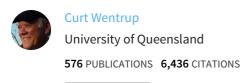
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Iminopropadienones, RN:C:C:C:O. Theory and experiment

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Iminopropadienones, RN=C=C=C=O: Theory and Experiment

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Ab initio molecular orbital calculations at the MP2/6-31G* level have been used to examine the structures and infrared spectra of a new class of compounds, the iminopropadienones, RN=C=C=C=O (R = H, CH₃, and Ph). The agreement between calculated and experimental IR spectra of PhNCCCO, Ph¹⁵NCCCO, and PhNCCC-¹³CO is excellent. Inclusion of electron correlation is found to be important for the correct prediction of the relative intensities of the cumulenic stretching vibrations. All three iminopropadienones are predicted to have a slightly bent NCCCO backbone (\angle CCC \approx 176°). As with carbon suboxide, these cumulenes are calculated to have an extremely flat CCC bending potential. The parent compound, HNCCCO (4a) is calculated to be thermodynamically stable toward dissociations into (i) HNCC + CO, (ii) HNC + CCO, (iii) HN + CCCO, and (iv) H°+NCCCO°. Rearrangement of 4a to the more stable cyanoketene isomer requires a sizeable barrier of 402 kJ mol⁻¹ [G2(MP2)]. The calculated stability of 4a is consistent with its experimental observation in neutralization—reionization mass spectrum. The adiabatic ionization energy and heat of formation of HNCCCO are predicted to be 9.51 eV and 175 kJ mol⁻¹, respectively.

1. Introduction

The first experimental evidence for the existence of iminopropadienones, RN=C=C=C(4), was reported recently.1 The phenyl derivative 4c was obtained from three different precursors (1-3) in flash vacuum pyrolysis (FVP) reactions, whereby 1 and 2 gave particularly good yields (Scheme 1). 4c was characterized by chemical trapping reactions. Both the methyl (4b) and phenyl (4c) derivatives have been observed by IR1 and mass spectrometry,2 and the unsubstituted compound, HN=C=C=C (4a), has been produced by dissociative ionization of methyl 5-aminoisoxazole-4-carboxylate in a tandem mass spectrometer and characterized by neutralization-reionization mass spectrometry.² Further work on the characterization of the related bisimines, RN=C=C=CNR', is currently in progress.³ All these compounds can be considered as imines of carbon suboxide, O=C=C=O, a reactive but isolable molecule.4

In this paper, we report high-level ab initio theoretical calculations on the structures and infrared spectra of these cumulenes. ¹⁵N- and ¹³C-labeled isotopomers of the phenyl derivative (4c) have been generated, and their IR spectra are in excellent agreement with the theoretical predictions. In addition, we have examined the stability with respect to rearrangement and fragmentation, ionization energy and thermochemistry for the parent compound, HN=C=C=C=O (4a).

2. Theory

Method and Results. Standard ab initio molecular orbital calculations⁵ were carried out with the GAUSSIAN 92 series of programs.⁶ Geometry optimizations for 4a-c were carried out with the split-valence polarized 6-31G* basis set⁵ at the Hartree-Fock (HF) and second-order Møller-Plesset perturbation (MP2) levels.⁵ Harmonic vibrational wavenumbers and infrared intensities were predicted at these equilibrium geometries using analytical second derivatives.⁷ The directly calculated HF/6-31G* wavenumbers were scaled by a factor of 0.89 to account for the overestimation of vibrational wavenumbers at this level of theory.⁸ For the unsubstituted iminopropadienone (4a) and related species, higher-level energy calculations were computed using the Gaussian-2 [G2(MP2)] theory,⁹ which is a composite

SCHEME 1

procedure based effectively on QCISD(T)/6-311+G(3df,2p)//MP2/6-31G* energies together with zero-point vibrational and isogyric corrections. Unless otherwise noted, these are the values referred to in the text. To investigate the bending potential function of HNC₃O (4a), we have carried out geometry optimizations at the HF and MP2 levels for a number of bent structures corresponding to various values of the CCC angle. These were performed with the 6-31G* basis set. Spin-restricted calculations were used for closed shell systems and spin-unrestricted for openshell systems. The frozen-core approximation was employed for all correlation calculations.

The results of geometry optimizations for the iminopropadienones (4a-c) are summarized in Tables 1 and 2. Figure 1 displays the planar (4a) and nonplanar, bent (4a') structures of HNC₃O.

^{*} Abstract published in Advance ACS Abstracts, December 15, 1993.

TABLE 1: Optimized Structural Parameters for Iminopropadienones (RN=C=C=C)

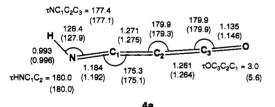
		HF/6-310	G*		MP2/6-31	G*
${\tt parameter}^b$	R = H	$R = CH_3$	$R = C_6H_5$	R = H	$R = CH_3$	$R = C_6H_5$
r(R-N)	0.996	1.420	1.382	1.011	1.437	1.390
$r(N-C_1)$	1.192	1.156	1.161	1.222	1.213	1.218
$r(C_1-C_2)$	1.275	1.296	1.291	1.287	1.292	1.288
$r(C_2-C_3)$	1.264	1.253	1.256	1.278	1.276	1.278
$r(C_3-O)$	1.146	1.155	1.153	1.188	1.192	1.190
$\angle RNC_1$	127.9	179.9	179.9	131.0	144.0	144.6
$\angle NC_1C_2$	175.1	180.0	180.0	171.4	172.7	173.2
$\angle C_1C_2C_3$	179.3	179.8	180.0	175.6	175.7	177.0
∠C ₂ C ₃ O	179.9	180.0	180.0	179.6	179.5	179.6
τRNC_1C_2	180.0	176.5	176.9	180.0	179.4	178.4
$\tau NC_1C_2C_3$	177.1	174.2	168.1	180.0	179.5	179.8
$\tau OC_3C_2C_1$	5.6	178.5	177.2	180.0	177.2	170.1

^a Bond lengths are given in angstroms and bond angles in degrees. ^b Atom labels are illustrated in Figure 1.

