

## Lösungshinweise Blatt 7

24a) Die Impulswellenfkt. des Zustands  $\hat{x}|\psi\rangle$  ist

$$\begin{aligned}\langle \tilde{\varphi}_h | \hat{x} | \psi \rangle &= \int dx \langle \tilde{\varphi}_h | \hat{x} | x \rangle \langle x | \psi \rangle \\ &= \int dx \, x \, e^{-ihx} \psi(x) \\ &= i \frac{\partial}{\partial h} \int dx \, e^{-ihx} \psi(x) = i \frac{\partial}{\partial h} \tilde{\psi}(h) \quad .\end{aligned}$$

b)

$$\begin{aligned}\langle x \rangle_{|\psi\rangle} &= \int \frac{dh}{2\pi} \underbrace{\langle \psi | \tilde{\varphi}_h \rangle}_{\parallel} \underbrace{\langle \tilde{\varphi}_h | x | \psi \rangle}_{\parallel} \\ &= \int \frac{dh}{2\pi} \tilde{\psi}^*(h) \, i \frac{\partial}{\partial h} \tilde{\psi}(h) \quad ,\end{aligned}$$

$$\begin{aligned}\langle p \rangle_{|\psi\rangle} &= \int \frac{dh}{2\pi} \underbrace{\langle \psi | p | \tilde{\varphi}_h \rangle}_{\hbar h \langle \tilde{\varphi}_h |} \langle \tilde{\varphi}_h | \psi \rangle \\ &= \int \frac{dh}{2\pi} \tilde{\psi}^*(h) \, \hbar h \, \tilde{\psi}(h) \quad .\end{aligned}$$

c) z.z.:  $e^{i p_0 \hat{x}} | \tilde{\varphi}_p \rangle \stackrel{!}{=} | \tilde{\varphi}_{p+p_0} \rangle \quad ; \quad (\hbar=1)$

$$\begin{aligned}e^{i p_0 \hat{x}} | \tilde{\varphi}_p \rangle &= \int dx \, e^{i p_0 \hat{x}} | x \rangle \langle x | \tilde{\varphi}_p \rangle \\ &= \int dx \, | x \rangle \underbrace{e^{i p_0 x} e^{i p x}}_{\stackrel{!}{=} \langle x | \tilde{\varphi}_{p+p_0} \rangle} = \int dx \, | x \rangle \langle x | \tilde{\varphi}_{p+p_0} \rangle \quad .\end{aligned}$$

$$d) \quad \tilde{T}^\dagger(p_0) = (e^{i p_0 \hat{x}})^\dagger = e^{-i p_0 \hat{x}} = \tilde{T}(p_0)^{-1},$$

$\uparrow$   
 $\hat{x}$  hermitesch

$$\begin{aligned} [\hat{p}, \tilde{T}(p_0)] |\varphi_p\rangle &= (\hat{p} \tilde{T}(p_0) - \tilde{T}(p_0) \hat{p}) |\varphi_p\rangle \\ &= (p + p_0 - p) |\varphi_{p+p_0}\rangle \\ &= p_0 \tilde{T}(p_0) |\varphi_p\rangle. \end{aligned}$$

