

$$\psi(n) = \langle n | \psi \rangle$$

Harmonic oscillator:

Compare: $\psi_n(x) = N_n H_n(x/\lambda_{osc}) e^{-\frac{1}{2}(\frac{x}{\lambda_{osc}})^2}$
 $= \langle x | \psi_n \rangle$

completeness of

- charge basis

$$1 = \sum_{n=-\infty}^{+\infty} |n\rangle \langle n|$$

- phase basis

$$1 = \int_0^{2\pi} d\varphi |\varphi\rangle \langle \varphi|$$

completeness of the position basis

$$1 = \int_{-\infty}^{+\infty} dx |x\rangle \langle x|$$

$$1 = \int dp |p\rangle \langle p|$$

$$\langle x | p \rangle = \frac{1}{\sqrt{v}} e^{ipx}$$

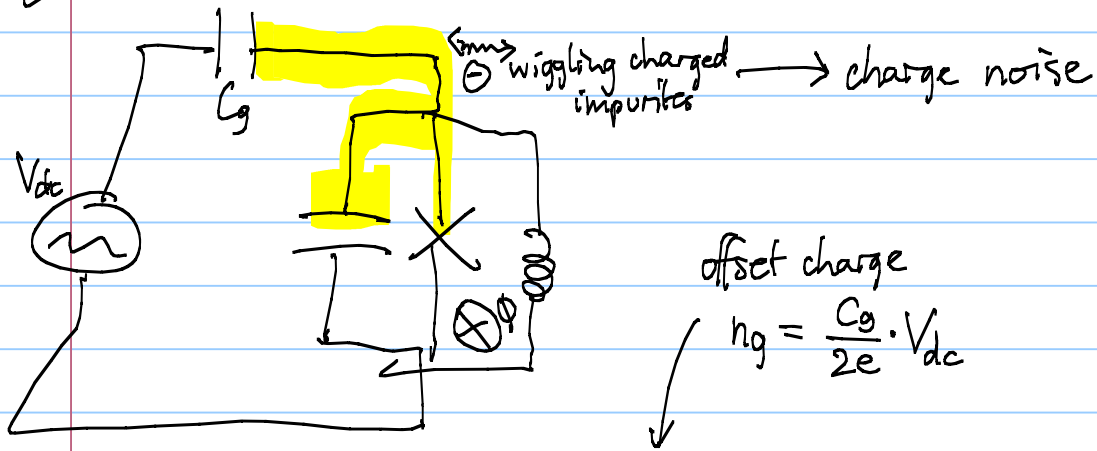
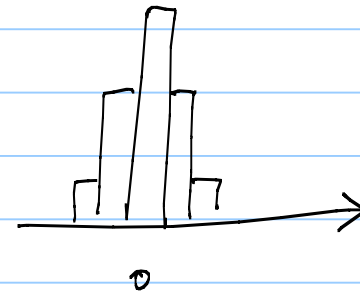
$$\psi(\varphi) = \langle \varphi | \psi \rangle = \underbrace{\langle \varphi | \left(\sum_{n=-\infty}^{+\infty} |n\rangle \langle n| \right) | \psi \rangle}_1 = \sum_{n=-\infty}^{+\infty} \underbrace{\langle \varphi | n \rangle}_{\frac{1}{\sqrt{2\pi}} e^{i\varphi n}} \underbrace{\langle n | \psi \rangle}_{\psi(n)}$$

$$f(\varphi)$$

To do

- * Plot wavefunctions for CPB
- * Numerics for the CPBL

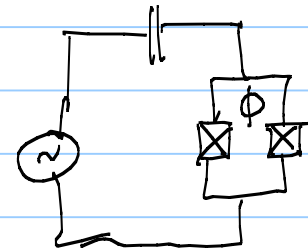
$$-i \frac{d}{d\varphi} f(\varphi) = n f(\varphi)$$



