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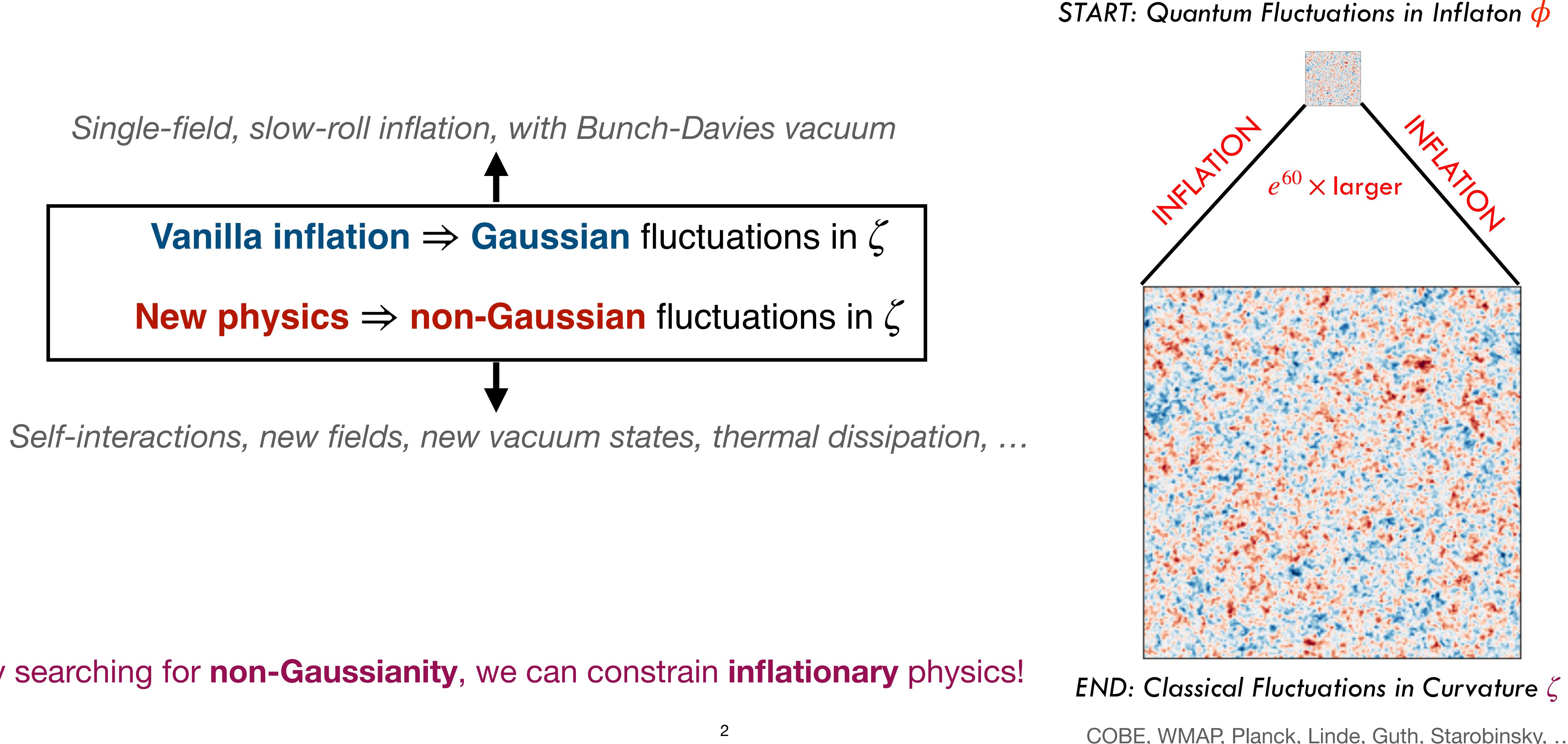
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# *Colliders in the Sky*

Constraining Primordial Non-Gaussianity  
with CMB and LSS Observations

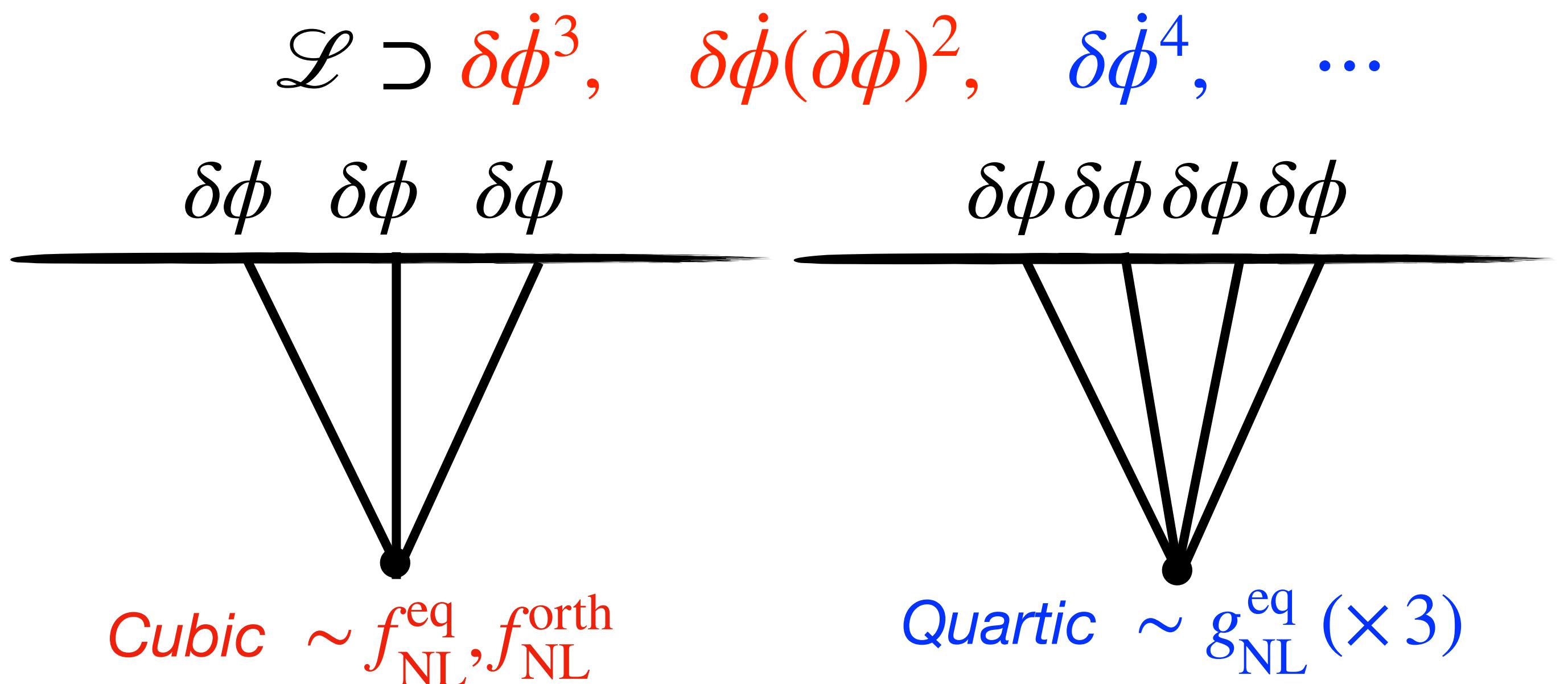
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# Primordial non-Gaussianity



