

University of Jyväskylä - Course TIEJ6003  
INTRODUCTION TO QUANTUM COMPUTING

Day-05a

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## the challenge: unstructured search

We would like to devise a method to locate known objects (‘needles’) in a large unordered set (‘haystack’) of  $N$  objects.

Suppose that  $f$  is a function from  $\{0, 1, \dots, N-1\} \rightarrow \{0, 1\}$ , and the target is to locate some  $x$  for which  $f(x) = 1$ .

For simplicity, assume a single *winner*  $w$  for which  $f(w) = 1$ .

Evaluating  $f(x)$  for random  $x$  will locate  $w$  in an expected  $N/2$  steps (alternatively,  $\mathcal{O}(\frac{N}{2})$  steps to find the object with probability 50%).

Grover developed a quantum algorithm to this end and published it in 1996. His method achieves a success-probability greater than 50% in  $\mathcal{O}(\sqrt{N})$  steps!

## introducing the oracle

We assume that we are provided with a black-box that has the ability to *recognize* solutions to the problem.

This assumption is not as strong as it may sound at first - there are many black-box problems in reality whose solutions may be recognized once located. We shall call it the *oracle*.

Intuitively, the oracle enables us to define a Hilbert subspace of winners versus non-winners, and to define a rotation in this subspace toward the winners.

It will be beneficial to work with the uniform superposition state  $|h\rangle = \frac{1}{\sqrt{N}} \sum_{0 \leq x < N} |x\rangle$  (prepared by  $H^{\otimes n} |0\rangle^{\otimes n}$ ).

Then, starting with  $|h\rangle$ , we will repeatedly apply rotations for  $T$  iterations and measure!

## the oracle formalized

The oracle's recognition is signalled via a qubit  $|q\rangle$ , and altogether the oracle's operation may be formally described as follows:

$$\mathcal{U}_{\text{oracle}} |x\rangle |q\rangle = |x\rangle |q \oplus f(x)\rangle. \quad (1)$$

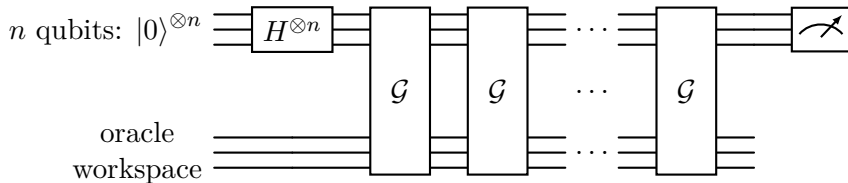
Importantly, it pays off to apply this operator on the superposition state  $(|0\rangle - |1\rangle)/\sqrt{2}$  (similar to Deutsch-Jozsa):

$$\mathcal{U}_{\text{oracle}} |x\rangle \left( \frac{|0\rangle - |1\rangle}{\sqrt{2}} \right) = (-1)^{f(x)} |x\rangle \left( \frac{|0\rangle - |1\rangle}{\sqrt{2}} \right)$$

In practice, the oracle *marks* the solution by *shifting its phase*.

Given an unstructured search with  $M$  solutions, it turns out that a solution can be located by calling the oracle  $\mathcal{O}(\sqrt{N/M})$  times.

## quantum search: schematic circuit



The Grover operator  $\mathcal{G}$  represents an iteration, which subsequently applies the Oracle followed by the so-called Diffusion operator  $\mathcal{D}$ :

$$\mathcal{G} := \mathcal{U}_{\text{oracle}} \mathcal{D}$$

## the Grover iteration

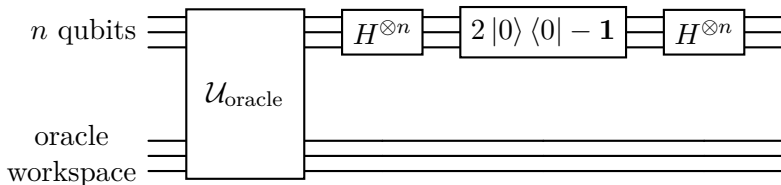
The Grover operator  $\mathcal{G}$  is applied repeatedly as a subroutine in each iteration. It can be broken down to four steps:

- 1 Call the oracle  $\mathcal{U}_{\text{oracle}}$
- 2 Apply the Hadamard transform  $H^{\otimes n}$
- 3 Conditionally apply phase shift:  
    **if**  $x > 0$  **then**  $|x\rangle \rightarrow -|x\rangle$   
    **else**  $|0\rangle \rightarrow |0\rangle$
- 4 Apply the Hadamard transform  $H^{\otimes n}$

The combined effect of steps 2+3+4 is the reflection on the uniform superposition  $|h\rangle$  (the Diffusion operator  $\mathcal{D}$ ):

$$H^{\otimes n} \left( 2|0\rangle^{\otimes n} \langle 0|^{\otimes n} - I^{\otimes n} \right) H^{\otimes n} = 2|h\rangle \langle h| - I^{\otimes n} = \mathcal{D}$$

# Grover iteration: concrete circuit per $\mathcal{G}$



$$\begin{aligned}\mathcal{G} &:= \mathcal{D}\mathcal{U}_{\text{oracle}} \\ \mathcal{U}_{\text{oracle}} &:= \mathbf{1} - 2|w\rangle\langle w| \\ \mathcal{D} &:= 2|h\rangle\langle h| - \mathbf{1}\end{aligned}$$



## subspace framework

Let  $S$  denote the set of objects we are searching for, and let  $M \geq 1$  be its cardinality.  $S$  is the *solution set* (“*winners*”), and its elements are called the *solutions*.