offset charge

$$n_g = \frac{C_g}{2e} \cdot V_{dc}$$

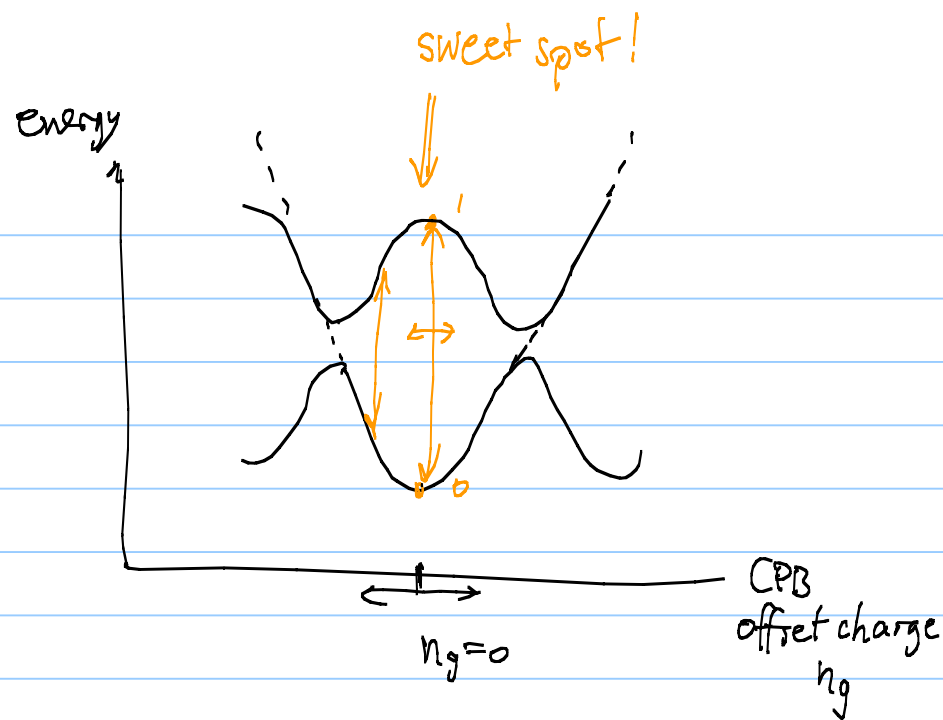
CPB with 2 junctions



$$H = 4E_c (n - n_g)^2 - E_J \cos \varphi + \frac{1}{2} E_L \left(\varphi - \frac{2\pi \Phi}{\Phi_0} \right)^2$$

skipping n_g

$$= -4E_c \frac{d^2}{d\varphi^2} - E_J \cos \varphi + \frac{1}{2} E_L \left(\varphi - \frac{2\pi \Phi}{\Phi_0} \right)^2$$



Variations in transition freq. (E_{01})
result in
"dephasing"

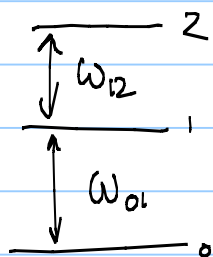
$$H|\psi\rangle = E|\psi\rangle$$

$$i\hbar \frac{d}{dt} |\psi\rangle = H|\psi\rangle$$

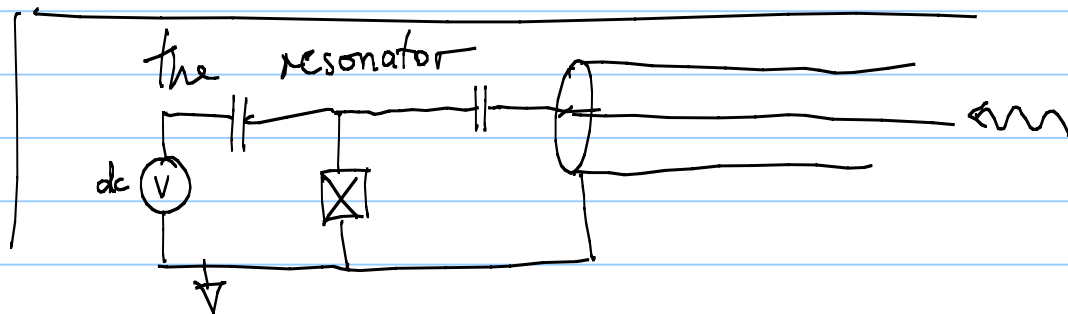
$$|\psi(t)\rangle = e^{-iEt/\hbar} |\psi\rangle$$

