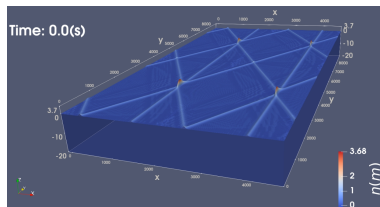


Modeling extreme water waves with Firedrake

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Leeds Institute for Fluid Dynamics; Totnes' Firedrake Workshop 2023



Summary with movies

Firedrake extreme wave modelling of *Benney-Luke (BL) eqns*:

- ▶ horizontally $\Phi(x, y, t), \eta(x, y, t)$: CG1, CG2-2-CG4;
space-time discrete variational principle (VP),
- ▶ discrete VP: wave-amplitude & phase-space conservation,
- ▶ x-periodic mesh, 3 weak forms,
- ▶ mesh refinement needed; other element types?

Summary

Firedrake modelling of (driven) **potential-flow** (PF) *water waves*:

- ▶ space-time discrete VP, variables $h(x, y, t)$, $\tilde{\phi}(x, y, t) = \phi(x, y, b(x, y) + h(x, y, t), t)$, $\phi(x, y, z, t)$ with mixed horizontal and vertical coordinates,
- ▶ transformation to fixed domain,
- ▶ **now**, vertical z : one vertical element with Lagrange/Chebyshev polynomials: user-arranged,
- ▶ space-time discrete VP in horizontal: CG1 polynomials; 3 to 5 weak forms.
- ▶ **In progress**: FD implementation of VPs for complicated moving domains in x, y, z .

Firedrake for extreme wave amplification in BL

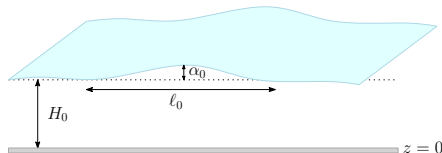
- ▶ Origin 2010 *bore-soliton-splash*:
- ▶ To what extent do exact but idealised extreme- or rogue-wave solutions survive in more realistic settings?
- ▶ **Rogue wave**: $A_r \geq 2.2 \times A_{\text{ambient}}$.
- ▶ Will such extreme waves fall apart due to dispersion or other mechanisms?
- ▶ Use $4\times$ & $9\times$ KP amplifications of interacting solitons/cnoidal waves.
- ▶ What do you think: will we be **able to reach the ninefold wave amplification** in more realistic calculations or in reality?



Mathematical hierarchy: BL and KP approximations

- ▶ Kadomtsev & Petviashvili (1970) eqn: unidirectional in 2DH
- ▶ Benney-Luke (1964) eqns –BL: bidirectional in 2DH

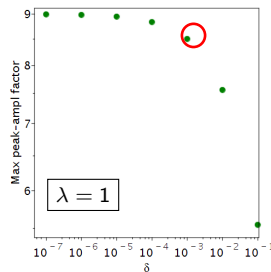
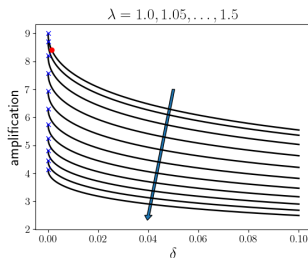
$$\begin{aligned}\partial_t \eta - \frac{\mu}{2} \partial_t \nabla^2 \eta + \nabla \cdot ((1 + \epsilon \eta) \nabla \Phi) - \frac{2\mu}{3} \nabla^4 \Phi &= 0 && \text{in } \Omega \\ \partial_t \Phi - \frac{\mu}{2} \partial_t \nabla^2 \Phi + \frac{\epsilon}{2} |\nabla \Phi|^2 + \eta &= 0 && \text{in } \Omega \\ \mathbf{n} \cdot \nabla \Phi = 0 \text{ on } \partial\Omega &\text{ and } \mathbf{n} \cdot \nabla (\nabla^2 \Phi) = 0 && \text{on } \partial\Omega\end{aligned}$$



$$\begin{aligned}\epsilon &= \alpha_0 / H_0 \ll 1 \\ \mu &= (H_0 / \ell_0)^2 \ll 1\end{aligned}$$

