

Oliver Gorton
 San Diego State University
 Department of Physics
 Physics 610 by Dr. Johnson
 Master Homework Problem

”For a free particle, compute the time evolution of a Gaussian packet and how its expectation values change in time.”

The Hamiltonian for a free particle is just the kinetic energy,

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \quad (1)$$

We start with the Gaussian wavepacket centered at $x = x_0$, which we write as

$$\Psi(x, t = 0) = \Psi_0(x) = (2\pi\sigma^2)^{-1/4} e^{-\frac{(x-x_0)^2}{4\sigma^2}} \quad (2)$$

1 Normalization

”For the general Gaussian wavefunction $\Psi(x) = Ae^{(-ax^2/2)}$, with complex valued a , find the normalization A .”

For normalization we require that $1 = \langle \Psi | \Psi \rangle$. In position space we have that:

$$\begin{aligned} 1 &= \int_{-\infty}^{+\infty} dx \Psi^*(x) \Psi(x) \\ &= \int_{-\infty}^{+\infty} dx A^* e^{-\frac{a^* x^2}{2}} A e^{-\frac{ax^2}{2}} \\ &= |A|^2 \int_{-\infty}^{+\infty} dx e^{-(a^*+a)x^2/2} \\ &= |A|^2 \int_{-\infty}^{+\infty} dx e^{-\text{Re}(a)x^2} \end{aligned} \quad (3)$$

To continue we need to solve the Gaussian integral $I = \int_{-\infty}^{+\infty} dx e^{-ax^2}$. To do this we consider the quantity I^2 , make a change of variables to polar coordinates, and solve by a

simple substitution.

$$\begin{aligned}
I^2 &= \left(\int_{-\infty}^{+\infty} dx e^{-ax^2} \right) \left(\int_{-\infty}^{+\infty} dy e^{-ay^2} \right) \\
&= \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} dx dy e^{-a(x^2+y^2)} \\
&= \int_0^{2\pi} \int_0^{+\infty} r d\theta dr e^{-ar^2} \\
&= 2\pi \int_0^{+\infty} \frac{1}{2} du e^{-au} \\
&= \pi \left(-\frac{1}{a} e^{-au} \Big|_0^{+\infty} \right) \\
&= \frac{\pi}{a}
\end{aligned} \tag{4}$$

We now have all we need to complete the normalization. Resuming from the last line of (3),

$$\begin{aligned}
1 &= |A|^2 \int_{-\infty}^{+\infty} dx e^{-Re(a)x^2} \\
&= |A|^2 \sqrt{\frac{\pi}{Re(a)}}
\end{aligned} \tag{5}$$

Thus, as long as $Re(A) > 0$ (otherwise the magnitude of A is imaginary),

$$\boxed{A = \pm \left(\frac{Re(a)}{\pi} \right)^{\frac{1}{4}}} \tag{6}$$

2 Expectation value $\langle x^2 \rangle$

”For the normalized Gaussian wavefunction in problem 1, find the expectation value $\langle x^2 \rangle$.”

For a general operator (observable) $\hat{\Omega}$ the expectation value in a state Ψ is calculated as $\langle \hat{\Omega} \rangle = \langle \Psi | \hat{\Omega} | \Psi \rangle$. Using the positive normalization found in problem 1, we arrive at the wavefunction

$$\Psi(x) = \left(\frac{Re(a)}{\pi} \right)^{\frac{1}{4}} e^{-\frac{ax^2}{2}} \tag{7}$$

Thus for the operator x^2 ,

$$\begin{aligned}
\langle x^2 \rangle &= \langle \Psi | x^2 | \Psi \rangle \\
&= \int_{-\infty}^{+\infty} dx \Psi^*(x) x^2 \Psi(x) \\
&= \left(\frac{Re(a)}{\pi} \right)^{\frac{1}{2}} \int_{-\infty}^{+\infty} dx x^2 e^{-Re(a)x^2}
\end{aligned} \tag{8}$$

The Guassian integral $I = \int_{-\infty}^{+\infty} dx \, x^2 \, e^{-ax^2}$ can be solved by the method of differentiation under the integral which is demonstrated here. To begin, first notice that

$$x^2 \, e^{-ax^2} = -\frac{\partial}{\partial a} e^{-ax^2} \quad (9)$$

Thus,

$$\begin{aligned} I &= \int_{-\infty}^{+\infty} dx \, x^2 \, e^{-ax^2} \\ &= - \int_{-\infty}^{+\infty} dx \, \frac{\partial}{\partial a} e^{-ax^2} \\ &= -\frac{\partial}{\partial a} \int_{-\infty}^{+\infty} dx \, e^{-ax^2} \\ &= -\frac{\partial}{\partial a} \sqrt{\frac{\pi}{a}} \\ &= \frac{1}{2a} \sqrt{\frac{\pi}{a}} \end{aligned} \quad (10)$$

Where the third equality follows from a being a constant with respect to x , and where the fourth equality comes from the results of problem 1 (4). Picking up from the last line of (8),

$$\begin{aligned} \langle x^2 \rangle &= \left(\frac{Re(a)}{\pi} \right)^{\frac{1}{2}} \int_{-\infty}^{+\infty} dx \, x^2 \, e^{-Re(a)x^2} \\ &= \left(\frac{Re(a)}{\pi} \right)^{\frac{1}{2}} \frac{1}{2Re(a)} \left(\frac{\pi}{Re(a)} \right)^{\frac{1}{2}} \end{aligned} \quad (11)$$

$$\boxed{\langle x^2 \rangle = \frac{1}{2Re(a)}}$$

It's worth verifying the dimensions of this solution. Note that in order for the argument of the *exp* function in Ψ to be dimensionless, a must have dimensions of L^{-2} . So our solution has the correct dimensions, L^2 .