

# Casey - Exercise 3

$$1. I[u] = \int_{\Omega} \left( \frac{1}{2} |\nabla u|^2 - u f \right) d\Omega$$

Variation of the functional  $I$  is defined as:

$$\delta I(u) = \frac{\partial I}{\partial \epsilon} = \lim_{\epsilon \rightarrow 0} \frac{I(u + \epsilon \delta u) - I(u)}{\epsilon} = 0$$

$$\Rightarrow I(u + \epsilon \delta u) = \frac{1}{2} \int_{\Omega} |\nabla(u + \epsilon \delta u)|^2 dA - \int_{\Omega} f(u + \epsilon \delta u) dA$$

( $dA = dx dy$ )

Expanding  $|\nabla(u + \epsilon \delta u)|^2$

$$= \nabla(u + \epsilon \delta u) \cdot \nabla(u + \epsilon \delta u)$$

$$= \nabla u \cdot \nabla u + 2\epsilon (\nabla u \cdot \nabla \delta u) + \epsilon^2 (\nabla \delta u \cdot \nabla \delta u)$$

$$= |\nabla u|^2 + 2\epsilon (\nabla u \cdot \nabla \delta u) + O(\epsilon^2)$$

$$\Rightarrow I(u + \epsilon \delta u) = \frac{1}{2} \int_{\Omega} [|\nabla u|^2 + 2\epsilon (\nabla u \cdot \nabla \delta u)] dA - \int_{\Omega} f(u + \epsilon \delta u) dA$$

$$= \left[ \frac{1}{2} \int_{\Omega} |\nabla u|^2 - \int_{\Omega} u f dA \right] + \epsilon \left[ \int_{\Omega} (\nabla u \cdot \nabla \delta u) dA - \int_{\Omega} \delta u f dA \right]$$

$$= I[u] + \epsilon \left[ \int_{\Omega} (\nabla u \cdot \nabla \delta u) dA - \int_{\Omega} \delta u f dA \right]$$

$$\text{Variation } \delta I[u] = \int_{\Omega} (\nabla u \cdot \nabla \delta u) dA - \int_{\Omega} \delta u f dA$$

Ritz-Galerkin principle - variation = 0 for all admissible  $\delta u$

$$\rightarrow \int_{\Omega} \nabla u \cdot \nabla \delta u \, dA - \int_{\Omega} \delta u f \, dA = 0$$

$$\rightarrow \int_{\Omega} \nabla u \cdot \nabla \delta u \, dA = \int_{\Omega} \delta u f \, dA$$

LHS can be written as:

$$- \int_{\Omega} (\nabla^2 u) \delta u \, dA + \oint_{\partial\Omega} (\nabla u \cdot \vec{n}) \delta u \, ds$$

Where  $d\Omega$  is the boundary

This yields:

$$- \int_{\Omega} (\nabla^2 u) \delta u \, dA + \oint_{\partial\Omega} (\nabla u \cdot \vec{n}) \delta u \, ds = \int_{\Omega} f \delta u \, dA$$

$$\Rightarrow \int_{\Omega} (-\nabla^2 u - f) \delta u \, dA + \oint_{\partial\Omega} (\nabla u \cdot \vec{n}) \delta u \, ds = 0$$

Bit more detail on how BC's applied.

Given the boundary conditions, the boundary integral = 0 for all boundaries

$$\Rightarrow \int_{\Omega} (-\nabla^2 u - f) \delta u \, dA = 0$$

This must hold for all admissible  $\delta u \Rightarrow -\nabla^2 u - f = 0$  to satisfy, recovering the system.