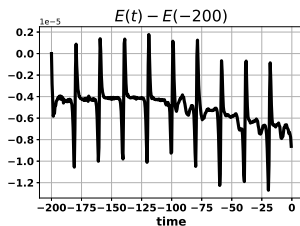
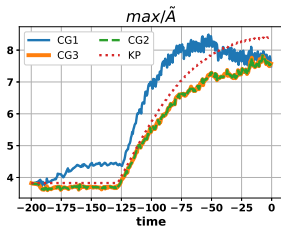
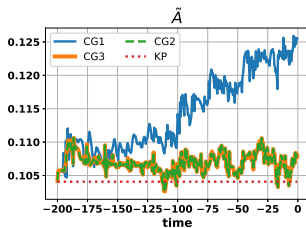
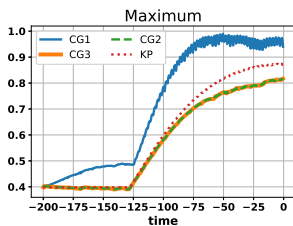
Maximum 9-fold amplification in KP & BL?

- Amplification: $\frac{u(X_*, Y_*, \tau_*)}{\tilde{A}} = \frac{(\sqrt{\lambda} + 2)^2}{\lambda} + \mathcal{O}(\sqrt{\delta}) \xrightarrow{\delta=0, \lambda \rightarrow 1} 9$



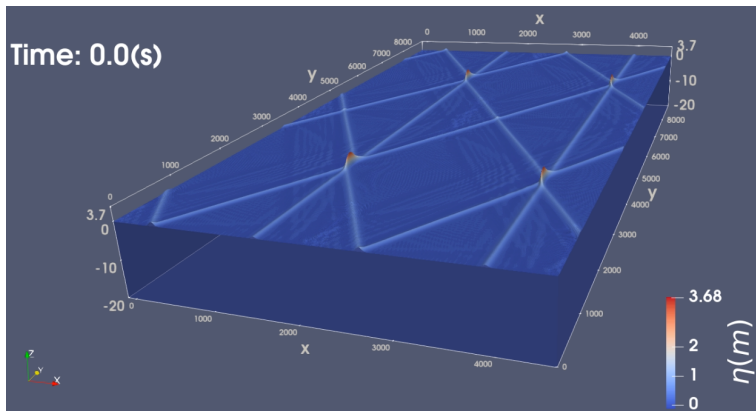
- Seed BL with KP-solution prior to maximum reached.
- Wave amplitude & phase space volume preserved via space-time discrete VP: no loss of amplitude.

Results BL-simulation three-soliton interaction



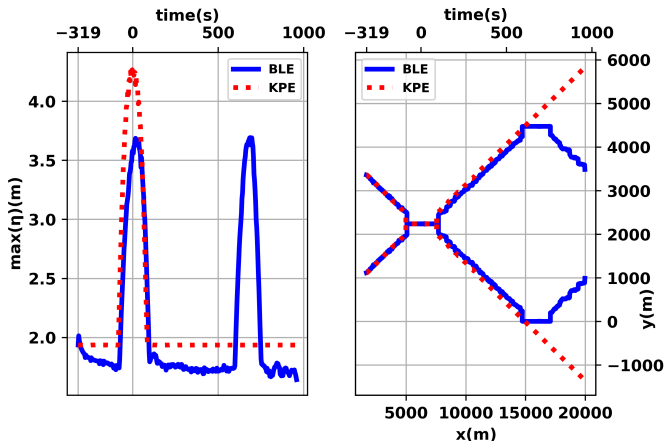
Results simulation three-soliton interaction (dimensional)

Crossing seas (4 or 8 domains combined – YouTube)



Results simulation three-soliton interaction (dimensional)

Cnoidal waves with periodicity in x, y, t (max. vs. t & x - y tracks):



Summary

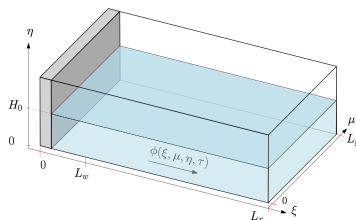
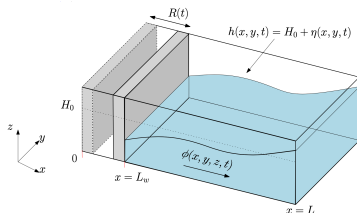
- ▶ Nine-fold soliton amplification shown theoretically, but only in limit $\delta \rightarrow 0$
- ▶ Web-soliton amplification $KP \approx 8.4$, simulated for $BL \approx 7.8$
- ▶ Amplifications achieved as simulated cnoidal crossing seas
- ▶ Local p or h mesh-refinement needed in x -periodic channel.
- ▶ Can amplifications survive in **potential-flow equations**?