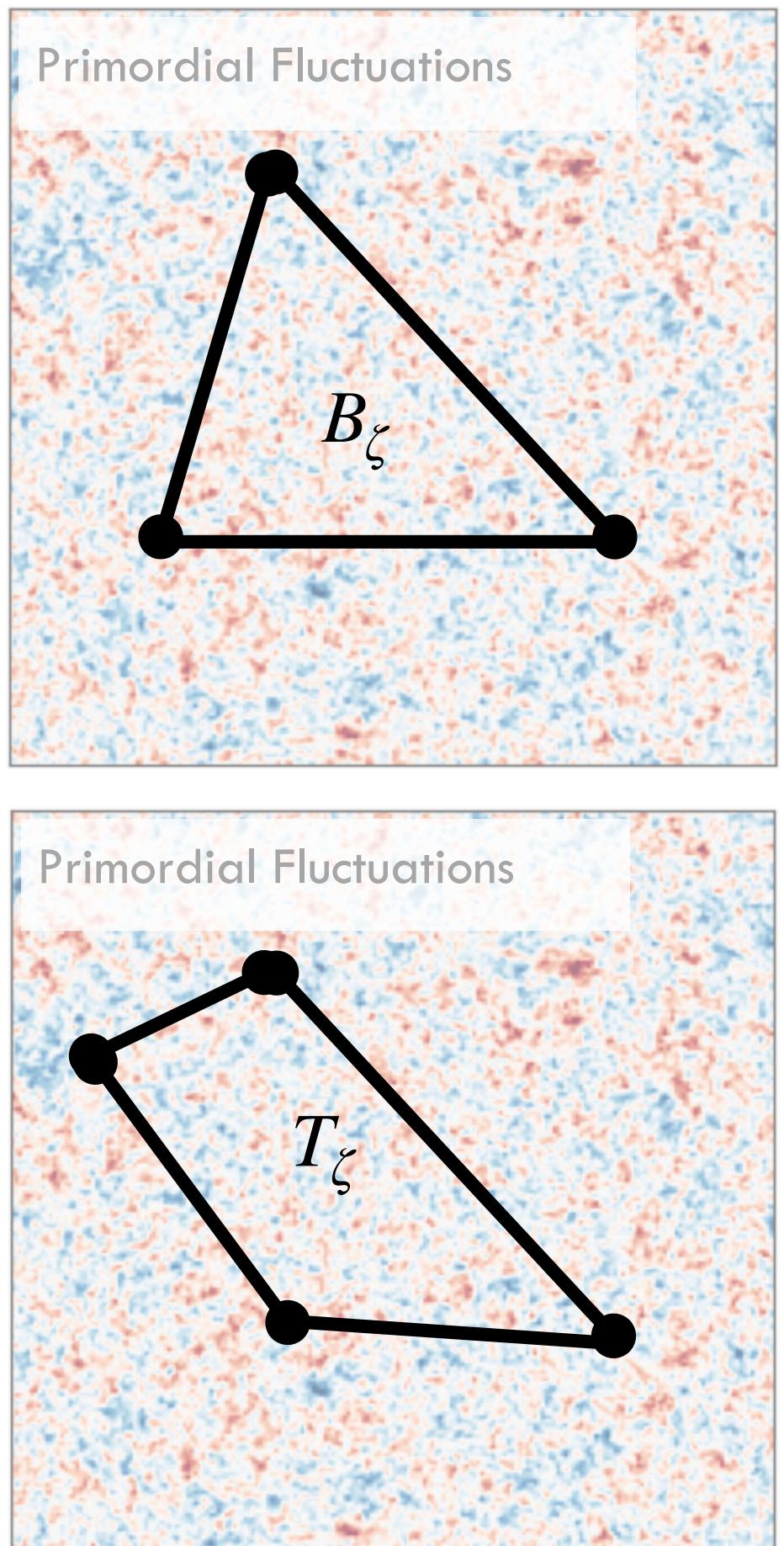
# Self-Interactions

- Many models of inflation feature **self-interactions**:



- These lead to **three-** and **four-point** functions at the end of inflation
- The **shape** encodes the **interaction vertex**, the **amplitude** encodes the microphysics