TABLE 2: Calculated Structural Parameters² and Rotational Constants^b for HNC₃O (4a)^a

2/ MP2/ G(2d) 6-311+G(2d	QCISD(T)/ lf.2p) 6-31G*
	11,2p) 0-31G
Geometry	
2 1.004	1.015
7 1.210	1.228
3 1.281	1.290
1.274	1.281
1.174	1.188
133.3	125.8
173.2	171.3
178.9	174.2
180.0	179.4
180.0	178.5
179.7	163.5
3 179.2	179.9
Rotational Constants	
.41 1133.02	814.61
58 2.1783	2.1471
24 2.1741	2.1416
	Geometry 1.004 1.210 1.210 1.281 1.274 1.174 1.133.3 1.73.2 1.78.9 1.80.0 1.80.0 1.80.0 1.79.7 1.79.2 Rotational Constants 1.41 1.133.02 1.41 1.133.02

^a Bond lengths are given in angstroms and bond angles in degrees. ^b In gigahertz.



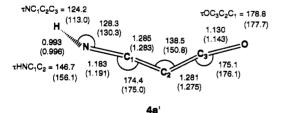


Figure 1. Optimized geometries (HF/6-311+G(2df,2p), with HF/6-31G* values in parentheses) of planar (4a) and bent (4a') structures of HNC₃O.

Calculated bending potential functions for HNC₃O is presented in Figure 2. Calculated charges (MP2/6-31G*), based on natural bond orbital (NBO) analysis, 10 of the iminopropadienones are given in Table 3. The G2(MP2) total and relative energies of 4a and related species are given in Table 4. Optimized MP2/ 6-31G* structures for cyanoketene (9), transition structure for the rearrangement of 4a to 9 (10) and iminopropadienone radical cation (HNC₃O⁺, 11) are shown in Figure 3. Vibrational wavenumbers and infrared intensities of the iminopropadienones

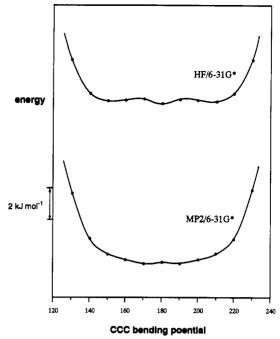


Figure 2. Calculated bending potential functions for HNC₃O.

TABLE 3: Calculated NBO Charges of Iminopropadienones $(RN=C=C=C=O)^s$

R = H (4a)	$R = CH_3 (4b)$	R = Ph (4c)
0.435	0.252	0.203
-0.613	-0.381	-0.374
0.440	0.434	0.438
-0.482	0.484	-0.463
0.629	0.607	0.611
-0.408	-0.428	-0.415
	0.435 -0.613 0.440 -0.482 0.629	0.435 0.252 -0.613 -0.381 0.440 0.434 -0.482 -0.484 0.629 0.607

^a MP2/6-31G* wave function. ^b Atom labels are illustrated in Figure

TABLE 4: Calculated Total Energies^a (hartrees) and Relative Energies^a (kJ mol⁻¹) of HNC₃O (4a) and Related Species

		k	
state	sym	total energy	rel energy
¹ A′	Cs	-244.448 67	0.0
$1\Sigma^{+}$	C∞v	-244.443 51	13.5
$^{1}A'$	C_s	-244.474 71	-68.4
^{1}A	C_1	-244.295 55	402.0
2A"	C_s	-244.099 08	917.9
2A"	C_s	-244.090 82	939.5
		-244.361 52	228.8
		-244.322 95	330.1
		-243.311 13	361.1
		-244.273 00	461.2
		-244.188 73	682.5
	¹ A' 1Σ ⁺ ¹ A' ¹ A 2A"	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$

^a G2(MP2) E₀ energies. ^b Corresponding to vertical ionization of HNC₃O. c Note that the lowest-energy singlet state of C₂O corresponds to a $^{1}\Delta$ state. The energy difference between $^{1}\Sigma^{+}$ and $^{1}\Delta$ states is 25.5 kJ mol-1 (complex MP2/6-31G*).

and their ¹⁵N and ¹³C isotopomers are presented in Table 5. The normal vibrations calculated for HNC₃O at the MP2/6-31G* level are illustrated in Figure 4. Comparison of MP2/3-21G and MP2/6-31G* cumulenicic stretching wavenumbers and IR intensities are given in Table 6. Finally, comparison of calculated and experimental cumulenic CO stretching wavenumbers are summarized in Table 7.

Equilibrium Structures of RNC₃O. To the best of our knowledge, there is no theoretical calculation reported on the structures of the iminopropadienones. The parent species, HNC₃O (4a), is an isoelectronic analogue of carbon suboxide (O=C=C=O), of which the spectra have been studied extensively in the millimeter-wave, far-infrared, mid-infrared, and ultraviolet regions.4 Strong experimental evidence has

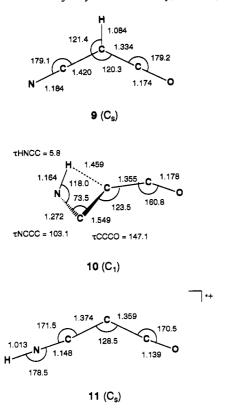


Figure 3. Optimized geometries (MP2/6-31G*) of cyanoketene (9), transition structure for the rearrangement of 4a to 9 (10), and HNC₃O⁻⁺ radical cation (11).

indicated that C₃O₂ is a semirigid bender, with a central CCC bond angle of 156°.4b,c Thus, the molecular structures of the iminopropadienones are of considerable interest in terms of their bending potentials.

