

# Casey - Exercise 3

$$1. I[u] = \iint_{\Omega} \left( \frac{1}{2} |\nabla u|^2 - uf \right) dA$$

Variation of the functional  $I$  is defined as:

$$\delta I(u) = \frac{\partial I}{\partial \epsilon} = \lim_{\epsilon \rightarrow 0} \frac{I(u + \epsilon \delta u) - I(u)}{\epsilon} = 0$$

$$\Rightarrow I(u + \epsilon \delta u) = \frac{1}{2} \int_{\Omega} |\nabla(u + \epsilon \delta u)|^2 dA - \int_{\Omega} f(u + \epsilon \delta u) dA$$

$(dA = dx dy)$

Expanding  $|\nabla(u + \epsilon \delta u)|^2$

$$\begin{aligned} &= \nabla(u + \epsilon \delta u) \cdot \nabla(u + \epsilon \delta u) \\ &= \nabla u \cdot \nabla u + 2\epsilon (\nabla u \cdot \nabla \delta u) + \epsilon^2 (\nabla \delta u \cdot \nabla \delta u) \\ &= |\nabla u|^2 + 2\epsilon (\nabla u \cdot \nabla \delta u) + O(\epsilon^2) \end{aligned}$$

$$\Rightarrow I(u + \epsilon \delta u) = \frac{1}{2} \int_{\Omega} [|\nabla u|^2 + 2\epsilon (\nabla u \cdot \nabla \delta u)] dA - \int_{\Omega} f(u + \epsilon \delta u) dA$$

$$\begin{aligned} &= \left[ \frac{1}{2} \int_{\Omega} |\nabla u|^2 dA - \int_{\Omega} uf dA \right] + \epsilon \left[ \int_{\Omega} (\nabla u \cdot \nabla \delta u) dA - \int_{\Omega} f \delta u dA \right] \\ &= I[u] + \epsilon \left[ \int_{\Omega} (\nabla u \cdot \nabla \delta u) dA - \int_{\Omega} f \delta u dA \right] \end{aligned}$$

$$\text{Variation } \delta I[u] = \int_{\Omega} (\nabla u \cdot \nabla \delta u) dA - \int_{\Omega} f \delta u dA$$

Ritz-Galerkin principle - variation = 0 for all

admissible  $\delta u$

$$\rightarrow \int_{\Omega} \nabla u \cdot \nabla \delta u \, dA - \int_{\Omega} \delta u f \, dA = 0$$

$$\rightarrow \int_{\Omega} \nabla u \cdot \nabla \delta u \, dA = \int_{\Omega} \delta u f \, dA$$

LHS can be written as:

$$- \int_{\Omega} (\nabla^2 u) \delta u \, dA + \oint_{\partial\Omega} (\nabla u \cdot \hat{n}) \delta u \, ds$$

where  $\partial\Omega$  is the boundary

This yields:

$$- \int_{\Omega} (\nabla^2 u) \delta u \, dA + \oint_{\partial\Omega} (\nabla u \cdot \hat{n}) \delta u \, ds = \int_{\Omega} f \delta u \, dA$$

$$\Rightarrow \int_{\Omega} (-\nabla^2 u - f) \delta u \, dA + \oint_{\partial\Omega} (\nabla u \cdot \hat{n}) \delta u \, ds = 0$$

Given the boundary conditions, the boundary integral

= 0 for all boundaries

$$\Rightarrow \int_{\Omega} (-\nabla^2 u - f) \delta u \, dA = 0$$

This must hold for all admissible  $\delta u \Rightarrow -\nabla^2 u - f = 0$  to

satisfy, recovering the system.

