

费曼3作业10

郑子诺，物理41

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80.4

(a) Suppose $\beta = a + ib$, use $|\beta|^2 + |1 + \beta|^2 = 1$ we have

$$b^2 = -a(1 + a)$$

and we also have

$$\frac{\beta}{1 + \beta} = \frac{a + ib}{1 + a + ib} = \frac{a(1 + a) + b^2 + ib}{(1 + a)^2 + b^2}$$

and therefore

$$\operatorname{Re}\left\{\frac{\beta}{1 + \beta}\right\} = \frac{a(1 + a) + b^2}{(1 + a)^2 + b^2} = 0$$

(b) Let $\frac{\beta}{1 + \beta} = i \tan \eta$, we have

$$\beta = \frac{i \tan \eta}{1 - i \tan \eta} = \frac{i \sin \eta}{\cos \eta - i \sin \eta} = i(\cos \eta + i \sin \eta) \sin \eta = ie^{i\eta} \sin \eta$$

83.4

(a) Use the Schrodinger equation, we have

$$i\hbar\left(-\frac{da}{dt}x^2 - \frac{dc}{dt}\right) = -\frac{\hbar^2}{2m}(4a^2x^2 - 2a)$$

and then

$$(i\hbar\frac{da}{dt} - \frac{2\hbar^2}{m}a^2)x^2 + i\hbar\frac{dc}{dt} + \frac{\hbar^2}{m}a = 0$$

Because this equation should be correct for all x , we have

$$\frac{da}{dt} = -\frac{2i\hbar}{m}a^2$$

and

$$\frac{1}{a(t)} = \frac{1}{a_0} + \frac{2i\hbar}{m}t$$

(b) Use the equation above, we also have

$$c(t) = \frac{1}{2} \ln\left(1 + \frac{2i\hbar a_0}{m}t\right) + c_0$$

(c) We have

$$\operatorname{Re}\{a(t)\} = \frac{a_0}{1 + \frac{4\hbar^2 a_0^2}{m^2}t^2}$$

Let $\frac{1}{\sqrt{a_0}} = 1\text{\AA}$, $t = 1\text{s}$, we have

$$d = \frac{1}{\sqrt{\operatorname{Re}\{a(t)\}}} = 2.3 \times 10^6 \text{m}$$

(d) We have

$$\phi(p) = \int_{-\infty}^{+\infty} \frac{1}{\sqrt{2\pi\hbar}} e^{-i\frac{p}{\hbar}x} \psi(x) dx = \frac{K}{\sqrt{2a\hbar}} e^{-(\frac{p^2}{4a\hbar^2} + c)}$$

(e) We have

$$\Delta p = \sqrt{4\operatorname{Re}\{a(t)\}\hbar^2} = \sqrt{\frac{4a_0\hbar^2}{1 + \frac{4\hbar^2 a_0^2}{m^2}t^2}}$$

(f) Obviously we have

$$\frac{d\frac{1}{\sqrt{a(t)}}}{dt} = i\hbar\sqrt{a(t)}$$

Then

$$2|i\hbar\sqrt{a(t)}| = 2\sqrt{|a(t)|}\hbar = \Delta p$$