

Letzte Vr(sq):

$$SO(3) \ni R \xrightarrow{\quad} U(R) \in \mathcal{L}(\mathcal{H})$$

$$\vec{r} \mapsto R\vec{r}$$

$$|\psi\rangle \mapsto U(R)|\psi\rangle$$

s.d. :

- $U(R)$ unitär
- $U(I_3) = \mathbb{1}_{\mathcal{H}}$
- $U(R' R) = U(R') U(R)$

R $\mapsto U(R)$

Darstellung
der $SO(3)$

- $U_{\text{Ort}}(R) |\vec{r}\rangle = |R\vec{r}\rangle$

\rightarrow Drehimpuls

$$J_e := i\hbar \frac{\partial}{\partial \varphi} U(R_{e,\varphi}) \Big|_{\varphi=0}$$

\rightarrow Vertauschungsrelationen:

$$[J_x, J_z] = i\hbar \epsilon_{xzy} J_y$$

ebenso für L_d, S_a
da $\vec{J} = \vec{L} + \vec{S}$

$$\rightarrow U(R_{\vec{n},\varphi}) = e^{-i\vec{n} \cdot \vec{J} \varphi/\hbar} = e^{-i\vec{n} \cdot \vec{L} \varphi/\hbar} \otimes e^{-i\vec{n} \cdot \vec{S} \varphi/\hbar}$$

$$\rightarrow$$
 Bahn-drehimpuls: $\hat{\vec{L}} = \hat{\vec{r}} \times \hat{\vec{p}}$

$$\underline{\text{Spinoperatoren}} \quad \vec{S} = \begin{pmatrix} S_x \\ S_y \\ S_z \end{pmatrix}$$

(= Drehimpulsoperatoren des "Eigenwertes des Drehimpulses")

- genügen

$$[S_x, S_y] = i\hbar \epsilon_{hern} S_z$$

- operieren auf $\mathcal{H}_{\text{spin}} = \text{Span}\{|\uparrow\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, |\downarrow\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}\} \cong \mathbb{C}^2$

→ d.h. $S_x \in M(2 \times 2, \mathbb{C})$

$$S_x = \frac{i\hbar}{2} \sigma_x$$

(eindeutig bis auf unitär äquivalente Darstellungen, vgl. Ülg.)

d.h. $\vec{S} = \frac{i\hbar}{2} \vec{\sigma}$ mit $\vec{\sigma} = \begin{pmatrix} \sigma_1 \\ \sigma_2 \\ \sigma_3 \end{pmatrix} :$

$$S_x = \frac{i\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$S_y = \frac{i\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$S_z = \frac{i\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Eigenzust.: $|x\pm\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ \pm 1 \end{pmatrix}$

$$|y\pm\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ \pm i \end{pmatrix}$$

$$\begin{aligned} |\uparrow\rangle &= |z+\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ |\downarrow\rangle &= |z-\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \end{aligned}$$

Eigenwerte: $\pm \hbar/2$

$$\pm \hbar/2$$

$$\pm \hbar/2$$

→ Transformation eines Spinzustandes

$$|\psi\rangle = \alpha |1\rangle + \beta |0\rangle = \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \quad (\in \mathbb{C}^2)$$

unter Rotation $R_{\vec{u}, \varphi}$ ($\in SO(3)$):

$$|\psi\rangle \mapsto |\psi'\rangle = U_{\text{spin}}(R_{\vec{u}, \varphi}) |\psi\rangle \text{ mit}$$

$$U_{\text{spin}}(R_{\vec{u}, \varphi}) = e^{-i \vec{u} \cdot \vec{s} \varphi / \hbar} = e^{-i \frac{\vec{u} \cdot \vec{\sigma} \varphi}{2}}$$