## Conditions for $\delta u(x,y)$

- Must belong to the same function space as the test function  $w(x,y)$ .
- Must satisfy the form of the boundary conditions.  
 $\delta u(0,y) = 0, \delta u(1,y) = 0$

Weak formulation for test function  $w(x,y)$

$$\rightarrow \int_{\Omega} \nabla u \cdot \nabla w \, dA = \int_{\Omega} wf \, dA$$

which will yield the same result as for  $\delta u$

$$\Rightarrow w(x,y) = \delta u(x,y)$$

$$2. u(x,y) \sim u_n(x,y) = \sum_{j=1}^N U_j \phi_j(x,y)$$

where  $U_j$  are nodal coefficients

Substituting into the function I

$$I(u_n) = \frac{1}{2} \int_{\Omega} |\nabla u_n|^2 \, dA - \int_{\Omega} u_n f \, dA$$

$$I(u) = \frac{1}{2} \int_{\Omega} \left| \nabla \left( \sum_{j=1}^N u_j \phi_j \right) \right|^2 dA - \int_{\Omega} f \left( \sum_{j=1}^N u_j \phi_j \right) dA$$

Since coefficients are constant w.r.t integration

$$\Rightarrow I(u) = \frac{1}{2} \sum_{i=1}^N \sum_{j=1}^N u_i u_j \left( \int_{\Omega} \nabla \phi_i \cdot \nabla \phi_j dA \right) - \sum_{j=1}^N u_j \left( \int_{\Omega} f \phi_j dA \right)$$

Ritz-Galerkin principle states solutions vector  $\vec{U}$  minimises the functional  $I$

Take partials of  $I$  to find minimum:

$$\frac{\partial I(\vec{U})}{\partial U_k} = 0 \quad k = 1, \dots, N$$

$$\text{LHS : } \frac{\partial}{\partial U_k} \left[ \sum_{i=1}^N \sum_{j=1}^N u_i u_j \int_{\Omega} \nabla \phi_i \cdot \nabla \phi_j dA \right] = \sum_{i=1}^N u_i \int_{\Omega} \nabla \phi_k \cdot \nabla \phi_i dA$$

How is RHS dealt with?

$$\text{RHS : } \frac{\partial}{\partial U_k} \left[ \sum_{j=1}^N u_j F_j \right] = \int_{\Omega} f \phi_k dA$$

$$\Rightarrow \frac{\partial I(u)}{\partial U_k} = \sum_{j=1}^N u_j \int_{\Omega} \nabla \phi_k \cdot \nabla \phi_j dA - \int_{\Omega} f \phi_k dA = 0 \quad \text{for } k=1, \dots, N$$

$$\Rightarrow \sum_{j=1}^N u_j \int_{\Omega} \nabla \phi_k \cdot \nabla \phi_j dA = \int_{\Omega} f \phi_k dA = 0$$

$\delta u_k = 0$  for  $k \neq k'$ .

Can be written as discrete algebraic system

$$\text{where : } \vec{K} \vec{U} = \vec{F} \quad \text{where } \vec{K} = \int_{\Omega} \nabla \phi_i \cdot \nabla \phi_k dA, \vec{F} = \int_{\Omega} f \phi_k dA$$

and  $\vec{U}$  the vector of unknown coefficients

Local Global

N - number of nodes in x  
M - number of nodes in y

	0	1	2	3	4	5	6	7	8	9
0	0	1	2	3	4	5	6	7	8	9
1	1	0	N <sub>1</sub> <sub>2</sub>	N <sub>1</sub> <sub>3</sub>	N <sub>1</sub> <sub>4</sub>	N <sub>1</sub> <sub>5</sub>	N <sub>1</sub> <sub>6</sub>	N <sub>1</sub> <sub>7</sub>	N <sub>1</sub> <sub>8</sub>	N <sub>1</sub> <sub>9</sub>
2	2	3	0	1	2	3	4	5	6	7
3	3	2	1	0	1	2	3	4	5	6
4	4	5	6	7	8	9	0	1	2	3
5	5	6	7	8	9	0	1	2	3	4
6	6	7	8	9	0	1	2	3	4	5
7	7	8	9	0	1	2	3	4	5	6
8	8	9	0	1	2	3	4	5	6	7
9	9	0	1	2	3	4	5	6	7	8