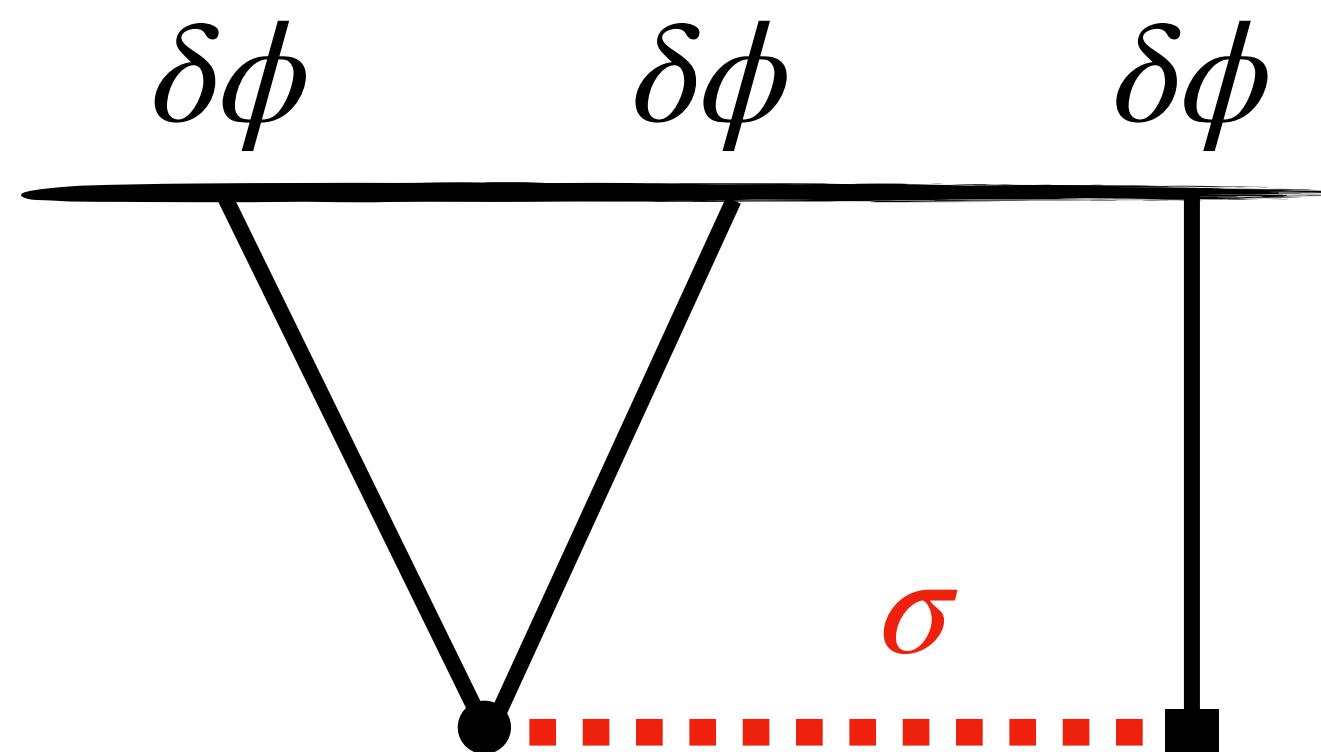
e.g.  $\langle \zeta(\mathbf{k}_1)\zeta(\mathbf{k}_2)\zeta(\mathbf{k}_3) \rangle \sim f_{\text{NL}}^{\text{eq}} \times \text{shape}$



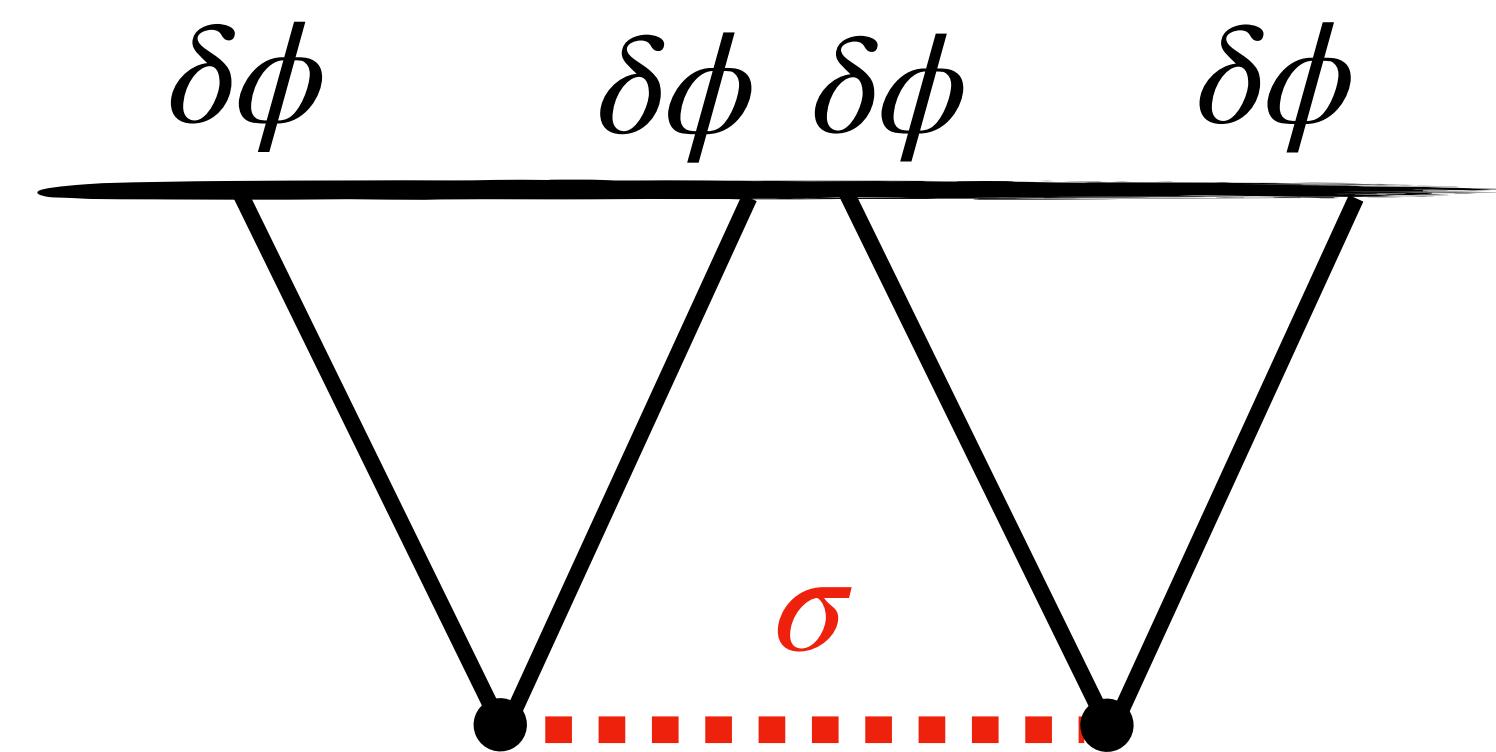
# Self-Interactions

- Other models feature **new particles**,  $\sigma$ :

$$\mathcal{L} \supset \delta\dot{\phi}\sigma, \quad \delta\dot{\phi}^2\sigma, \quad \delta\dot{\phi}^3\sigma, \dots$$



