

Casey - Exercise 3

$$1. I[u] = \int_{\Omega} \left(\frac{1}{2} |\nabla u|^2 - u f \right) d\Omega$$

Variation of the functional I is defined as:

$$\delta I(u) = \frac{\partial I}{\partial \epsilon} = \lim_{\epsilon \rightarrow 0} \frac{I(u + \epsilon \delta u) - I(u)}{\epsilon} = 0$$

$$\Rightarrow I(u + \epsilon \delta u) = \frac{1}{2} \int_{\Omega} |\nabla(u + \epsilon \delta u)|^2 dA - \int_{\Omega} f(u + \epsilon \delta u) dA$$

($dA = dx dy$)

Expanding $|\nabla(u + \epsilon \delta u)|^2$

$$= \nabla(u + \epsilon \delta u) \cdot \nabla(u + \epsilon \delta u)$$

$$= \nabla u \cdot \nabla u + 2\epsilon (\nabla u \cdot \nabla \delta u) + \epsilon^2 (\nabla \delta u \cdot \nabla \delta u)$$

$$= |\nabla u|^2 + 2\epsilon (\nabla u \cdot \nabla \delta u) + O(\epsilon^2)$$

$$\Rightarrow I(u + \epsilon \delta u) = \frac{1}{2} \int_{\Omega} [|\nabla u|^2 + 2\epsilon (\nabla u \cdot \nabla \delta u)] dA - \int_{\Omega} f(u + \epsilon \delta u) dA$$

$$= \left[\frac{1}{2} \int_{\Omega} |\nabla u|^2 - \int_{\Omega} u f dA \right] + \epsilon \left[\int_{\Omega} (\nabla u \cdot \nabla \delta u) dA - \int_{\Omega} \delta u f dA \right]$$

$$= I[u] + \epsilon \left[\int_{\Omega} (\nabla u \cdot \nabla \delta u) dA - \int_{\Omega} \delta u f dA \right]$$

$$\text{Variation } \delta I[u] = \int_{\Omega} (\nabla u \cdot \nabla \delta u) dA - \int_{\Omega} \delta u f dA$$

Ritz-Galerkin principle - variation = 0 for all admissible δu

$$\rightarrow \int_{\Omega} \nabla u \cdot \nabla \delta u \, dA - \int_{\Omega} \delta u f \, dA = 0$$

$$\rightarrow \int_{\Omega} \nabla u \cdot \nabla \delta u \, dA = \int_{\Omega} \delta u f \, dA$$

LHS can be written as:

$$- \int_{\Omega} (\nabla^2 u) \delta u \, dA + \oint_{\partial\Omega} (\nabla u \cdot \vec{n}) \delta u \, ds$$

Where $d\Omega$ is the boundary

This yields:

$$- \int_{\Omega} (\nabla^2 u) \delta u \, dA + \oint_{\partial\Omega} (\nabla u \cdot \vec{n}) \delta u \, ds = \int_{\Omega} f \delta u \, dA$$

$$\Rightarrow \int_{\Omega} (-\nabla^2 u - f) \delta u \, dA + \oint_{\partial\Omega} (\nabla u \cdot \vec{n}) \delta u \, ds = 0$$

Given the boundary conditions, the boundary integral = 0 for all boundaries

$$\Rightarrow \int_{\Omega} (-\nabla^2 u - f) \delta u \, dA = 0$$

This must hold for all admissible $\delta u \Rightarrow -\nabla^2 u - f = 0$ to satisfy, recovering the system.

Conditions for $\delta u(x, y)$

- Must belong to the same function space as the test function $w(x, y)$.
- Must satisfy the form of the Boundary Conditions.
 $\delta u(0, y) = 0$, $\delta u(1, y) = 0$

Weak formulation for test function $w(x, y)$

$$\rightarrow \int_{\Omega} \nabla u \cdot \nabla w \, dA = \int_{\Omega} w f \, dA$$

Which will yield the same result as for δu
 $\Rightarrow w(x, y) = \delta u(x, y)$

$$2. \quad u(x, y) \sim u_n(x, y) = \sum_{j=1}^N U_j \phi_j(x, y)$$

where U_j are nodal coefficients

Substituting into the functional \mathcal{I}

$$\mathcal{I}(u_n) = \frac{1}{2} \int_{\Omega} |\nabla u_n|^2 \, dA - \int_{\Omega} u_n f \, dA$$

$$I(u) = \frac{1}{2} \int_{\Omega} \left| \nabla \left(\sum_{j=1}^N U_j \phi_j \right) \right|^2 dA - \int_{\Omega} f \left(\sum_{j=1}^N U_j \phi_j \right) dA$$

Since coefficients are constant wrt integration

$$\Rightarrow I(u) = \frac{1}{2} \sum_{i=1}^N \sum_{j=1}^N U_i U_j \left(\int_{\Omega} \nabla \phi_i \cdot \nabla \phi_j dA \right) - \sum_{j=1}^N U_j \left(\int_{\Omega} f \phi_j dA \right)$$

Ritz - Galerkin principle states solution vector \vec{U} minimises the functional I

Take partials of I to find minimum:

$$\frac{\partial I(\vec{U})}{\partial U_k} = 0 \quad k = 1, \dots, N$$

$$\text{LHS: } \frac{\partial}{\partial U_k} \left[\frac{1}{2} \sum_{i=1}^N \sum_{j=1}^N U_i U_j \int_{\Omega} \nabla \phi_i \cdot \nabla \phi_j dA \right] = \sum_{j=1}^N U_j \int_{\Omega} \nabla \phi_k \cdot \nabla \phi_j dA$$

$$\text{RHS: } \frac{\partial}{\partial U_k} \left[\sum_{j=1}^N U_j F_j \right] = \int_{\Omega} f \phi_k dA$$

$$\Rightarrow \frac{\partial I(u)}{\partial U_k} = \sum_{j=1}^N U_j \int_{\Omega} \nabla \phi_k \cdot \nabla \phi_j dA - \int_{\Omega} f \phi_k dA = 0 \quad \text{for } k=1, \dots, N$$

$$\Rightarrow \sum_{j=1}^N U_j \int_{\Omega} \nabla \phi_k \cdot \nabla \phi_j dA = \int_{\Omega} f \phi_k dA = 0$$

