

# Indirect measurements of nuclear cross sections: tempering experimental results with theory

Oliver Gorton<sup>1</sup>, Jutta Escher<sup>2</sup>, Orlando Olivas-Gomez<sup>3</sup>

<sup>1</sup>Computational Science (PhD)  
San Diego State University | UC Irvine

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<sup>2</sup>Lawrence Livermore National Laboratory  
Nuclear Data & Theory  
PLS/NACS

<sup>3</sup>University of Notre Dame

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National Laboratory

# I will give a summary of the major results of my summer projects

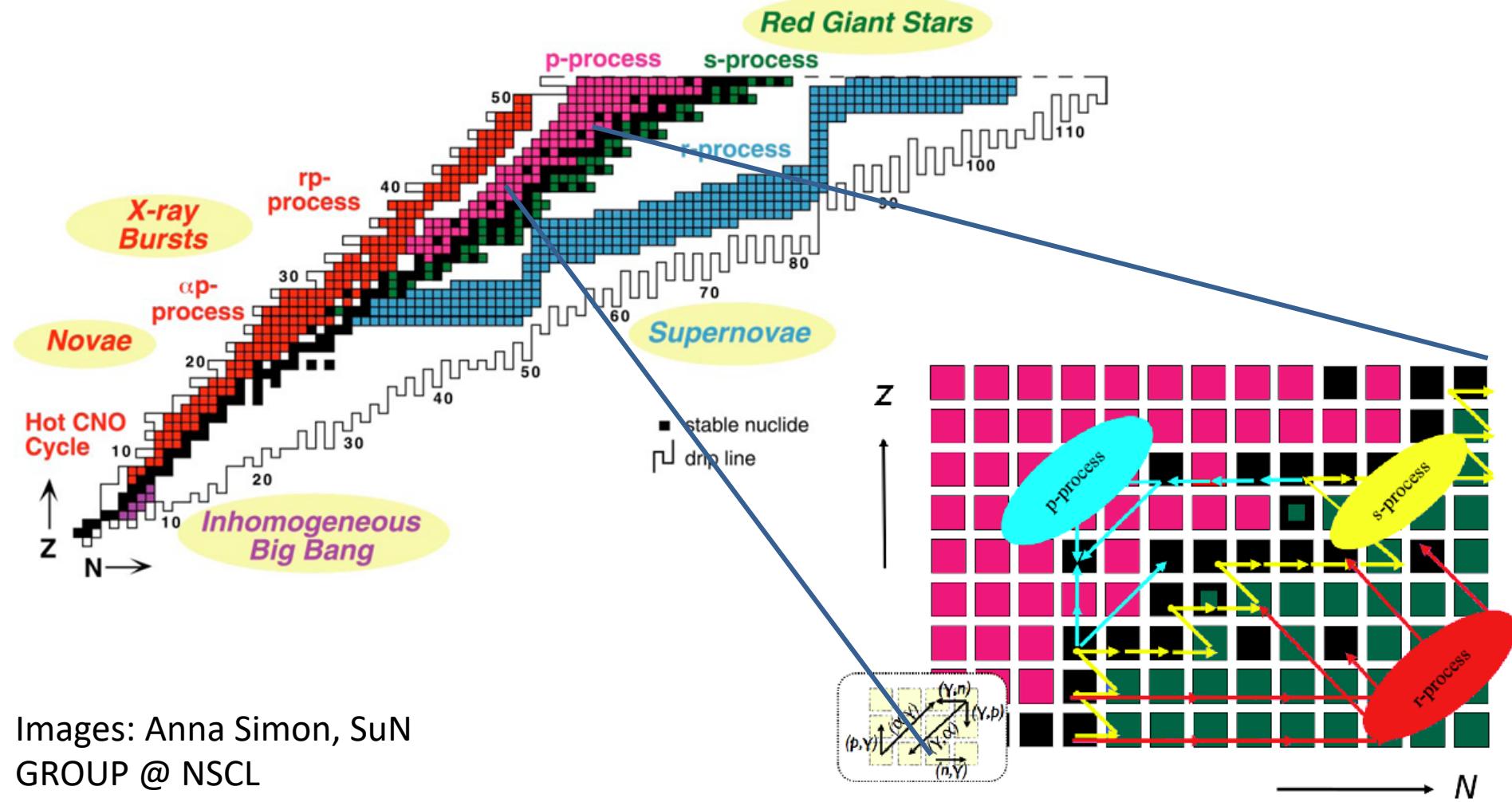
## PART 1:

- Use of inverse-reactions for indirect measurements at University of Notre Dame
- Applying our theory capabilities to improve predictive power of indirect measurements

## PART 2:

- Sensitivity of an approximation method used for indirect measurements
- Applying this surrogate method to recent data

# All of the elements in the universe came from a handful of astrophysical processes

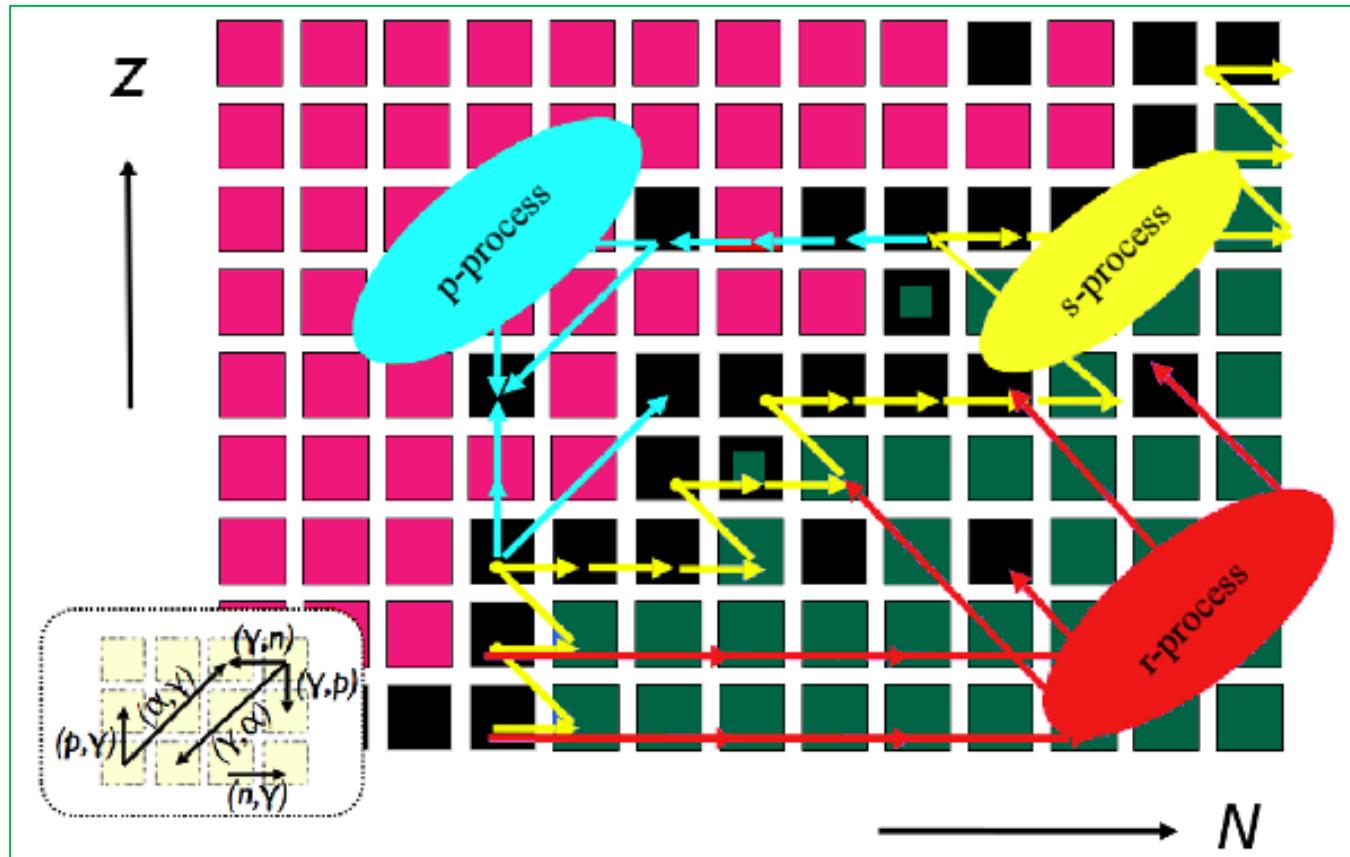


Images: Anna Simon, SuN

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# Proton-rich elements formed in p-process



Images: Anna Simon, SuN

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# The details of the p-process(es) remain a mystery: unable to predict abundances

Z				103Sb	104Sb	105Sb	106Sb	107Sb	108Sb	109Sb	110Sb	111Sb	112Sb	113Sb	114Sb	115Sb	
	99Sn	100Sn	101Sn	102Sn	103Sn	104Sn	105Sn	106Sn	107Sn	108Sn	109Sn	110Sn	111Sn	112Sn	113Sn	114Sn	
49	97In	98In	99In	100In	101In	102In	103In	104In	105In	106In	107In	108In	109In	110In	111In	112In	113In
	96Cd	97Cd	98Cd	99Cd	100Cd	101Cd	102Cd	103Cd	104Cd	105Cd	106Cd	107Cd	108Cd	109Cd	110Cd	111Cd	112Cd
47	95Ag	96Ag	97Ag	98Ag	99Ag	100Ag	101Ag	102Ag	103Ag	104Ag	105Ag	106Ag	107Ag	108Ag	109Ag	110Ag	111Ag
	94Pd	95Pd	96Pd	97Pd	98Pd	99P	( $\gamma, \alpha$ )	101Pd	102Pd	103Pd	104Pd	105Pd	106Pd	107Pd	108Pd	109Pd	110Pd
45	93Rh	94Rh	95Rh	96Rh	97Rh	98Rh	99Rh	100Rh	101Rh	102Rh	103Rh	104Rh	105Rh	106Rh	107Rh	108Rh	109Rh
	92Ru	93Ru	94Ru	95Ru	96Ru	97Ru	98Ru	99Ru	100Ru	101Ru	102Ru	103Ru	104Ru	105Ru	106Ru	107Ru	108Ru
43	91Tc	92Tc	93Tc	94Tc	95Tc	96Tc	97Tc	98Tc	99Tc	100Tc	101Tc	102Tc	103Tc	104Tc	105Tc	106Tc	107Tc
	48	50	52	54	56	58	60	62	N								

Images: NNDC chart of nuclides

# Measuring $(p,\gamma)$ can tell us about $(\gamma,p)$

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- Astrophysics application need better cross sections
- Orlando Olivas-Gomez from University of Notre Dame
- Indirect measurement of  $(\gamma,p)$  cross sections from  $(p,\gamma)$  experiment

$(p,\gamma)$  = Proton capture

$(\gamma,p)$  = photo-disintegration of a photon

# UND group measured proton capture cross sections

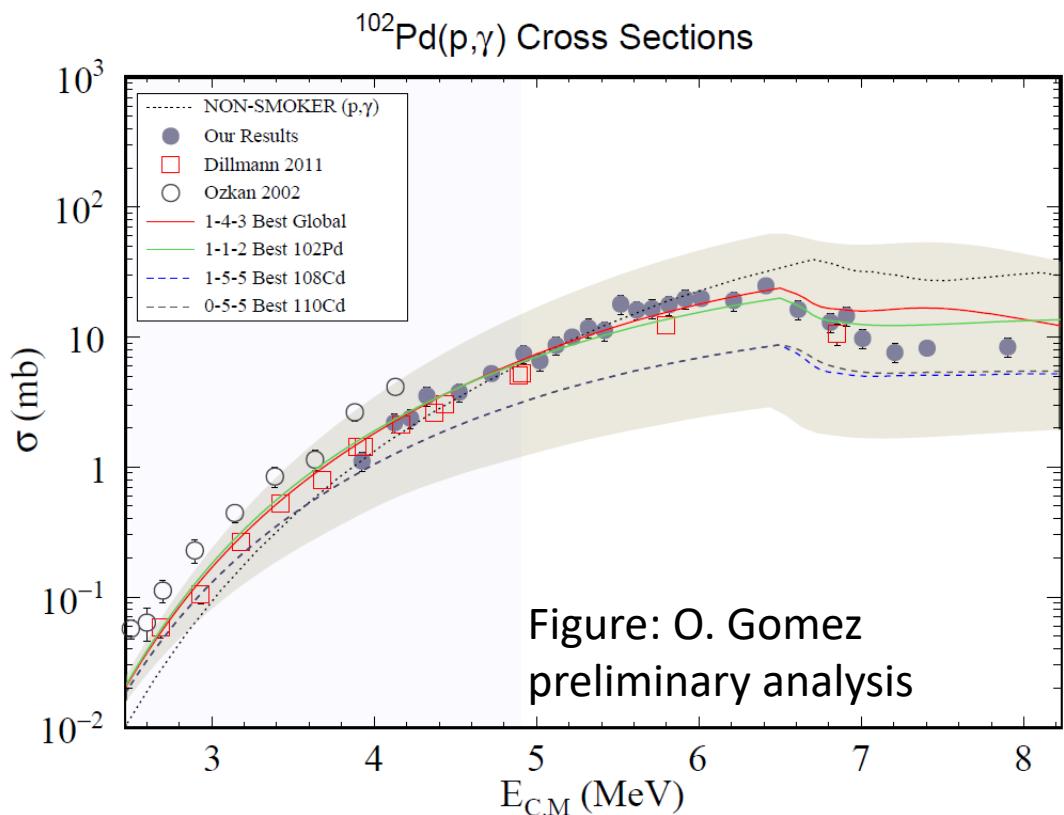
- Original idea:
  - Find the best theoretical model to describe  $(p, \gamma)$  data
  - Use that model to predict  $(\gamma, p)$  cross sections

Targets:

$^{102}\text{Pd}$

$^{108}\text{Cd}$

$^{110}\text{Cd}$  (Never done before!)



This is where I come in...

