

Chap 6 Ex 4

The potential energy is given by

$$\begin{aligned} V &= -m_1 gl \cos \theta_1 - m_2 gl (\cos \theta_1 + \cos \theta_2) \\ V &= -lg ((m_1 + m_2) \cos \theta_1 + m_2 \cos \theta_2), \end{aligned}$$

where m_1 is the mass of the upper weight, m_2 the mass of the lower weight θ_1 is the angle that the upper bar makes with the vertical axis, and θ_2 the angle the lower bar makes. The position of the lower weight is

$$\vec{r}_2 = l(\cos \theta_1 + \cos \theta_2)\hat{j} + l(\sin \theta_1 + \sin \theta_2)\hat{i}.$$

So, the kinetic energy is

$$\begin{aligned} T &= \frac{l^2}{2} \left(\dot{\theta}_1^2 (m_1 + m_2) + \dot{\theta}_2^2 m_2 + 2\dot{\theta}_1 \dot{\theta}_2 m_2 (\sin \theta_1 \sin \theta_2 + \cos \theta_1 \cos \theta_2) \right) \\ T &= \frac{l^2}{2} \left(\dot{\theta}_1^2 (m_1 + m_2) + \dot{\theta}_2^2 m_2 + 2\dot{\theta}_1 \dot{\theta}_2 m_2 \cos(\theta_1 - \theta_2) \right). \end{aligned}$$

If we use the small angle approximation, then the kinetic and potential energies become.

$$\begin{aligned} V &= -\frac{lg}{2} ((m_1 + m_2)(2 - \theta_1^2) + m_2(2 - \theta_2^2)) \\ T &= \frac{l^2}{2} \left(\dot{\theta}_1^2 (m_1 + m_2) + \dot{\theta}_2^2 m_2 + 2\dot{\theta}_1 \dot{\theta}_2 m_2 \right). \end{aligned}$$

So, the T and V tensors are

$$\begin{aligned} V &= lg \begin{bmatrix} m_1 + m_2 & 0 \\ 0 & m_2 \end{bmatrix} \\ T &= l^2 \begin{bmatrix} m_1 + m_2 & m_2 \\ m_2 & m_2 \end{bmatrix} \end{aligned}$$

So, we need to solve $\det(V - \omega^2 T) = 0$, let $\lambda = \omega^2$

$$\begin{vmatrix} (g - l\lambda)(m_1 + m_2) & -l\lambda m_2 \\ -l\lambda m_2 & (g - l\lambda)m_2 \end{vmatrix} = 0$$

This has the following solutions

$$\lambda = \frac{2g}{l} \left(1 - \sqrt{\frac{m_2}{m_1}} \right), \frac{2g}{l} \left(1 + \frac{2m_2}{m_1} + \sqrt{\frac{m_2}{m_1}} \right).$$

Notice that the units of ω^2 are Hz^2 , as expected. Also if $m_1 \gg m_2$, then we get that

$$\frac{m_2}{m_1} \approx 0$$

So, if $m_1 \gg m_2$ the two frequencies are approximately equal to $\frac{2g}{l}$. The normal modes are

$$\left(\sqrt{\frac{m_2}{m_1 + m_2}}, 1 \right) \text{ and } \left(-\sqrt{\frac{m_2}{m_1 + m_2}}, 1 \right)$$

where the first one corresponds to the eigenvalue $\frac{2g}{l} \left(1 - \sqrt{\frac{m_2}{m_1}} \right)$. Notice that the eigenvectors are approximately equal if $m_1 \gg m_2$. Let ω_1 , and ω_2 denote the two eigenvalues, and \vec{c}_1 , and \vec{c}_2 denote the two eigenvectors. The most general motion this system can have is given by

$$\vec{r}(t) = f_1 c_1 \cos(\omega_1 t + \delta_1) + f_2 \cos(\omega_2 t + \delta_2)$$

If we have the initial condition

$$\begin{aligned} \vec{r}(0) &= (\theta_0, 0) \\ \dot{\vec{r}}(0) &= (0, 0) \end{aligned}$$

(i.e. the upper mass is slightly pulled away from the vertical, and the lower mass is allowed to hang free), then we have

$$\begin{aligned} \vec{r}(0) &= (\theta_0, 0) \\ f_1 c_1 \cos(\delta_1) + f_2 c_2 \cos(\delta_2) &= (\theta_0, 0) \\ \dot{\vec{r}}(0) &= (0, 0) \\ f_1 c_1 \omega_1 \sin(\delta_1) + f_2 c_2 \omega_2 \sin(\delta_2) &= (0, 0) \\ f_1 \cos(\delta_1) - f_2 \cos(\delta_2) &= \theta_0 \sqrt{\frac{m_1 + m_2}{m_2}} \\ f_1 \cos(\delta_1) + f_2 \cos(\delta_2) &= 0 \\ f_1 \omega_1 \sin(\delta_1) - f_2 \omega_2 \sin(\delta_2) &= 0 \\ f_1 \omega_1 \sin(\delta_1) + f_2 \omega_2 \sin(\delta_2) &= 0 \end{aligned}$$

