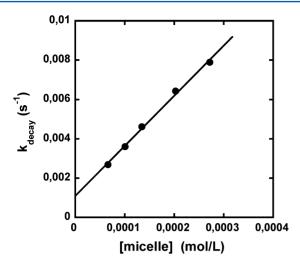


Figure 2. Time-scan experiment monitoring the decrease in the excimer emission ( $\lambda_{\rm em}=480$  nm) and the increase of the monomer emission ( $\lambda_{\rm em}=375.5$  nm) after mixing 1 mL of a P103 solution (20 g/L) with 0.05 mL of P103 solution (20 g/L) containing PyC<sub>18</sub>. The solid line represents the fit to a single-exponential expression. Inset: residuals from the exponential fit of the excimer decay.

480 nm,  $I_{\rm E}$ ) and the increase of the monomer intensity ( $\lambda_{\rm em}$  = 105 375.5 nm,  $I_{\rm M}$ ) following the mixing of P103 (20 g/L) 106 containing PyC<sub>18</sub> with P103 solution at 20g/L. Particularly 107 for the case of small  $\langle n \rangle$  ( $\langle n \rangle$  < 0.5) the  $I_{\rm E}$  and  $I_{\rm M}$  are 108 proportional to the fraction of micelles bearing two probes  $I_{\rm E} \propto 109$ P(t) and  $I_{\rm M} \propto P(t)$ . Unlike the probe exchange in common 110 surfactant micelles (Triton X100), which shows rigorous single 111 exponential behavior,  $^6$  the  $I_{\rm E}$  and  $I_{\rm M}$  decays exhibit a small but 112 noticeable deviation from the single exponential (Figure 2b). 113 This is most likely due to the effect of polymer polydispersity 114 on fusion and fragmentation. The relaxation time  $(\tau)$  from 115 the single exponential is similar to the average value  $\langle \tau \rangle$  116 calculated from the fit to two exponentials. When the kinetics is 117 repeated at different copolymer concentration, we observe a 118 strong dependence of the exchange rate  $k_{\rm decay} = 1/ au$  on the 119 concentration of empty micelles. In Figure 3 we show that  $k_{
m decay}$  120 f3 exhibits a linear dependence on the micelle concentration  $k_{
m decay}$  121 =  $k_1 + k_2$ [micelle], with [micelle] = ([P103] - cmc)/ $N_{agg}$ . The 122 linear dependence of  $k_{\text{decay}}$  vs [micelle] with a finite intercept 123 reflects the existence of two mechanisms: a first-order process 124 with a rate  $k_1$  independent of the empty micelles ( $k_1 = 1.07 \times 125$ 10<sup>-3</sup> s<sup>-1</sup>) and a second-order process with a linear dependence 126 of  $k_{\text{decay}}$  on [micelle]  $(k_2 = 25.5 \text{ s}^{-1} \text{ M}^{-1})$ .

The exchange of solutes between micelles shares the same 128 pathways with the exchange of copolymer chains (Chart 1). 129 The transfer of a probe molecule (P) from a full to an empty 130 micelle proceeds through three main paths (Chart 1): (a) one 131 probe exits the full micelle to the aqueous phase and then 132 enters an empty one, (b) the full micelles fragment into two 133 micelles, each bearing one probe followed by the growth of the 134 fragments via insertion of copolymer chains or fusion with an 135

Macromolecules Article



**Figure 3.** Relaxation rate  $k_{\text{decay}}$  calculated from the fits of the exchange decays of PyC<sub>18</sub> in P103 to the single exponential, plotted against the concentration of empty micelle.

136 empty micelle, and (c) the full micelles can fuse with empty 137 micelles to form large ones that break into two proper micelles 138 containing each one probe. The exit—entry and fragmentation—139 growth are first-order processes, which should lead to a 140 constant exchange rate  $k_{\rm exit}$  and  $k_{\rm fra}$ , respectively. The fusion—141 fragmentation is a bimodal process, which leads to second-142 order kinetics. In the case where the concentration of empty 143 micelles is much larger than the full ones, the fusion—144 fragmentation gives pseudo-first-order kinetics with a rate 145  $k_{\rm fus}$ [micelle]. The exchange rate can be written as  $k_{\rm decay} = k_{\rm exit} + 146$   $k_{\rm fra} + k_{\rm fus}$ [micelle]. One can also imagine other side processes 147 that are not considered in this analysis such as fragmentation of 148 a full micelle in two fragments followed by their fusion. Micelles 149 can also collide, adhere, exchange their contents, and fragment 150 without fusing.