# I wrote a code for this!

## My work with the theory side of indirect measurements

- Fit nuclear reaction model parameters with MCMC
- Indirect measurement
  - > data
  - > parameter estimation
  - > predict desired reaction
- Surrogate method

## Tech specs

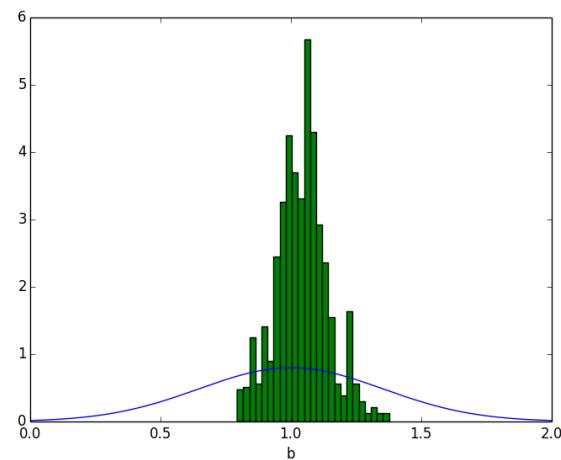
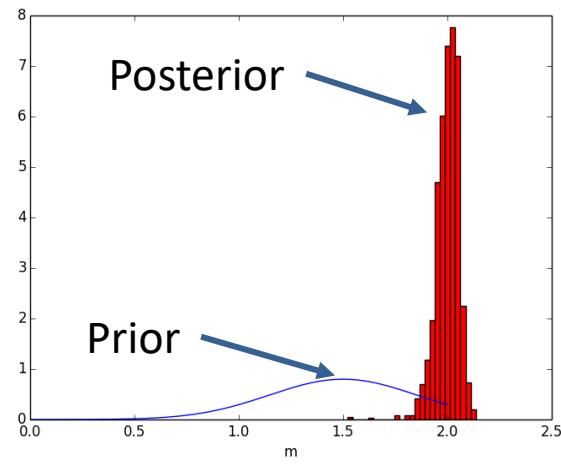
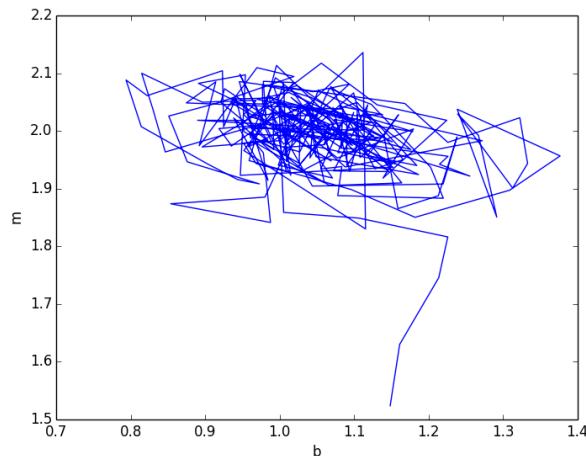
- Python, multiple Hauser-Feshbach reaction codes, multiple observables, simultaneous fitting, parallel sampling, 2000+ LOC

## Notable improvements this summer

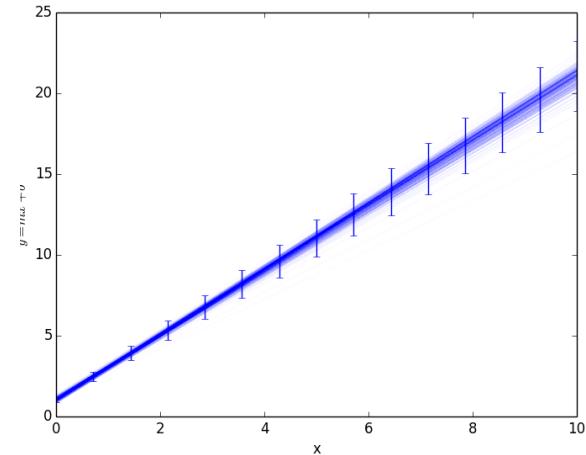
- Ability to include prior distributions for parameters
- Compatibility with open-source Hauser-Feshbach model code, EMPIRE

# Applying MCMC to linear regression

Works by censoring a random walk



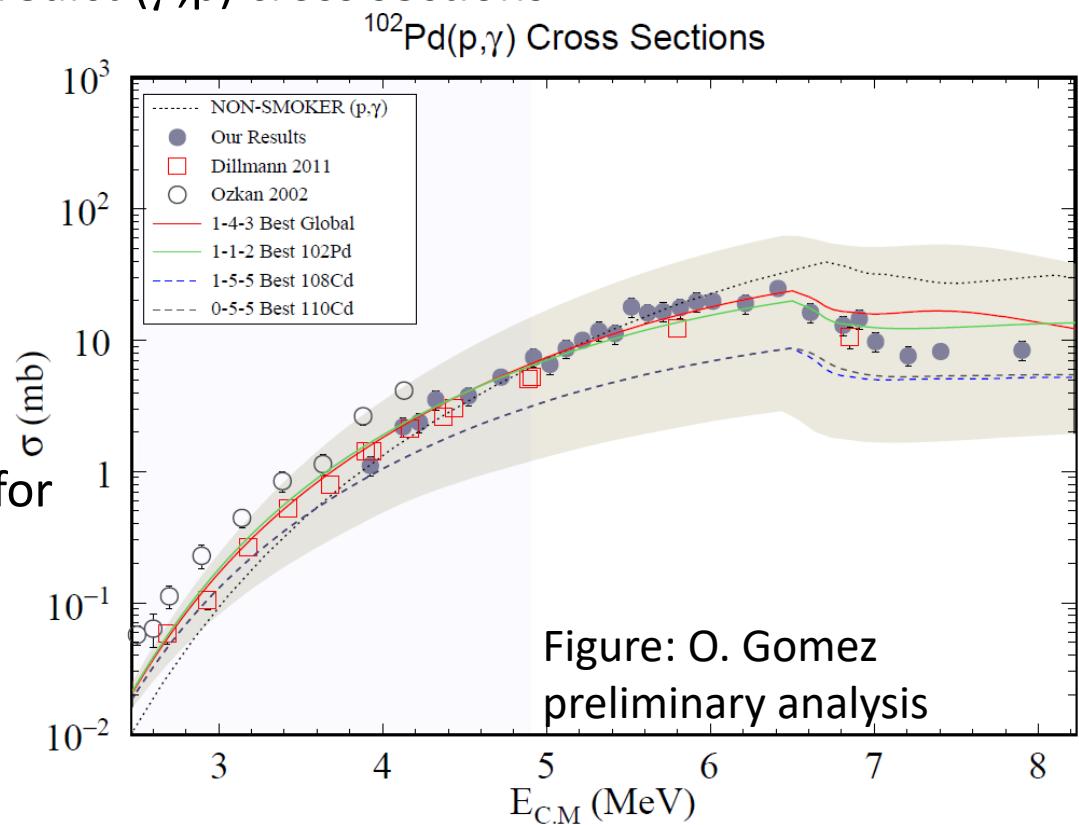
$$y = m x + b$$



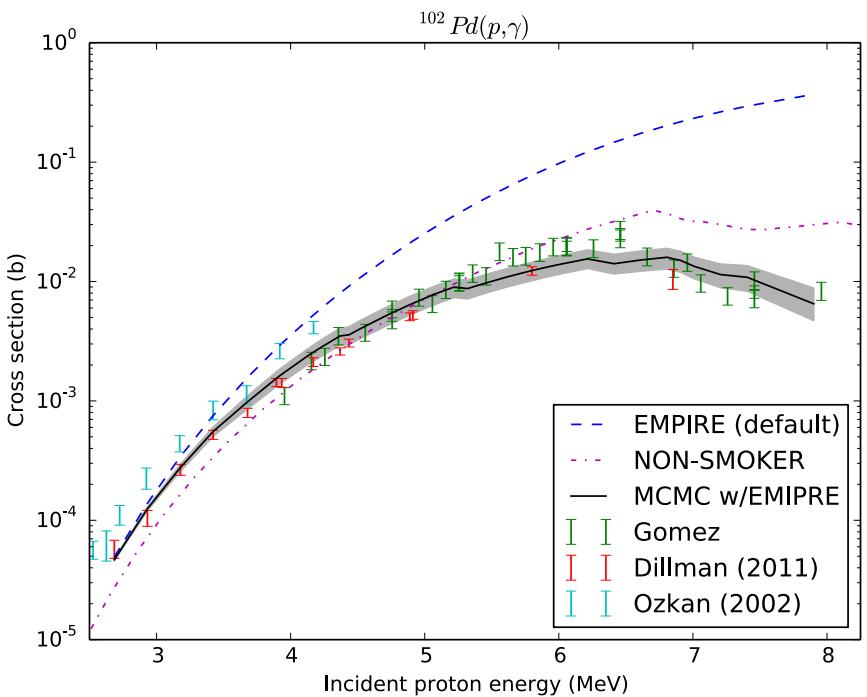
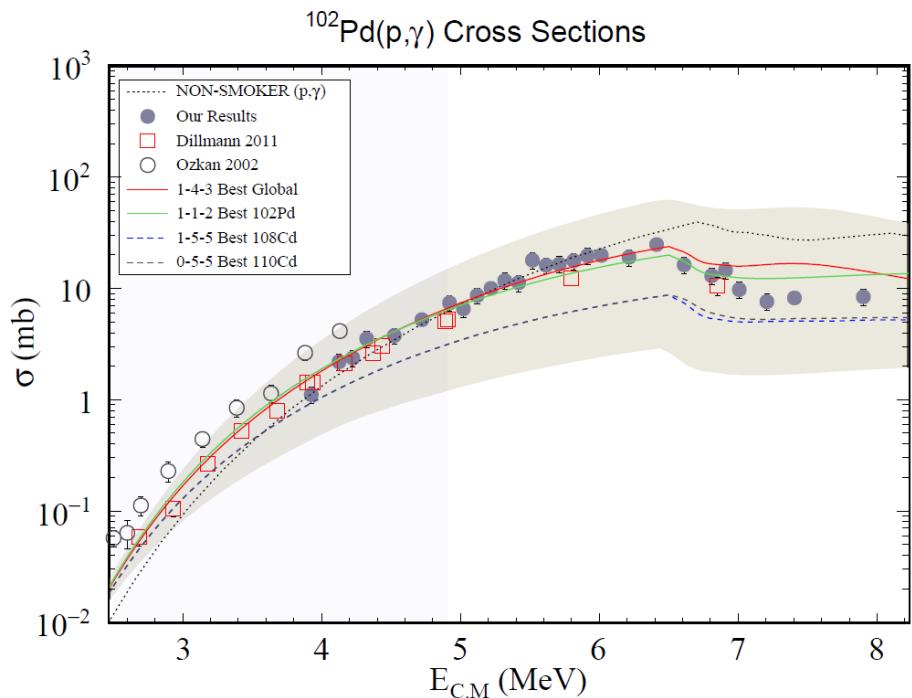
# I can improve the resulting theory with MCMC

- Original idea:
  - Find the best theoretical model to describe  $(p, \gamma)$  data
  - Use that best model to predict  $(\gamma, p)$  cross sections

- What this will improve:
  - Extract more information from the data
  - Probability distributions for parameters
  - Uncertainty bars



# From nearest-facsimile to probability distribution



Selecting the best ‘default’ option

Tuning a model to constraints given by the data

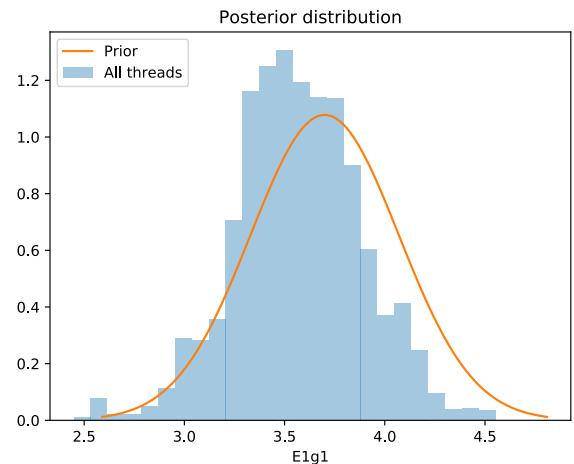
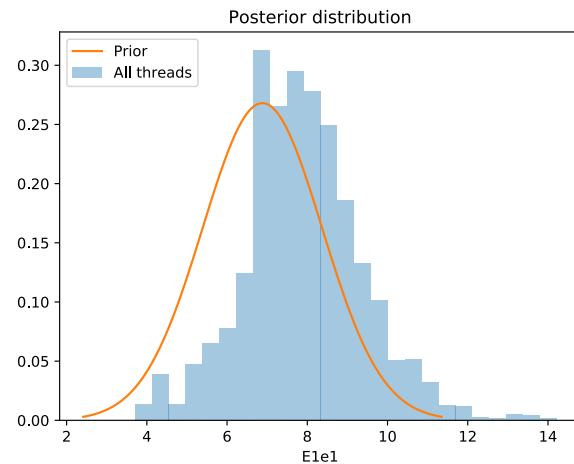
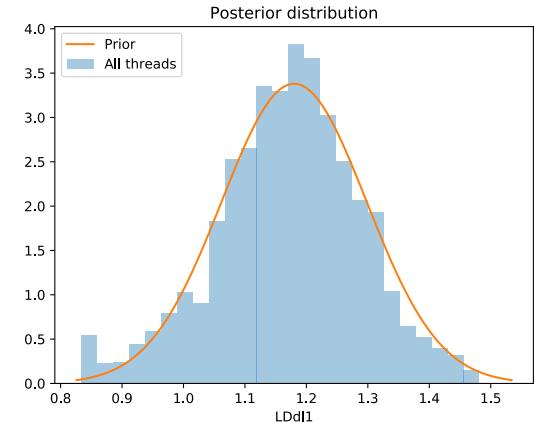
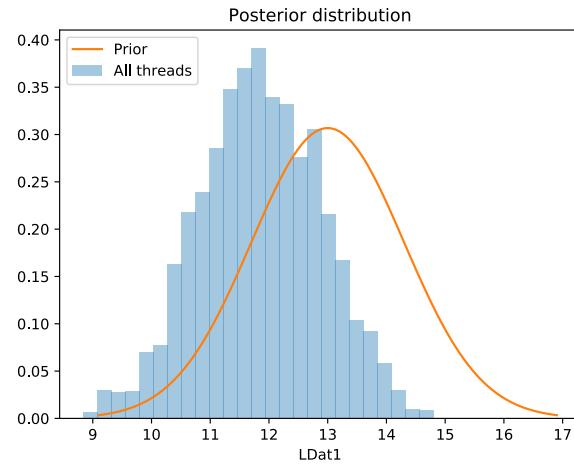
# I examined nuclear structure model parameters