To investigate the possible conformations of iminopropadienone, geometry optimizations for HNC₃O (4a) were carried out systematically with a sequence of basis sets, namely, 3-21G, 6-31G*, 6-311G*, 6-31+G*, 6-31G(2d), 6-31G(df), and 6-311+G(2df,2p) basis sets,8 at the Hartree-Fock level. The effect of electron correlation was investigated at the MP25 and QCISD(T)¹¹ levels, using the 6-31G* basis set. The calculated structure of HNC₃O is very sensitive to the theoretical level employed. With the smallest basis set, 3-21G, there is only one stable equilibrium structure, with a linear HNCCCO framework. At the HF/6-31G* level, there exists two stable conformations, the planar, essentially linear (4a) and the nonplanar, bent (4a') structures (Figure 1). Both structures have almost identical energy. The central CCC bond angle of the bent structure (4a') is 151°, similar to that of C₃O₂.4b,c Both conformations are true local minima (with all real frequencies) on the HNC₃O potential energy surface. Expansion of the basis set from double-\(\zeta \) valence to triple- ζ valence (6-31G* \rightarrow 6-311G*) and inclusion of f functions on heavy atoms (6-31G* \rightarrow 6-31G(df)) have little effect on the structures and relative energies of 4a and 4a'. Inclusion of diffuse sp functions on heavy atoms (6-31G* \rightarrow 6-31+G*) favors the "planar" structure. In fact, at this level of theory (HF/6-31+G*), the "planar" structure (4a) is the only stable equilibrium structure of HNC₃O. On the other hand, expansion of the basis set from one set of d-polarization functions to two $(6-31G^* \rightarrow 6-31G(2d))$ strongly favors the bent structure (4a'), with a significant increase of the degree of bending (∠CCC = 132°). 4a' is the only stable HNCCCO structure at the HF/ 6-31G(2d) level. With the largest 6-311+G(2df,2p) basis set, both 4a and 4a' exist and the energy difference between the conformations is just 3 kJ mol⁻¹, in favor of the bent structure. Inclusion of electron correlation at the MP2 level strongly favors the "planar" conformation. No stable bent conformation of HNC₃O (4a') was found at the MP2/6-31G* level. This is confirmed by a larger-basis-set [6-31+G(2d)] optimization at the MP2 level (Table 2). To establish a more definitive theoretical estimate of the molecular structure of the cumulene, we have carried out geometry optimizations of 4a and 4a' at the MP2/ 6-311+G(2df,2p) and QCISD(T)/6-31G* levels (Table 2). As with the MP2/6-31G* and MP2/6-31+G(2d) levels, the "planar" structure (4a) is the only stable form of HNCCCO. In other words, the improvement of basis set $[6-31+G(2d) \rightarrow 6-311+$ G(2df,2p)] as well as higher-level treatment of electron correlation [QCISD(T)] favor the "planar" structure. This result lends confidence to our prediction of the molecular structures of the iminopropadienones at the MP2/6-31G* level. Calculated rotational constants of 4a are given in Table 2.

Geometry optimizations were carried out for all three derivatives of iminopropadienones (4a-c) with the 6-31G* basis set at both the Hartree-Fock and MP2 levels (Table 1). At the HF level, the iminopropadienones are predicted to be essentially planar, with almost linear CCC angles. However, inclusion of electron correlation (MP2) increases significantly the degree of bending for all three derivatives (by 3-4°). A similar finding has been reported for the structurally related species, propadienone (CH₂=C=C).¹² It is worthwhile to note that the calculated structural parameters for propadienone^{10a} at the MP2/6-31G* level agree well with experimental values. 13 At the MP2/6-31G* level, the iminopropadienones are predicted to be have a slightly bent NCCCO framework (\angle CCC = 176–177°). For the parent compound (4a), we have examined the CCC bending potential functions, obtained by optimizing a number of bent structures, with a range of CCC angles, at both the HF/6-31G* and MP2/ 6-31G* levels. The result (Figure 3) clearly indicates that the CCC bending potential is extremely flat. Constraining the HNCCCO skeleton to be linear (i.e., ${}^{1}\Sigma^{+}$ state) increases the energy by 14 kJ mol-1 (Table 4). Thus, the calculated barrier to completely linearity is 14 kJ mol⁻¹. This result suggests that the HNC bending potential is significantly steeper than the CCC curve. A similar finding has been reported for the quasi-linear molecule HNCO.14 In summary, the iminopropadienones are best considered as quasi-linear molecules.

Interestingly, all three iminopropadienones display rather short C=C bond lengths (1.276-1.292 Å at the MP2/6-31G* level), similar to that of C₃O₂ (1.279 Å, experimental estimate). ^{4e,f} For comparison, the experimental C=C bond length in ethylene, ketene, and propadienone are 1.339, 1.314, and 1.322 Å, respectively. 15 For the methyl (4b) and phenyl (4c) derivatives, strong delocalization of the π orbitals are predicted. The (Me)C-N and (Ph)C-N bond lengths (1.437 and 1.390 Å, respectively) are surprisingly short.

Charge Distributions. We have examined the charge distributions of the iminopropadienones using natural bond orbital (NBO) analysis¹⁰ (MP2/6-31G*). As can be seen in Table 3, the NCCCO skeletons of all three iminopropadienones display an alternating charge distribution. It thus appears that the relatively short C=C bond lengths (1.276-1.292 Å) and the stabilities of the iminopropadienones may in part be due to strong electrostatic attraction. The large negative charge calculated for the terminal nitrogen atom suggests that an electronwithdrawing substituent will have a strong effect at the N terminus. The calculated dipole moments of 4a-c at the MP2/ 6-31G* level are 3.02, 4.61, and 4.06 D, respectively. In view of their large dipole moments, the iminopropadienones are predicted to be significantly stabilized in solution. The charge distributions of these cumulenes also reveals that protonation at the central carbon (C₂) is highly favorable. This is readily confirmed by a recent combined experimental and theoretical study of protonated iminopropadienones.¹⁶

Stability of HNC₃O. To assess the stability of the unsubstituted iminopropadienone (4a), several fragmentation and rearrangement pathways have been examined. We have considered the dissociation reactions of 4a to CO, HNC, HN, and H. (Table 4). The lowest energy fragmentation corresponding to the loss

TABLE 5: Calculated IR Spectra of Iminopropadienones (RN=C=C=C=O) and Their Isotopomers