This equation can be solved for the amplitudes and phase shifts, but even without solving it, it can be seen that this equation will give rise to beats, since when $\cos(\omega_1 t + \delta_1) = 1$, and $\cos(\omega_2 t + \delta_2) = -1$ the first pendulum will have maximum amplitude, and $\sin(\omega_1 t + \delta_1) = \sin(\omega_2 t + \delta_2) = 0$, so the second pendulum will be at rest.

Chap 6 Ex 5

For the tri-atomic molecule, the eigenvectors and eigenvalues are

$$\omega_1 = 0$$

$$\vec{c}_1 = \frac{1}{\sqrt{2m+M}}(1, 1, 1)$$

$$\omega_2 = \sqrt{\frac{k}{m}}$$

$$\vec{c}_2 = \frac{1}{\sqrt{2m}}(1, 0, -1)$$

$$\omega_3 = \sqrt{\frac{k}{m} \left(1 + \frac{2m}{M}\right)}$$

$$\vec{c}_3 = \left(\frac{1}{\sqrt{2m \left(1 + \frac{2m}{M}\right)}}, \frac{-2}{\sqrt{2M \left(2 + \frac{M}{m}\right)}}, \frac{1}{\sqrt{2m \left(1 + \frac{2m}{M}\right)}} \right)$$

Assuming that there is no translational motion ($f_1=0$), then the general solution is

$$x_1(t) = \frac{1}{\sqrt{2m}} f_2 \cos(\omega_2 t + \delta_2) + \frac{1}{\sqrt{2m \left(1 + \frac{2m}{M}\right)}} f_3 \cos(\omega_3 t + \delta_3)$$

$$x_2(t) = \frac{-2}{\sqrt{2M \left(2 + \frac{M}{m}\right)}} f_3 \cos(\omega_3 t + \delta_3)$$

$$x_3(t) = -\frac{1}{\sqrt{2m}} f_2 \cos(\omega_2 t + \delta_2) + \frac{1}{\sqrt{2m \left(1 + \frac{2m}{M}\right)}} f_3 \cos(\omega_3 t + \delta_3)$$

a)

The initial conditions are that everyone starts at rest, and $x_2(0) = a_0$, with the

other masses being at equilibrium. The equations for x_2 gives

$$x_2(0) = a_0$$

$$\dot{x}_2(0) = 0$$

$$\frac{-2}{\sqrt{2M(2 + \frac{M}{m})}} f_3 \cos(\delta_3) = a_0$$

$$\frac{2}{\sqrt{2M(2 + \frac{M}{m})}} \omega_3 f_3 \sin(\delta_3) = 0$$

$$f_3 = \frac{a_0 \sqrt{2M(2 + \frac{M}{m})}}{-2}$$

$$d_3 = 0.$$

From the other equations, we get

$$f_2 = \frac{-a_0}{2} \sqrt{2M \left(\frac{2 + \frac{M}{m}}{1 + \frac{2m}{M}} \right)}$$

$$d_2 = 0.$$

b)

If the middle mass had an initial velocity of v_0 , then

$$x_2(0) = a_0$$

$$\dot{x}_2(0) = v_0$$

$$\frac{-2}{\sqrt{2M(2 + \frac{M}{m})}} f_3 \cos(\delta_3) = a_0$$

$$\frac{2}{\sqrt{2M(2 + \frac{M}{m})}} \omega_3 f_3 \sin(\delta_3) = v_0$$

for simplicity of notation, let $a = \frac{a_0 \sqrt{2M(2 + \frac{M}{m})}}{-2}$ and $b = \frac{v_0 \sqrt{2M(2 + \frac{M}{m})}}{2\omega_3}$, so the

equations become

$$\begin{aligned}f_3 \cos(\delta_3) &= a \\f_3 \sin(\delta_3) &= b\end{aligned}$$

$$\begin{aligned}f_3 &= \sqrt{a^2 + b^2} \\d_3 &= \cos^{-1} \left(\frac{a}{\sqrt{a^2 + b^2}} \right).\end{aligned}$$

Similarly, f_2 and d_2 can be solved for using the other initial condition equations

$$\begin{aligned}f_2 &= \frac{1}{\sqrt{1 + \frac{2m}{M}}} a \\d_2 &= \cos^{-1} \left(\frac{a}{\sqrt{a^2 + b^2}} \right).\end{aligned}$$