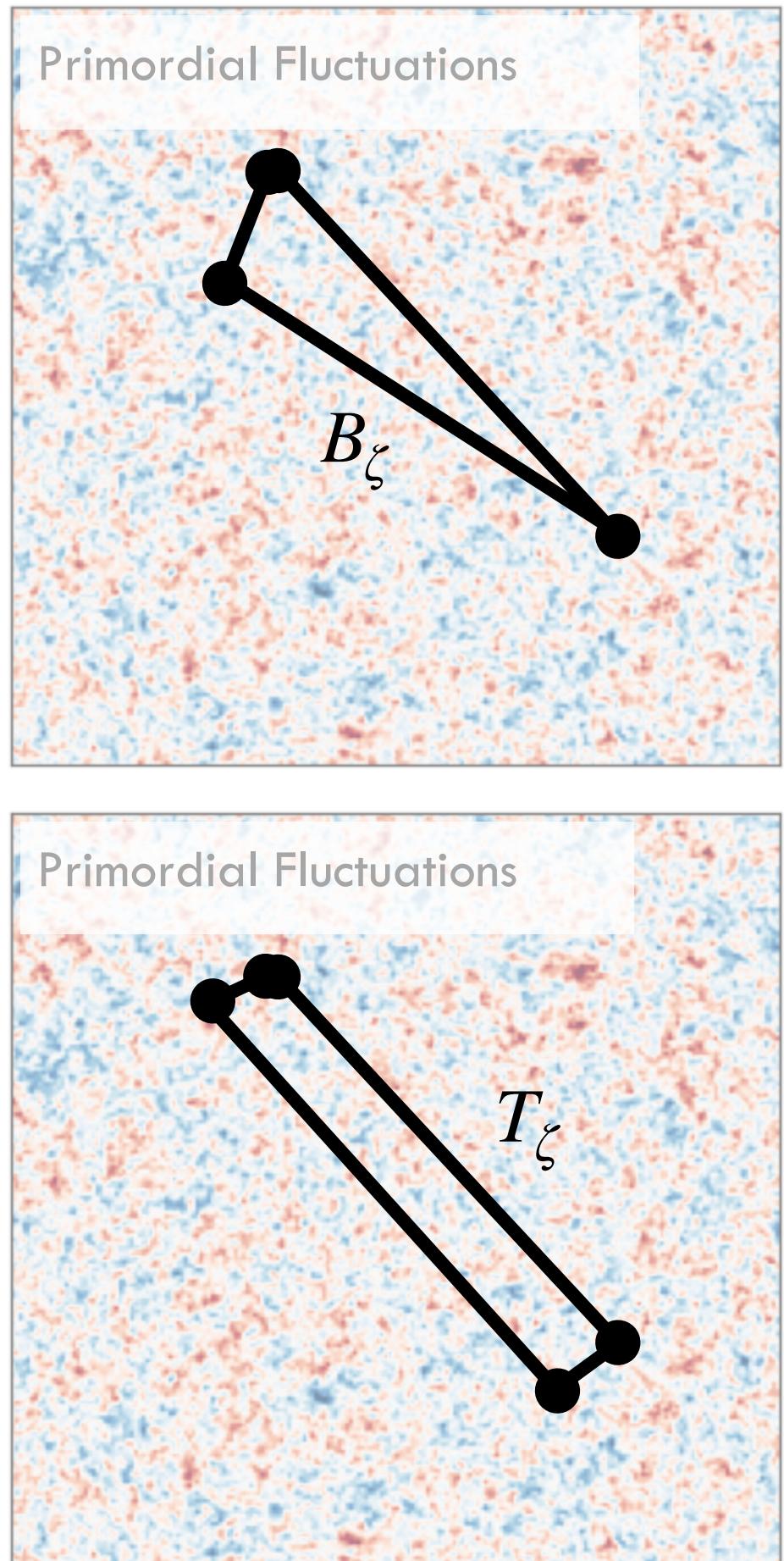
$$\text{Linear-Quadratic} \sim f_{\text{NL}}^{\text{loc}}$$



$$\text{Quadratic}^2 \sim \tau_{\text{NL}}^{\text{loc}} \quad (+ \text{linear-cubic}, \sim g_{\text{NL}}^{\text{loc}})$$

- These lead to **three-** and **four-point** functions at the end of inflation
- The **shape** encodes the **interaction vertex**, the **amplitude** encodes the microphysics

$$\text{e.g. } \langle \zeta(\mathbf{k}_1)\zeta(\mathbf{k}_2)\zeta(\mathbf{k}_3) \rangle \sim f_{\text{NL}}^{\text{loc}} \times \text{shape}$$



# The Cosmological Collider

- The **three-** and **four-point** functions track the **exchange** of a particle  $\sigma_{\mu_1 \dots \mu_s}$  of mass  $m_\sigma \sim H$  and spin  $s = 0, 1, 2, \dots$
- In the **collapsed limit** (low exchange momentum), the inflationary signatures are set by **symmetry** and depend **only** on the mass  $m_\sigma$ , the spin,  $s$ , and the speed  $c_\sigma$

$$\langle \zeta^4 \rangle \sim \tau_{\text{NL}} \times \frac{1}{k_1^3 k_3^3 k_{12}^3} \left[ \left( \frac{k_{12}}{k_1 k_3} \right)^{3/2+i\mu_s} + \left( \frac{k_{12}}{k_1 k_3} \right)^{3/2-i\mu_s} \right] \mathcal{L}_s(\hat{\mathbf{k}}_1 \cdot \hat{\mathbf{k}}_3)$$