this phase factor
becomes unpredictable

Anharmonicity



$$\alpha = \omega_{12} - \omega_{01}$$



$$H = 4E_c (n - n_g)^2 - E_J \cos \varphi + \frac{1}{2} E_L \left(\varphi - \frac{2\pi\Phi}{\Phi_0} \right)^2$$

skipping n_g

$$= -4E_c \frac{d^2}{d\varphi^2} - E_J \cos \varphi + \frac{1}{2} E_L \left(\varphi - \frac{2\pi\Phi}{\Phi_0} \right)^2$$

Step 0: Shift φ :

$$\varphi' = \varphi - \frac{2\pi\Phi}{\Phi_0}$$

↓
call it φ ...

Step 1: Harmonic oscillator

$$H_0 = -4E_c \frac{d^2}{d\varphi^2} + \frac{1}{2} E_L \left(\varphi - \frac{2\pi\Phi}{\Phi_0} \right)^2$$

for later: $V = -E_J \cos \left(\varphi + \frac{2\pi\Phi}{\Phi_0} \right)$

$$a \sim \left(\frac{d}{d\varphi} + \varphi \right)$$

diagonalize H_0

$$\hookrightarrow H_0 = \hbar\omega_0 \left(a^\dagger a + \frac{1}{2} \right)$$

you should get: $\hbar\omega_0 = \sqrt{8E_L E_c}$, I hope!

In the h.o. basis:

$$H_0 = \hbar\omega_0 \begin{pmatrix} \frac{1}{2} & & & 0 \\ & \frac{3}{2} & & \\ & & \frac{5}{2} & \\ 0 & & & \ddots \end{pmatrix}$$

$$H_0 = \hbar\omega_0 (a^\dagger a + \frac{1}{2})$$

Step 2 : What to do with V ?

Let's call the eigenstates of the harm. osc. $|m\rangle$

$$\text{So } H_0 |m\rangle = E_m |m\rangle = \hbar\omega_0 (m + \frac{1}{2}) |m\rangle$$

$$(H_0)_{mm'} = \langle m | H_0 | m' \rangle = \delta_{mm'} \hbar\omega_0 (m + \frac{1}{2})$$

$$(H)_{mm'} = (H_0 + V)_{mm'} = \underbrace{\langle m | H_0 | m' \rangle}_{\text{this we know!}} + \langle m | V | m' \rangle$$

$$\langle m | V | m' \rangle = -E_J \int_{-\infty}^{+\infty} d\varphi \mathcal{N}_m \mathcal{N}_{m'} \cos\left(\varphi + \frac{2\pi\Phi}{\Phi_0}\right) H_m(\varphi/\varphi_{\text{osc}}) H_{m'}(\varphi/\varphi_{\text{osc}}) \times \exp\left[-\left(\frac{\varphi}{\varphi_{\text{osc}}}\right)^2\right]$$

like m, m'

$$\langle i | \cos \varphi | i + 2j \rangle$$

$$= (-2)^m \left[\frac{n!}{(n+2m)!} \right]^{1/2} \varphi_0^{2m} e^{-\varphi_0^2/4} L_n^{2m}(\varphi_0^2/2)$$

$$\langle i | \sin \varphi | i + 2j + 1 \rangle \quad (A3)$$

$$= (-2)^m \left[\frac{n!}{2(n+2m+1)!} \right]^{1/2} \varphi_0^{2m+1} e^{-\varphi_0^2/4} L_n^{2m+1}(\varphi_0^2/2),$$

$$\varphi_0 = \varphi_{\text{osc}} = (8E_C/E_L)^{1/4}$$

generalized Laguerre polynomial
(A2)