## Conditions for $\delta u(x, y)$

- Must belong to the same function space as the test function  $w(x, y)$ .
- Must satisfy the form of the Boundary Conditions.  
 $\delta u(0, y) = 0$  ,  $\delta u(1, y) = 0$

Weak formulation for test function  $w(x, y)$

$$\rightarrow \int_{\Omega} \nabla u \cdot \nabla w \, dA = \int_{\Omega} w f \, dA$$

Which will yield the same result as for  $\delta u$   
 $\Rightarrow w(x, y) = \delta u(x, y)$

$$2. \quad u(x, y) \sim u_n(x, y) = \sum_{j=1}^N U_j \phi_j(x, y)$$

where  $U_j$  are nodal coefficients

Substituting into the functional  $I$

$$I(u_n) = \frac{1}{2} \int_{\Omega} |\nabla u_n|^2 \, dA - \int_{\Omega} u_n f \, dA$$

$$I(u) = \frac{1}{2} \int_{\Omega} \left| \nabla \left( \sum_{j=1}^N U_j \phi_j \right) \right|^2 dA - \int_{\Omega} f \left( \sum_{j=1}^N U_j \phi_j \right) dA$$

Since coefficients are constant wrt integration

$$\Rightarrow I(u) = \frac{1}{2} \sum_{i=1}^N \sum_{j=1}^N U_i U_j \left( \int_{\Omega} \nabla \phi_i \cdot \nabla \phi_j dA \right) - \sum_{j=1}^N U_j \left( \int_{\Omega} f \phi_j dA \right)$$

Ritz - Galerkin principle states solution vector  $\vec{U}$  minimises the functional  $I$

Take partials of  $I$  to find minimum:

$$\frac{\partial I(\vec{U})}{\partial U_k} = 0 \quad k = 1, \dots, N$$

$$\text{LHS: } \frac{\partial}{\partial U_k} \left[ \frac{1}{2} \sum_{i=1}^N \sum_{j=1}^N U_i U_j \int_{\Omega} \nabla \phi_i \cdot \nabla \phi_j dA \right] = \sum_{j=1}^N U_j \int_{\Omega} \nabla \phi_k \cdot \nabla \phi_j dA$$

How is IHS dealt with?

$$\text{RHS: } \frac{\partial}{\partial U_k} \left[ \sum_{j=1}^N U_j F_j \right] = \int_{\Omega} f \phi_k dA$$

$$\Rightarrow \frac{\partial I(u)}{\partial U_k} = \sum_{j=1}^N U_j \int_{\Omega} \nabla \phi_k \cdot \nabla \phi_j dA - \int_{\Omega} f \phi_k dA = 0 \quad \text{for } k=1, \dots, N$$

$$\Rightarrow \sum_{j=1}^N U_j \int_{\Omega} \nabla \phi_k \cdot \nabla \phi_j dA = \int_{\Omega} f \phi_k dA = 0$$

$\Delta U_k = 0$  for  $k \neq k'$

Can be written as discrete algebraic system

where:  $\vec{K} \vec{U} = \vec{F}$  where  $\vec{K} = \int_{\Omega} \nabla \phi_i \cdot \nabla \phi_k dA$ ,  $\vec{F} = \int_{\Omega} f \phi_k dA$

and  $\vec{U}$  the vector of unknown coefficients

local  
K values  
global

$N$  - number of nodes in  $x$   
 $M$  - number of nodes in  $y$

	0	1	2	3	4	5	...	$N-2$	$N-1$	$N$
0	1	2	3	4	5	...	...	$N-1$	$N$	$2N$
1	6	7	8	9	10	...	...	$2N-1$	$2N$	$2N+1$
2	11	12	13	14	15	...	...	$2N+1$	$2N+2$	$2N+3$
...	...	...	...	...	...	...	...	...	...	...
$N-1$	$(N-1)$	$(N)$	$(N+1)$	$(N+2)$	$(N+3)$	...	...	$(2N-1)$	$(2N)$	$(2N+1)$
$N$	$(2N)$	$(2N+1)$	$(2N+2)$	$(2N+3)$	$(2N+4)$	...	...	$(4N-1)$	$(4N)$	$(4N+1)$
$2N+1$	$(4N+1)$	$(4N+2)$	$(4N+3)$	$(4N+4)$	$(4N+5)$	...	...	$(6N-1)$	$(6N)$	$(6N+1)$
...	...	...	...	...	...	...	...	...	...	...
$4N+1$	$(6N+1)$	$(6N+2)$	$(6N+3)$	$(6N+4)$	$(6N+5)$	...	...	$(8N-1)$	$(8N)$	$(8N+1)$
$5N+1$	$(8N+1)$	$(8N+2)$	$(8N+3)$	$(8N+4)$	$(8N+5)$	...	...	$(10N-1)$	$(10N)$	$(10N+1)$
...	...	...	...	...	...	...	...	...	...	...
$MN-N$	$(M-1)N$	$(M)N$	$(M+1)N$	$(M+2)N$	$(M+3)N$	...	...	$(M-2)N$	$(M-1)N$	$MN$
$MN-N+1$	$MN$	$(M)N+1$	$(M)N+2$	$(M)N+3$	$(M)N+4$	...	...	$MN-1$	$MN$	$(M+1)N$