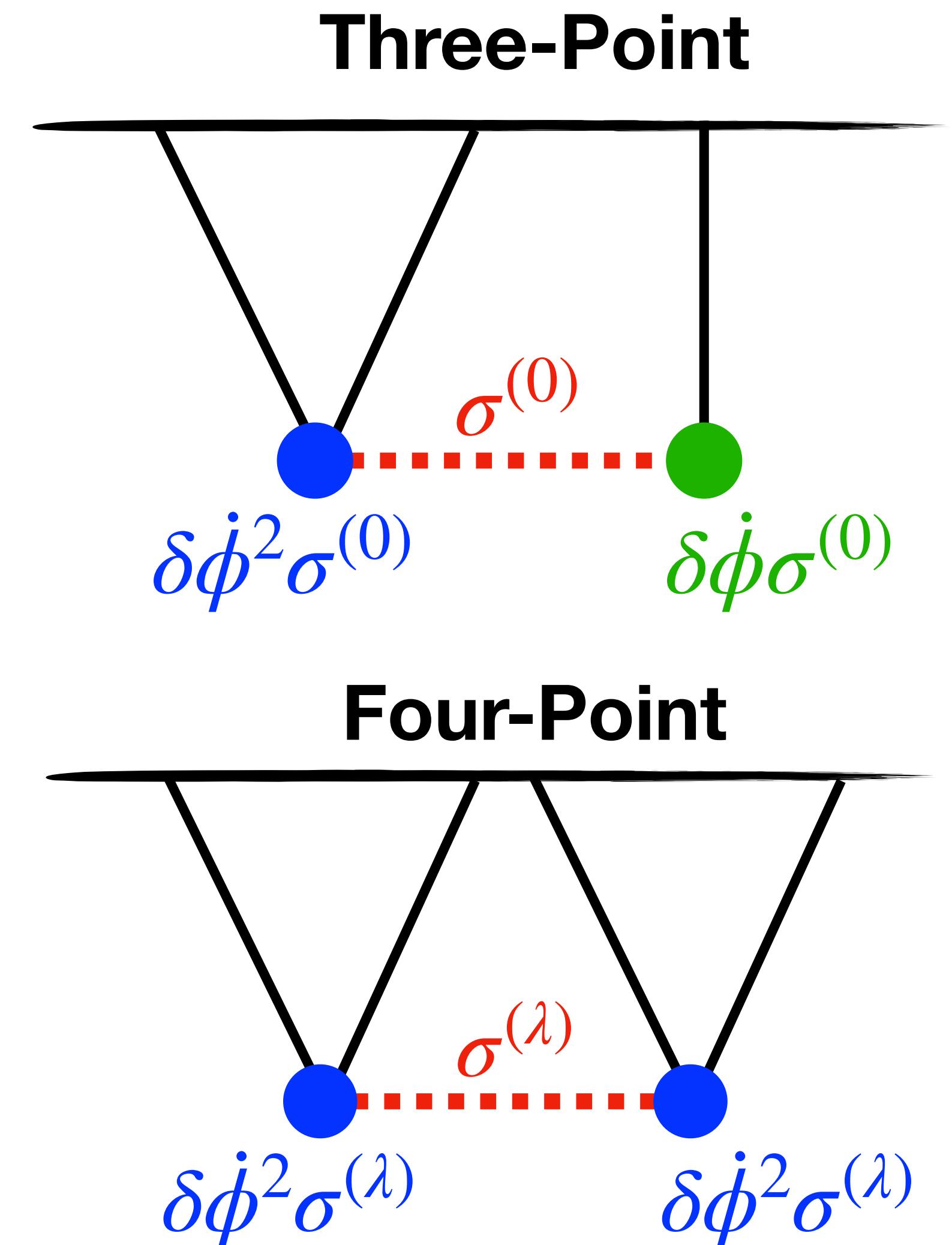
Amplitude  $\sim e^{-m_\sigma/H}$

Shape (mass dependent)

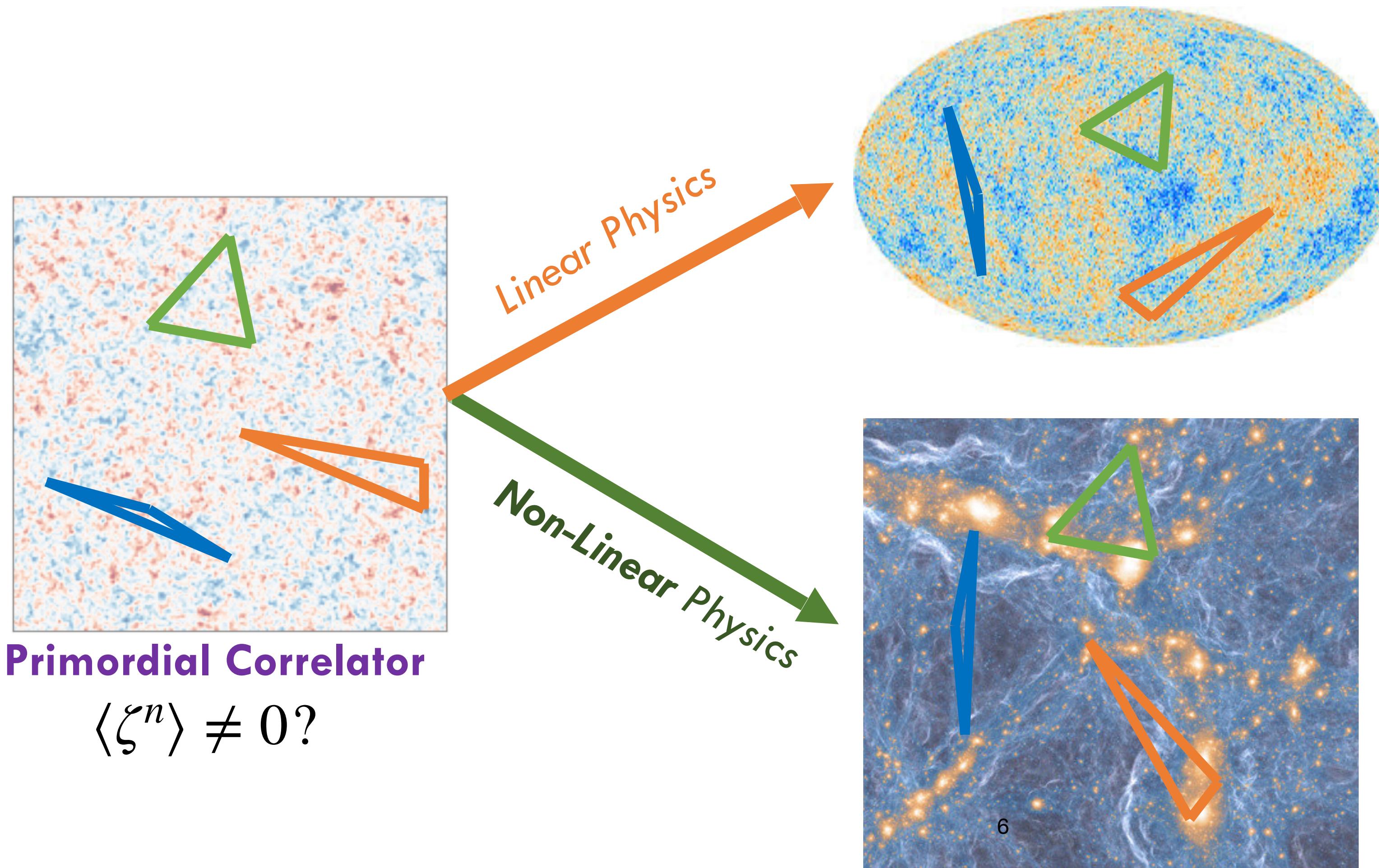
Angle (spin-dependent)