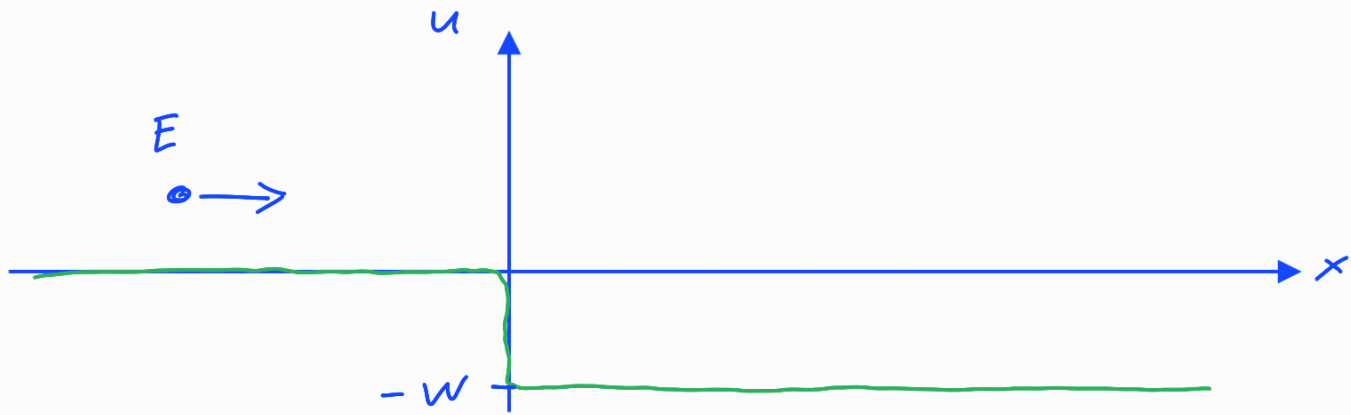
e)

$$\begin{aligned} \langle p \rangle_{\tilde{T}_{p_0}|\psi} &= \langle \psi | \tilde{T}_{p_0}^\dagger \hat{p} \tilde{T}_{p_0} | \psi \rangle \\ &\stackrel{d)}{=} \langle \psi | \tilde{T}_{p_0}^\dagger (\tilde{T}_{p_0} \hat{p} + p_0 \tilde{T}_{p_0}) | \psi \rangle \\ &= \langle \psi | \hat{p} | \psi \rangle + p_0 \langle \psi | \psi \rangle = \langle p \rangle_{|\psi\rangle} + p_0. \end{aligned}$$

$$\begin{aligned} f) \quad \tilde{T}_{\hbar_0} |\psi\rangle &= \int dx \tilde{T}_{\hbar_0} |x\rangle \langle x | \psi \rangle \\ &\stackrel{a)}{=} \int dx e^{i \hbar_0 x} |x\rangle \langle x | \psi \rangle \\ &= \int dx \underbrace{e^{i \hbar_0 x} \psi(x)}_{\hat{=} \text{Wellenfunkt. von } \tilde{T}_{\hbar_0} |\psi\rangle} |x\rangle \\ &\hat{=} \text{Wellenfunkt. von } \tilde{T}_{\hbar_0} |\psi\rangle. \end{aligned}$$

$$\begin{aligned} e) \quad T_{x_0} |\psi\rangle &= \int \frac{d\hbar}{2\pi} e^{-i x_0 \hat{p} / \hbar} |\tilde{\varphi}_\hbar\rangle \langle \tilde{\varphi}_\hbar | \psi \rangle \\ &= \int d\hbar \underbrace{e^{-i x_0 \hbar} \tilde{\psi}(\hbar)}_{\hat{=} \text{Impulswellenfunkt. von } T_{x_0} |\psi\rangle} |\tilde{\varphi}_\hbar\rangle \\ &\hat{=} \text{Impulswellenfunkt. von } T_{x_0} |\psi\rangle \end{aligned}$$

25)



Streuansatz:

$$\psi(x) = \begin{cases} e^{ikx} + r e^{-ikx} & : x < 0 \\ t e^{ik'x} & : x \geq 0 \end{cases}$$