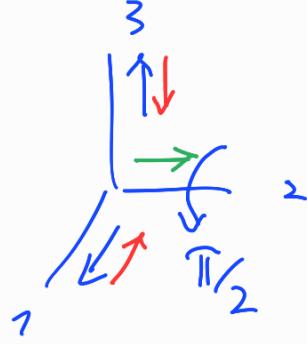
$\vec{s} = \frac{\hbar}{2} \vec{\sigma}$

$$\vec{u} \cdot \vec{\sigma} = u_1 \sigma_1 + u_2 \sigma_2 + u_3 \sigma_3 = \begin{pmatrix} u_3 & u_1 - i u_2 \\ u_1 + i u_2 & -u_3 \end{pmatrix}$$

$$\rightarrow (\vec{u} \cdot \vec{\sigma})^2 = \mathbb{1}_2 \quad (|\vec{u}| = r)$$

$$\rightarrow U_{\text{spin}}(R_{\vec{u}, \underline{\varphi}}) = \mathbb{1} \cos \frac{\underline{\varphi}}{2} - i \vec{u} \cdot \vec{\sigma} \sin \frac{\underline{\varphi}}{2}$$

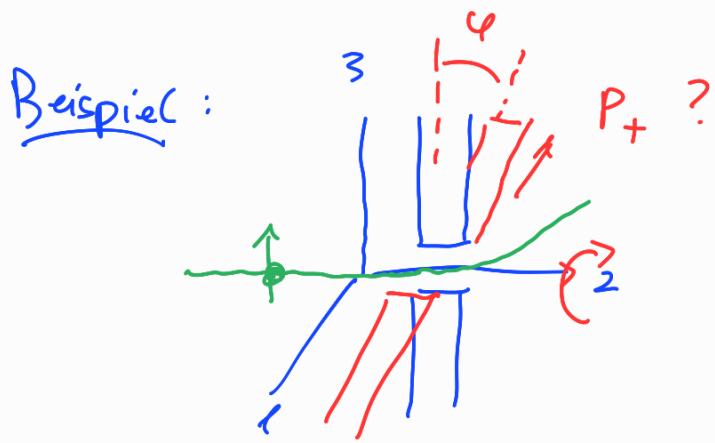
(Lec 6): $U_{2,\pi_2} := U_{sp}(R_{2,\pi_2}) \quad , \quad -i\pi_2 = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$



$$= \underbrace{\frac{1}{2} \cos \pi_4}_{1/\sqrt{2}} + \underbrace{\begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}}_{i/\sqrt{2}} \underbrace{\sin \pi_4}_{i/\sqrt{2}}$$

$$U_{2,\pi_2} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} ;$$

- $U_{2,\pi_2} |z+\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = |x+\rangle \checkmark$
- " " $|z- \rangle = " " \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} -1 \\ 1 \end{pmatrix} = -|x-\rangle \checkmark$
- $U_{2,\pi_2} |x+\rangle = \frac{1}{2} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} = |z-\rangle \checkmark$
- " " $|x-\rangle = " " \begin{pmatrix} 1 \\ -1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} = |z+\rangle \checkmark$
- $U_{2,\pi_2} |y+\rangle = \frac{1}{2} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ i \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1-i \\ 1+i \end{pmatrix}$
- $= \frac{1-i}{2} \begin{pmatrix} 1 \\ 1+i \end{pmatrix} = e^{-i\pi_4} \underbrace{\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix}}_{|y+\rangle} \checkmark$
- " " $|y-\rangle = " " = e^{+i\pi_4} |y-\rangle \checkmark$



$$|z^+\rangle = U(R_{2,\phi}) |z+\rangle$$

$$= \underbrace{1}_{\text{Z}} \cos \varphi_2 + \underbrace{\begin{pmatrix} 0 & -1 \\ 1 & 1 \end{pmatrix}}_{\text{Z} - i\Omega_2} \sin \varphi_2 = \begin{pmatrix} \cancel{\cos \varphi_2} & -\sin \varphi_2 \\ \sin \varphi_2 & \cos \varphi_2 \end{pmatrix}$$



$$P_+ = \left| \langle z^+ | z^+ \rangle \right|^2 = \left| \langle z^+ | U_{2,\phi} | z^+ \rangle \right|^2$$

$$= \left(\cos \varphi_2 \right)^2$$

$$\Gamma \quad \varphi = 30^\circ = \frac{\pi}{6} \quad \rightarrow \quad P_+ = \left(\cos \frac{\pi}{6} \right)^2$$

$$= 0,93 \quad . \quad \boxed{1}$$

Bemerkung: $U(R_{\vec{u}, \varphi}) = \underbrace{1}_{=} \cos \frac{\varphi}{2} - i \underbrace{\vec{0} \cdot \vec{u}}_{=} \sin \frac{\varphi}{2}$