for mass parameter  $\mu_s = \sqrt{m_\sigma^2/H^2 - 9/4}$

- We get **oscillations** for particles with  $m_\sigma \gtrsim H$



# How to Measure Primordial Non-Gaussianity



# How to Measure a Three-Point Function

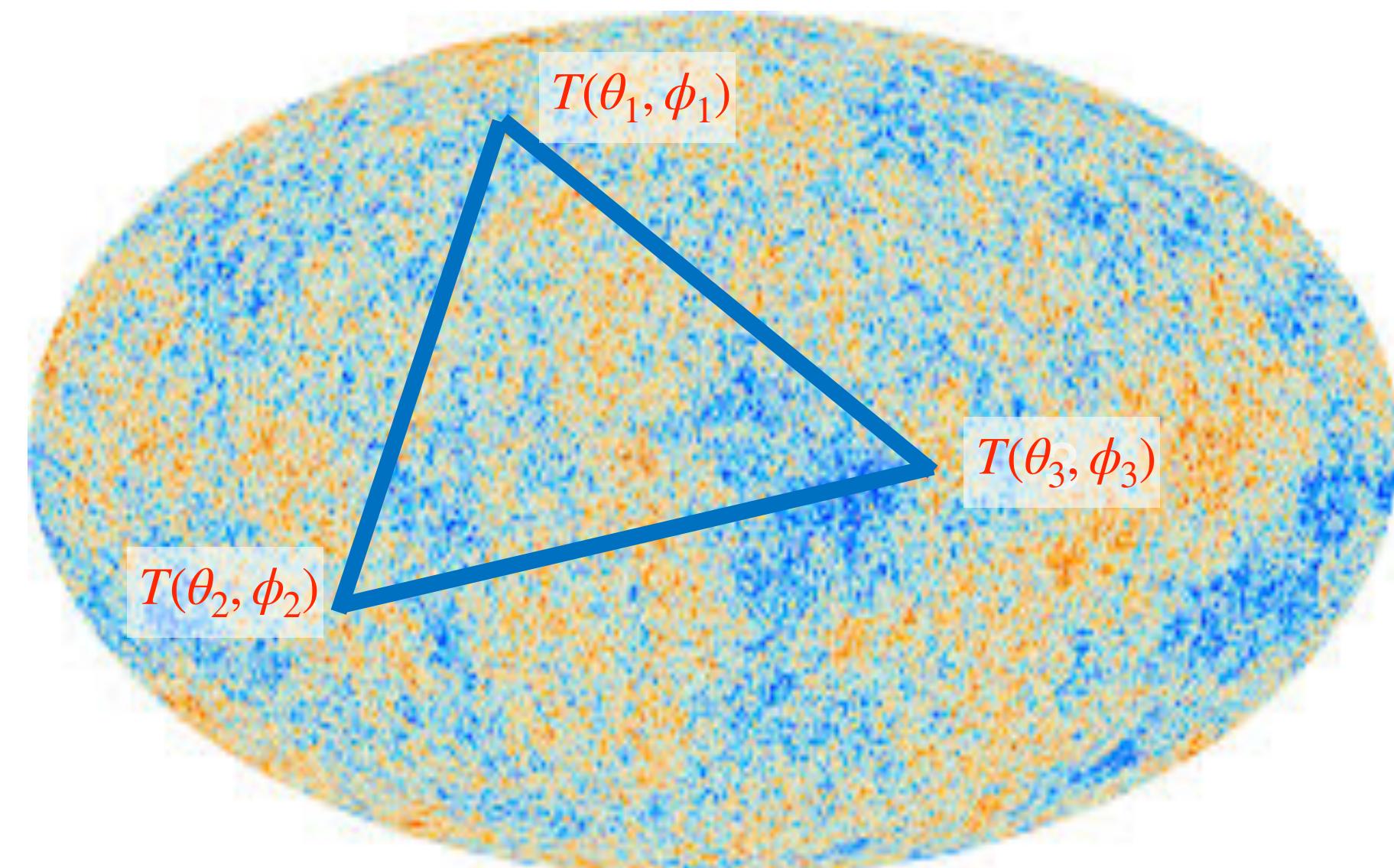
- CMB experiments measure the **temperature** and **polarization** across the whole sky

$$T(\theta, \phi), \quad E(\theta, \phi) \quad \leftrightarrow \quad a_{\ell m}^T, \quad a_{\ell m}^E$$

- Since the physics is **linear** we just need to correlate the CMB at **three** angles

$$\langle T(\theta_1, \phi_1) T(\theta_2, \phi_2) T(\theta_3, \phi_3) \rangle \leftrightarrow \langle a_{\ell_1 m_1}^T a_{\ell_2 m_2}^T a_{\ell_3 m_3}^T \rangle$$

- This is computationally **expensive**:
  - The bispectrum is **3-dimensional** [after symmetries]
  - There's  $N_{\text{pix}}^3 \sim 10^{21}$  combinations of points!



# How to Measure a Three-Point Function

Most CMB analyses use two tricks:

## **1. Compression:**

- We compress all  $10^{21}$  elements into a single number, encoding the amplitude of a specific model, e.g.,  $f_{\text{NL}}^{\text{loc}}$

$$\widehat{f}_{\text{NL}}^{\text{loc}} \sim \sum_{\ell_1 m_1 \ell_2 m_2 \ell_3 m_3} \langle a_{\ell_1 m_1} a_{\ell_2 m_2} a_{\ell_3 m_3} \rangle_{\text{theory}}^\dagger \times (C^{-1} a)_{\ell_1 m_1} (C^{-1} a)_{\ell_2 m_2} (C^{-1} a)_{\ell_3 m_3}$$