We denote the set of objects that are not a solution (“*non-winners*”) by

$$S^\perp := \{0, 1, \dots, N-1\} \setminus S$$

and accordingly construct two vectors:

$$|\Psi_S\rangle := \frac{1}{\sqrt{M}} \sum_{x \in S} |x\rangle \quad |\Psi_{S^\perp}\rangle := \frac{1}{\sqrt{N-M}} \sum_{x \in S^\perp} |x\rangle \quad (2)$$

## projections

By considering the two projection operators onto the two subspaces

$$P_S := \sum_{x \in S} |x\rangle \langle x| \quad P_{S^\perp} := \sum_{x \in S^\perp} |x\rangle \langle x| = \mathbf{1} - P_S, \quad (3)$$

and by recalling that  $S \cup S^\perp = \{0, 1, \dots, N-1\}$ , every state can be expanded in the computational basis as follows

$$|\psi\rangle = (P_S + P_{S^\perp}) |\psi\rangle = \sum_{x \in S} \psi_x |x\rangle + \sum_{x \in S^\perp} \psi_x |x\rangle = \alpha |\Psi_S\rangle + \beta |\Psi_{S^\perp}\rangle \quad (4)$$

In particular, the uniform superposition state  $|h\rangle$  can be expanded as:

$$|h\rangle = \sqrt{\frac{M}{N}} |\Psi_S\rangle + \sqrt{\frac{N-M}{N}} |\Psi_{S^\perp}\rangle \quad (5)$$

## interpretation: reflections

By observing the 2D subspace spanned by  $|\Psi_S\rangle$  (“the good”) and  $|\Psi_{S^\perp}\rangle$  (“the bad”), Grover’s iteration may be interpreted as two *reflection* operations when starting with  $|h\rangle$ :

- 1  $\mathcal{U}_{\text{oracle}}$  performs a *reflection* about the vector  $|\Psi_{S^\perp}\rangle$ .
- 2  $\mathcal{D}$  performs a *reflection* about the vector  $|h\rangle$ .

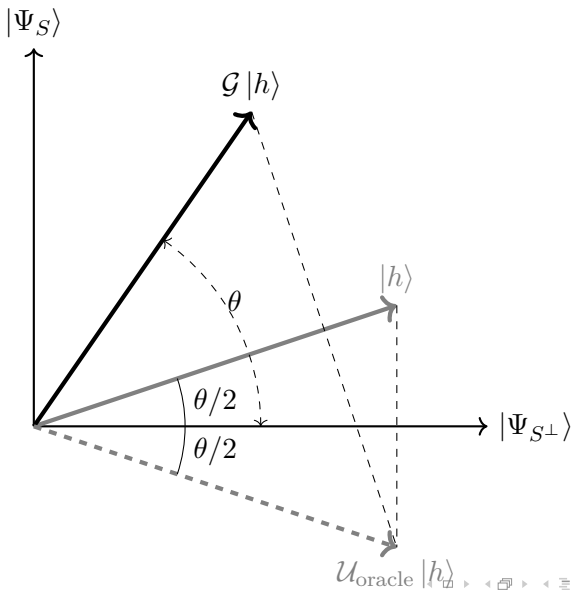
Importantly, the repeated application of  $\mathcal{G}$  always keeps the states in the plane defined by  $|\Psi_S\rangle$  and  $|\Psi_{S^\perp}\rangle$ .

Altogether,  $\mathcal{G}$  yields a rotation defined as

$$\mathcal{G} := \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}, \quad (6)$$

with  $\theta$  satisfying  $\sin \theta = 2\sqrt{M(N-M)}/N$ .

# geometric visualization



# Grover formalized per $M = 1$

**Inputs:** (1) a black-box oracle  $\mathcal{U}_{\text{oracle}}$  which performs the operation  $\mathcal{U}_{\text{oracle}} |x\rangle |q\rangle = |x\rangle |q \oplus f(x)\rangle$ .

(2)  $(n + 1)$  qubits in the state  $|0\rangle$ .

**Outputs:**  $x_0$ .

**Runtime:**  $\mathcal{O}(\sqrt{2^n})$  operations. Succeeds with probability  $\mathcal{O}(1)$ .

**Procedure:**  $|\psi_0\rangle = |0\rangle^{\otimes n} |0\rangle$

$$|\psi_1\rangle = \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} |x\rangle \left[ \frac{|0\rangle - |1\rangle}{\sqrt{2}} \right]$$

$$\rightarrow [\mathcal{D}\mathcal{U}_{\text{oracle}}]^T |\psi_1\rangle \approx |x_0\rangle \left[ \frac{|0\rangle - |1\rangle}{\sqrt{2}} \right]$$

$$\rightarrow x_0$$

## Grover: complexity and aftermath

Now that the search is reduced to rotations in the plane of  $\{|\Psi_S\rangle, |\Psi_{S^\perp}\rangle\}$ , the algorithm's complexity is reduced to the question

*“how many radians are needed to approach  $|\Psi_S\rangle$ ?”*

Given the starting point  $|h\rangle = \sqrt{\frac{M}{N}} |\Psi_S\rangle + \sqrt{\frac{N-M}{N}} |\Psi_{S^\perp}\rangle$ , rotating it through  $\arccos(\sqrt{M/N})$  radians will drive the system to  $|\Psi_S\rangle$ :

$$T = \text{round} \left( \frac{\arccos(\sqrt{M/N})}{\theta} \right) \leq \left\lceil \frac{\pi}{4} \sqrt{\frac{N}{M}} \right\rceil. \quad (7)$$

$\mathcal{O}(\sqrt{N/M})$  Grover iterations (and thus oracle calls) are required, versus  $\mathcal{O}(N/M)$  oracle calls that are required classically.