$\frac{2\pi}{2}$ -Rot.: $U(R_{\vec{u}, 2\pi}) = -1$ | $R_{\vec{u}, 2\pi} = \underline{1}$
(II)

$\frac{4\pi}{4}$ -Rot.: $U(R_{\vec{u}, 4\pi}) = +1$ | $R_{\vec{u}, 4\pi} = \underline{1}$

(\rightarrow Aqu., QFT.)

Drehimpuls-eigenwerte und -zustände

Ausgangspkt.: $[J_x, J_z] = i\hbar \epsilon_{QCM} J_3$ (*)

\rightarrow u.Ug. fassen über Eigenwerte/-zustände:

J^2 und $\underline{J_3}$
 Γ $\quad [J^2 = J_1^2 + J_2^2 + J_3^2]$

(*) $\rightarrow [J^2, J_x] = \dots = 0$

insb. $[J^2, J_3] = 0$

\rightarrow gemeinsame Eigenbasis von J^2, J_3 :

Eigenbasis $|\alpha, \beta\rangle$:

$$\hat{J}^2 |\alpha, \beta\rangle = \alpha |\alpha, \beta\rangle$$

$$\hat{J}_z |\alpha, \beta\rangle = \beta |\alpha, \beta\rangle$$

$\alpha, \beta \in \mathbb{Z}$

hilfreich:

$$\begin{aligned}\hat{J}_+ &:= \hat{J}_x + i \hat{J}_y & \uparrow, + \\ \hat{J}_- &:= \hat{J}_x - i \hat{J}_y & \downarrow, -\end{aligned}$$

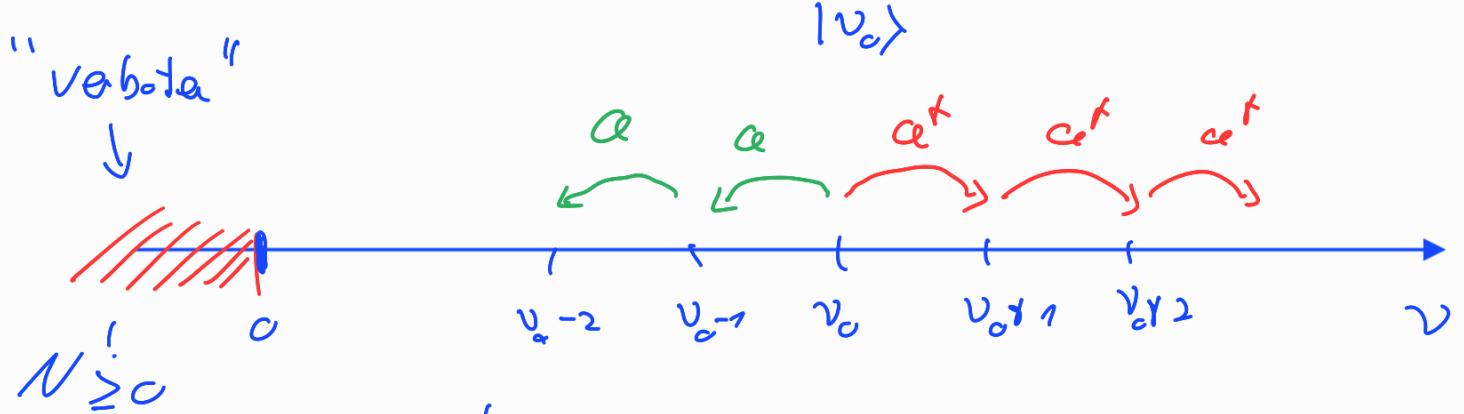
(*) \rightarrow

- $[\underline{\hat{J}_3}, \underline{\hat{J}_+}] = \frac{\hbar}{2} \hat{J}_+$ ||
- $[\underline{\hat{J}_3}, \underline{\hat{J}_-}] = \frac{\hbar}{2} \hat{J}_-$!
- (• $[\hat{J}^2, \hat{J}_{\pm}] = 0$)