		wave	number ^b			int	ensity	-			wave	number ^b			in	tensity ^c	
	RN	C ₃ O	R ¹⁵ N-	RNC ₂ -	RN	C ₃ O	R15N-	RNC ₂ -		RN	C ₃ O	R15N-	RNC ₂ -	RN	IC₃O	R15N-	RNC ₂ -
	HF	MP2	C_3O^d	$^{13}CO^d$	HF	MP2	C ₃ O ^e	13COe		HF	MP2	C_3O^d	¹³ CO ^d	HF	MP2	C ₃ O ^e	13COe
								HNC ₃	O (4a))							
ν_1	3885	3695	-10	0	303	288	-10	1	ν7	688	568	0	-16	72	30	0	-2
ν_2	2451	2446	-1	-30	4237	2538	-18	-55	ν_8	646	551	0	-11	57	16	1	-5
ν_3	2368	2204	-9	-30	1298	21	3	-20	ν_9	629	505	-1	-5	44	10	1	4
ν_4	1760	1595	-12	-3	116	99	0	7	ν_{10}	523	498	-1	0	315	63	-3	0
ν ₅	843	785	-10	-4	6	9	-2	1	ν_{11}	108	114	-1	0	4	3	0	0
ν_6	748	683	-4	0	578	647	1	0	ν_{12}	69	80	-1	0	4	0	0	0
	2212	2222	•	0	1.0	0	0	CH ₃ NC			1007	0		•	0	•	^
ν_1	3313	3223 3197	0 0	0 0	16 16	9 10	0 0	0	ν_{12}	1151	1086 664	-9 -5	0 -3	3 69	8 7	0	0
ν_2	3313 3233	3197	0	0	40	50	0	0 1	ν_{13}	665 665	557	-5 0	-3 -15	69	22	0	0
ν3	2568	2453	-3	-26	2023	3131	-47	- 7 9	ν ₁₄	657	557 541	0	-15 -15	69 45		0	0 -2
ν4	2434	2245	-3 -15	-26 -31	5062	22	11	-19 -13	ν ₁₅	603	509	-2	-13 -1	43 66	11 59	-1	-2 -1
ν ₅	1849	1669	-13 -16	-51 -5	45	146	5	12	ν ₁₆	603	457	-2 -2	-1 -1	66	39	-1 -1	-1 1
ע ₆ דע	1636	1557	0	0	6	6	0	0	ν ₁₇ ν ₁₈	78	174	-2 -2	0	1	41	-1 -1	Ó
ν8	1636	1548	0	ő	6	6	Ô	ő	ν ₁₉	78	81	- <u>2</u> -1	0	1	0	0	0
ν ₈	1602	1500	0	0	206	139	2	2	ν ₂₀	42	76	0	0	7	4	ŏ	ő
ν ₁₀	1260	1190	-2	Ö	1	59	0	õ	ν_{21}	42	51	0	ŏ	7	2	ŏ	ő
ν ₁₁	1260	1155	-2	ŏ	î	Ó	ŏ	ő	- 21			ŭ	ŭ	,	-	ŭ	Ŭ
								PhNC:	O (4c	y							
ν_1	3403	3228	0	0	4	7	0	0	ν22	1091	923	0	0	2	0	0	0
ν_2	3399	3219	0	0	16	9	0	0	ν_{23}	1040	888	-7	-1	8	1	-1	0
ν_3	3390	3210	0	0	27	7	0	0	ν_{24}	972	869	0	0	19	2	0	0
ν_4	3378	3201	0	0	8	1	0	0	ν_{25}	947	819	0	0	0	0	0	0
ν5	3368	3194	0	0	1	2	0	0	ν_{26}	853	738	0	0	78	108	0	0
ν_6	2532	2423	-4	-22	1978	4049	-82	-50	ν_{27}	762	694	-2	-1	19	16	1	-1
ν_7	2421	2205	-16	-32	8488	55	19	-53	ν_{28}	726	661	0	0	69	4	0	0
ν8	1866	1643	-16	-7	176	352	29	28	ν_{29}	680	652	-1	0	1	6	-1	0
ν9	1798	1594	0	0	74	50	-12	-5	ν_{30}	666	628	0	-12	69	40	-1	10
ν_{10}	1782	1579	0	0	6	9	0	0	ν_{31}	663	603	0	-15	69	33	0	9
ν_{11}	1660	1526	0	0	90	49	-5	-3	ν_{32}	604	589	-4	-6	58	41	-2	-10
ν_{12}	1619	1497	0	0	5	4	0	0	ν_{33}	600 554	572 502	−3 −1	-2	49	18	1	-10
ν_{13}	1470	1389	0 0	0 0	0 1	1 10	0 0	0 0	V34	554 457	302 437	-1 -5	0 0	7 0	5 16	0 0	0
V14	1333	1311	0 ·	0	52	2	0	0	V35	437 438	437 417	-3 0	0	0	16	0	0 0
ν ₁₅	1328 1292	1245 1241	0	0	2	0	0	0	ν36	438 375	346	0	-2	2	9	0	0
ν ₁₆	1292	1241	-1	0	1	75	5	0	V37	278	251	_3	-2 0	0	0	0	0
ν ₁₇	1184	1105	-1 0	0	3	9	0	0	V38	92	90	-3 0	0	1	6	0	0
ν ₁₈	1128	1052	0	0	1	6	0	0	ν ₃₉ ν ₄₀	63	71	-1	0	1	Ö	0	0
ν ₁₉ ν ₂₀	1128	1032	0	0	0	0	0	0	ν40 ν41	35	24	0	0	0	1	Ö	Ö
ν ₂₀ ν ₂₁	1101	928	Ö	0	0	ŏ	0	0	V41	11	18	Ö	0	0	1	Ö	ŏ

^a 6-31G* basis set, unless otherwise noted. ^b In cm⁻¹. ^c In km mol⁻¹. ^d Isotopomer wavenumber shifts (MP2 values). ^e Intensity changes due to isotopic substitutions (MP2 values). f 3-21G basis set at the MP2 level.

of CO is enthothermic by 330 kJ mol⁻¹. Loss of the hydrogen atom is also unfavorable (by 361 kJ mol-1). Note that the product of the reaction 4a -> HNCC + CO, can rearrange to a more stable isomer, i.e., HCCN. However, the "nonlinear" dissociations to HCCN is also calculated to be endothermic (by 229 kJ mol^{-1}). The energy requirement for other dissociation reactions are significantly higher (Table 4). Hence, iminopropadienone (4a) is predicted to be a thermodynamically stable species. This finding is consistent with the experimental observation of HNC₃O by neutralization-reionization mass spectrometry.²

Cyanoketene (9) is calculated to be the global minimum on the C₃HNO potential energy surface. The iminopropadienone isomer (4a) lies 68 kJ mol⁻¹ above 9. Rearrangement of 4a to 9 can, in theory, take place via a 1,3-hydrogen shift, involving a fourcentre transition structure 10 (Figure 3). The barrier for this process is substantial, 402 kJ mol-1, even larger than those required for the fragmentations 4a to HNCC + CO and H' + NCCCO'. There is strong experimental evidence that the tautomerization of iminopropadienone to cyanoketene in high temperature FVP experiments² occurs via wall collision or by intermolecular proton transfer.3b