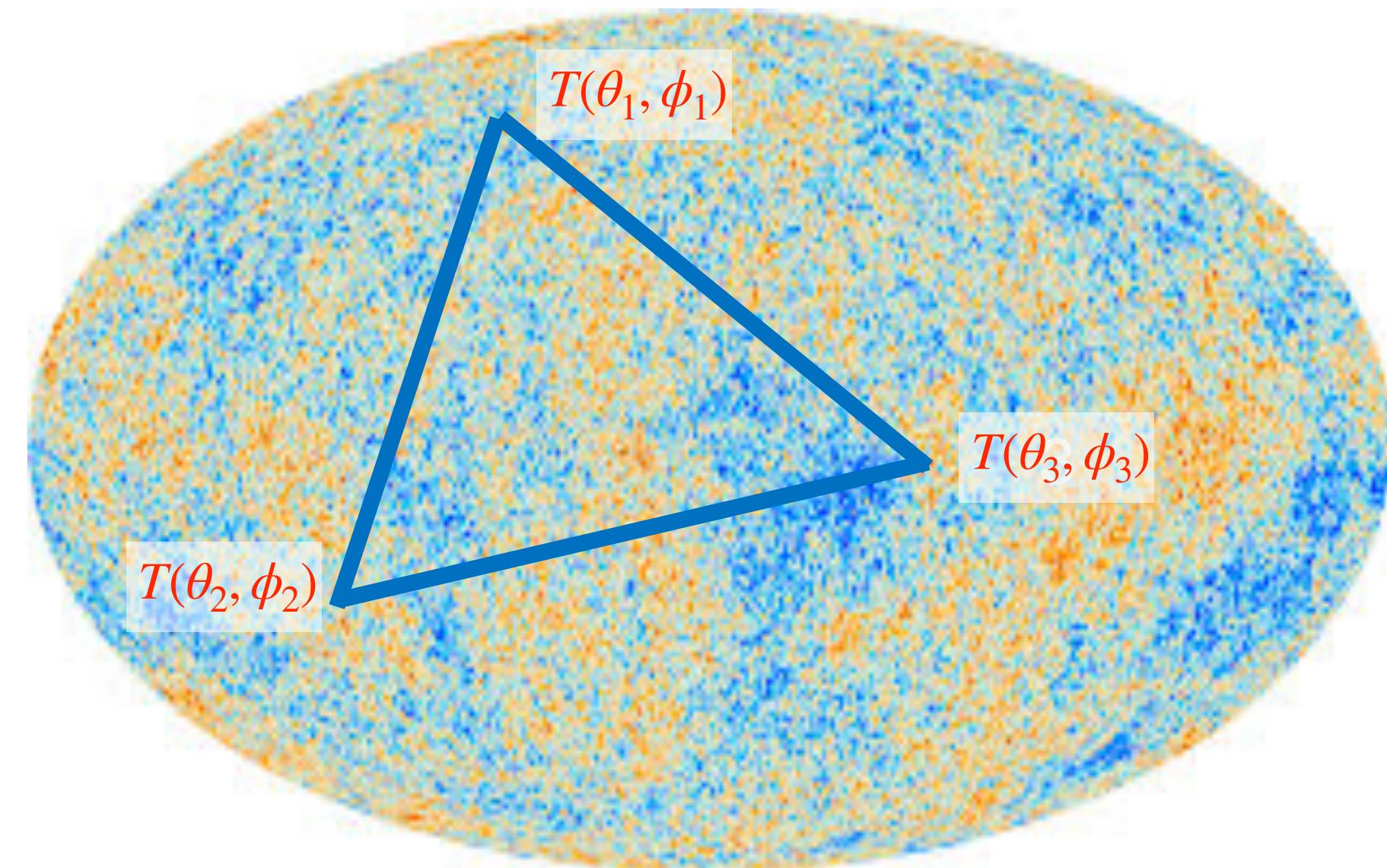
- This is an **optimal estimator** for  $f_{\text{NL}}$ , i.e. it is **lossless**

## 2. Separability:

- If the **theory model** is **separable**, we can rewrite the  $\ell, m$  sum using spherical harmonic transforms!

$$B_\zeta(k_1, k_2, k_3) \sim \sum_n \alpha_n(k_1) \beta_n(k_2) \gamma_n(k_3)$$

- This reduces the complexity from  $\mathcal{O}(N_{\text{pix}}^3)$  to  $\mathcal{O}(N_{\text{pix}} \log N_{\text{pix}})$



**(Note:** binned/modal analyses use a similar trick, but compress to a lower-dimensional basis, rather than a single amplitude)

# CMB Bispectrum Constraints

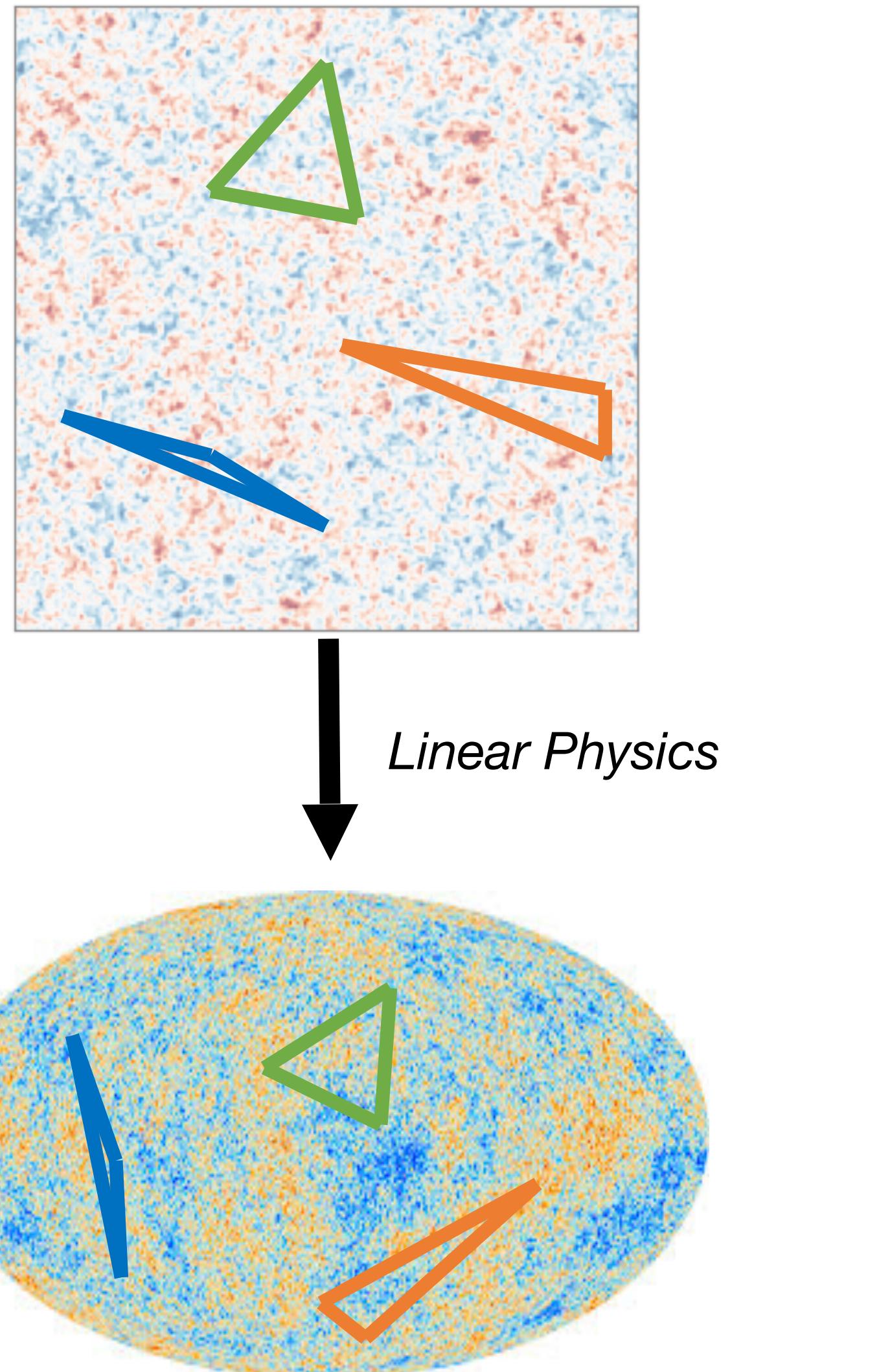
- Planck placed **strong** constraints on scalar **three-point** functions, e.g.,