Uniform Mesh

3.

Matrix assembly:

For element  $k$ :

$\alpha$  and  $\beta$  representative

of local index

$k$  representative of

global index

$$\Rightarrow i = \text{index}(k, \alpha)$$

$$j = \text{index}(k, \beta)$$

$$\Rightarrow A_{ij} = A_{ij} + \hat{A}_{\alpha\beta}$$

(determined by element)

+ locally determined

$$b_i = b_i + \hat{b}_\alpha$$

Based solely on local

index.

Giving algebraic system

$\bar{A}\bar{b}$  representing

the mesh

element wise.

## Part 2

1. Weak formulation derivation:  $\times q(y)$  and integrate

$$\rightarrow \int_0^{L_y} q \partial_y h_m dy - \int_0^{L_y} q \alpha g dy (h_m \partial_y h_m) dy = \int_0^{L_y} \frac{q R}{\eta_{gr} \delta \epsilon} dy$$

Integrating by parts (second term)

$$\int_0^{L_y} q \alpha g dy (h_m \partial_y h_m) dy$$

$$\rightarrow \int_0^{L_y} \alpha g h_m \partial_y q \partial_y h_m dy + [\alpha g q h_m \partial_y h_m]_0^{L_y}$$

At  $y = L_y$ ,  $\partial_y h_m = 0 \Rightarrow$  term  $\rightarrow 0$  (Eq. 12)

At  $y = 0$ , can be evaluated as  $\alpha g q(0) h_m \partial_y h_m|_{y=0}$   
which can be rewritten to give  $\alpha g \frac{1}{2} \partial_y (h_m)^2|_{y=0} = 0$

Using canal level ODE (Eq. 14)  $\rightarrow$  (Eq. 28)

Substituting back to weak form yields  
(Eq. 29) the weak formulation.

Using the FEM expansion for  $u_m$   
 $u_m(y, t) \approx \sum h_j(t) \varphi_j(y)$  and  $z = \varphi_j(y)$   
 and considering time stepping:  $\Delta t = \frac{u^{n+1} - u^n}{\Delta t}$   
 Algebraic system is as seen (Eq. 33).

### Time step restriction

Considering system as:

$$\bar{M} \dot{\vec{u}}^{n+1} = \bar{M} \vec{u}^n - \Delta t \bar{K} (\vec{u}^n) \vec{u}^n$$

$\Delta t$  restriction can be derived from stability  
 of the algebraic system

$$\Rightarrow \vec{u}^{n+1} = (\bar{I} - \Delta t \bar{M}^{-1} \bar{K}) \vec{u}^n$$

Eigenvalues of  $(\bar{I} - \Delta t \bar{M}^{-1} \bar{K})$  must be  $\leq 1$

$$\Rightarrow \Delta t \leq \frac{2}{\lambda_{\max}(\bar{M}^{-1} \bar{K})}$$

This can be expressed in terms of the parameters  
 of the problem

$$\Rightarrow \Delta t \leq \frac{\max \sigma_e(\Delta y)^2}{\alpha g h_m}$$

where  $h_m = \max(h_m)$  in the domain.