	M <sub>N-1</sub>	M <sub>N</sub>	M <sub>N+1</sub>	M <sub>N+2</sub>	M <sub>N-1</sub>	M <sub>N</sub>	M <sub>N+1</sub>	M <sub>N+2</sub>	M <sub>N-1</sub>	M <sub>N</sub>
	0	1	2	3	4	5	6	7	8	9
0	0	1	2	3	4	5	6	7	8	9
1	1	0	N <sub>1</sub> <sub>2</sub>	N <sub>1</sub> <sub>3</sub>	N <sub>1</sub> <sub>4</sub>	N <sub>1</sub> <sub>5</sub>	N <sub>1</sub> <sub>6</sub>	N <sub>1</sub> <sub>7</sub>	N <sub>1</sub> <sub>8</sub>	N <sub>1</sub> <sub>9</sub>
2	2	3	0	1	2	3	4	5	6	7
3	3	2	1	0	1	2	3	4	5	6
4	4	5	6	7	8	9	0	1	2	3
5	5	6	7	8	9	0	1	2	3	4
6	6	7	8	9	0	1	2	3	4	5
7	7	8	9	0	1	2	3	4	5	6
8	8	9	0	1	2	3	4	5	6	7
9	9	0	1	2	3	4	5	6	7	8

## Uniform Mesh

3.

Matrix assembly:

For element k:

α and β representation

of local index

k representation of

global index

$$\Rightarrow i = \text{index}(k, \alpha)$$

$$j = \text{index}(k, \beta)$$

$$\Rightarrow A_{ij} = A_{ij} + \hat{A}_{\alpha\beta}$$

determined by element

+ locally determined

$$b_i = b_i + \hat{b}_{\alpha}$$

Based solely on local

index

Giving algebraic system

$\vec{AB}$  representing

the mesh

element wise

## Part 2

1. Weak formulation derivation:  $\times q(y)$  and integrate

$$\rightarrow \int_0^{Ly} q \frac{\partial}{\partial y} h_m dy - \int_0^{Ly} q \alpha \delta_y(h_m) dy = \int_0^{Ly} \frac{q^k}{\alpha_{per} \beta_e} dy$$

Integrating by parts (Second term)

$$\int_0^{Ly} q \alpha \delta_y(h_m) dy$$

$$\rightarrow \int_0^{Ly} \alpha g h_m \delta_y q \delta_y h_m dy + [\alpha g h_m \delta_y h_m]_0^{Ly}$$

At  $y = Ly$ ,  $\delta_y h_m = 0 \Rightarrow$  term  $\rightarrow 0$  (Eq. 12)

At  $y = 0$ , can be evaluated as  $\alpha g q(0) h_m \delta_y h_m = 0$

which can be rewritten to give  $\alpha g^2 \delta_y(h_m)^2 |_{Ly} = 0$

Using canal level ODE (Eq. 14)  $\rightarrow$  (Eq. 28)

Substituting back to weak form yields  
(Eq. 29) the weak formulation.

Using the FEM expansion for  $u_m$

$$u_m(y, t) \approx \sum u_i(t) \varphi_i(y) \text{ and } q = \varphi_i(y)$$

and considering time stepping:  $\Delta t h = \frac{u^{n+1} - u^n}{\Delta t}$

Algebraic system is as seen (Eq. 33).

### Time step restrictions

Considering system as:

$$\bar{\bar{M}}\vec{u}^{n+1} = \bar{\bar{M}}\vec{u}^n - \Delta t \bar{\bar{K}}(\vec{u}^n)\vec{u}^n$$

$\Delta t$  restriction can be derived from stability of the algebraic system

$$\rightarrow \vec{u}^{n+1} = (\bar{\bar{I}} - \Delta t \bar{\bar{M}}^{-1} \bar{\bar{K}}) \vec{u}^n$$

Eigenvalues of  $(\bar{\bar{I}} - \Delta t \bar{\bar{M}}^{-1} \bar{\bar{K}})$  must be  $\leq 1$

$$\Rightarrow \Delta t \leq \frac{2}{\lambda_{\max}(\bar{\bar{M}}^{-1} \bar{\bar{K}})}$$

This can be expressed in terms of the parameters of the problem

$$\rightarrow \Delta t \leq \frac{\text{Max } \sigma_e(\Delta y)^2}{\deg h_m}$$

where  $h_m = \max(h_m)$  in the domain.