Firedrake for potential flow water waves

Water-wave equations

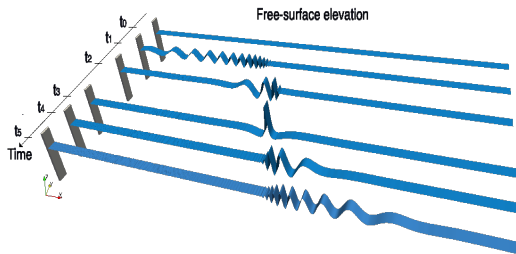
$$\begin{aligned}\nabla^2 \phi &= 0 && \text{in } \Omega \\ \partial_t \eta + \nabla \phi \cdot \nabla \eta - \partial_z \phi &= 0 && \text{at } z = H_0 + \eta \\ \partial_t \phi + \frac{1}{2} |\nabla \phi|^2 + g\eta &= 0 && \text{at } z = H_0 + \eta \\ \mathbf{n} \cdot \nabla \phi &= 0 && \text{on } z = 0 \text{ and } \partial\Omega\end{aligned}$$



Current strategy: Firedrake in horizontal domain

Current strategy (Gidel 2018; Gidel et al. 2022/23)

- ▶ Time-dependent free surface & piston wavemaker, transformed to fixed domain.
- ▶ One element in vertical Lagrange/Chebychev $p = 4, 9$; user-arranged.
- ▶ Firedrake: space-time discrete VP in horizontal with CG1/CG2, Störmer-Verlet.
- ▶ Time-dependent VP, non-autonomous Hamiltonian.
- ▶ Good comparison with experimental data (Gidel et al 2022/23); soliton amplification simulation in progress.



Future strategy: Firedrake via VPs

Goal: implement (time-discrete) VP & derive weak forms automatically.

- ▶ **Why:** VPs in moving domains complicated, e.g. wave tank with waveflap, wave-beam FSI, numerical mesh motion.
- ▶ Simple case; nonlinear PF in $\{x, z\}$, $\partial_y = 0$, $W = L_w \leq L_x$, no wavemaker:

$$0 = \delta \int_0^T \int_0^{L_x} \int_0^{H_0} - \left[\frac{1}{2} \frac{L_w^2}{W} h(\phi_x + (z/h)h_x)\phi_z \right]^2 + \frac{1}{2} W \frac{H_0^2}{h} (\phi_z)^2 \Big] dz dx \\ + \int_0^{L_x} -gH_0 W h \left(\frac{1}{2} h - H_0 \right) + H_0 W \phi|_{z=H_0} h_t dx dt$$

- ▶ Time-discrete version, coupled surface-interior system, use *fd.derivative*: $\delta h^{n+1}, \delta \varphi^{n+1}$, explicit fderivative ϕ^{n+1} :

$$0 = \delta \int_0^{L_x} \int_0^{H_0} - \left[\frac{1}{2} \frac{L_w^2}{W} h^n (\phi_x^{n+1} + (z/h^n)h_x^n) \phi_z^{n+1} \right]^2 + \frac{1}{2} W \frac{H_0^2}{h^n} (\phi_z^{n+1})^2 \Big] dz dx \\ + \int_0^{L_x} -gH_0 W h^n \left(\frac{1}{2} h^n - H_0 \right) + H_0 W \phi^{n+1}|_{z=H_0} \frac{(h^{n+1} - h^n)}{\Delta t} + H_0 W \phi^n|_{z=H_0} \frac{h^n}{\Delta t} dx$$

Future strategy: Firedrake for complicated VPs

Why? Complexity of eqns associated with VP for waveflap-driven water waves:

$$z^{n+1} = \eta h^{n+1}/H_0, \quad x_\xi^{n+1,n} = \frac{(L_w - W^{n+1,n})}{L_w} + \frac{\eta}{H_0} W_z^{n+1,n} \frac{(L_w - \xi)}{L_w} h_\xi^{n+1}, \quad (2.69a)$$

$$x_\eta^{n+1,n} = \frac{(L_w - \xi)}{L_w} W_z^{n+1,n} \frac{h^{n+1}}{H_0}, \quad (2.69b)$$

$$z_\xi^{n+1} = \eta h_\xi^{n+1}/H_0, \quad z_\eta^{n+1} = h^{n+1}/H_0, \quad (2.69c)$$

$$W^{n+1,n} = W(\eta h^{n+1}/H_0, \tau^n), \quad W_\tau^{n+1,n} = \partial_\tau W(z, \tau^n)|_{z=\eta h^{n+1}/H_0}, \quad (2.69d)$$

