

Considering $\rho \equiv qx$, and $[J] \equiv \sqrt{2J+1}$, for M_J , according to Eq.(3a) of Ref [1]

$$\begin{aligned} \langle n'(l'1/2)j' || M_J(q\vec{x}) || n(l1/2)j \rangle &= \frac{1}{(4\pi)^{1/2}} (-1)^{J+j+1/2} [l'] [l] [j'] [j] [J] \begin{Bmatrix} l' & j' & \frac{1}{2} \\ j & l & J \end{Bmatrix} \begin{pmatrix} l' & J & l \\ 0 & 0 & 0 \end{pmatrix} \\ &\times \langle n'l'j' | j_J(\rho) | nlj \rangle \end{aligned} \quad (1)$$

For Δ_J , according to Eq.(3c) of Ref. [1] and notice that $L = J$ here,

$$\begin{aligned} \langle n'(l'1/2)j' || \Delta_J(q\vec{x}) || n(l1/2)j \rangle &= \frac{1}{(4\pi)^{1/2}} (-1)^{J+j+1/2} [l'] [j'] [j] [J] [J] \begin{Bmatrix} l' & j' & \frac{1}{2} \\ j & l & J \end{Bmatrix} \\ &\times \left\{ -(l+1)^{1/2} [l+1] \begin{Bmatrix} J & 1 & J \\ l & l' & l+1 \end{Bmatrix} \begin{pmatrix} l' & J & l+1 \\ 0 & 0 & 0 \end{pmatrix} \langle n'l'j' | j_J(\rho) \left(\frac{d}{d\rho} - \frac{l}{\rho} \right) | nlj \rangle \right. \\ &\left. + l^{1/2} [l-1] \begin{Bmatrix} J & 1 & J \\ l & l' & l-1 \end{Bmatrix} \begin{pmatrix} l' & J & l-1 \\ 0 & 0 & 0 \end{pmatrix} \langle n'l'j' | j_J(\rho) \left(\frac{d}{d\rho} + \frac{l+1}{\rho} \right) | nlj \rangle \right\} \end{aligned} \quad (2)$$

For Σ'_J and Σ''_J , they relate to $M_{JL} \cdot \sigma$ according to Eq. (1e) and (1f) of Ref. [1], therefore we first derive $M_{JL} \cdot \sigma$. It can be obtained from Eq.(3b) of Ref. [1],

$$\begin{aligned} \langle n'(l'1/2)j' || M_{JL}(q\vec{x}) \cdot \sigma || n(l1/2)j \rangle &= \frac{1}{(4\pi)^{1/2}} (-1)^{l'} 6^{1/2} [l'] [l] [j'] [j] [L] [J] \begin{Bmatrix} l' & l & L \\ \frac{1}{2} & \frac{1}{2} & 1 \end{Bmatrix} \begin{pmatrix} l' & L & l \\ 0 & 0 & 0 \end{pmatrix} \\ &\times \langle n'l'j' | j_L(\rho) | nlj \rangle \end{aligned} \quad (3)$$

Then we can get Σ'_J and Σ''_J ,

$$\begin{aligned} \langle n'(l'1/2)j' || \Sigma'_J(q\vec{x}) || n(l1/2)j \rangle &= -\frac{\sqrt{J}}{[J]} \langle n'(l'1/2)j' || M_{JJ+1}(q\vec{x}) \cdot \sigma || n(l1/2)j \rangle \\ &+ \frac{\sqrt{J+1}}{[J]} \langle n'(l'1/2)j' || M_{JJ-1}(q\vec{x}) \cdot \sigma || n(l1/2)j \rangle, \end{aligned} \quad (4)$$

$$\begin{aligned} \langle n'(l'1/2)j' || \Sigma''_J(q\vec{x}) || n(l1/2)j \rangle &= \frac{\sqrt{J+1}}{[J]} \langle n'(l'1/2)j' || M_{JJ+1}(q\vec{x}) \cdot \sigma || n(l1/2)j \rangle \\ &+ \frac{\sqrt{J}}{[J]} \langle n'(l'1/2)j' || M_{JJ-1}(q\vec{x}) \cdot \sigma || n(l1/2)j \rangle. \end{aligned} \quad (5)$$

For Φ''_J , the nuclear operator can be given by

$$\Phi''(q\vec{x}) = i \left(\frac{1}{q} \vec{\Delta} M_J(q\vec{x}) \right) \quad (6)$$

Now the question is how to derive $\langle n'l'j'|j_L(\rho)|nlj\rangle$, $\langle n'l'j'|j_L(\rho)\left(\frac{d}{d\rho} - \frac{l}{\rho}\right)|nlj\rangle$, and $\langle n'l'j'|j_L(\rho)\left(\frac{d}{d\rho} + \frac{l+1}{\rho}\right)|nlj\rangle$. It is defined by

$$\langle n'l'j'|\theta(\rho)|nlj\rangle = \int x^2 dx R_{n'l'j'}^*(x)\theta(\rho)R_{nlj}(x), \quad (7)$$

for

$$\theta(\rho) = j_L(\rho), j_L(\rho)\left(\frac{d}{d\rho} - \frac{l}{\rho}\right), \text{ and } j_L(\rho)\left(\frac{d}{d\rho} + \frac{l+1}{\rho}\right) \quad (8)$$

