

$$\text{flux } F(A, s)|_{s_{k+1}} = F_{k+1}(\bar{A}_k^n, \bar{A}_{k+1}^n) \approx \frac{1}{\Delta t} \int_{t_n}^{t_{n+1}} F(t, s_{k+1}) dt$$

$$\text{so } \bar{A}_k^{n+1} = \bar{A}_k^n - \frac{1}{\Delta t} \int_{t_n}^{t_{n+1}} \cancel{F_{k+1}(\bar{A}_k^n, \bar{A}_{k+1}^n)} - \cancel{F_{k+1}(\bar{A}_k^n, \bar{A}_{k+1}^n)} dt \\ F(A, s)|_{s_{k+1}} - F(A, s)|_{s_{k+1}} \Delta t$$

$$\Rightarrow \bar{A}_k^{n+1} = \bar{A}_k^n - \frac{\Delta t}{\Delta x} (F_{k+1}(\bar{A}_k^n, \bar{A}_{k+1}^n) - F_k(\bar{A}_k^n, \bar{A}_{k+1}^n))$$

- 4) From question 2, we saw that the characteristic 'speed' $F'(t) > 0$, so the characteristics move to the right. So, the flux taken at a boundary $t + \frac{1}{2}$ is determined by the volume to the left of the boundary. Thus, the Godunov flux $F_{k+\frac{1}{2}} = F(\bar{A}_k^n)$ which is equal to ~~the flux from the Riemann problem solution~~ $F(A_k)$, using A_k ~~is~~ from the Riemann problem solution.

- 5) The CFL stability condition for the Godunov scheme above defines a timestep, Δt , that must satisfy:

$$\Delta t < \text{CFL} \cdot \min_k \frac{h_k}{|F'(A_k)|}$$

For $\text{CFL} \in (0, 1]$, and $\lambda_k = \cancel{F'(A_k)} > 0$ from question 2.

$$\text{so } \Delta t < \cancel{\text{CFL}} \cdot \min_k \frac{h_k}{F'(A_k)} \Leftrightarrow \text{our time-step restriction.}$$

- 6) At $s=0$, define our boundary condition, \bar{A}_{-1}^n , as our inflow area. Since $F'(1) > 0$, our characteristics move ~~backward~~ from left to right in the domain $s = [0, 1, \dots, L]$.

$$\text{Fix the flux at } s=0 \quad F_{-1+\frac{1}{2}}(\bar{A}_{-1}^n) = Q_0$$

∅

Since $\lambda(A) > 0$, giving us an upwind scheme, the characteristics simply leave the domain at $s=L$, so fixing the flux at the end of the domain would not have any effect on the rest of the domain. Thus, no boundary condition is needed at the downstream boundary.