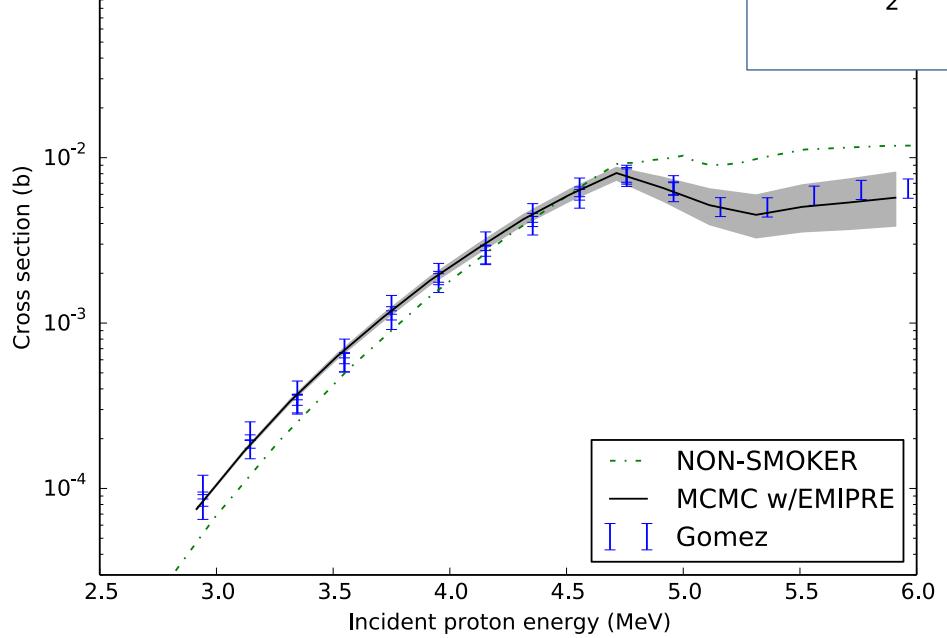
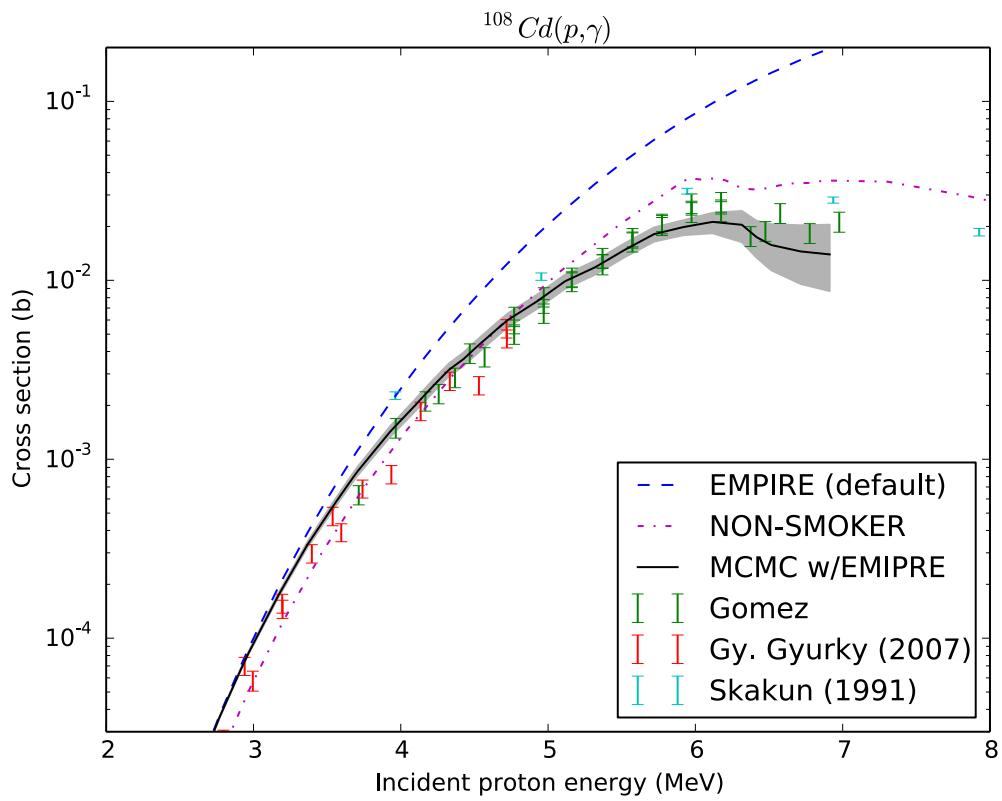
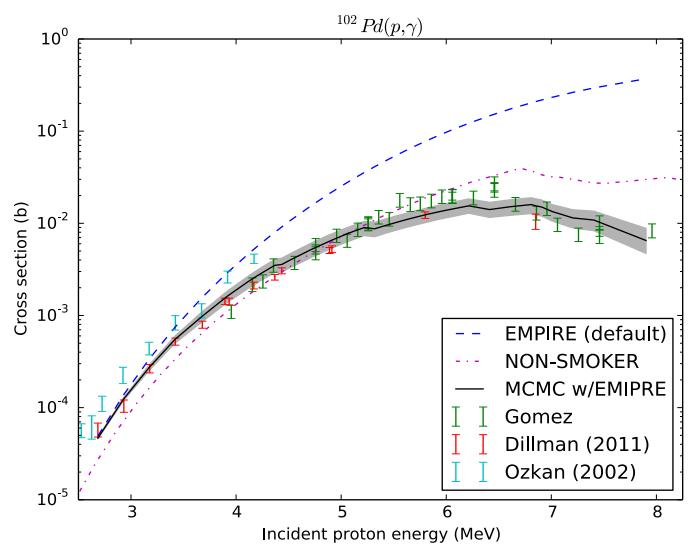
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- Gamma-ray strength functions
- Nuclear level densities
- Priors mostly from Reference Input Parameter Library (RIPL-3)

# Nuclear reaction model parameter estimation

Parameter distributions  
from the  $^{102}\text{Pd}$  cross  
section fitting

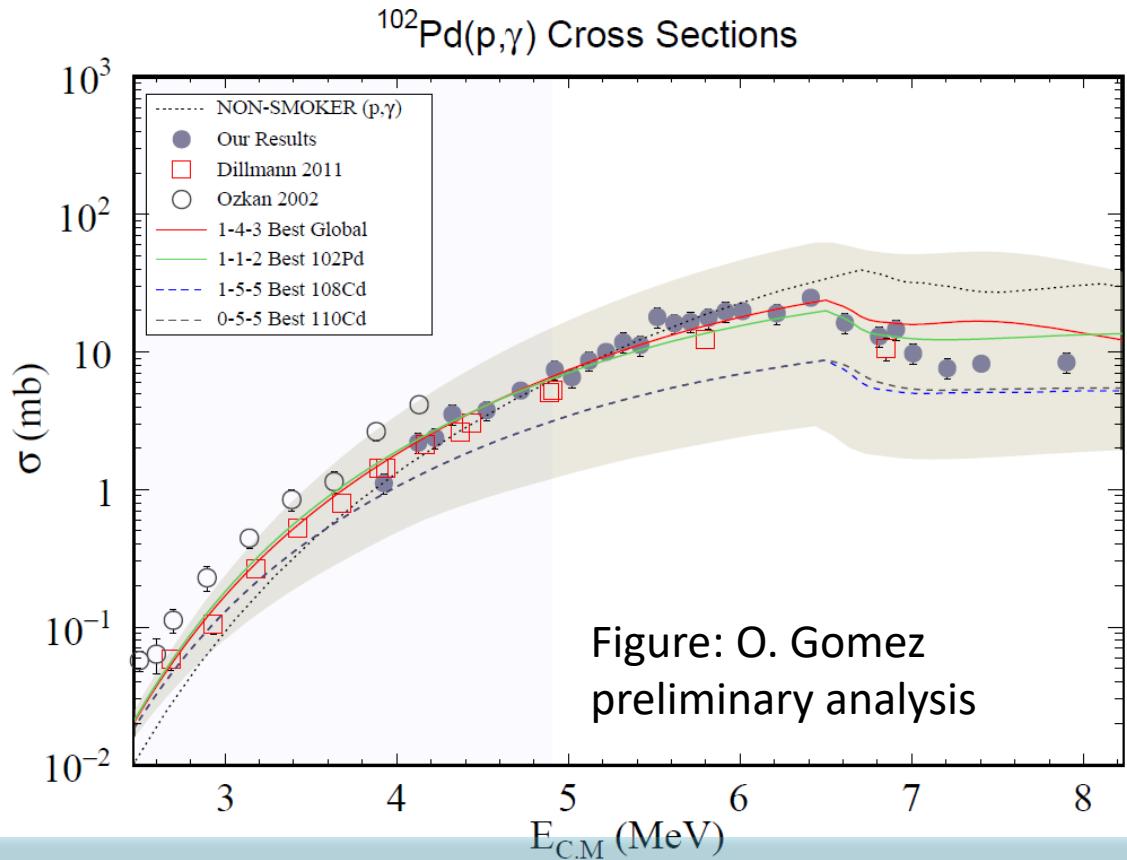




# With parameters that reflect the data, now we can find the desired cross sections

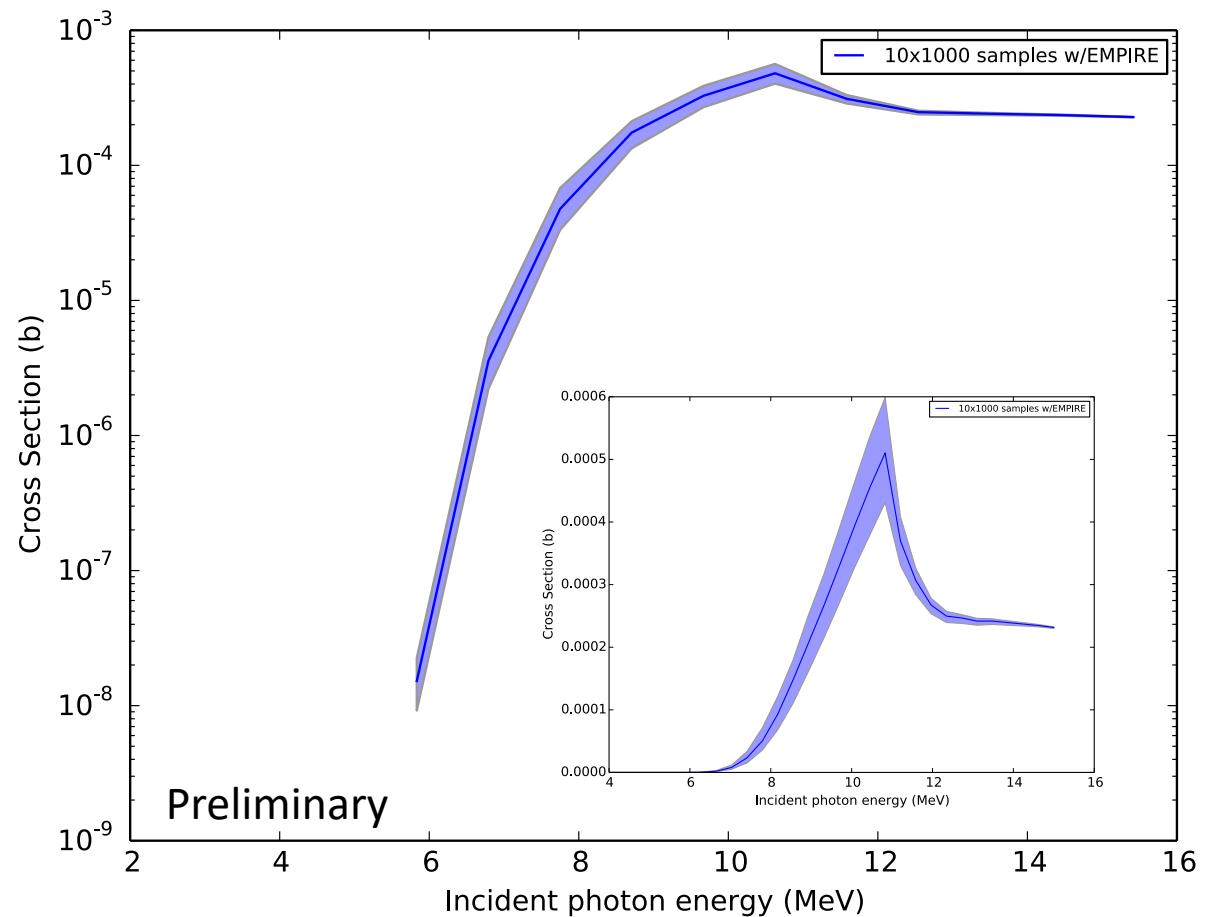
## Steps

1. Find the best theoretical model to describe  $(p, \gamma)$  data
2. Use that model to predict  $(\gamma, p)$  cross sections



# Ag103 ( $\gamma$ ,p) cross section from (p, $\gamma$ )-constrained parameters

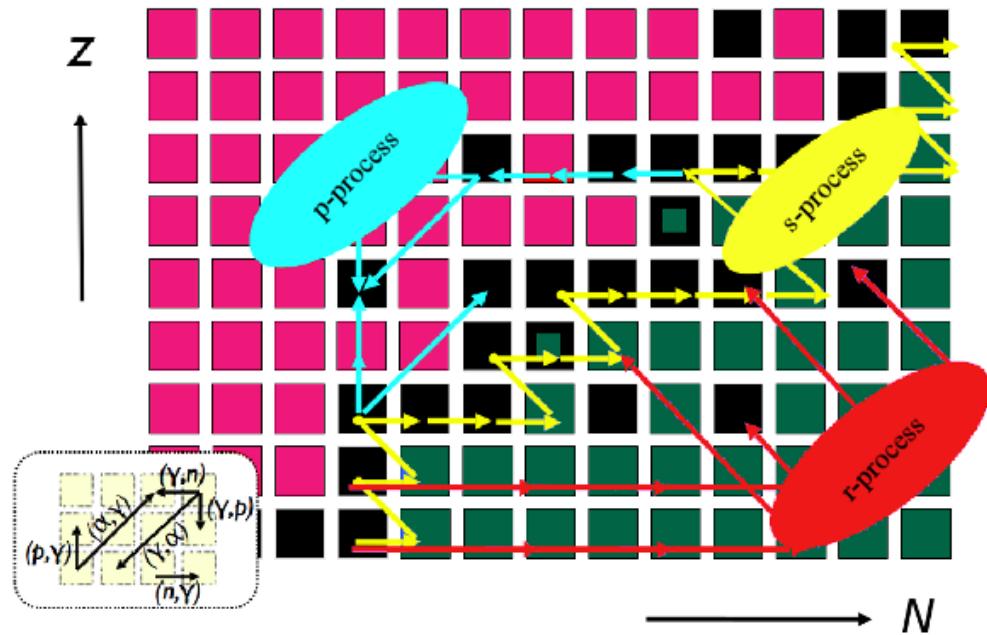
Distribution (blue band)  
reflects information  
encoded in experimental  
error bars