To completely evaluate Eq. (7), we employ harmonic oscillators in here and hence drop the lable j now,

$$R_{nl}(x) = \left(\frac{2e^z}{b^3(n-1)!\Gamma(n+l+1/2)z^{l+1}} \right)^{1/2} \times \frac{d^{n-1}}{dz^{n-1}}(z^{n+l-1/2}e^{-z}), \quad (9)$$

where $z \equiv (x/b)^2$ and b is the oscillator parameter. Harmonic oscillator recursion relations give

$$R_{nl}(x) = \sqrt{(n-1)!\Gamma(n+l+1/2)} \sum_{m=0}^{n-1} \frac{(-1)^m}{m!(n-1-m)!} \frac{\sqrt{\Gamma(l+2m+3/2)}}{\Gamma(l+m+3/2)} R_{1l+2m}(x). \quad (10)$$

so that the matrix elements in Eq. (7) can be reduced to linear combination of terms having only $n' = n = 1$. In addition, we have [1]

$$\begin{aligned} \left(\frac{d}{d\rho} - \frac{l}{\rho}\right)R_{1l}(x) &= -(8y)^{-1/2}[l+1]R_{1l+1}(x), \\ \left(\frac{d}{d\rho} + \frac{l+1}{\rho}\right)R_{1l}(x) &= (8y)^{-1/2}\{2[l]R_{1l-1}(x) - [l+1]R_{1l+1}(x)\}, \end{aligned} \quad (11)$$

where $y \equiv (qb/2)^2$. In Ref. [1], we can get

$$\langle 1l'|j_L(\rho)|1l\rangle = \frac{(2y)^{L/2}e^{-y}(L+l'+l+1)!!}{(2L+1)!!\{(2l'+1)!!(2l+1)!!\}^{1/2}} \times {}_1F_1[(L-l'-l; L+3/2; y], \quad (12)$$

where ${}_1F_1(a; b; z)$ is the confluent hypergeometric function.

Insert Eq. (12), Eq. (11), Eq. (10) into Eq. (7), and keep the relation between gamma function and double factorial in mind,

$$\Gamma(n + \frac{1}{2}) = \frac{(2n-1)!!}{2^n} \sqrt{\pi}, \quad (13)$$

we can evaluate $\langle n'l'j'|j_L(\rho)|nlj\rangle$, $\langle n'l'j'|j_L(\rho)\left(\frac{d}{d\rho} - \frac{l}{\rho}\right)|nlj\rangle$, and $\langle n'l'j'|j_L(\rho)\left(\frac{d}{d\rho} + \frac{l+1}{\rho}\right)|nlj\rangle$

as follow

$$\begin{aligned}
\langle n'l'j'|j_L(\rho)|nlj\rangle &= \frac{2^L}{(2L+1)!!} y^{L/2} e^{-y} \sqrt{(n'-1)!(n-1)!} \\
&\times \sqrt{\Gamma(n'+l'+1/2)\Gamma(n+l+1/2)} \\
&\times \sum_{m=0}^{n-1} \sum_{m'=0}^{n'-1} \frac{(-1)^{m+m'}}{m!m'!(n-m-1)!(n'-m'-1)!} \\
&\times \frac{\Gamma[(l+l'+L+2m+2m'+3)/2]}{\Gamma(l+m+3/2)\Gamma(l'+m'+3/2)} \\
&\times {}_1F_1[(L-l'-l-2m'-2m)/2; L+3/2; y]
\end{aligned} \tag{14}$$

$$\begin{aligned}
\langle n'l'j'|j_L(\rho)\left(\frac{d}{d\rho} - \frac{l}{\rho}\right)|nlj\rangle &= \frac{2^L}{(2L+1)!!} y^{\frac{L-1}{2}} e^{-y} \sqrt{(n'-1)!(n-1)!} \\
&\times \sqrt{\Gamma(n'+l'+1/2)\Gamma(n+l+1/2)} \\
&\times \sum_{m=0}^{n-1} \sum_{m'=0}^{n'-1} \frac{(-1)^{m+m'}}{m!m'!(n-m-1)!(n'-m'-1)!} \\
&\times \frac{\Gamma[(l+l'+L+2m+2m'+2)/2]}{\Gamma(l+m+3/2)\Gamma(l'+m'+3/2)} \times \\
&\left\{ -\frac{l+l'+L+2m+2m'+2}{2} {}_1F_1[(L-l'-l-2m'-2m-1)/2; L+3/2; y] \right. \\
&\left. + 2m \times {}_1F_1[(L-l'-l-2m'-2m+1)/2; L+3/2; y] \right\}
\end{aligned} \tag{15}$$

$$\begin{aligned}
\langle n'l'j'|j_L(\rho)\left(\frac{d}{d\rho} + \frac{l+1}{\rho}\right)|nlj\rangle &= \frac{2^{L-1}}{(2L+1)!!} y^{\frac{L-1}{2}} e^{-y} \sqrt{(n'-1)!(n-1)!} \\
&\times \sqrt{\Gamma(n'+l'+1/2)\Gamma(n+l+1/2)} \\
&\times \sum_{m=0}^{n-1} \sum_{m'=0}^{n'-1} \frac{(-1)^{m+m'}}{m!m'!(n-m-1)!(n'-m'-1)!} \\
&\times \frac{\Gamma[(l+l'+L+2m+2m'+2)/2]}{\Gamma(l+m+3/2)\Gamma(l'+m'+3/2)} \\
&\left\{ -\frac{l+l'+L+2m+2m'+2}{2} {}_1F_1[(L-l'-l-2m'-2m-1)/2; L+3/2; y] \right. \\
&\left. + (2l+2m+1) \times {}_1F_1[(L-l'-l-2m'-2m+1)/2; L+3/2; y] \right\}
\end{aligned} \tag{16}$$

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- [1] T. Donnelly and W. Haxton, Atomic Data and Nuclear Data Tables **23**, 103 (1979), ISSN 0092-640X, URL <http://www.sciencedirect.com/science/article/pii/0092640X79900032>.