[Erinnerung: harm. Oszill.: α, α^+ , $[\alpha, \alpha^+] = 0$]

$$N = \underline{\underline{\alpha \alpha^+}}, \quad N |\nu\rangle = \nu |\nu\rangle$$

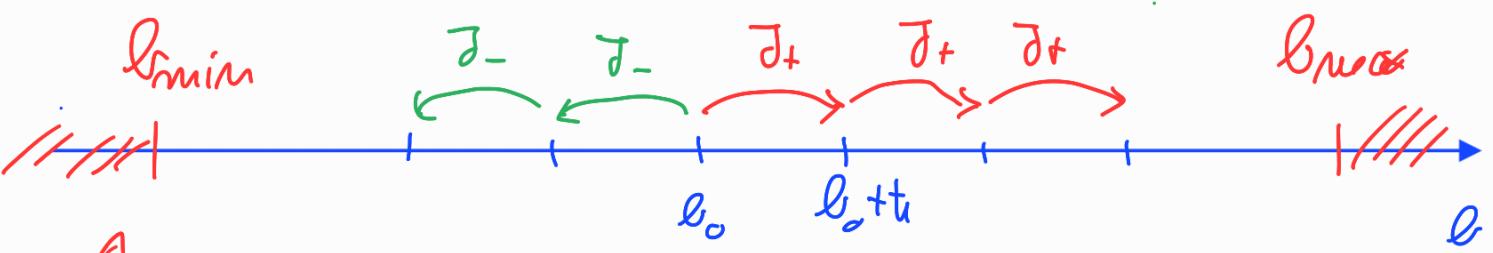
- $[\underline{\underline{N}}, \underline{\underline{\alpha^+}}] = \underline{\underline{\alpha^+}}$ ||
- $[\underline{\underline{N}}, \underline{\underline{\alpha}}] = \underline{\underline{\alpha}}$ ||



$$\alpha|v_o> = 0 \rightarrow v_o \in N$$

$\rightarrow |v>, v = 0, 1, 2, 3, \dots$

mit gleicher Rechnung (aufgrund gleicher Kom.-Relativierung):



(*)

$$0 \leq J^2 - J_3^2$$

$$\left[\langle J^2 - J_3^2 \rangle_{\psi} = \langle J_1^2 \rangle_{\psi} + \langle J_2^2 \rangle_{\psi} \geq c! \right]$$

für beliebiges $|>$

\rightarrow es ex. Zust. $|\alpha, b_{\max}\rangle, |\alpha, b_{\min}\rangle$ s. d.

$$+ \quad J_+ |\alpha, b_{\max}\rangle = 0 !$$

$$- \quad J_- |\alpha, b_{\min}\rangle = c !$$

$$\begin{aligned} \rightarrow & \underbrace{J_- J_+}_{!!} |\alpha, b_{\max}\rangle = 0 \\ & (J_1 - i J_2)(J_1 + i J_2) = \underbrace{J_1^2}_{J^2} + \underbrace{J_2^2}_{-J_3^2} + i [J_1, J_2] \\ & = J^2 - J_3^2 - t J_3 \end{aligned}$$