# Improved cross sections means improved astrophysics models

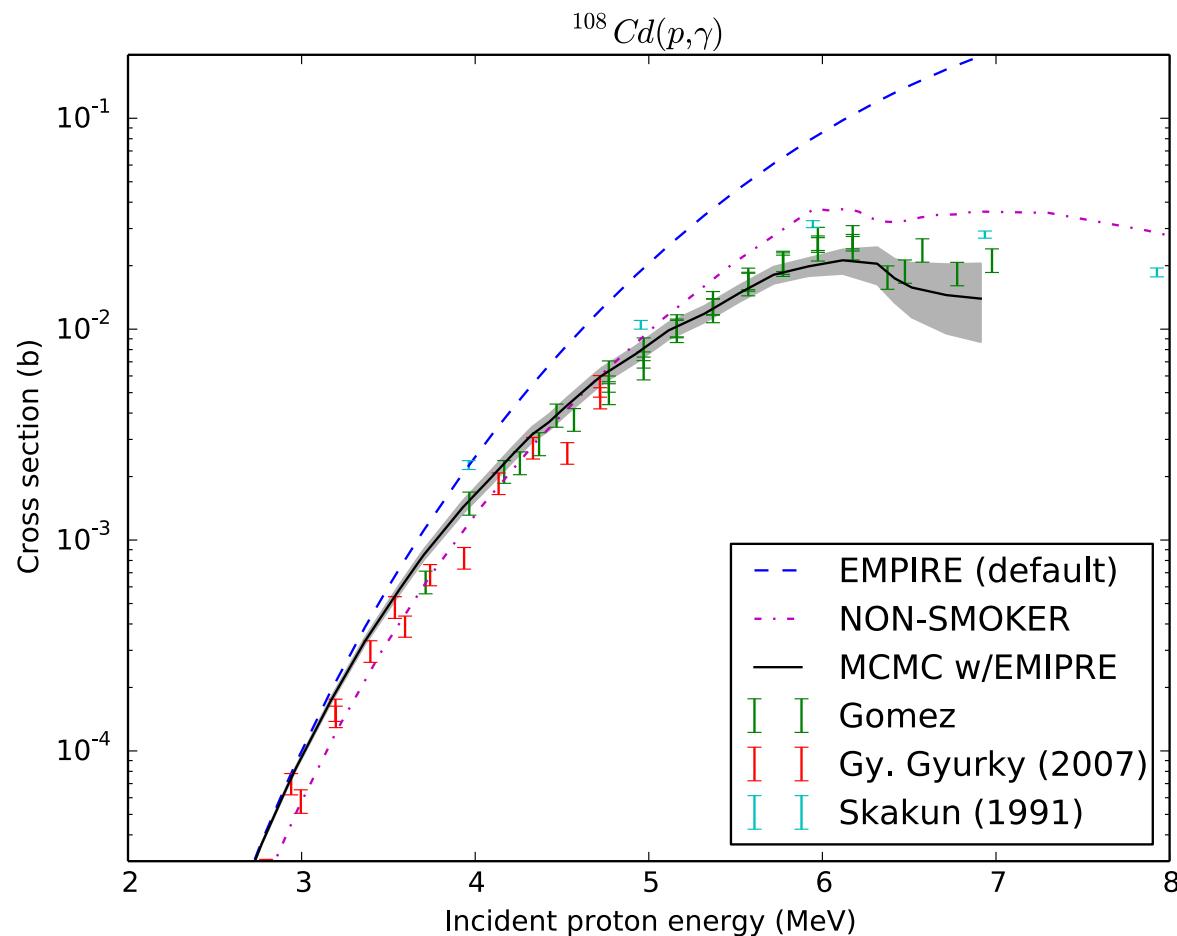
## Take aways

- Re-evaluated cross sections may change understanding of the p-process
- Error bars tell us how much we can infer from the new measurements

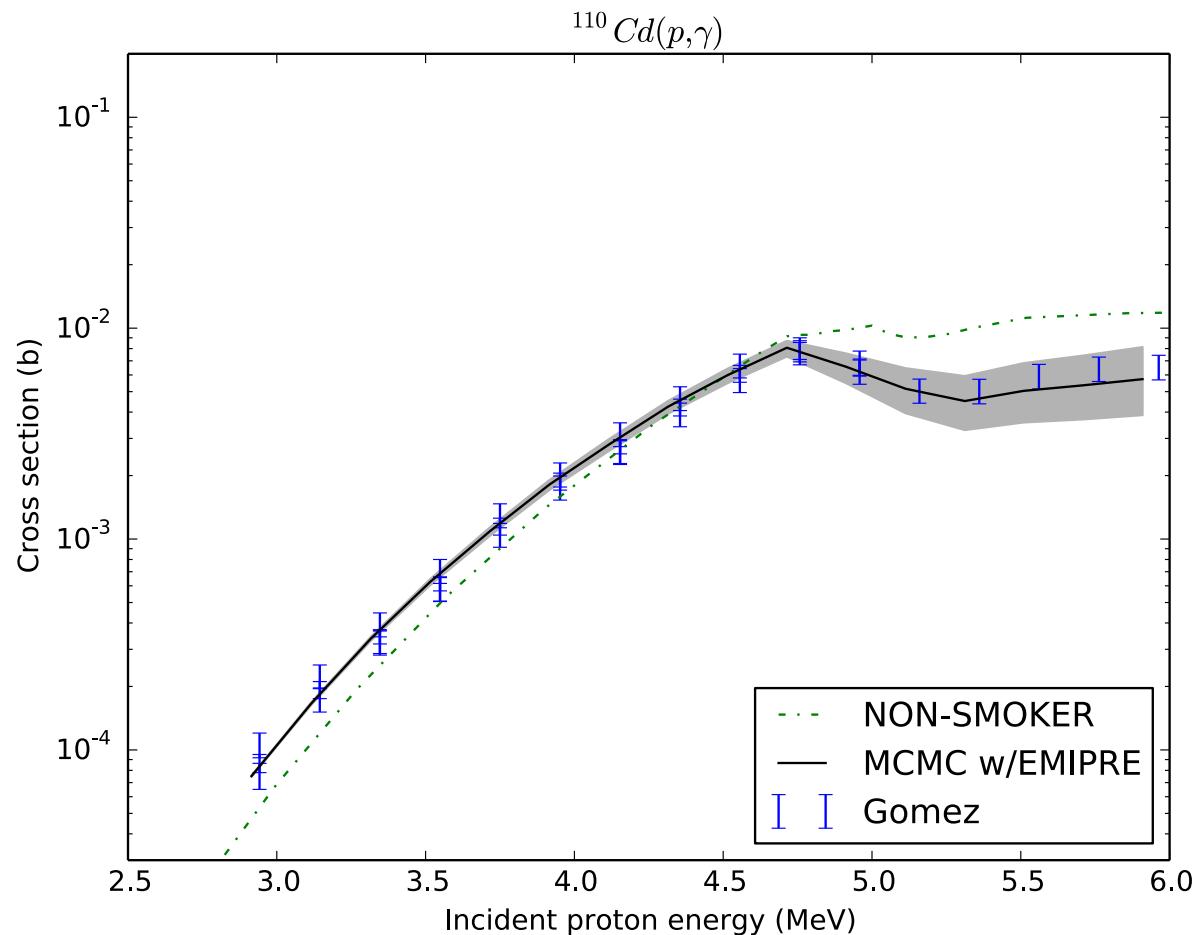


Images: Anna Simon, SuN  
GROUP @ NSCL  
[groups.nscl.msu.edu/SuN/](http://groups.nscl.msu.edu/SuN/)

# $p + {}^{108}\text{Cd} \rightarrow {}^{109}\text{In} + \gamma$



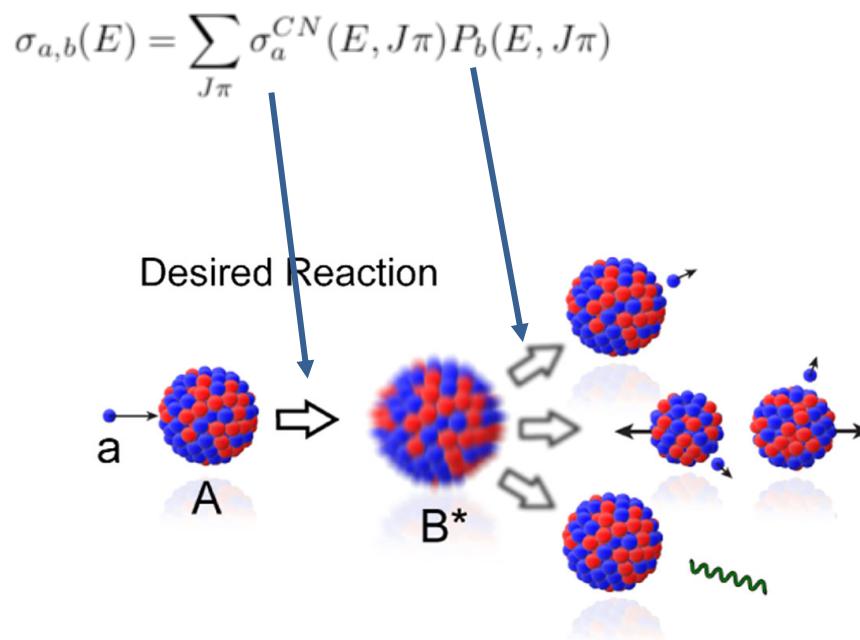
# $p + {}^{110}\text{Cd} \rightarrow {}^{111}\text{In} + \gamma$



# Hauser-Feshbach Model

Full Hauser-Feshbach formulation:

$$\frac{d\sigma_{\alpha\chi}^{HF}(E_\alpha)}{dE_{\chi'}} = \pi \frac{\lambda_\alpha^2}{2\pi} \sum_{J\pi} \omega_\alpha^J \sum_{lsl's'I'} \frac{T_{\alpha ls}^J(E_\alpha) T_{\chi' l's'}^J(E_{\chi'}) \rho_{I'}(U')}{\sum_{\chi''l''s''} T_{\chi''l''s''}^J(E_{\chi''}) + \sum_{\chi''l''s''I''} \int T_{\chi''l''s''}^J(E_{\chi''}) \rho_{I''}(U'') dE_{\chi''}} \quad (8)$$

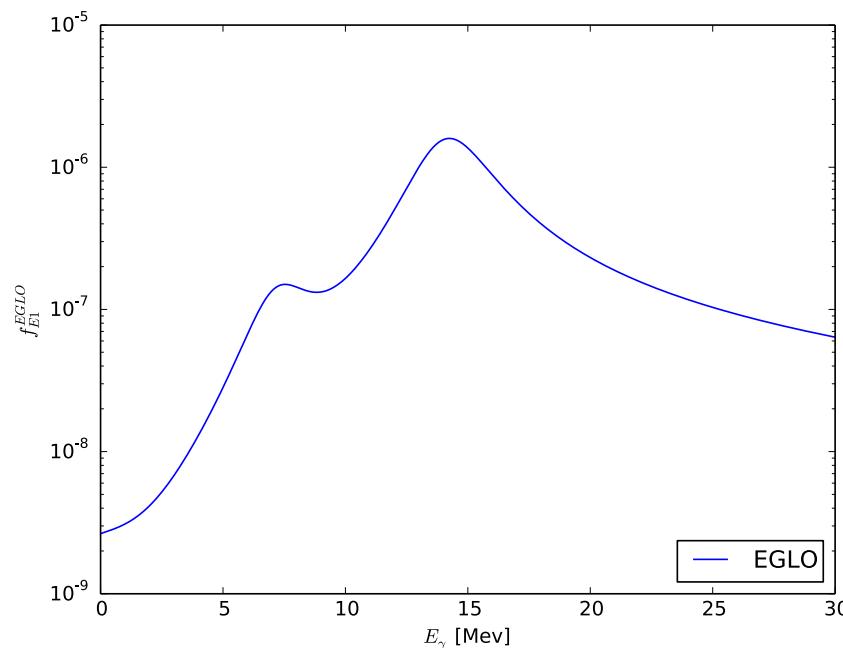


# Gamma-ray strength functions

$$f_{E1}^{EGLO}(E_\gamma) = 8.674 \times 10^{-8} \left\{ \sum_{i=1}^2 \Gamma_{E1i} \sigma_{E1i} \frac{E_\gamma \bar{\Gamma}(E_\gamma, T)}{(E_{E1i}^2 - E_\gamma^2)^2 + (E_\gamma \bar{\Gamma}(E_\gamma, T))^2} + \frac{0.7}{E_{E1i}^3} \bar{\Gamma} \right\}$$

The energy dependent width is

$$\bar{\Gamma}(E_\gamma, T) = \left[ k_0 + (1 - k_0) \frac{(E_\gamma - \epsilon)}{(E_{E1i} - \epsilon)} \right] \frac{\Gamma_{E1i} (E_\gamma^2 + 4\pi T^2)}{E_{E1i}^2}.$$



# Nuclear level density

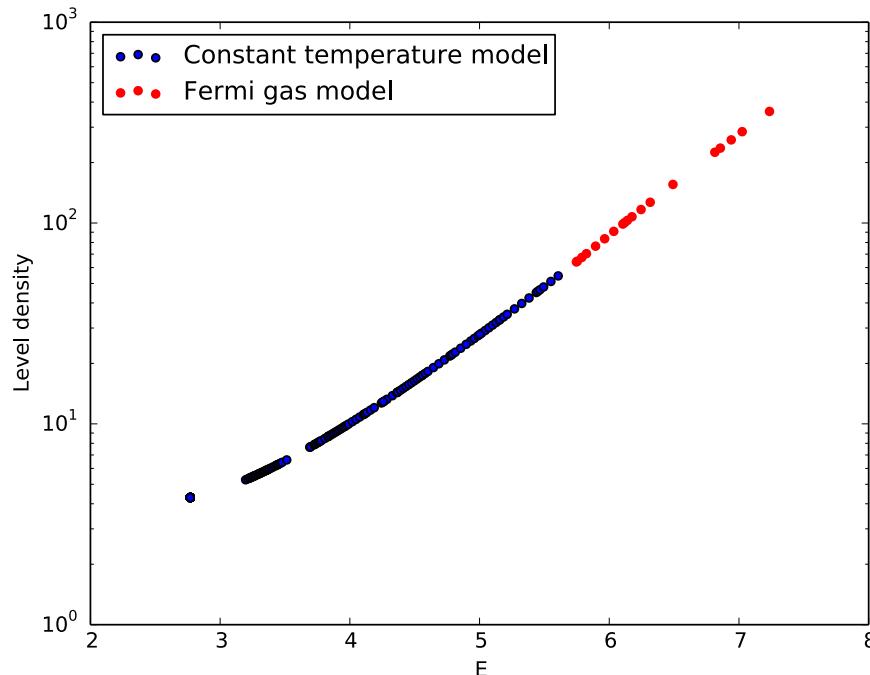
$$\rho(E_x) = \begin{cases} \rho^{CT}(E_x) = \frac{1}{T} \exp\left(\frac{E_x - E_0}{T}\right) & E_x \geq U_x \\ \rho^{FG}(E_x) \propto \frac{1}{\sigma^{3/2} a^{1/4} U^{5/4}} \exp\left(2\sqrt{aU} - \frac{(J+1/2)^2}{2\sigma^2}\right) & E_x > U_x \end{cases}$$

where

$$\sigma^2 = \lambda \sqrt{aU} A^{2/3} = 0.146 \sqrt{aU} A^{2/3}$$

$$a = \tilde{a}[1 + (1 - \exp(-\gamma U))\delta W/U] = E_x/T^2$$

$$U = E_x - \Delta$$



## PART 1:

- Collaboration with University of Notre Dame experimentalists using in inverse-reactions for indirect measurements
- Applying our theory capabilities to improve predictive power of their measurements

## PART 2:

- Sensitivity of an approximation method used for indirect measurements
- Applying this surrogate method to recently new data

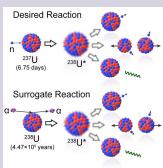
# Indirect Measurements are Necessary for Radioactive Target Measurements



Reactions on *radioactive targets* are difficult or impossible to measure.