The rate liming for the probe exit-entry is either the water 152 solubility  $(C_w)$  or the diffusion through the core/corona. In the 153 case, where water solubility is the dominant barrier, the exit rate 154 can be estimated from partitioning equilibrium  $k_{\text{exit}} = k_{\text{entry}} C_{\text{w}} /$ 155  $n_{\rm m}$ , where  $n_{\rm m}$  is the average number of probes per micelle at 156 equilibrium. The diffusion through the viscous PPO core 157 should give  $k_{\text{entry}}$  smaller than the diffusion controlled rate in 158 water  $(k_{\text{entry}} < 3 \times 10^9 \text{ M}^{-1} \text{ s}^{-1})$ . Probe solubility in P103 159 micelle infers that  $n_{\rm w} > 1$ .  $C_{\rm w}$  of PyC<sub>18</sub> is too small to be 160 detected easily but can be estimated using the energy for the 161 transfer of one methylene group from water to the micelles  $(\Delta \mu_{\rm CH_2})$ ,  $\ln(C_{\rm w}) = \ln(C_{\rm w}^0) - N\Delta \mu_{\rm CH_2}/RT$ , where  $C_{\rm w}^0$  can be 163 taken as the water solubility of 1-acetylpyrene (N = 2) ( $10^{-6}$  $_{164}$  mol/L). The different literature values of  $\Delta\mu_{\rm CH}$ , give  $C_{\rm w} \approx 4.3$  $165 \times 10^{-15} \text{ mol/L}, ^{18} C_{\text{w}} \approx 4.7 \times 10^{-17} \text{ mol/L}, ^{19,20} \text{ and } C_{\text{w}} \approx 2.5 \times 10^{-15} \text{ mol/L}, ^{18} C_{\text{w}} \approx 2.5 \times 10^{ 166 \ 10^{-16} \ mol/L.^{21}$  Even if we omit the role of the core viscosity on 167  $k_{\rm entry}$  ( $k_{\rm entry} \approx 3 \times 10^9~{\rm M}^{-1}~{\rm s}^{-1}$ ) and we take  $n_{\rm m} = 1$ , we find  $k_{\rm exit}$  168  $1.3 \times 10^{-5}$ ,  $7.5 \times 10^{-6}$ , and  $1.4 \times 10^{-7}~{\rm s}^{-1}$ , which is 2–4 orders 169 of magnitude lower than the measured exchange rate. 170 Moreover, an upper limit for  $k_{\rm exit}$  can be estimated from the 171 exchange of PyC<sub>18</sub> between spherical micelles of sodium 172 dodecyl sulfate (SDS) because (i) fragmentation and fusion are 173 extremely slow in SDS in the absence of added salt, 23 (ii) the 174 exchange is dominated by water solubility, 23 and (iii) the exit 175 rate in SDS is expected to be faster than in P103 since the SDS 176 core is smaller and less viscous than PPO core. 15,22,23 When the

exchange of  $PyC_{18}$  is carried out in SDS in absence of salt, 177 following the same procedure as in ref 23, we found the 178 exchange rate to be negligible compared to that in P103, which 179 suggests that  $k_{\rm exit}$  can be safely neglected in P103. One might 180 imagine an alternative exit mechanism where the probe exit is 181 assisted by chain expulsion ( $k^-$ ). When the P103 solution 182 containing  $PyC_{18}$  is diluted below the cmc, the  $PyC_{18}$  does not 183 dissolve in the PPO of the free chains but rather forms large 184 aggregates. This confirms that the PPO of the free chains 185 cannot solubilize the  $PyC_{18}$ . This result rejects the model of 186 probe exit assisted by chain expulsion. Therefore, the first-order 187 process involves mainly the fragmentation—growth mechanism. 188

The energy barrier to fragmentation is estimated from the 189 combination of surface tension energy and core elastic energy 190 as  $E_{\rm fission} \sim (N_{\rm PPO})^{2/3} (N_{\rm agg})^{2/3} x^{2/3}$ , where  $N_{\rm PPO}$  is the chain 191 length of the PPO and  $x = N_{\rm agg}^1/N_{\rm agg}$ , with  $N_{\rm agg}^1$  is the size of the 192 fragment. This favors the expulsion rate of single chains  $N_{\rm agg}^1 = 193$  1. In the case of P103, the fragmentation rate measured with 194 PyC<sub>18</sub> is less than  $10^{-6}$  the estimated value of the expulsion rate 195 of single chains  $(k^- \approx 2000 \ {\rm s}^{-1}),^{24} k_{\rm fra}/k^- \approx 10^{-6}$ . However this 196 argument is not sufficient to exclude the contribution of the 197 fission in copolymer dynamics. For example, even in a system 198 of small nonionic surfactant like Triton X-100, where fusion 199 fragmentation dominates several aspects of their dynamics, the 200  $k_{\rm fra}/k^-$  is similar to that reported here  $k_{\rm fra}/k^- = 5.5 \times 10^{-6}$ .

