## Coordinates for Tip-Semiconductor System in SEMITIP VERSION 3

In SEMITIP VERSION 3 and higher, a new set of coordinates in the vacuum is used that are a generalization of the usual prolate spheroidal coordinates. The new coordinates  $\xi$  and  $\eta$  in the vacuum are related to the cylindrical coordinates r and z by

$$\frac{r^2}{\xi^2 - 1} + \frac{(z - ac\eta)^2}{\xi^2} = a^2,$$
 (1a)

$$-\frac{r^2}{1-\eta^2} + \frac{(z - ac\eta)^2}{\eta^2} = a^2$$
 (1b)

where  $c \equiv z_0/(a\eta_T)$  with  $\eta_T$  being the  $\eta$  value defining the hyperboloid that corresponds to the boundary of the probe tip and  $z_0$  being the center point of this hyperboloid. The values of  $\eta$  thus run from 0 on the surface to  $\eta_T$  at the tip. The value of  $\xi$  is 1 on the central axis and increase with distance away from that axis. For  $z_0=0$  these equations reduce to the standard definition of prolate spheroidal coordinates. With specified values for R, b, and s we have  $\eta_T=1/\sqrt{1+b^{-2}}$ ,  $a=Rb^2/\eta_T$ , and  $z_0=s-a\eta_T$ . The inverse equations to (1a) and (1b) are

$$z = a\xi\eta + ac\eta = a(\xi + c)\eta, \qquad (2a)$$

$$r = a[(\xi^2 - 1)(1 - \eta^2)]^{1/2}.$$
 (2b)

Laplace's equation for the electrostatic potential energy  $\phi$  in the vacuum is found to be

$$f_{1}(\xi,\eta)\frac{\partial^{2}\phi}{\partial\xi^{2}} + f_{2}(\xi,\eta)\frac{\partial^{2}\phi}{\partial\eta^{2}} + f_{3}(\xi,\eta)\frac{\partial^{2}\phi}{\partial\theta^{2}} + f_{4}(\xi,\eta)\frac{\partial^{2}\phi}{\partial\xi\partial\eta} + f_{5}(\xi,\eta)\frac{\partial\phi}{\partial\xi} + f_{6}(\xi,\eta)\frac{\partial\phi}{\partial\eta} = 0$$
(3a)

where

$$f_1(\xi,\eta) = \frac{(\xi^2 - 1)[(\xi + c)^2 - \eta^2 (2c\xi + c^2 + 1)]}{\xi(\xi + c) - \eta^2 (c\xi + 1)},$$
 (3b)

$$f_2(\xi, \eta) = \frac{(1 - \eta^2)(\xi^2 - \eta^2)}{\xi(\xi + c) - \eta^2(c\xi + 1)},$$
(3c)

$$f_3(\xi, \eta) = \frac{\xi(\xi + c) - \eta^2(c\xi + 1)}{(\xi^2 - 1)(1 - \eta^2)},$$
(3d)

$$f_4(\xi, \eta) = \frac{-2c\eta(\xi^2 - 1)(1 - \eta^2)}{\xi(\xi + c) - \eta^2(c\xi + 1)},$$
(3e)

$$f_5(\xi, \eta) = \frac{g_5(\xi, \eta)}{\left[\xi(\xi + c) - \eta^2(c\xi + 1)\right]^2},$$
(3f)

and

$$f_6(\xi, \eta) = \frac{g_6(\xi, \eta)}{\left[\xi(\xi + c) - \eta^2(c\xi + 1)\right]^2},$$
 (3g)

with

$$g_{5}(\xi,\eta) = c^{3} + 3c^{2}\xi + c(2+c^{2})\xi^{2} + 3c^{2}\xi^{3} + 4c\xi^{4} + 2\xi^{5} + \eta^{4}[c^{3} + (2+3c^{2})\xi + c(6+c^{2})\xi^{2} + 3c^{2}\xi^{3}] - (3h)$$

$$2\eta^{2}[c(c^{2} - 1) + 3c^{2}\xi + c(6+c^{2})\xi^{2} + (2+3c^{2})\xi^{3} + c\xi^{4}]$$

and

$$g_{6}(\xi,\eta) = -\eta \{c^{2} + 4c\xi + c^{2}\xi^{2} + 2\xi^{4} + \eta^{4}[2 + c^{2} + 4c\xi + c^{2}\xi^{2}] - 2\eta^{2}[c^{2} + 4c\xi + (2 + c^{2})\xi^{2}]\}.$$
(3i)

In VERSIONS 3, 4, and 5 there is not  $\theta$  dependence, so the second derivative with respect to  $\theta$  [third term of Eq. (3a)] is zero.

Grid points are labeled by (i,j) with i referring to the radial direction and j referring to the direction perpendicular to the surface. In the new coordinate system for the vacuum the grid points are given by  $(\xi_i, \eta_j)$ . The spacing between the  $\eta_j$  values is chosen to be simply  $\Delta \eta = \eta_T / \text{NV}$ , so that  $\eta_j = j (\eta_T / \text{NV})$ . The  $\xi_i$  values can be obtained by noting that the radial grid values at the surface are matched between the vacuum and the grid. Thus, using the  $R_i$  values output by the program and with j = 0 referring to the surface, we have  $r_{i0} = R_i$ . On the surface,  $\eta = 0$ , so that  $r_{i0} = R_i = a(\xi_i^2 - 1)^{1/2}$  so that  $\xi_i = \{1 + [R_i / a]^2\}^{1/2}$ . The particular values of  $r_{ij}$  and  $z_{ij}$  corresponding to the grid points can therefore be obtained by:

- (i) for  $z_{ij}$  we have  $z_{ij} = a(\xi_i + c)\eta_j = ([a^2 + R_i^2]^{1/2} + ac)j\eta_T / NV$ . The program defines DELV<sub>i</sub> as  $([a^2 + R(I)^2]^{1/2} + ac)\eta_T / NV$ , so that  $z_{ij} = j \times DELV_i$ .
- (ii) for the  $r_{ij}$  values, we again note that the radial values at the surface are  $R_i = a(\xi_i^2 1)^{1/2}$ . Then, away from the surface we have  $r_{ij} = a[(\xi_i^2 1)(1 \eta_j^2)]^{1/2} = R_i \{1 [j(\eta_T/NV)]^2\}^{1/2}$  with  $\eta_T = 1/\sqrt{1 + b^{-2}}$  and where b is the user specified shank slope.