mit  $k = \sqrt{2mE}/\hbar$ ,  $k' = \sqrt{2m(E+W)}/\hbar$

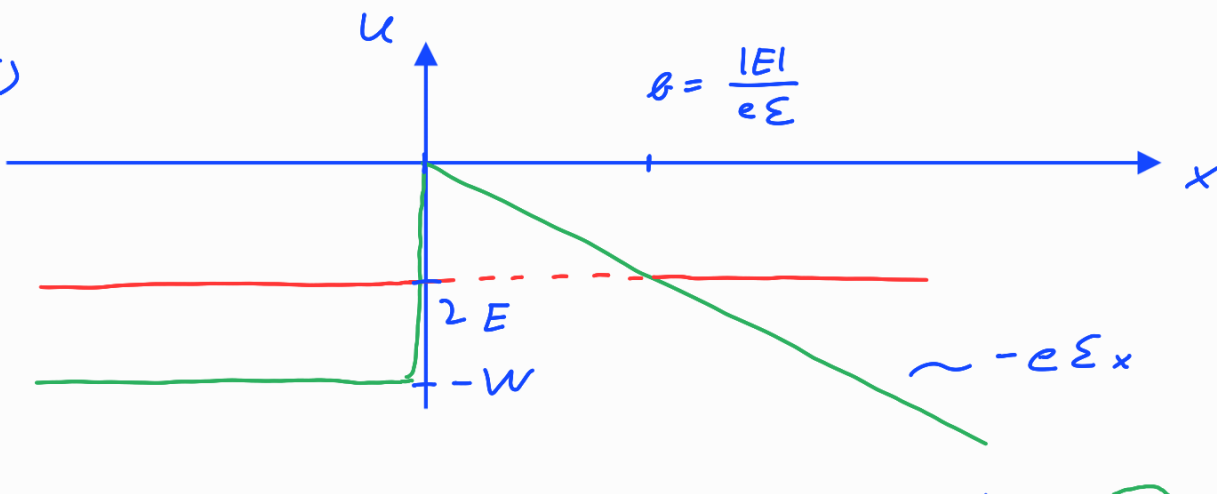
Stetigkeit von  $\psi$  in  $x=0$ :  $1+r = t$   
 " " " " " " :  $k(1-r) = k't$  }  $\rightarrow 1-r = \frac{k'}{k}(1+r)$

$$\rightarrow r = \frac{1 - k'/k}{1 + k'/k} = \frac{1 - \sqrt{1+W/E}}{1 + \sqrt{1+W/E}} = -8/10$$

$W/E = 80$

$$\rightarrow R = |r|^2 = 0.64$$

26)



$$\frac{1}{\hbar} \int_0^L (8m(U_0 - E))^{1/2} dx = \frac{\sqrt{8m}}{\hbar} \int_0^L (|E| - eEx)^{1/2} dx$$

$$= - \frac{\sqrt{8m}}{\hbar} \frac{2}{3eE} (|E| - eEx)^{3/2} \Big|_0^L = \frac{4}{3} \frac{\sqrt{2m}}{\hbar} \frac{|E|^{3/2}}{eE}$$

$$\rightarrow I(E, E) \simeq I_0 \exp\left(-\frac{4}{3} \frac{\sqrt{2m}}{\hbar} \frac{|E|^{3/2}}{eE}\right)$$

signifikanter Tunnelstrom bei Feldstärke

$$\Sigma_c \simeq \frac{4}{3} \frac{\sqrt{2me}}{\hbar} \frac{|E|^{3/2}}{e} \simeq 10^{10} \text{ V/m}.$$

$$27) |\psi(t)\rangle = \frac{e^{-i\omega t/2}}{\sqrt{2}} (|0\rangle + e^{-i\omega t} |1\rangle)$$

$$\rightarrow \langle a \rangle_t = \frac{1}{2} (\langle 0| + e^{i\omega t} \langle 1|) a (|0\rangle + e^{-i\omega t} |1\rangle)$$

$$= \frac{e^{-i\omega t}}{2} \underbrace{\langle 0| a |1\rangle}_{=1} = \frac{e^{-i\omega t}}{2}$$

$$\hookrightarrow \langle a^\dagger \rangle_t = \langle a \rangle_t^* = \frac{e^{+i\omega t}}{2}$$

$$\hookrightarrow \langle x \rangle_t = \left(\frac{\hbar}{2m\omega}\right)^{1/2} \langle a^\dagger + a \rangle_t = \left(\frac{\hbar}{2m\omega}\right)^{1/2} \cos \omega t$$

$$\langle p \rangle_t = \left(\frac{\hbar m \omega}{2}\right)^{1/2} i \langle a^\dagger - a \rangle_t = -\left(\frac{\hbar m \omega}{2}\right)^{1/2} \sin \omega t$$