The second-order process  $k_2$  is likely to be dominated by 202 fusion-fragmentation, which involves several steps: collision of 203 a full and an empty micelle, adhesion of these micelles, fusion 204 of the two micelles to form a large one, exchange of the solute 205 within the large micelle, and fragmentation of the large micelle 206 into two proper micelles containing one probe each. Because 207 the diffusion-controlled rate is more than 109 times the 208 measured  $k_2$ , the fusion-fragmentation process cannot be 209 dominated by the collision step. The second-order fusion rate is 210 found to be  $25.5 \text{ s}^{-1} \text{ M}^{-1}$ . Because fusion can either yield solute  $_{211}$ exchange or not,  $k_{\text{fus}} = 2k_2 = 51 \text{ s}^{-1} \text{ M}^{-1}$ . The linear increase of 212  $k_{\rm obs}$  vs [micelle] infers that fragmentation of the large aggregate 213  $2N_{\rm agg}$ , resulting from the fragmentation of two micelles is much 214 faster than  $k_{\rm fra}$  of proper micelles. This is expected since the 215 energy resulting from the fission of micelles of size  $2N_{\rm agg}$  is 216 negative.<sup>8</sup> If the fission rate of  $2N_{agg}$  were similar to that of the 217 proper micelle  $N_{\rm agg}$ , the  $k_{\rm obs}$  will level off at high concentration. 218 Thus, the measured  $k_2$  describes the rate of fusion  $k_{\text{fus}}$ . It has 219 been shown in Triton X-100 and synperonic surfactants that 220 fusion rate is independent of the polarity of the probe,6 which 221 infers that the second-order fusion rate reflects the rate of 222 fusion. The energy barrier to fusion is the coronal energy 223 resulting from the steric repulsion or the elastic energy of the 224 corona. In the case studied here of crew-cut micelle (short 225 corona), the elastic energy barrier is described as  $E_{\rm fusion} \sim$  226  $N_{\rm PEO}^{2/3}/N_{\rm PPO}(N_{\rm agg})^{2.7}$  The barrier energy to insertion is  $E_{\rm fusion}$  227  $\sim N_{\rm PEO}/(N_{\rm PPO})^{-4/9}(N_{\rm agg})^{2/9}$ , which makes fusion less probable 228 than insertion. The measured fusion rate is found to be less 229 than  $10^{-6}$  times the expulsion rate  $(k^+ \approx 5 \times 10^6 \text{ s}^{-1} \text{ M}^{-1}).^{24}$  230 Yet the fusion process still controls several aspects of the P103 231 dynamic such as the sphere-to-rod transition. 11

In this article we show that fusion and fragmentation take 233 place between proper micelles at equilibrium in the triblock 234 copolymer PEO<sub>17</sub>PPO<sub>60</sub>PEO<sub>17</sub>with a rate 10<sup>6</sup> slower than the 235 rate of chain expulsion and insertion.

Macromolecules Article

#### AUTHOR INFORMATION

#### Corresponding Author

\*E-mail rharbi@uif-grenoble.fr.

#### ACKNOWLEDGMENTS

- 241 This work was supported by the joint program ECOS-Nord
- 242 M05P02E, MO6-P03 of the ministry of research and education
- 243 (France). We acknowledge Prof Armando Soltero for providing
- 244 the copolymer. We acknowledge M. Karrouch, E. Fevre, F.
- 245 Hugenel, and Dr. Hélène Galliard for their technical support.