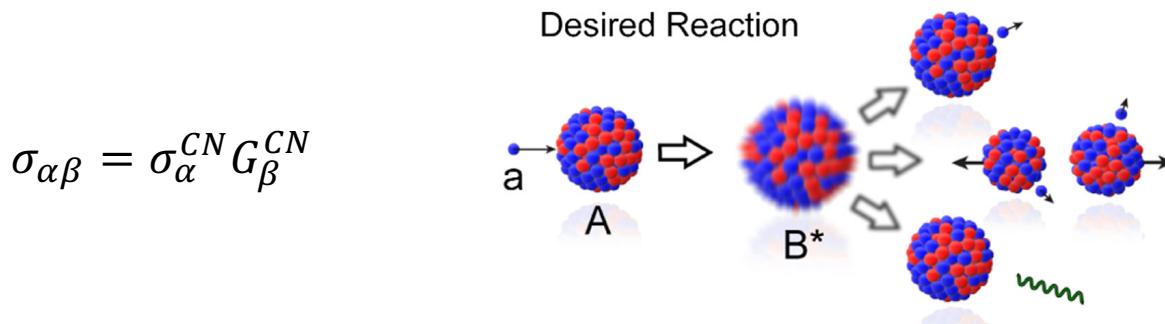
Theory of compound nuclear reactions limited by available nuclear structure.



*The Surrogate Method* allows indirect measurements of cross sections by combining surrogate data and theory.

# Hauser-Feshbach Theory models compound nuclear reactions in two stages

1. Formation of the compound nucleus (CN)
2. Decay of the CN



\* Conservation of energy, spin and parity

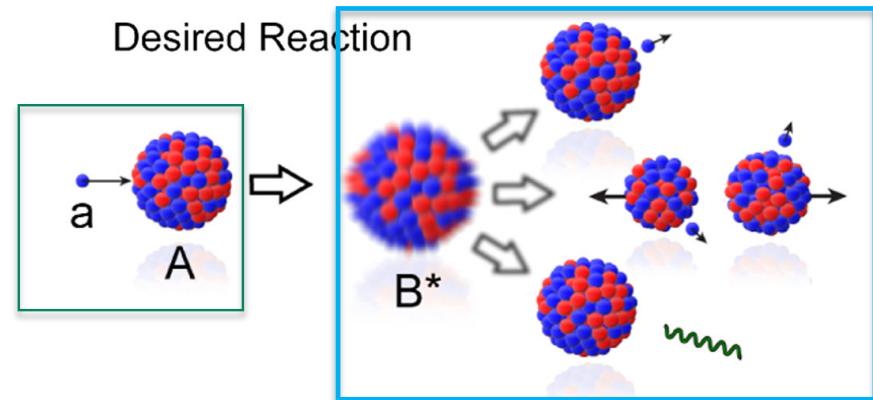
$$\sigma_{\alpha\beta}(E) = \sum_{J\pi} \sigma_{\alpha}^{CN}(E, J\pi) G_{\beta}^{CN}(E, J\pi)$$

# The Surrogate Idea: requires a lot of theory

Desired Reaction

$$\sigma_{\alpha\beta}(E) = \sum_{J\pi} \sigma_{\alpha}^{CN}(E, J\pi) G_{\beta}^{CN}(E, J\pi)$$

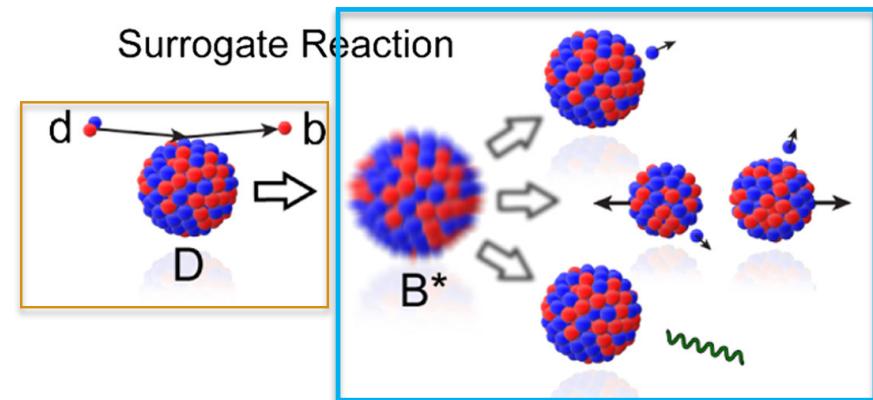
Desired Reaction



Surrogate Reaction

$$P_{\delta\beta}(E) = \sum_{J\pi} F_{\delta}^{CN}(E, J\pi) G_{\beta}^{CN}(E, J\pi)$$

Surrogate Reaction



# The Weiskopf-Ewing Limit applies under special conditions

$$\sigma_{\alpha\beta}(E) = \sum_{J\pi} \sigma_{\alpha}^{CN}(E, J\pi) G_{\beta}^{CN}(E, J\pi)$$

$E$  (CN) must be high ->  
decay into continuum of states

- Works for neutron induced fission ( $n,f$ )
- Fails for neutron capture



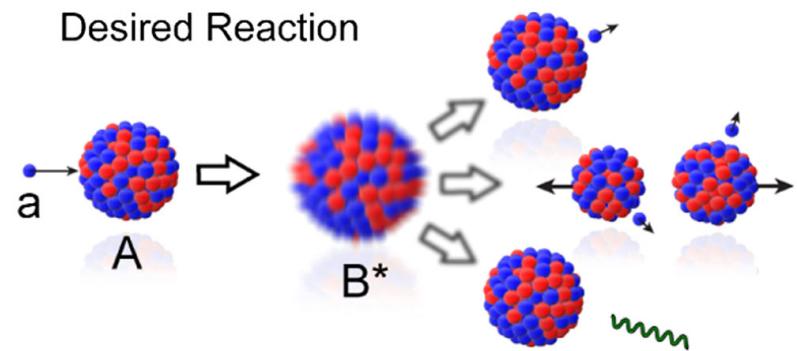
$$\sigma_{\alpha\beta}^{WE}(E) = \sigma_{\alpha}^{\prime CN}(E) G_{\beta}^{CN}(E)$$

# The Weiskopf-Ewing Limit Greatly Simplifies the Surrogate Method

Desired Reaction

$$\sigma_{\alpha\beta}^{WE}(E) = \sigma_{\alpha}'^{CN}(E) G_{\beta}^{CN}(E)$$

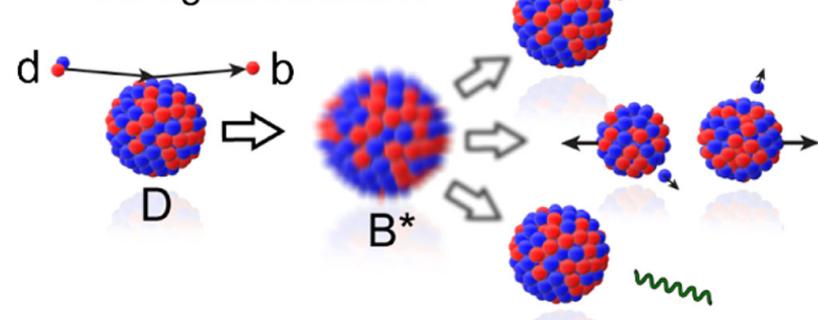
Desired Reaction



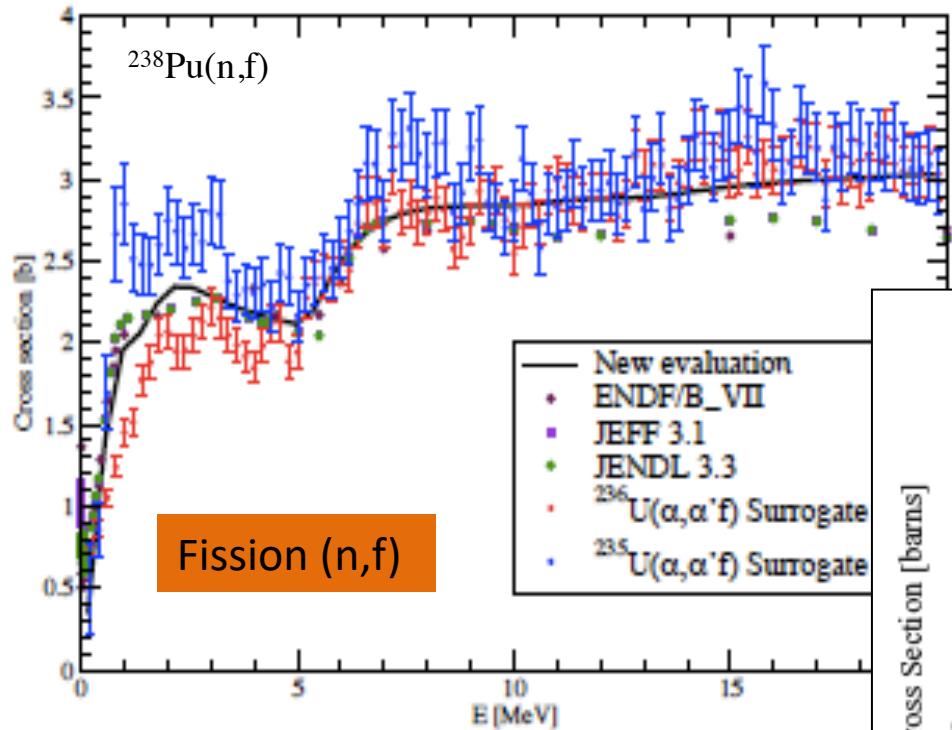
Surrogate Reaction

$$P_{d\beta}^{WE}(E) = G_{\beta}^{CN}(E)$$

Surrogate Reaction

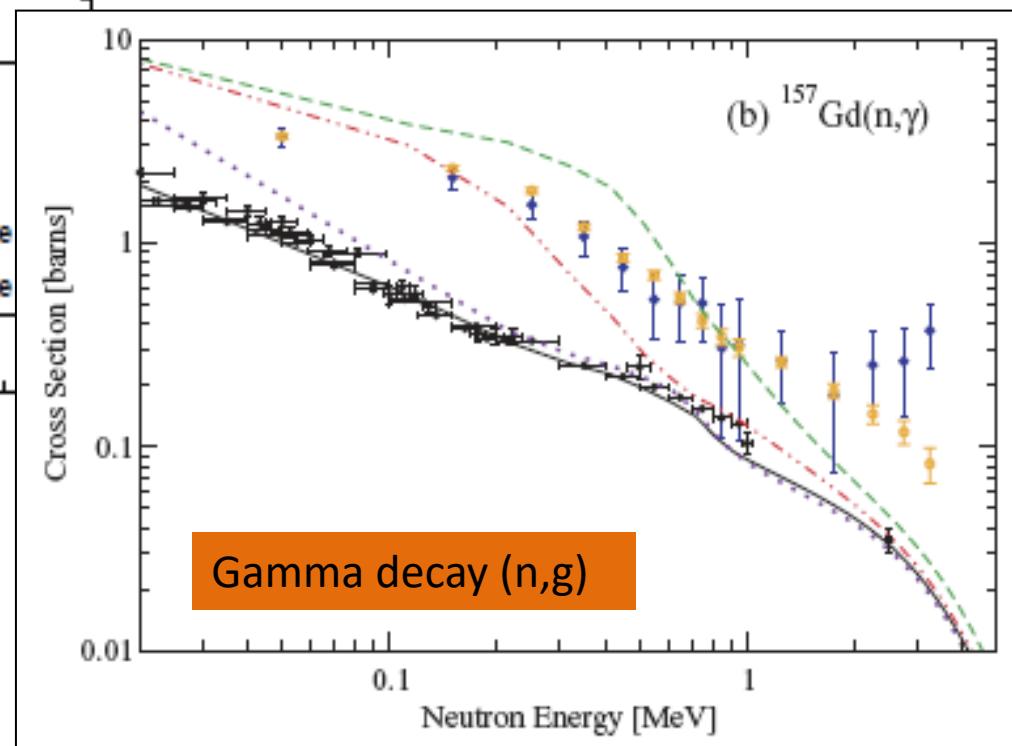


# Does the Weisskopf-Ewing approximation work? Yes, for fission ( $n,f$ ), not for gamma decay ( $n,g$ ).



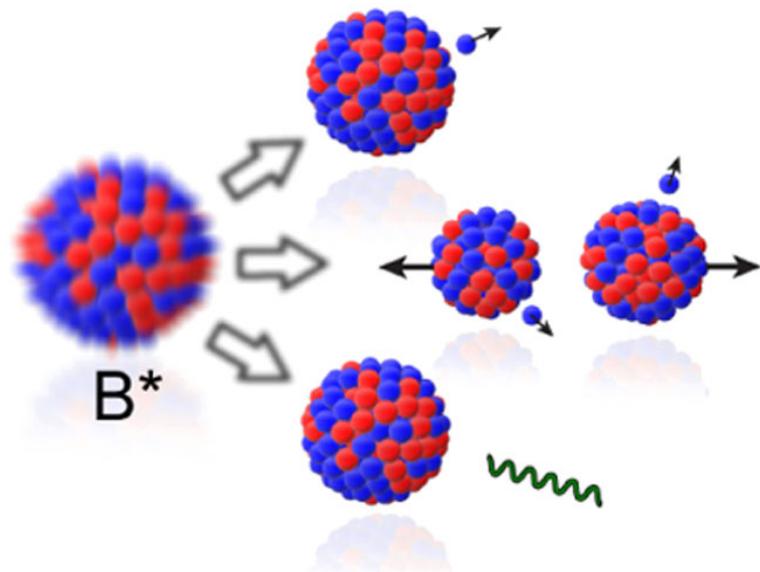
Ressler, Burke, Escher,...  
PRC 83 (2011) 054610