Ionization Energy and Heat of Formation of HNC₃O. Adiabatic (IE_a) and vertical (IE_v) ionization energies of HNC₃O (4a) were calculated at the G2(MP2) level of theory, which in general can reproduce ionization energies to within 0.1 eV.17 For the isoelectronic analogue C₃O₂, the calculated adiabatic ionization

energy (10.54 eV) is in very good agreement with the experimental value (10.60 eV).18 The predicted IE_a and IE_v values of iminopropadienone 4a are 9.51 and 9.74 eV, respectively (Table 4). The large difference between IE_a and IE_v (0.23 eV) is attributed to the large change in molecular geometry upon ionization. The iminopropadienone radical cation is calculated to have a strongly bent structure (11, Figure 3), in contrast to the essentially linear structure of the corresponding neutral. The calculated central CCC angle of 11 is 128.5°, significantly smaller than the experimental ∠CCC value (156°) for carbon suboxide. 4b,c

Two independent methods (I and II) were used to determine the heat of formation (ΔH_1°) of HNC₃O (4a). In method I, $\Delta H_1^{\circ}_{298}$ (HNC₃O) was evaluated from the theoretical energy change (-66 kJ mol-1 at 298 K)19 for the isodesmic reaction (eq 1) together with accepted literature $\Delta H_{\rm f}^{\circ}_{298}$ values for formal-

$$HN=C=C=C=O + CH_2=O \rightarrow O=C=C=C=O + CH_2=NH (1)$$

dehyde (-109 kJ mol⁻¹)¹⁸ and carbon suboxide (-94 kJ mol⁻¹),¹⁸ and a recently revised $\Delta H_1^{\circ}_{298}$ value for formaldimine (86 kJ mol⁻¹).²⁰ The predicted $\Delta H_{\rm f}^{\circ}_{298}$ value for 4a using method I is 167 kJ mol⁻¹. In method II, $\Delta H_{\rm f}^{\circ}$ was obtained from the calculated heat of atomization (eq 2)in conjunction with the

$$HN=C=C=C=O \to 3C + N + O + H$$
 (2)

Figure 4. Calculated normal vibrations of HNC₃O (4a) at the MP2/6-31G* level. Vibrations ν_{1} - ν_{9} are in the x, y plane (plane of paper), while vibrations ν_{10} - ν_{12} are in the x, z plane (out of the plane of paper).

TABLE 6: Calculated Wavenumbers (cm⁻¹) and IR Intensities (km mol⁻¹) of Cumulenic Stretching Vibrations of Iminopropadienones

	wave	number	inte	ensity
	MP2/3-21G	MP2/6-31G*	MP2/3-21G	MP2/6-31G*
		HNC ₃ O	(4a)	
ν_2	2434	2446	2066	2538
ν_3	2175	2204	2	21
		CH ₃ NC ₃ C) (4b)	
ν_4	2432	2452	2668	3131
ν5	2216	2244	15	22

TABLE 7: Comparison of Calculated Cumulenic CO Stretching Wavenumbers (cm⁻¹) with Experimental Values

molecule	theory ^a	experiment	$ratio^b$
0=C=0	2448	2349 ^c	0.96
O=C=C=C=O	2485	2290°	0.92
$CH_2 = C = O$	2235	2152^{c}	0.96
$CH_2=C=C=O$	2207	2125^{d}	0.96
$CH_3N=C=C=C=O$	2453	2279°	0.93
HN=C=C=C=O	2446		

^a MP2/6-31G* values. ^b Experiment value/calculated value. ^c From ref 23. ^d From ref 13. ^e From ref 1.

accurate heats of formation of atomic C, N, O, and $H^{.18}$ This leads to a $\Delta H_{\rm f}^{\circ}{}_{298}$ value of 175 kJ mol $^{-1}$ for ${\bf 4a}$, reasonably close to the value computed by method I. Using a similar approach (method II), the G2(MP2) $\Delta H_{\rm f}^{\circ}{}_{298}$ for C₃O₂ is -98 kJ mol $^{-1}$, in excellent agreement with the experimental value of -94 kJ mol $^{-1}$. ¹⁸ Because of the uncertainty of $\Delta H_{\rm f}^{\circ}{}_{298}$ (CH₂NH) involved in method I, the method II value (175 kJ mol $^{-1}$) constitutes our final best theoretical estimate, to which an error bar of approximately 10 kJ mol $^{-1}$ is to be associated. ¹⁷

Vibrational Wavenumbers and Infrared Spectra. The infrared spectra of all three iminopropadienones (4a-c), calculated at the HF/6-31G* and MP2/6-31G* (MP2/3-21G for the phenyl derivative) levels are reported in Table 5. The normal vibrations calculated for the parent compound, HNC₃O (4a) at the MP2/6-31G* level are illustrated in Figure 4. As can be seen in Table 5, the characteristic features of the spectra of the iminopropa-

dienones are the strong absorption bands of the cumulenic²¹ N=C=C=O stretching wavenumbers, ν_2 and ν_3 for 4a, ν_4 and ν_5 for 4b, and ν_6 and ν_7 for 4c. In the first approximation, the two cumulenic vibrations can be considered as a mixing of the N=C=C and C=O vibrations. The higher-wavenumber "antisymmetric" ketenic²¹ stretch is associated predominantly with the C=C=O moiety, while the lower-wavenumber "symmetric" stretch has a higher proportion of C=N vibration. At the Hartree-Fock level, both the antisymmetric and symmetric vibrations are predicted to yield strong absorptions, with the symmetric stretch considerably stronger than the antisymmetric stretch for 4b and 4c. Inclusion of electron correlation at the MP2 level leads to a drastic change in the relative intensities of the two cumulenic vibrations. In fact, the symmetric cumulenic stretch (e.g., ν_3 in 4a) becomes a very weak band at the MP2 level. Since MP2 wavenumber calculation employing the 6-31G* basis set was not feasible for the phenyl derivative (4c), a smaller 3-21G basis set was used instead. To assess the likely accuracy of the MP2/3-21G level, we have determined also the IR spectra of 4a and 4b at the MP2/3-21G level. As can be seen in Table 6, the calculated MP2/3-21G harmonic wavenumbers of the cumulenic stretches are reasonably close to the MP2/6-31G* values. More importantly, the relative intensities are essentially the same as when using the larger 6-31G* basis set. This lends confidence to our prediction of the IR intensities for the phenyl derivative (4c). Indeed, the calculated MP2 relative intensities of v_6 and v_7 for 4a is in excellent accord with the experimental observations, as described in section 3. Hence, inclusion of electron correlation is essential for the correct prediction of relative intensities of the cumulenic stretching wavenumbers of iminopropadienones. Significant wavenumber shifts are observed for these vibrations when going from R = H to $R = CH_3$ to R = Ph (Table 5).