$$W_z^{n+1,n} = \partial_z W(z, \tau^n)|_{z=\eta h^{n+1}/H_0}, \quad (2.69e)$$

$$\psi^* = \psi^{n,n+1} \equiv \frac{\Pi^n}{1 - W(\eta h^{n+1}/H_0, \tau^n)/L_w}, \quad (2.69f)$$

$$J^{n+1,n} = x_\xi^{n+1,n} z_\eta^{n+1} - x_\eta^{n+1,n} z_\xi^{n+1}, \quad (2.69g)$$

$$\begin{aligned} 0 = & \delta \int_0^{L_x} \left(\Pi^n \frac{h^{n+1} - h^n}{\Delta \tau} - \Pi^{n+1} \frac{h^{n+1}}{\Delta \tau} + (1 - \xi/L_w) h_\xi^{n+1} W_\tau^{n+1,n} \psi^* \right. \\ & \left. + x_\xi^{n+1,n} g \left(\frac{1}{2} (h^{n+1})^2 - H_0 h^{n+1} \right) \right) |_{\eta=H_0} d\xi \\ & + \int_0^{H_0} \left(z_\eta^* W_\tau^{n+1,n} (\psi^* \eta / H_0 + \varphi^*) + x_\eta^{n+1,n} g \left(\frac{1}{2} (z^{n+1})^2 - H_0 z^{n+1} \right) \right) |_{\xi=0} d\eta \\ & + \int_0^{L_x} \int_0^{H_0} \frac{1}{2 J^{n+1,n}} \left(((z_\eta^{n+1,n})^2 + (z_\xi^{n+1})^2) ((\eta/H_0) \partial_\xi \psi^* + \partial_\xi \varphi^*)^2 \right. \\ & \left. - 2(x_\xi^{n+1,n} z_\eta^{n+1,n} + z_\xi^{n+1} z_\eta^{n+1}) ((\eta/H_0) \partial_\xi \psi^* + \partial_\xi \varphi^*) (\psi^* / H_0 + \partial_\eta \varphi^*) \right. \\ & \left. + ((x_\xi^{n+1,n})^2 + (z_\xi^{n+1,n})^2) |\psi^* / H_0 + \partial_\eta \varphi^*|^2 \right) d\xi d\eta. \quad (2.69h) \end{aligned}$$

In the above, the definitions in the first lines are short-hands for substitution in the VP in the last line. The two weak formulations following from the variations $\delta \Pi^n, \delta \varphi^*$ of (2.69h) are solved in unison, solving h^{n+1} and φ^* , followed by the explicit solver step in the weak formulation resulting from the variation δh^{n+1} , yielding Π^{n+1} of (2.69h). These variations will be undertaken automatically within *Firedrake* as well as the associated iterative solver for the coupled two weak formulations and the explicit solver for the last

Challenges Firedrake via VPs

Computational-science challenges within Firedrake?

- ▶ Coupled interior-free-surface system, two sets of mixed partially canonical variables.
- ▶ Weak forms automatically generated: via *fd.derivatives* wrt $\{h^{n+1}(x), \varphi^{n+1}(x, z)\}$ for $\{\phi^{n+1}(x, H_0), \varphi(x, z)^{n+1}\}$, and via
- ▶ *fd.derivative* wrt $\phi^{n+1}(x, H_0)$ for $h^{n+1}(x)$.
- ▶ Use linearised system as stepping stone, then add nonlinearity, add piston wavemaker, add waveflap wavemaker,

Outlook

- ▶ Firedrake simulation for extreme 3-soliton amplifications: speed-increase needed & MPI.
- ▶ Efficiency gain via Firedrake implementation via VPs & *fd.derivative*.
- ▶ Moving meshes for water waves within VPs (B 2022), overturning & breaking waves, FSI.

References

- ▶ Gidel et al. 2022/23: *EarthArxiv link* (has GitHub link).
- ▶ Choi, B, Kalogirou, Kelmanson (2022) *Water Waves* **4** (has GitHub/Zenodo links).
- ▶ Gidel 2018: PhD thesis Leeds.
<https://etheses.whiterose.ac.uk/21730/>
- ▶ B, Kalogirou 2016: Variational Water Wave Modelling: from Continuum to Experiment. *Theory of Water Waves, London Math. Soc.* **426**. *preprint link*.
- ▶ Choi et al 2022: [Crossing seas](#) *YouTube movie –link*.