J. Escher and F.S. Dietrich, PRC 81 (2010) 024612  
N. Scielzo, J. Escher, et al., PRC 81 (2010) 034608

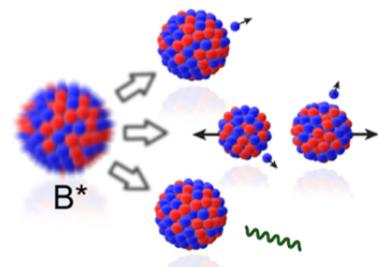


# What about other neutron induced reactions?

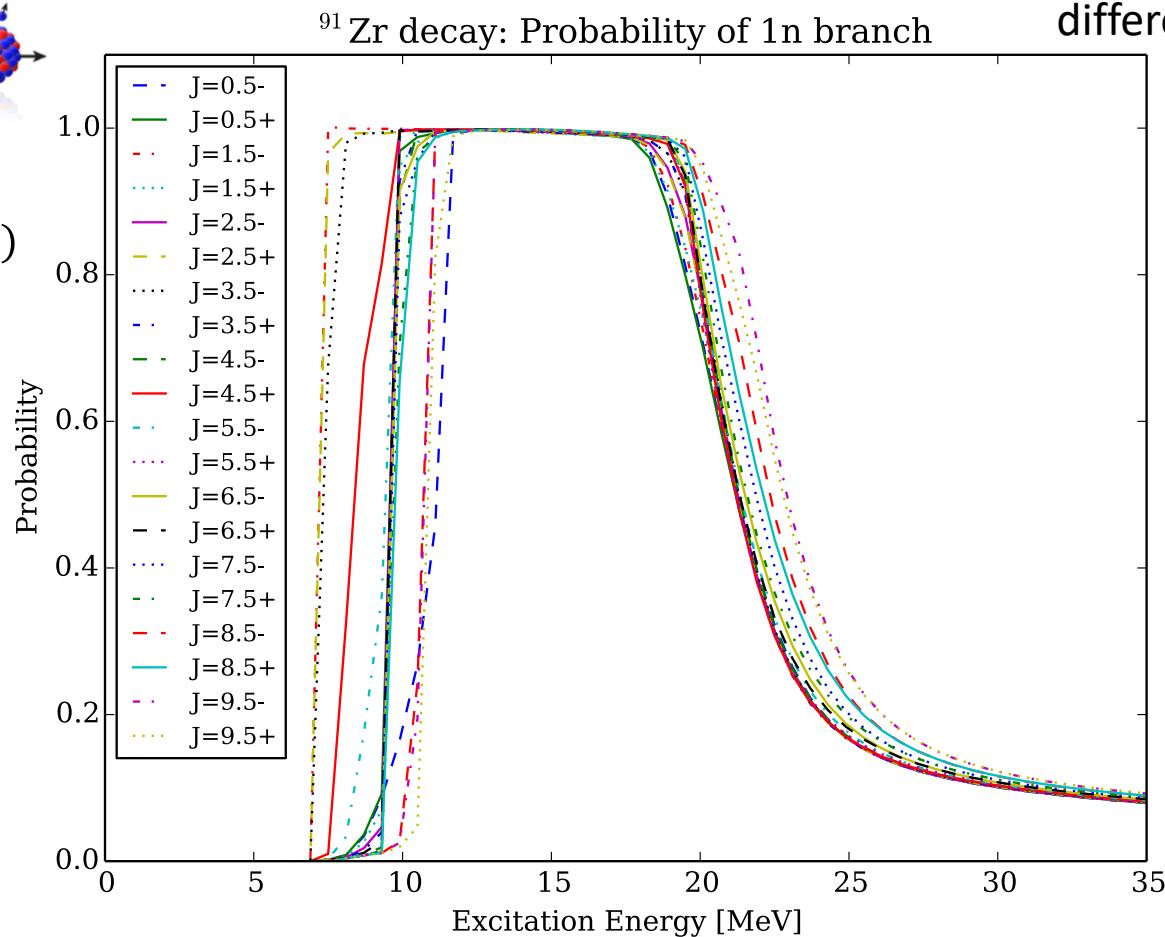
- Remains untested for  $(n,n')$  and  $(n,2n)$  reactions
- We can test the model sensitivity to Spin and Parity



# I tested the Weisskopf-Ewing approximation for the (n,n') reaction on 90Zr targets

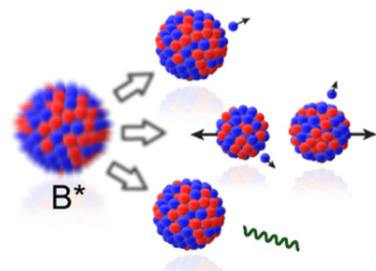


$$P_{d\beta}^{WE}(E) = G_\beta^{CN}(E)$$

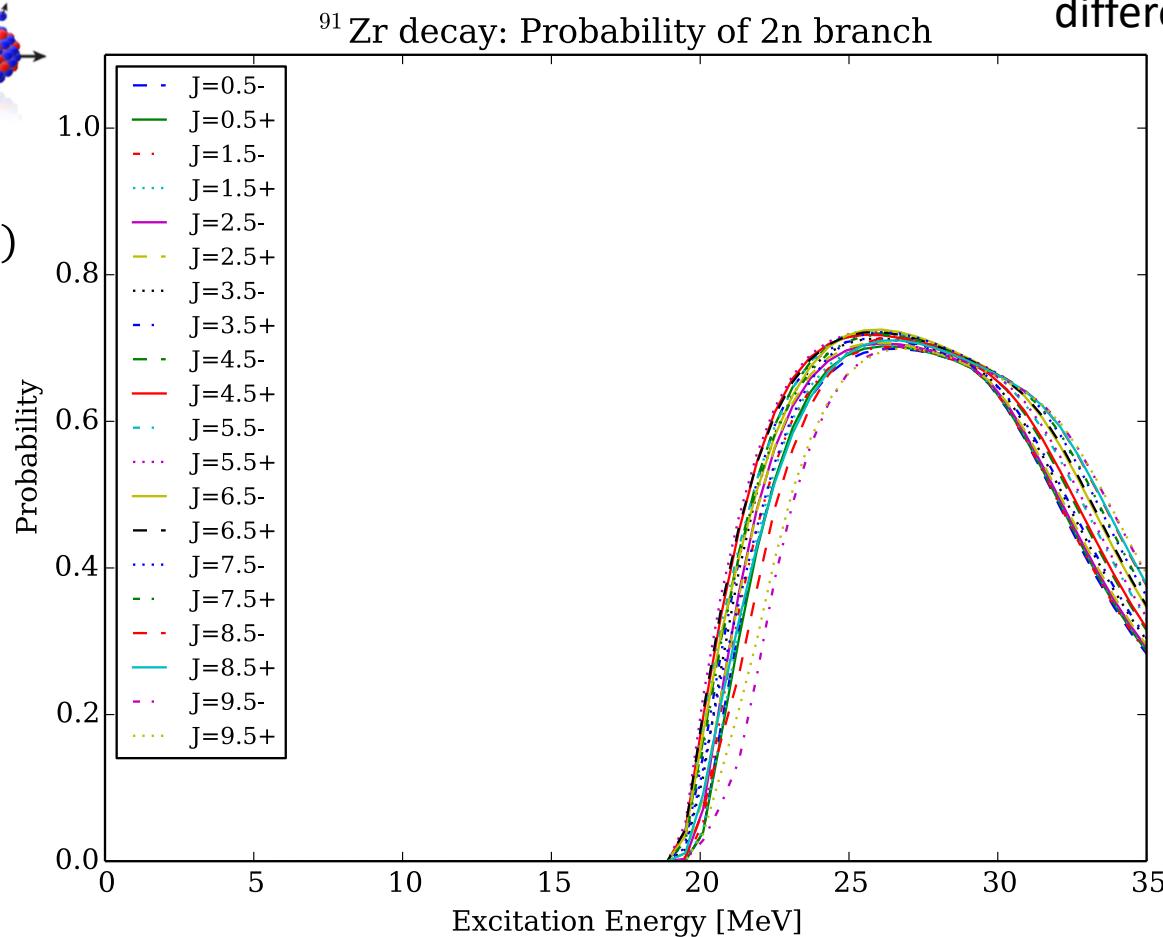


How does a nucleus with different spins decay?

# I tested the Weisskopf-Ewing approximation for the (n,2n) reaction on $^{90}\text{Zr}$ targets



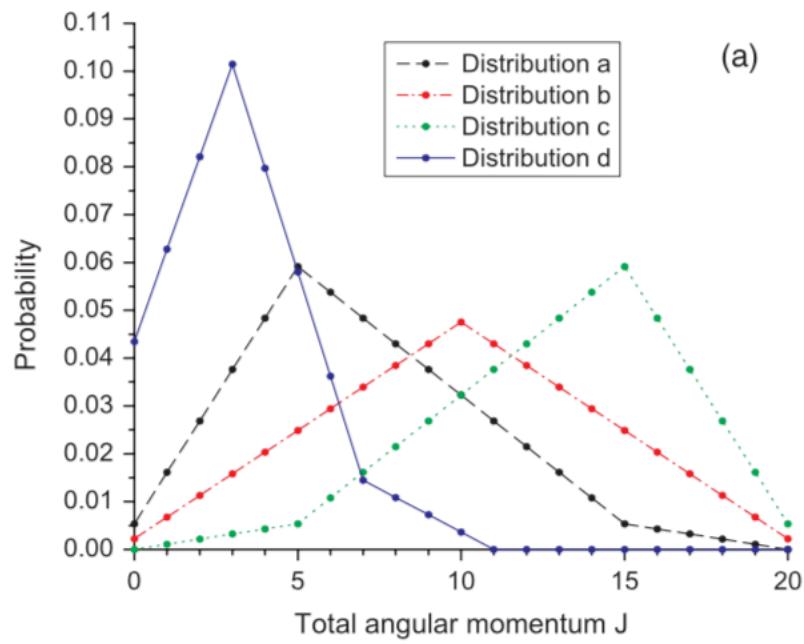
$$P_{d\beta}^{WE}(E) = G_\beta^{CN}(E)$$



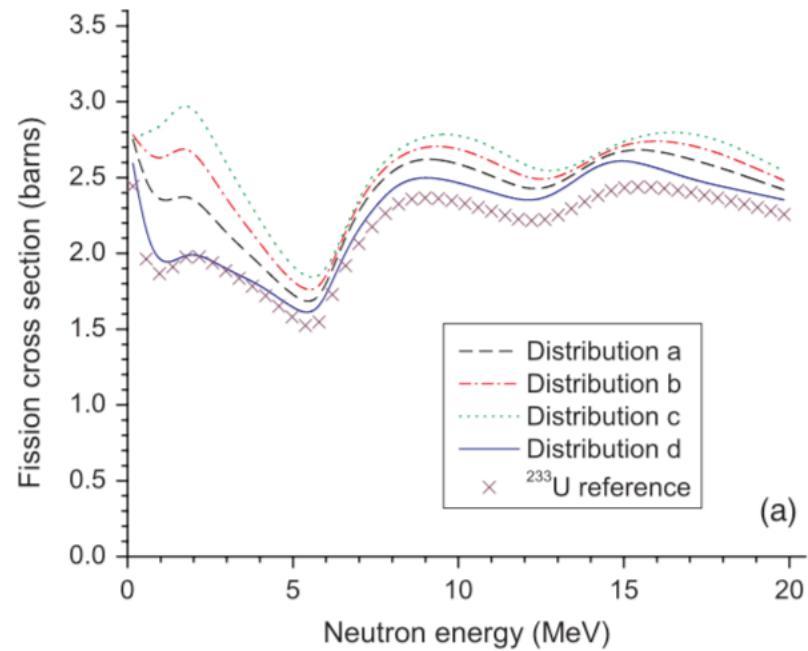
How does a nucleus with different spins decay?

# A reaction produces nuclei with a range of spins and parities!

JUTTA E. ESCHER AND FRANK S. DIETRICH

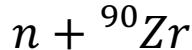
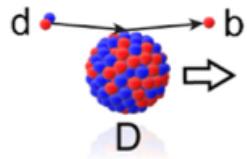


JUTTA E. ESCHER AND FRANK S. DIETRICH

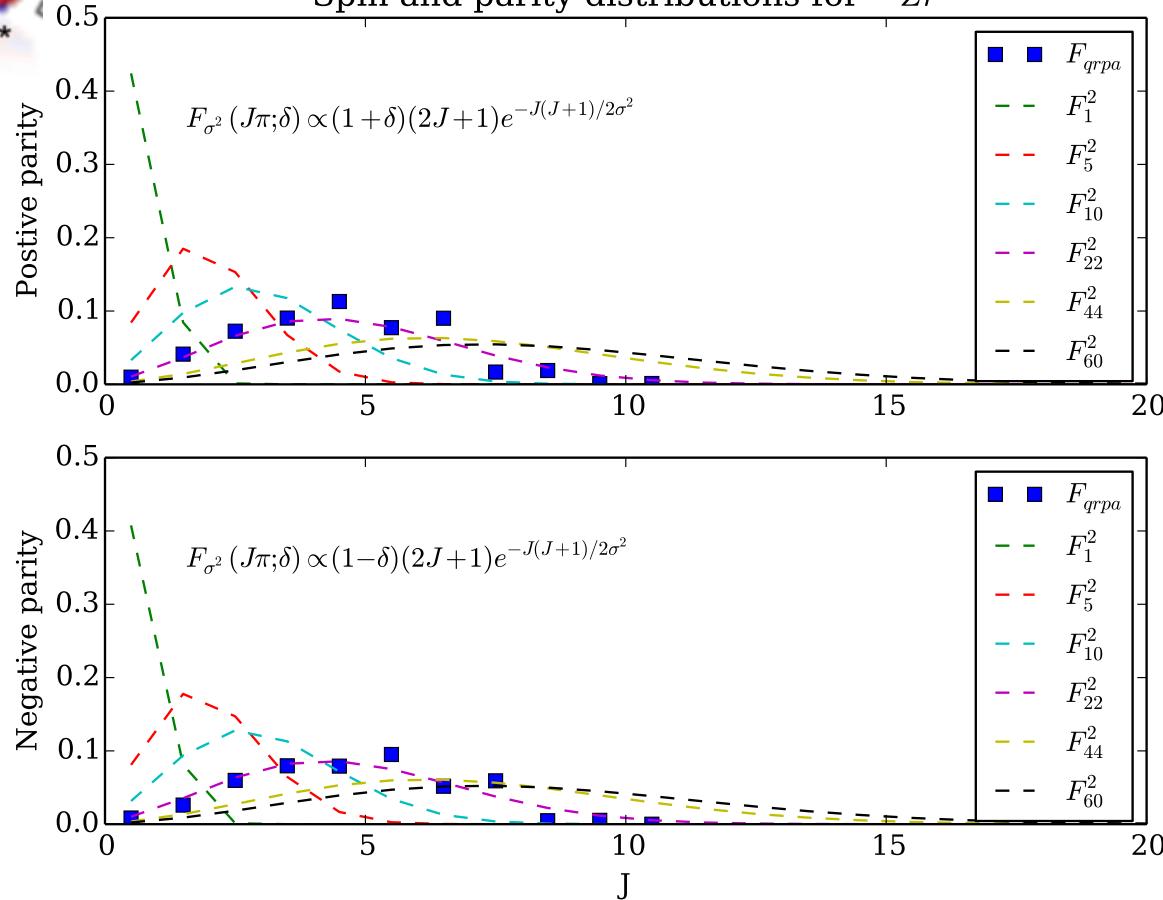


Escher, Dietrich, Phys. Rev. C 74, 054601 (2006)