Can be written as discrete algebraic system

where: $\vec{K} \vec{U} = \vec{F}$ where $\vec{K} = \int_{\Omega} \nabla \phi_i \cdot \nabla \phi_k dA$, $\vec{F} = \int_{\Omega} f \phi_k dA$
and \vec{U} the vector of unknown coefficients

local
K values
global

N - number of nodes in x
 M - number of nodes in y

	0	1	2	3	4	5	...	$N-2$	$N-1$	N
0	1	2	3	4	5	$N-1$	N	$2N$
1	6	7	8	9	10	$2N-1$	$2N$	$2N+1$
2	11	12	13	14	15	$2N+1$	$2N+2$	$2N+3$
...
$N-1$	N	$N+1$	$N+2$	$N+3$	$N+4$	$2N-1$	$2N$	$2N+1$
N	$N+1$	$N+2$	$N+3$	$N+4$	$N+5$	$2N$	$2N+1$	$2N+2$
...
$M-1$	$(M-1)N$	$(M-1)N+1$	$(M-1)N+2$	$(M-1)N+3$	$(M-1)N+4$	$(M-2)N-1$	$(M-2)N$	$(M-2)N+1$
M	$(M-1)N+1$	$(M-1)N+2$	$(M-1)N+3$	$(M-1)N+4$	$(M-1)N+5$	$(M-2)N$	$(M-2)N+1$	$(M-2)N+2$
...
$MN-1$	$(M-1)N$	$(M-1)N+1$	$(M-1)N+2$	$(M-1)N+3$	$(M-1)N+4$	$(M-2)N-1$	$(M-2)N$	$(M-2)N+1$
MN	$(M-1)N+1$	$(M-1)N+2$	$(M-1)N+3$	$(M-1)N+4$	$(M-1)N+5$	$(M-2)N$	$(M-2)N+1$	$(M-2)N+2$
...
$MN-1$	$(M-1)N$	$(M-1)N+1$	$(M-1)N+2$	$(M-1)N+3$	$(M-1)N+4$	$(M-2)N-1$	$(M-2)N$	$(M-2)N+1$
MN	$(M-1)N+1$	$(M-1)N+2$	$(M-1)N+3$	$(M-1)N+4$	$(M-1)N+5$	$(M-2)N$	$(M-2)N+1$	$(M-2)N+2$

Uniform Mesh

3.

Matrix assembly:

For element k :

α and β representative

of local index

k representative of

global index

$$\Rightarrow i = \text{index}(k, \alpha)$$

$$j = \text{index}(k, \beta)$$

$$\Rightarrow A_{ij} = A_{ij} + \hat{A}_{\alpha\beta}$$

determined by element

+ locally determined

$$b_i = b_i + \hat{b}_\alpha$$

Based solely on local

index.

Giving algebraic system

$\bar{A}\bar{b}$ representing

the mesh

element wise.

Part 2

1. Weak formulation derivation: $\times q(y)$ and integrate

$$\rightarrow \int_0^{L_y} q \partial_y h_m dy - \int_0^{L_y} q \alpha g dy (h_m \partial_y h_m) dy = \int_0^{L_y} \frac{q R}{\rho g \beta_e} dy$$

Integrating by parts (second term)

$$\int_0^{L_y} q \alpha g dy (h_m \partial_y h_m) dy$$

$$\rightarrow \int_0^{L_y} \alpha g h_m \partial_y q \partial_y h_m dy + [\alpha g q h_m \partial_y h_m]_0^{L_y}$$

At $y = L_y$, $\partial_y h_m = 0 \Rightarrow$ term $\rightarrow 0$ (Eq. 12)

At $y = 0$, can be evaluated as $\alpha g q(0) h_m \partial_y h_m|_{y=0}$
which can be rewritten to give $\alpha g \frac{1}{2} \partial_y (h_m)^2|_{y=0} = 0$

Using canal level ODE (Eq. 14) \rightarrow (Eq. 28)

Substituting back to weak form yields
(Eq. 29) the weak formulation.

Using the FEM expansion for u_m
 $u_m(y, t) \approx \sum h_j(t) \varphi_j(y)$ and $z = \varphi_j(y)$
 and considering time stepping: $\Delta t = \frac{u^{n+1} - u^n}{\Delta t}$
 Algebraic system is as seen (Eq. 33).

Time step restriction

Considering system as:

$$\bar{M} \dot{\vec{u}}^{n+1} = \bar{M} \vec{u}^n - \Delta t \bar{K} (\vec{u}^n) \vec{u}^n$$

Δt restriction can be derived from stability
 of the algebraic system

$$\Rightarrow \vec{u}^{n+1} = (\bar{I} - \Delta t \bar{M}^{-1} \bar{K}) \vec{u}^n$$

Eigenvalues of $(\bar{I} - \Delta t \bar{M}^{-1} \bar{K})$ must be ≤ 1

$$\Rightarrow \Delta t \leq \frac{2}{\lambda_{\max}(\bar{M}^{-1} \bar{K})}$$

This can be expressed in terms of the parameters
 of the problem

$$\Rightarrow \Delta t \leq \frac{\max \sigma_e(\Delta y)^2}{\alpha g h_m}$$

where $h_m = \max(h_m)$ in the domain.