To provide a more definitive assignment of the vibrational spectra of the iminopropadienones (4a-c), the calculated IR spectra of the ¹⁵N (R¹⁵NCCCO) and ¹³C(RNCC¹³CO) isotopomers are reported in Table 5. Significant wavenumber shifts of the cumulenic bands are predicted for the ¹⁵N and ¹³C isotopomers for all three cumulenes. The shifts calculated for the ¹³C isotopomers are larger than those for the ¹⁵N isotopomers. Enormous intensity changes are calculated for the symmetric cumulenic stretches (the smaller peaks) of the ¹³C isotopomers. For the methyl derivative of iminopropadienone (4b), the intensity of the ν_5 vibration is calculated to decrease by 60%, while for the parent compound (ν_3 in 4a) and the phenyl derivative (ν_7 in 4c), the symmetric cumulenic band is predicted to almost disappear in their ¹³C spectra.

As pointed out in the structure section, the iminopropadienones are predicted to be "flexible" molecules. This is readily confirmed by the existence of some very low-wavenumber vibrations. In particular, we note the chain-bending vibrations, e.g., ν_{11} in **4a** (Table 5 and Figure 4). It is noteworthy that the three lowest bending fundamentals of HNCO are severely perturbed by Coriolis mixing. Likewise, the importance of including the Coriolis coupling term to describe the bending δ_{NS} of N_2S has been reported. Hence, the six frequencies in **4a** lying between 500 and 850 cm⁻¹ (ν_5 – ν_{10}) may require Coriolis correction for reliable theoretical prediction.

Finally, we have determined the CO stretching wavenumbers of several molecules containing a cumulenic CO moiety at the MP2/6-31G* level (Table 7). There is excellent agreement between theory and experiment.²³ The calculated cumulenic CO stretching wavenumbers are 4-8% larger than the experimental values and suggest a scaling factor of 0.94 may be appropriate for the cumulenic stretching wavenumbers of the unsubstituted iminopropadienone (4a). Thus, our predicted values of the cumulenic vibrations ν_2 and ν_3 of 4a are 2299 and 2072 cm⁻¹, respectively.

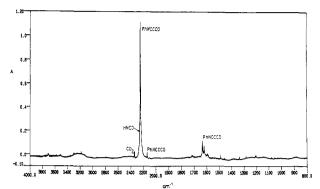


Figure 5. Ar matrix IR spectrum (12 K) of PhNC₃O (4c), generated by FVP of 1 at 800 °C.

TABLE 8: Experimental IR Spectra of PhN=C=C=C=O (4c) and Its Isotopomers (Ar Matrix, 12 K)

PhNC ₃ O ^b	Ph15NC3O	PhNC ₂ ¹³ CO
2247 vs (2253)	2247 vs	2222 vs
2141 w (2154)	2126 w	
1633 m (1661)	1623 m	1630 m
1620 m (1600)	1608 m	1617 m
1490 w (1478)	1490 w	1490 w
1284 w (1308)	1284 w	1284 w
1210 w (1182)	1208 w	1210 w
754 w (759)	754 w	
685 w (678)	685 w	
565 w (589)	565 w	

^a Wavenumbers in cm⁻¹. ^b HF/6-31G* values (scaled by a factor of 0.89)8 are in parentheses.

3. Experimental IR Spectra of PhNC₃O and Its Isotopomers

The Ar matrix IR spectrum of the phenyl derivative (4c), generated by FVP of 1c at 800 °C, is shown in Figure 5.24 The same spectrum is obtained from 2c and 3c except that CO2 and acetone together with dimethylamine or methanethiol will, of course, also be seen in the latter spectra. By far the most prominent feature is the very strong band at 2247 cm⁻¹. As shown below, the very weak band at 2141 cm⁻¹ is also due to the cumulenic moiety in PhN=C=C=C. Since the calculated Hartree-Fock wavenumbers generally exceed the experimental values by 12%,8 a scaling factor of 0.89 has been applied to the calculated harmonic wavenumbers at the HF/6-31G* level. The scaled calculated values are in excellent accord with the observed IR wavenumbers of 4c (Table 8).

The ¹⁵N and ¹³C labeled cumulenes 6 and 8 were generated by FVP of the Meldrum's acid derivatives 5 and 7, respectively (see Experimental Section for the syntheses). The infrared spectrum of 6 reveals a significant shift of the small peak at 2141 cm⁻¹, whereas the major peak at 2247 cm⁻¹ is hardly affected (Figure 6 and Table 8). In contrast, the ¹³C-labeled compound 8 shows a large shift of the ketenic band at 2247 cm⁻¹, while the small peak at 2141 cm⁻¹ virtually disappears (Figure 6 and Table 8).

These observations are nicely confirmed by the calculations at the MP2/3-21G level (Table 5). While both vibrations in the cumulene region (ν_6 and ν_7) involve all the cumulated atoms, ν_7 (the "symmetric" cumulenic vibration) is dominated by the C=N stretching. As a consequence, the band ν_7 is strongly affected by ¹⁵N labeling in 4c (2141–2126 cm⁻¹). The theoretical calculations reproduce the experimental isotopic shifts almost quantitatively (16 cm⁻¹ for ν_7 and 4 cm⁻¹ for ν_6 in 6). ¹³C labeling as in 8 is calculated to have a large effect on both ν_6 and ν_7 (22 and 32 cm⁻¹, respectively). The computed isotopic shift for ν_6 is in good accord with the measured value (25 cm⁻¹). Most importantly, the calculated intensity of ν_7 decreases 30-fold in 8, thereby explaining the virtual disappearance of this band in the experimental spectrum (Figure 6 and Table 8).

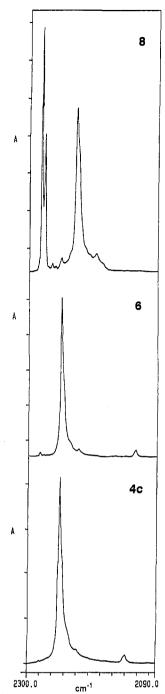


Figure 6. Partial IR spectra of 4c, 6, and 8 (Ar, 12 K). The cumulenic signals at 2247 and 2141 cm⁻¹ (4c), 2247 and 2126 cm⁻¹ (6), and 2222 cm⁻¹ (8) are shown. The strong signals near 2278 and 2273 cm⁻¹ in the spectrum of 8 are due to ¹³CO₂.