# I chose an analytic distribution with two parameters to occupy the compound nucleus



Spin and parity distributions for  ${}^{91}\text{Zr}$

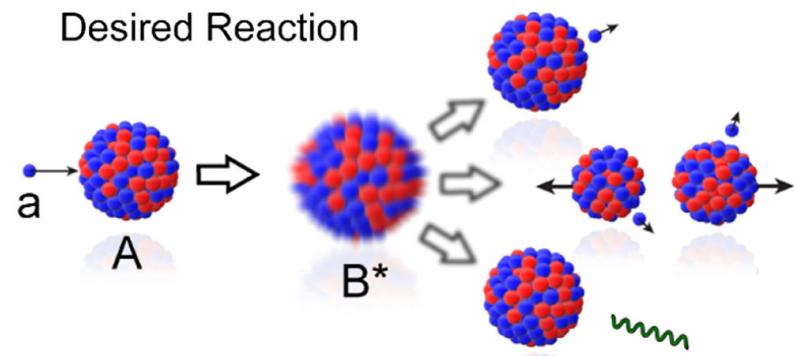


# Full picture: Weisskopf-Ewing approximation

Desired Reaction

$$\sigma_{\alpha\beta}^{WE}(E) = \sigma_{\alpha}'^{CN}(E) G_{\beta}^{CN}(E)$$

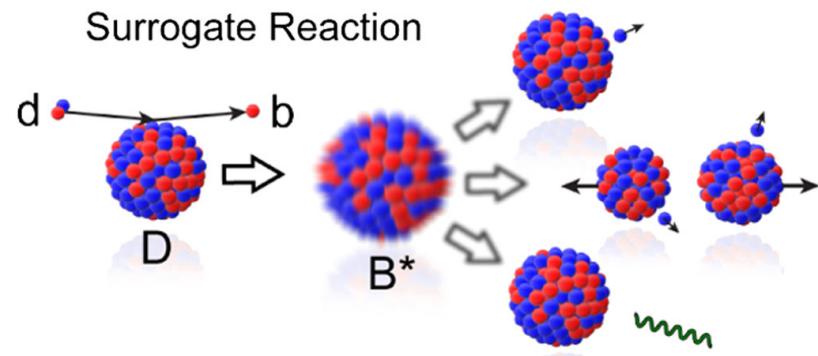
Desired Reaction



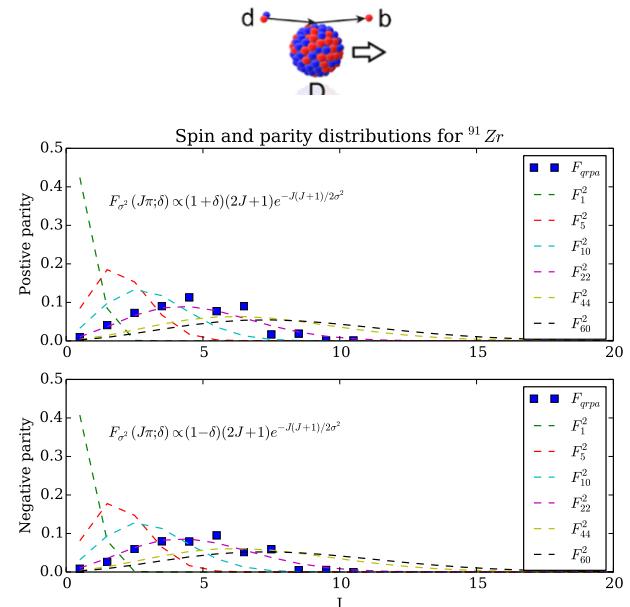
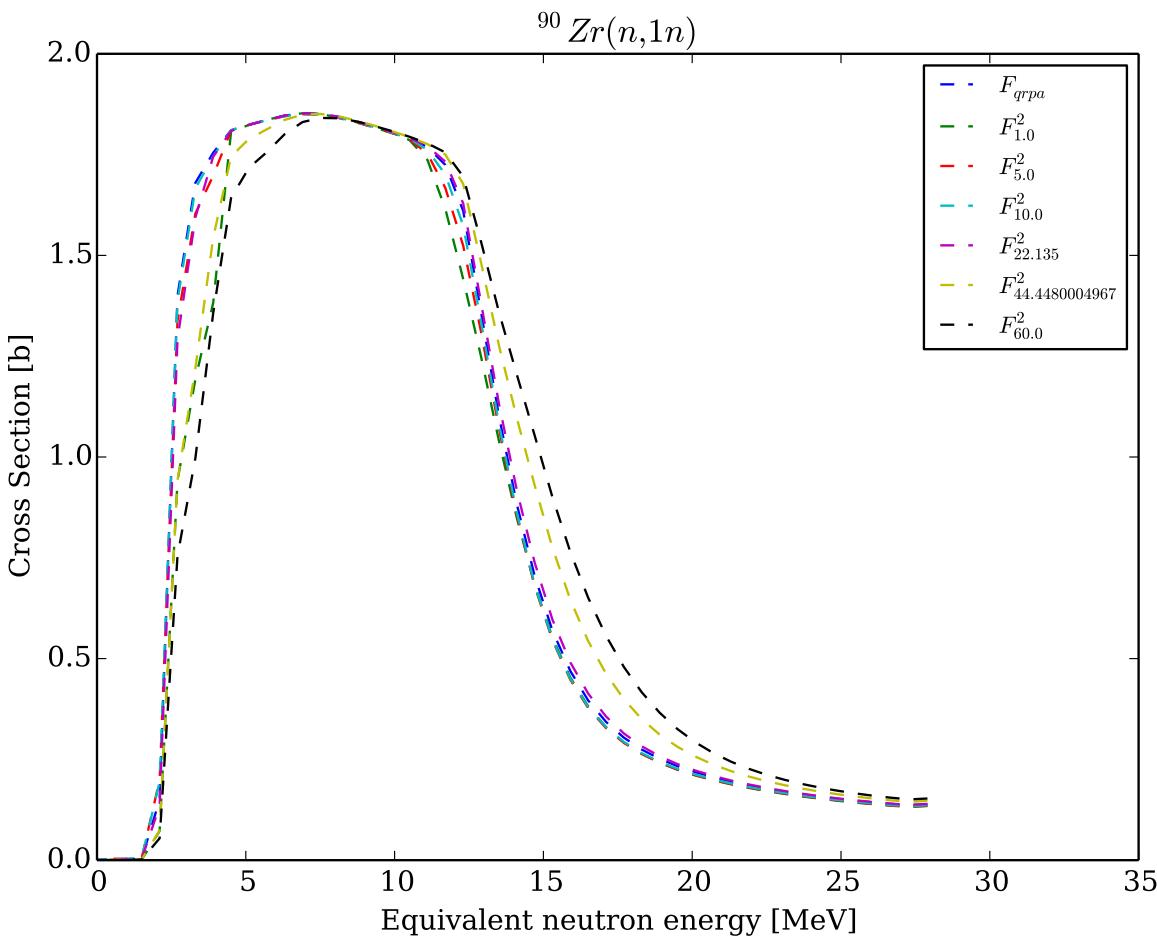
Surrogate Reaction

$$P_{d\beta}^{WE}(E) = G_{\beta}^{CN}(E)$$

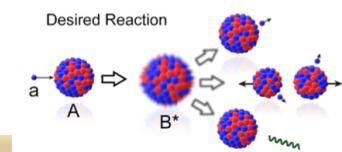
Surrogate Reaction



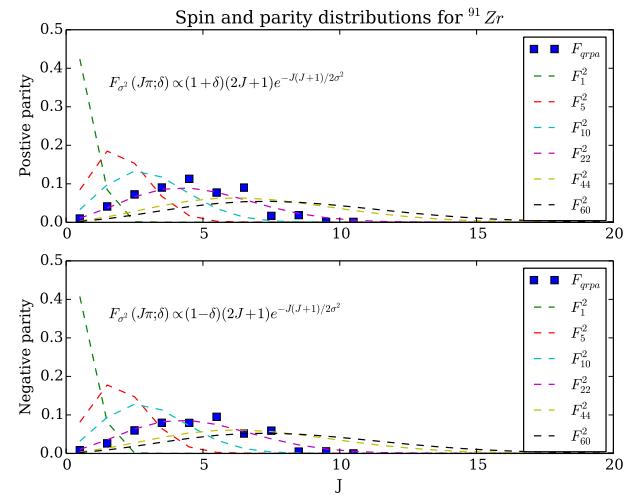
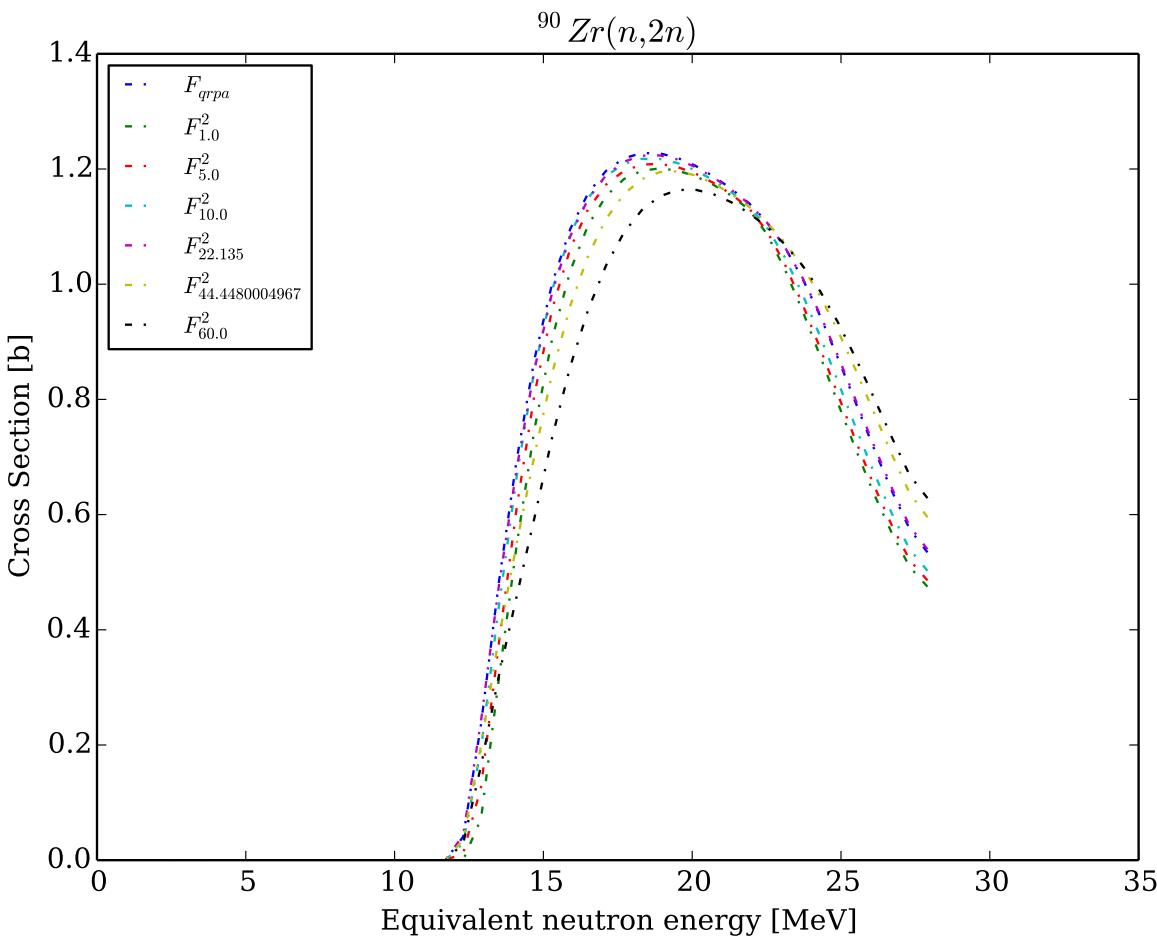
# The tests I preformed suggest that the WE approximation is valid for these reactions



$$\sigma_{\alpha\beta}^{WE}(E) = \sigma_{\alpha}'^{CN}(E) G_{\beta}^{CN}(E)$$

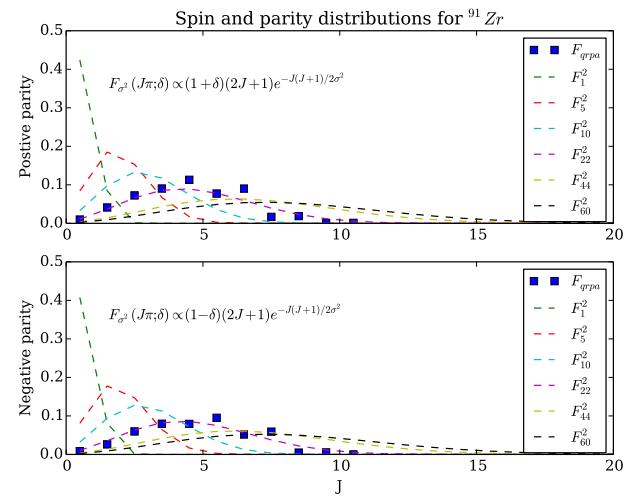
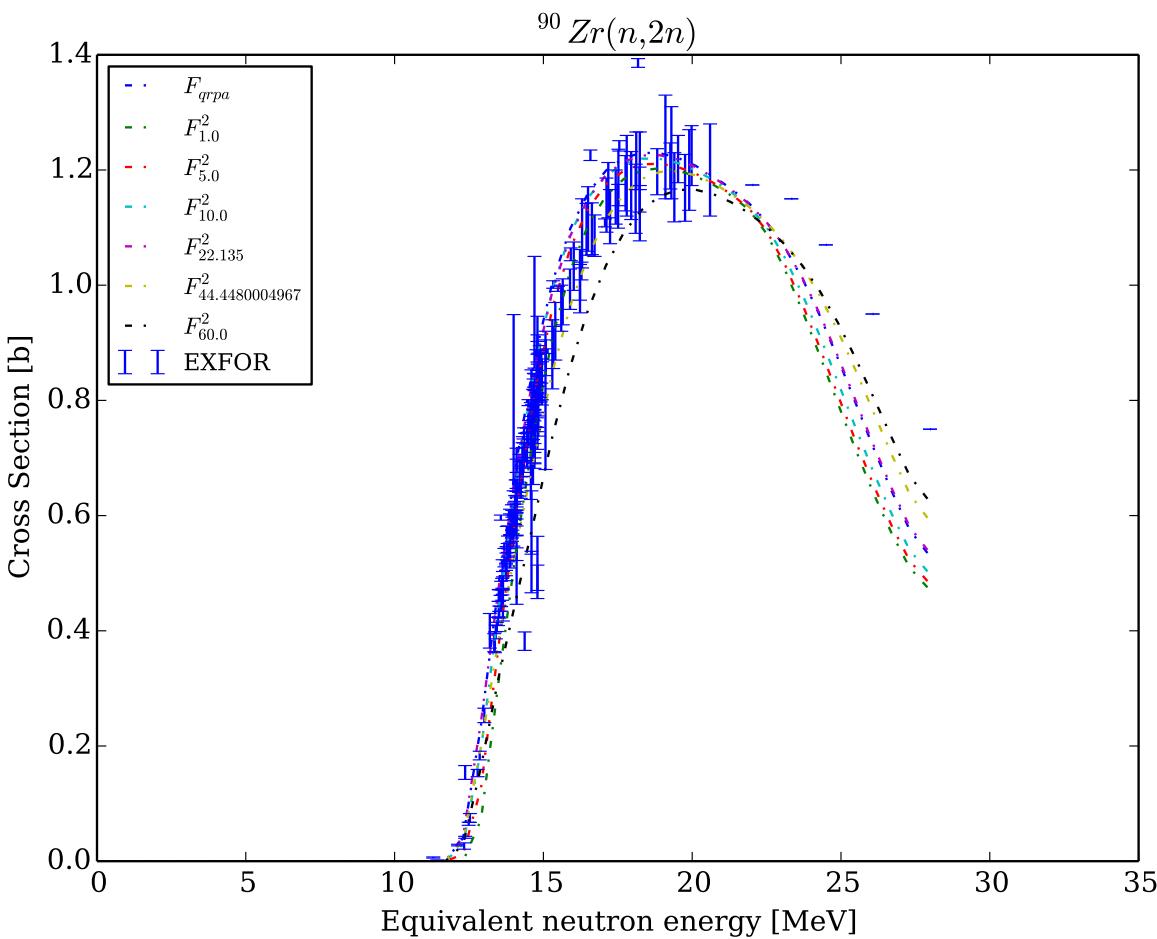