$$(J^2 - J_3^2 - t J_3) |\underline{\alpha}, \underline{b_{\max}}\rangle = 0$$

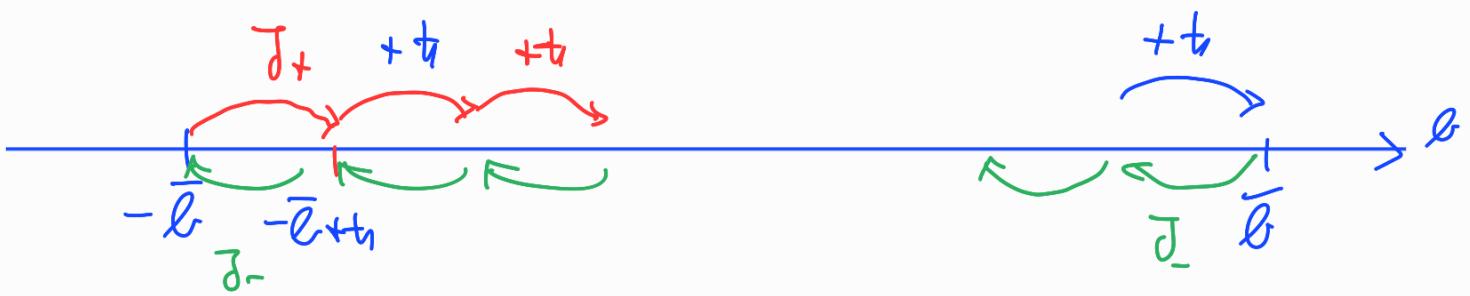
$$(\alpha - b_{\max}^2 - t b_{\max}) |\alpha, b_{\max}\rangle = 0$$

$$\text{d. h. } \cdot \quad \alpha = b_{\max} (b_{\max} + t)$$

$$\text{analog: } \cdot \quad \alpha = b_{\min} (b_{\min} - t)$$

$$\rightarrow b_{\min} = -b_{\max} \equiv -\bar{b}$$

$$b_{\max} = \bar{b}$$



$$\rightarrow 2 \bar{l} = \underline{m} \bar{t}, \quad m = 0, 1, \dots$$

$$\bar{l} = \bar{t} \cdot \underline{\frac{m}{2}} = (0, \frac{1}{2}, 1, \frac{3}{2}, \dots)$$

$$\rightarrow a = \bar{l} (\bar{l} + \bar{t})$$

$$= \bar{t}^2 \underline{\frac{m}{2}} \left(\underline{\frac{m}{2}} + 1 \right) \rightarrow j$$

Drehimpulsquantenzahlen

$$\hookrightarrow j = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots, \quad j \in \frac{1}{2} \mathbb{N}$$

$$m = -j, -j+1, \dots, j-1, j$$

Drehimpuls eiger 2D-Kreis

$$|j, m\rangle \quad (|\alpha, \beta\rangle)$$

$$\cancel{J^2(j, m)} = \bar{t}^2 j(j+1) |j, m\rangle, \quad \cancel{J_3} |j, m\rangle = \underline{m} |j, m\rangle$$

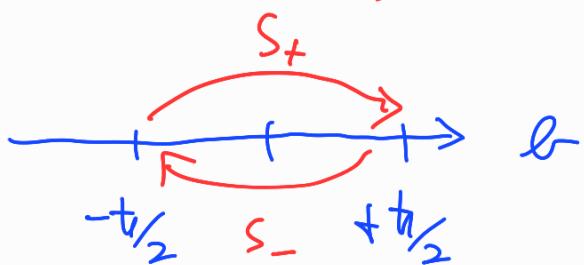
Beispiele : Spin des Elektrons :

$$S_3 (\geq \pm) = \pm \frac{\hbar}{2} (\geq \pm)$$



$$\underline{m} = \pm \frac{1}{2},$$

$$\boxed{j = s = \frac{1}{2} !}$$



$$\left(|z+> = |\frac{1}{2}, \frac{1}{2}>, |z-> = |\frac{1}{2}, -\frac{1}{2}> \right)$$

$$S_+ = S_1 + i S_2 = \frac{\hbar}{2} (0_1 + i 0_2) = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

$$S_- = S_1 - i S_2 = \dots = \frac{\hbar}{2} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

(↓)

$$S_+ \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{\hbar}{2} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{\hbar}{2} |1>$$

$$S_- \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \dots = \frac{\hbar}{2} |1>$$

|↑>

$$S_+ \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{\hbar}{2} \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = 0 !$$



, Elektron ist Spin $\frac{1}{2}$ Teilchen"

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j

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