#### REFERENCES 246

- (1) (a) Hamley, W. In Block Copolymers in Solution; Hamley, I. W., 247 248 Ed.; John Wiley & Sons: San Francisco, 2005. (b) Amphiphilic Block 249 Copolymers: Self-Assembly and Applications; Alexandridis, P., Lindmann, 250 B., Eds.; Elsevier: Amsterdam, 2000.
- (2) (a) Nicolai, T.; Colombani, O.; Chassenieux, C. Soft Matter 2010, 2.51
- 252 6, 3111. (b) Lund, R.; Willner, L.; Richter, D. Macromolecules 2006,
- 253 39, 4566. (c) Won, Y. Y.; Davis, H. T.; Bates, F. S. Macromolecules 254 2003, 36, 953-955. (d) Johnson, B. K.; Prud'homme, R. K. Phys. Rev.
- 255 Lett. 2003, 91, 118302.
- (3) (a) Aniansson, E. A. G.; Wall, S. N. J. Phys. Chem. 1974, 78, 1024;
- 257 1975, 75, 857. (b) Aniansson, E. A. G.; Wall, S. N.; Almgren, M.;
- 258 Hoffmann, H.; Kielmann, H.; Ulbricht, W.; Zana, R.; Lang, J.; Tondre, 259 C. J. Phys. Chem. 1976, 80, 905. (c) Wall, S. N.; Aniansson, E. A. J.
- 260 Phys. Chem. 1980, 84, 727.
- (4) (a) Kahlweit, M. J. Colloid Interface Sci. 1982, 90, 92. (b) Lessner, 262 E.; Teubner, M.; Kahlweit, M. J. Phys. Chem. 1981, 85, 3167.
- (5) Waton, G.; Michels, B.; Zana, R. Macromolecules 2000, 34, 907. 263
- (6) (a) Rharbi, Y.; Li, M.; Winnik, M. A.; Hahn, K. G. J. Am. Chem. 264
- 265 Soc. 2000, 122, 6242. (b) Rharbi, Y.; Winnik, M. A.; Hahn, K. G.
- 266 Langmuir 1999, 15, 4697. (c) Rharbi, Y.; Bechthold, N.; Landfester,
- 267 K.; Salzman, A.; Winnik, M. A. Langmuir 2003, 19, 10.
- (7) Halperin, A.; Alexander, S. Macromolecules 1989, 22, 2403. 2.68
- (8) Dormidontova, E. E. Macromolecules 1999, 32, 7630.
- (9) Burke, S. E.; Eisenberg, A. Langmuir 2001, 17, 6714. 270
- (10) Denkova, A. G.; Mendes, E.; Coppens, M.-O. J. Phys. Chem. B 2.71 272 **2009**, 113, 989.
- (11) Landazuri, G.; Fernandez, V. V. A.; Soltero, J. F.; Rharbi, Y. J. 274 Phys. Chem. B 2012, 116, 11720.
- (12) Friedel-Crafts and Related Reactions; Olah, G., Ed.; Wiley-275 276 Interscience: New York, 1963; Vol. 3, p 78.
- (13) Alexandridis, P.; Holzwarth, J. F.; Hatton, T. A. Macromolecules 278 1994, 27, 2414.
- (14) Fernandez, V. V. A.; Soltero, J. F. A.; Puig, J. E.; Rharbi, Y. J. 280 Phys. Chem. B 2009, 113, 3015.
- (15) Kadam, Y.; Yerramilli, U.; Bahadur, A.; Bahadur, P. Colloids Surf., 282 B **2011**, 83, 49
- (16) Infelta, P. P.; Gratzel, M. J. Chem. Phys. 1979, 70, 179. 283
- (17)Hilczer, M.; Barzykin, A. V.; Tachiya, M. Langmuir 2001, 14, 284 4196 285
- 286 (18) Kozlov, M. Y.; Melik-Nubarov, N. S.; Batrakova, E. V.; Kabanov, 287 A. V. Macromolecules 2000, 33, 3305-3313.
- (19) Tolls, J.; van Dijk, J.; Verbruggen, E. J. M.; Hermens, J. L. M.; 288
- 289 Loeprecht, B.; Schuurmann, G. J. Phys. Chem. A 2002, 106, 2760-290 2765.
- (20) Ferguson, A. L.; Debenedetti, P. G.; Panagiotopoulos, A. Z. J. 291 292 Phys. Chem. B 2009, 113, 6405-6414.
- (21) Taisne, L.; Walstra, P.; Cabane, B. J. Colloid Interface Sci. 1996, 294 184, 378
- (22) Nivaggioli, T.; Tsao, B.; Alexandridis, P.; Hatton, T. A. Langmuir 295 296 1996, 11, 119.
- (23) Rharbi, Y.; Winnik, M. A. J. Phys. Chem. B 2003, 107, 1491.
- (24)  $k^-$  and  $k^+$  are extrapolated from Figures 1 and 2 of the reference:
- 299 Zana, R.; Marques, C.; Johner, A. Adv. Colloid Interface Sci. 2006, 345, 300 123-126.