4. Conclusion

The iminopropadienones, RN=C=C=C (R = H, CH_3 , and Ph) are predicted to have an almost linear NCCCO backone (∠CCC≈ 176°). The calculated CCC bending potential, however, is extremely flat. The very good agreement between the calculated and experimental IR spectra of PhNC₃O (4a) and its ¹⁵N and ¹³C isotopomers, 6 and 8, strongly supports the identification of these compounds based largely on chemical and mass spectrometric evidence.^{1,2} This also gives cause to believe that the calculated structures of iminopropadienones are close to reality. The calculated stability of HNC₃O is consistent with its experimental observation in the gas phase.2 We are currently extending this fruitful interplay between theory and experiment to the characterization of bisimines RN=C=C=C=NR'.3

5. Experimental Section

Apparatus and cryostats used for flash vacuum pyrolyses were as previously described.²⁵ IR spectra were recorded on a Perkin-Elmer 1720X FTIR spectrometer at a resolution of 1 cm⁻¹.

Materials. Aniline[15N] (99 atom % 15N) and barium carbonate[13C] (>90 atom % 13C) were obtained from Cambridge Isotope Laboratories, Woburn, MA.

2,2-Dimethyl-5-[dimethylamino([15N]phenylamino)methylene]-1,3-dioxane-4,6-dione (5) was synthesized from 2,2-dimethyl-5-[bis(imethylthio)methylene]-1,3-dioxane-4,6-dione (0.263 g; 1.06 mmol) in 5 mL of EtOH by adapting the method published for the unlabeled compound1 with aniline[15N] (0.100 g; 1.06 mmol) under reflux for 2 h. The solvent was removed at 15 Torr (25 °C). THF (10 mL) was added, followed by 0.5 mL of 40%aqueous dimethylamine and 287 mg (1.06 mmol) of HgCl₂. The mixture was stirred for 30 min at 25 °C, filtered; and the clear filtrate was extracted with CH₂Cl₂ (13 × 10 mL). The organic layer was washed with H_2O (2 × 5 mL), dried over Na_2SO_4 , filtered, and concentrated at 17 Torr (25 °C). The resulting solid (0.24 g; 77%) was recrystallized from THF/pentane to give the colorless compound, mp 192-194 °C; ¹H NMR δ 1.71 (s, 6 H; 2 CH₃), 2.92 (s, 6 H; N(CH₃)₂), 7.1-7.4 (M, 5H; aromatic), 9.13 (s, 1 H; NH); 13 C NMR (DMSO- d_6) δ 26.26 (C(CH₃)₂), $39(N(CH_3)_2)$, 72.37(C-5), 100.80(C-2), 123.00 (phenyl-ortho), 125.00 (phenyl-para), 128.60 (d, $J_{15N} = 1.4$ Hz; phenyl-meta), 139.53 (d, J_{15} _N = 15.8 H z; phenyl-C1) 161.76 (CO), 162.91 (d, $J_{15N} = 19.0 \text{ Hz}; ^{15}N-C-N).$

2,2-Dimethyl-5-[dimethylamino(phenylamino)methylene]-1,3dioxane-4,6-dione-13C2 (7). Barium carbonate [13C] was converted to ¹³CO₂ using a published procedure.²⁶ ¹³CO₂ was converted into malonic acid[1,3-13C2] and then into the doubly labeled Meldrum's acid (2,2-dimethyl-1,3-dioxane-4,6-dione[4,6-13C₂]) using literature methods.²⁷ The Meldrum's acid was converted to the 5-[bis(imethylthio)methylene] derivative using the standard method,28 and the title compound 7 was then prepared in the same way as the unlabeled compound1 and as described for the ¹⁵N-labeled analogue (6) above. Mp 192-194 °C; ¹H NMR (CDCl₃) δ 1.61 (s, 6 H; (CH₃)₂C), 2.86 (s, 6 H; N(CH₃)₂), 7.0– 7.3 (m, 5 H; aromatic), 9.27 (br s, 1 H; NH); ¹³C NMR (CDCl₃) δ 26.30 ((CH₃)₂C), 41.67 (N(CH₃)₂), 72.15 (C-5), 102.78 (C-2), 122.86 (phenyl-ortho), 125.99 (phenyl-para), 129.47 (phenylmeta), 138.83 (phenyl-C1), 164.3 (CO), 166.7 (NCN).

General Procedure for Flash Vacuum Pyrolysis. The Meldrum's acid precursor (10-15 mg, 0.03-0.05 mmol) was gently sublimed (80–100 °C) with a continuous stream of Ar ((6–8) \times 10⁻⁵ mbar; 50 mbar of Ar in 15 min) into the pyrolysis apparatus²⁵ (10-cm quartz tube; 8-mm i.d. 500 °C). The product was collected on a BaF₂ (or KBr) disk at 16-18 K and FTIR spectra were recorded periodically to monitor the formation of the matrix. A typical pyrolysis/deposition took 3 h to completion. Reference spectra of the starting materials, acetone, and dimethylamine were obtained under similar conditions, without pyrolysis.

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References and Notes

- (1) Mosandl, T.; Kappe, C. O.; Flammang, R.; Wentrup, C. J. Chem. Soc., Chem. Commun. 1992, 1571.
- (2) Flammang, R.; Laurent, S.; Flammang-Barbieux, L.; Wentrup, C. Rapid. Commun. Mass Spectrom. 1992, 6, 667
- (3) (a) Stadtmüller, S.; White, R.; Flammang, R.; Laurent, S.; Wong, M. W.; Wentrup, C. To be published. (b) Flammang, R.; Laurent, S.; Flammang-Barbieux, M.; Wentrup, C. Org. Mass Spectrom. 1993, 28, 1161. (4) (a) Kappe, T. Angew. Chem., Int. Ed. Engl. 1974, 13, 491. (b) Jensen, P.; Johns, J. W. C. J. Mol. Spectrosc. 1986, 118, 248. (c) Jensen, P. J. Mol.
- Spectrosc. 1984, 104, 59. (d) Loses, R. L.; Sabin, J. R.; Oddershede, J. J. Mol. Spectrosc. 1981, 86, 357. (e) Bunker, P. R. J. Mol. Spectrosc. 1980, 80, 422. (f) Webber, W. H. J. Mol. Spectrosc. 1980, 79, 396. (g) Duckett, J. A.; Mills, I. M.; Robiette, A. G. J. Mol. Spectrosc. 1976, 63, 249. (h) Smith, W. M.; Leroy, G. E. J. Chem. Phys. 1966, 45, 1767.