Chap 6 Ex 8

The potential energy is given by

$$V = \frac{k}{2} \left((x_1 - x_2)^2 + (x_3 - x_2)^2 + (x_3 - x_1)^2 + (y_1 - y_2)^2 + (y_3 - y_2)^2 + (y_3 - y_1)^2 \right).$$

So, the potential and kinetic energy tensors are given by

$$V = k \begin{bmatrix} 2 & -1 & -1 & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 & 0 \\ -1 & -1 & 2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 2 & -1 & -1 \\ 0 & 0 & 0 & -1 & 2 & -1 \\ 0 & 0 & 0 & -1 & -1 & 2 \end{bmatrix}$$
$$T = \begin{bmatrix} m & 0 & 0 & 0 & 0 & 0 \\ 0 & m & 0 & 0 & 0 & 0 \\ 0 & 0 & m & 0 & 0 & 0 \\ 0 & 0 & 0 & m & 0 & 0 \\ 0 & 0 & 0 & 0 & m & 0 \\ 0 & 0 & 0 & 0 & 0 & m \end{bmatrix}$$

where the first three columns are the coordinates x_1, x_2, x_3 , and the last three

are y_1, y_2, y_3 . So, the eigenvalues and eigenvectors are

Translation in x

$$\begin{aligned}\omega_1 &= 0 \\ \vec{c}_1 &= (1, 1, 1, 0, 0, 0)\end{aligned}$$

Translation in y

$$\begin{aligned}\omega_2 &= 0 \\ \vec{c}_2 &= (0, 0, 0, 1, 1, 1)\end{aligned}$$

Oscillation between mass 1 and 2 in y

$$\begin{aligned}\omega_3 &= \sqrt{\frac{3k}{m}} \\ \vec{c}_2 &= (0, 0, 0, 1, -1, 0)\end{aligned}$$

Oscillation between mass 1 and 3 in y

$$\begin{aligned}\omega_3 &= \sqrt{\frac{3k}{m}} \\ \vec{c}_2 &= (0, 0, 0, 1, 0, -1)\end{aligned}$$

Oscillation between mass 1 and 2 in x

$$\begin{aligned}\omega_3 &= \sqrt{\frac{3k}{m}} \\ \vec{c}_2 &= (1, -1, 0, 0, 0, 0)\end{aligned}$$

Oscillation between mass 1 and 3 in x

$$\begin{aligned}\omega_3 &= \sqrt{\frac{3k}{m}} \\ \vec{c}_2 &= (1, 0, -1, 0, 0, 0).\end{aligned}$$

So, I found a double root instead of a triple root. This must mean that either there is a mistake in my expression for the kinetic or potential energy, or that I did not set up the system properly. **Chap 6 Ex 9**

The most general solution to the equation of motion of the last problem is

$$x_j = f_i c_{ij} \cos(\omega_i t + \delta_i).$$

Where there is a sum over i , and ω_i are the normal frequencies, and c_i are the corresponding eigenvectors. So, any solution that can be written in this way will satisfy the equations of motion. Clearly translation in x and y is a solution since those are the first two eigenvectors. I believe that rotation about the z should

probably be the third eigenvector, but either there is a mistake in my expression for the kinetic or potential energy, or my system is not set up correctly.

Chap 6 Ex 14

Chap 6 Ex 18

Rectangular plate problem

The kinetic energy is given by

$$T = \frac{1}{2}m \left(\dot{z}^2 + \frac{1}{6}w^2\dot{\phi}^2 + \frac{1}{6}h^2\dot{\theta}^2 \right).$$

The potential energy is given by

$$V = \frac{1}{2}k (4z^2 + w^2 \sin^2 \phi + h^2 \sin^2 \theta)$$

$$V = \frac{1}{2}k (4z^2 + w^2\phi^2 + h^2\theta^2)$$

So, the kinetic and potential energy tensors are given by

$$V = \begin{bmatrix} 4 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$T = m \begin{bmatrix} 1 & 0 & 0 \\ 0 & \frac{1}{6} & 0 \\ 0 & 0 & \frac{1}{6} \end{bmatrix}$$

where the coordinates are $z, w\phi, h\theta$ in that order. The coordinates have been chosen so that the frequency has units of Hz. The normal modes and frequencies

are

Oscillation of plate up and down

$$\omega_1 = \sqrt{\frac{4k}{m}}$$

$$\vec{c}_1 = (1, 0, 0, 0))$$

Oscillation of plate right and left

$$\omega_2 = \sqrt{\frac{6k}{m}}$$

$$\vec{c}_2 = (0, 1, 0, 0))$$

Oscillation of plate forward and backward

$$\omega_3 = \sqrt{\frac{6k}{m}}$$

$$\vec{c}_3 = (0, 0, 1, 0))$$

remember that the second two coordinates are $w\phi$ and $h\theta$ not ϕ and θ .