# The tests I preformed suggest that the WE approximation is valid for these reactions



$$\sigma_{\alpha\beta}^{WE}(E) = \sigma_{\alpha}'^{CN}(E) G_{\beta}^{CN}(E)$$

# Variance of the predictions fall within current experimental error



$$\sigma_{\alpha\beta}^{WE}(E) = \sigma_{\alpha}'^{CN}(E) G_{\beta}^{CN}(E)$$

# I tested the WE approximation for (n,n) and (n,2n) on 90Zr

1. Generate artificial spin and parity distributions
2. Create simulated surrogate data
3. Compute the WE limit using the simulated surrogate data

$$F(J, \pi)$$

$$P_{d\beta}^{WE}(E) = G_\beta^{CN}(E)$$

$$\sigma_{\alpha\beta}^{WE}(E) = \sigma_{\alpha}^{\prime CN}(E) P_{d\beta}^{CN}(E)$$

The simulated predictions  $\sigma_{\alpha\beta}^{WE}$  are sufficiently similar for reasonable  $F(J, \pi)$



The WE limit is valid

# Results are promising

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Conclusions:

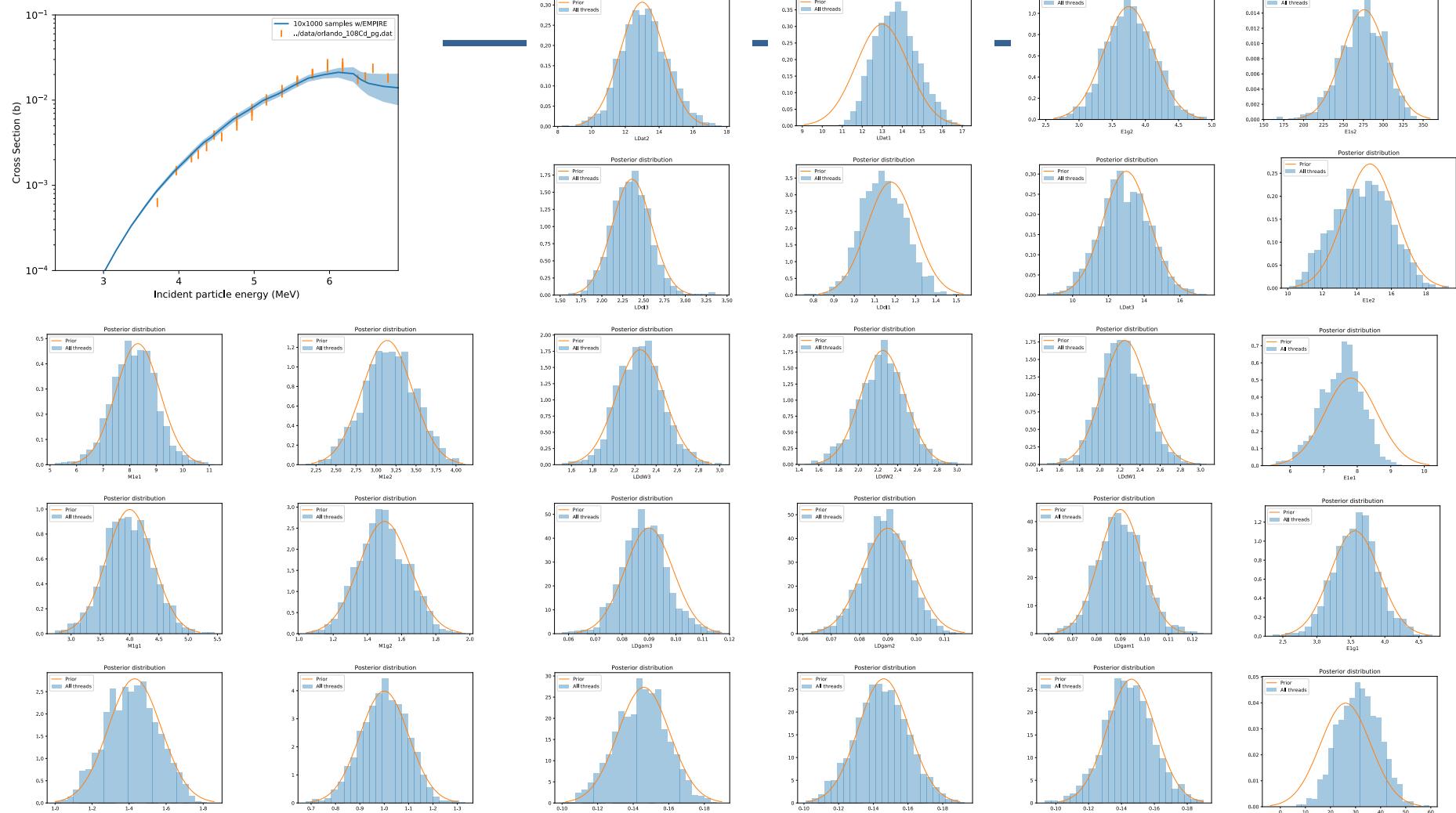
- $(n,n')$  and  $(n,2n)$  reactions on Zr are much less sensitive to spin and parity than capture reactions
- Whether this is sufficient depends on the application
- Further investigations are required, but results are promising

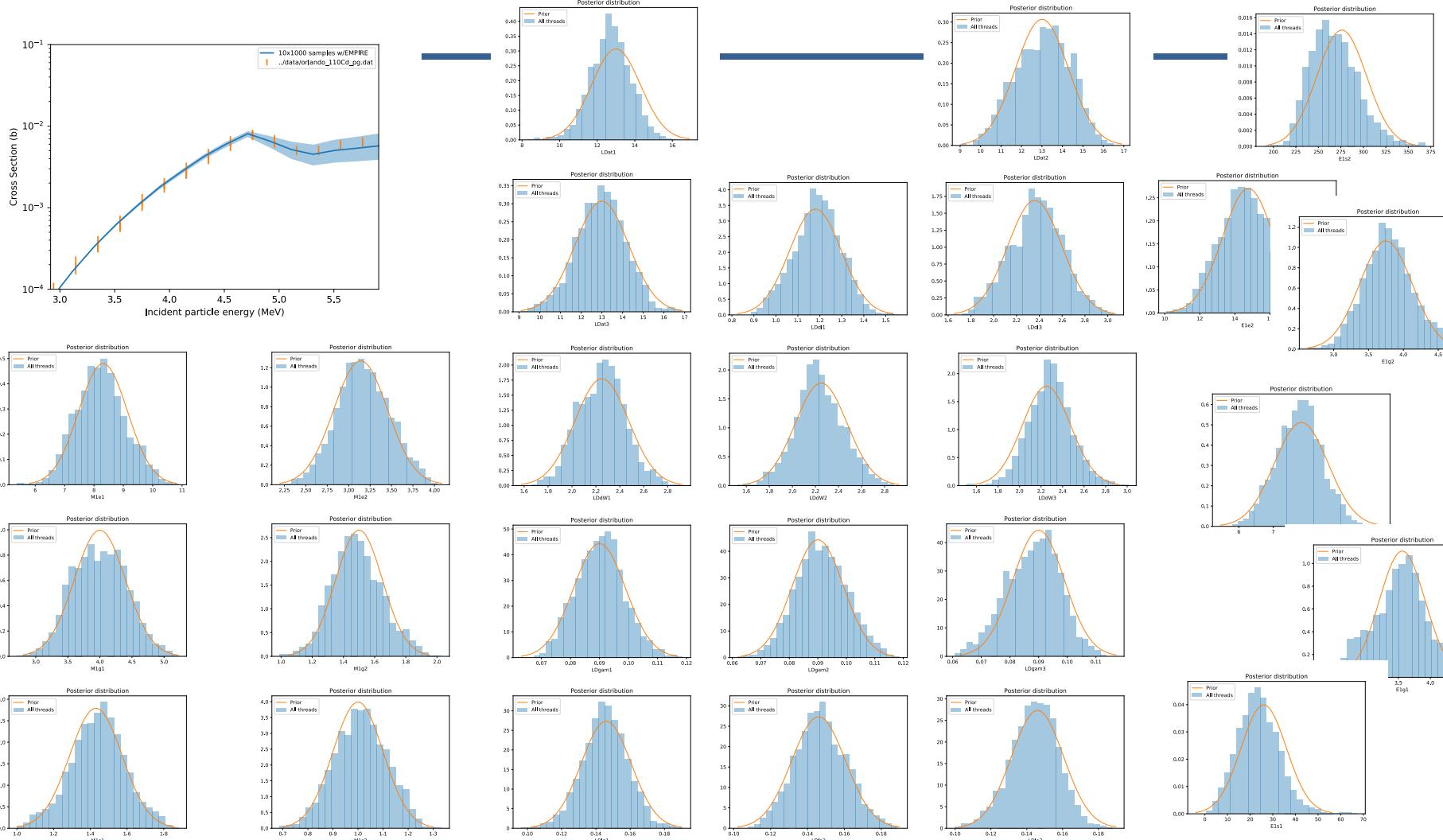
Next steps:

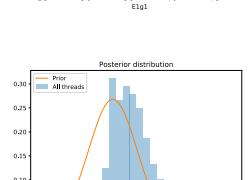
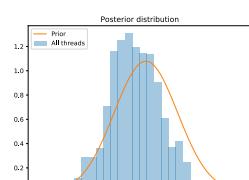
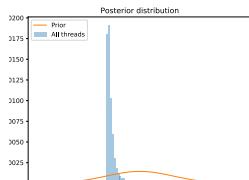
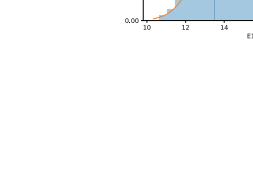
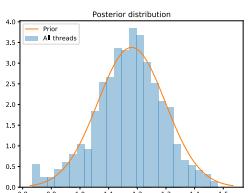
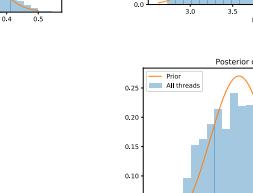
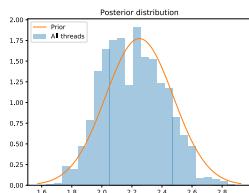
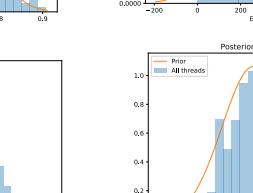
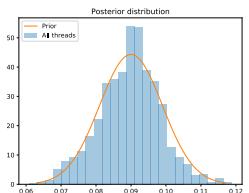
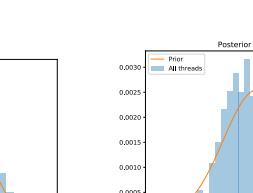
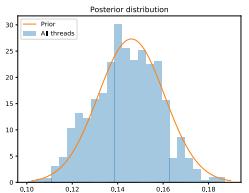
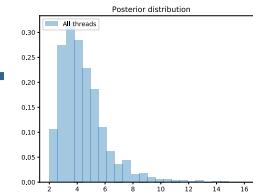
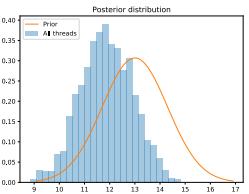
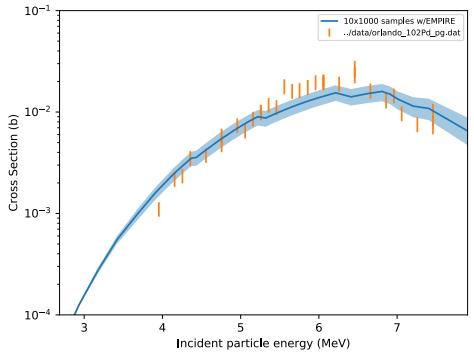
- Investigate effects of pre-equilibrium
- Apply this theory to new Gadolinium data



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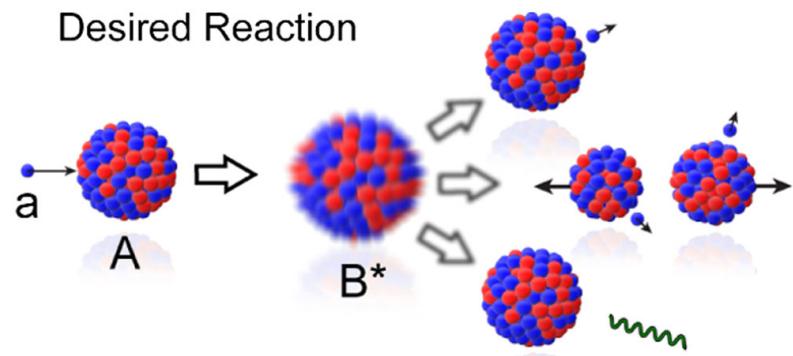
# The Surrogate Idea

Desired Reaction

$$\sigma_{\alpha\beta}(E) = \sum_{J\pi} \sigma_{\alpha}^{CN}(E, J\pi) G_{\beta}^{CN}(E, J\pi)$$

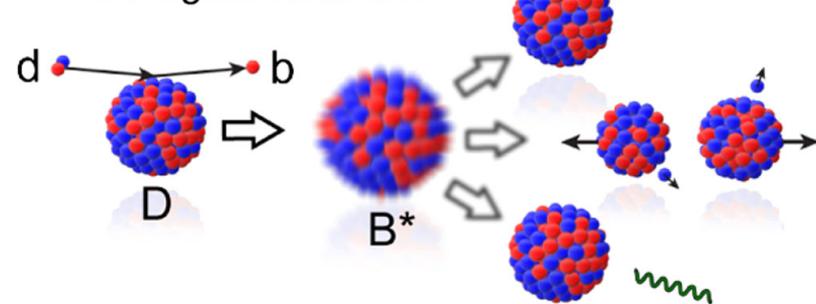
Hard part

Desired Reaction



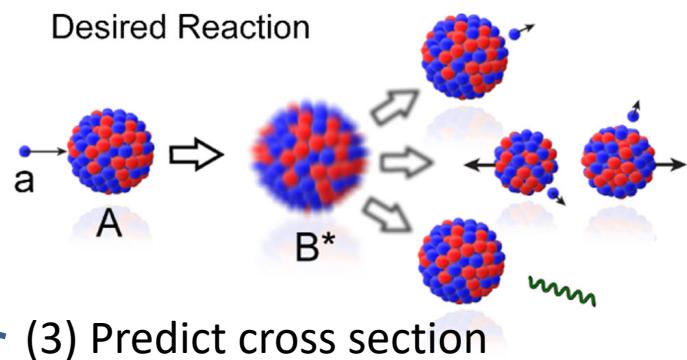
Surrogate Reaction

Surrogate Reaction



# The (General) Surrogate Method Combines Surrogate Data and Theory

- $\sigma_{\alpha\beta}(E) = \sum_{J\pi} \sigma_{\alpha}^{CN}(E, J\pi) G_{\beta}^{CN}(E, J\pi)$ 
  - $\sigma_{\alpha}^{CN}(E, J\pi)$ : well constrained
  - $G_{\beta}^{CN}(E, J\pi)$ : lack of nuclear structure info



- $P_{\delta\beta}(E) = \sum_{J\pi} F_{\delta}^{CN}(E, J\pi) G_{\beta}^{CN}(E, J\pi)$

(2) Constrain parameters