 (5) Hehre, W. J.; Radom, L.; Schleyer, P. v. R.; Pople, J. A. Ab Initio
- Molecular Orbital Theory; Wiley: New York, 1986.

 (6) Frisch, M. J.; Trucks, G. W.; Head-Gordon, M.; Gill, P. M. W.; Wong, M. W.; Foresman, J. B.; Johnson, B. G.; Schlegel, H. B.; Robb, M. A.; Replogle, E. S.; Gomperts, R., Andres, J. L.; Raghavachari, K.; Binkley, J. S.; Gonzalez, C.; Martin, R. L.; Fox, D. J.; DeFrees, D. J.; Baker, J Stewart, J. J. P.; Pople, J. A. GAUSSIAN 92; Gaussian Inc.: Pittsburgh PA,
- (7) (a) Handy, N. C.; Amos, R. D.; Gaw, J. F.; Rice, J. E.; Simandiras, E. D. Chem. Phys. Lett. 1985, 120, 151. (b) Pople, J. A.; Krishnan, R.; Schlegel, H. B.; Binkley, J. S. Int. J. Quantum Chem. Symp. 1979, 13, 225.
 (8) (a) Pople, J. A.; Scott, A. P.; Wong, M. W.; Radom, L. Isr. J. Chem.,
- in press. (b) Pople, J. A.; Schlegel, H. B.; Raghavachari, K.; DeFrees, D. J.; Binkley, J. S.; Frisch, M. J.; Whitesides, R. F.; Hout, R. F.; Hehre, W. J. Int. J. Quantum Chem. Symp. 1981, 15, 269
- (9) (a) Curtiss, L. A.; Raghavachari, K.; Pople, J. A. J. Chem. Phys. 1993, 98, 1293. (b) Curtiss, L. A., Raghavachari, K.; Trucks, G. W.; Pople, A. J. Chem. Phys. 1991, 94, 7221
- (10) (a) Glendening, E. D.; Reed, A. E.; Carpenter, J. E.; Weinhold, F. NBO Version 3.1. For reviews, see (b) Reed, A. E.; Curtiss, L. A.; Weinhlod, F. Chem. Rev. 1988, 88, 899. (c) Weinhlod, F.; Carpenter, J. E. In The Structure of Small Molecules and Ions; Naaman, R., Vager, Z., Eds.; Plenum: New York, 1988
- (11) Pople, J. A.; Head-Gordon, M.; Raghavachari, K. J. Chem. Phys. 1987, 87, 5968.
- (12) (a) Scott, A. P.; Radom, L. Chem. Phys. Lett. 1992, 200, 15. (b) Farnell, L.; Radom, L. J. Am. Chem. Soc. 1984, 106, 25. (c) Farnell, L.; Radom, L. Chem. Phys. Lett. 1982, 91, 373. (d) Taylor, P. R. J. Comput. Chem. 1984, 6, 589. (e) Brown, R. D.; Godfrey, P. D.; Ball, M. J.; Godfrey, S.; McNaughton, D.; Rodler, M.; Kleibomer, B.; Champion, R. J. Am. Chem. Soc. 1986, 108, 6534
- (13) Chapman, O. L.; Miller, M. D.; Pitzenberger, S. M. J. Am. Chem. Soc. 1987, 109, 6867.
- (14) East, A. L.; Johnson, C. S.; Allen, W. D. J. Chem. Phys. 1993, 98, 1299.
- (15) Harmony, M. D.; Laurie, V. W.; Kuczkowski, R. L.; Schwendeman, R. H.; Ramsay, D. A.; Lovas, F. L.; Lafferty, W. J.; Maki, A. G. J. Phys. Chem. Ref. Data 1979, 8, 619.
- (16) Flammang, R.; van Haverbeke, Y.; Wong, M. W.; Rühmann, A.; Wentrup, C. J. Phys. Chem., submitted.
 (17) Wong, M. W.; Radom, L. To be published.
 (18) Lias, S. G.; Bartmess, J. E.; Liebman, J. F.; Holmes, J. L.; Levin, R.
- D.; Mallard, W. G. J. Phys. Chem. Ref. Data 1988, 17, Suppl. 1.
- (19) The G2(MP2) H₂₉₈ energies of HNC₃O, C₃O₂, CH₂CO and CH₂-NH are -244.442 36, -264.342 48, -114.332 26, and -94.457 26 hartrees, respectively.
- (20) Smith, B. J.; Pople, J. A.; Curtiss, L. A.; Radom, L. Aust. J. Chem. 1992, 45, 285
- (21) In this paper, the term "ketenic" denotes the C=C=O moiety while "cumulenic" refers to the whole cumulene skeleton including the CCO moiety. (22) Collins, C. L.; Yamaguchi, Y.; Schaefer, H. F. J. Chem. Phys. 1993,
- 98, 4777 (23) (a) Shimanouchi, T. Tables of Molecular Vibrational Frequencies
- I; Publication NSRDS-NBS 39; National Bureau of Standards, Washington, (b) Shimanouchi, T. J. Phys. Chem. Ref. Data. 1977, 6, 993
- (24) It is not necessary to isolate 4a in an Ar matrix. It can be observed spectroscopically up to room temperature for brief periods of tim
- (25) Wentrup, C.; Blanch, R.; Briel, H.; Gross, G. J. Am. Chem. Soc. 1988, 110, 1874.
 - (26) Mayor, C.; Wentrup, C. J. Am. Chem. Soc. 1975, 97, 7467.
- (27) (a) Thies, H.; Franke, W.; Schwarz, H. Synthesis 1982, 587. (b) Brown, R. F. C.; Browne, N. R.; Coulston, K. J.; Eastwood, F. W.; Irvine, M. J.; Pullin, D. E.; Wiersum, U. E. Aust. J. Chem. 1989, 42, 1321
- (28) Ye, F.-C.; Chen, B.-C.; Huang, X. Syntheses 1989, 317. (b) Huang, X.; Chen, B.-C. Ibid. 1986, 967. (c) Ben Cheikh, A.; Chuche, J.; Manisse, ; Pommelet, J. C.; Netsch, K.-P.; Lorencak, P.; Wentrup, C. J. Org. Chem. 1991, 56, 970.