Planck 2018	Local . . . . .	$-0.9 \pm 5.1$	New light scalars
	Equilateral . . . . .	$-26 \pm 47$	
	Orthogonal . . . . .	$-38 \pm 24$	Self-interactions

- These span **many** phenomenological templates
- Planck uses both **separable** shapes and **modal/binned** approximations
- Recent work has also constrained **tensor** three-point functions, e.g.,  $\langle \zeta \zeta h \rangle$  and **cosmological collider** bispectra

**Conclusion:** Scalar primordial non-Gaussianity is **small**:  $10^{-5} |f_{\text{NL}}| \ll 1$

However, we are still far from the (rough) theory targets:  $\sigma(f_{\text{NL}}) \sim 1$



# What's Next for PNG?

## 1. More models

- Folded NG? Excited states? Slow colliders? Strongly-mixed colliders?

## 2. Higher-orders

- Four-point functions? Five-point functions? Non-perturbative effects?

## 3. Other datasets

- Next-generation CMB? Galaxy clustering? Weak lensing? 21cm observations?

# The CMB Trispectrum

**Very few** previous works have considered **four-point** functions!

- Are they worth investigating?

**Yes!**

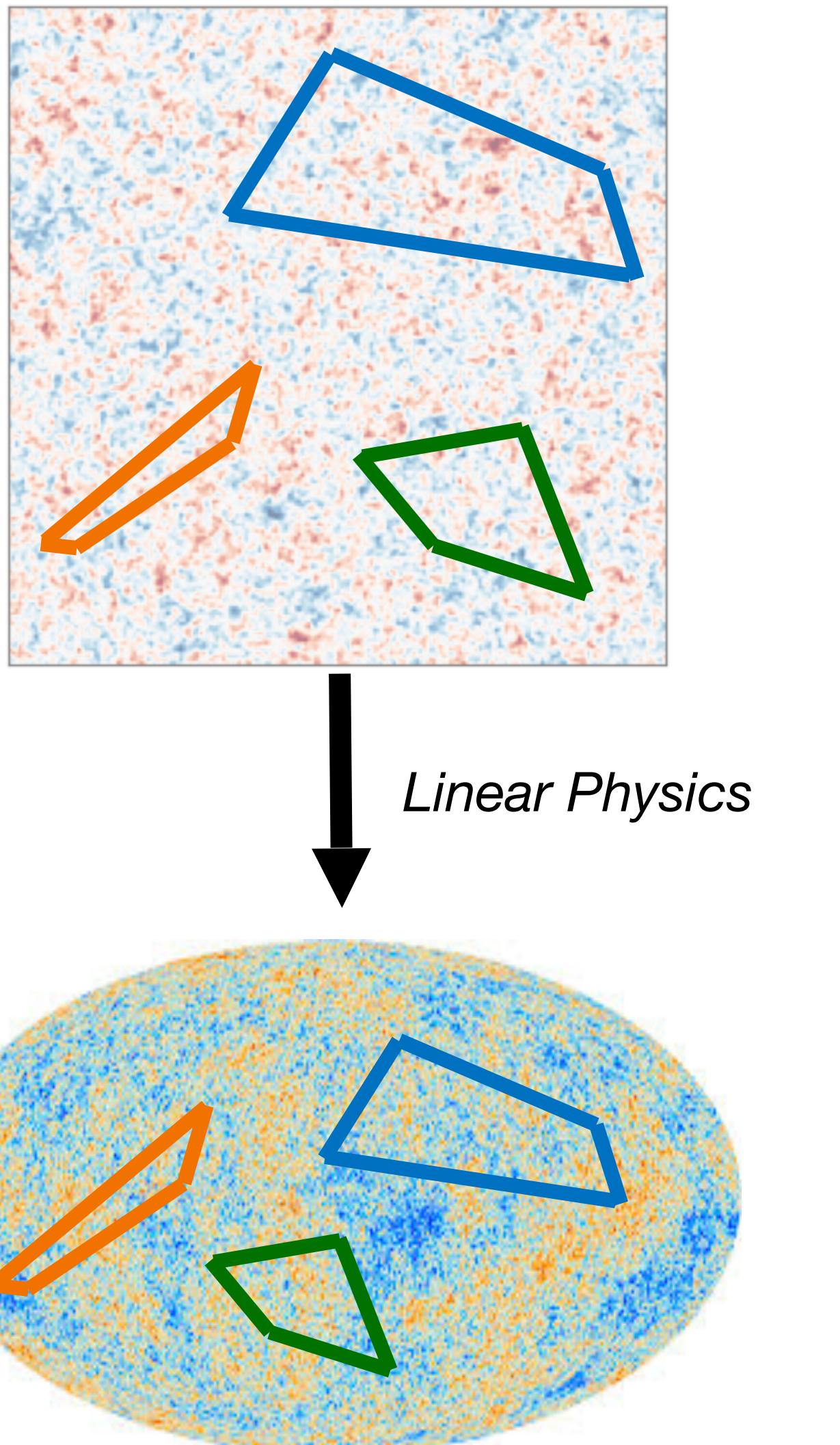
- **Cubic-terms** in the Lagrangian could be protected by **symmetry**

$$\mathcal{L} \sim \frac{1}{2}(\partial\sigma)^2 + \cancel{\dot{\sigma}^3} + \cancel{\dot{\sigma}(\partial\sigma)^2} + \delta\sigma^4 + \dots$$

(for a general light scalar  $\sigma$ , ignoring coupling amplitudes)

Killed by  $\mathbb{Z}_2$  symmetry ( $\sigma \rightarrow -\sigma$ ), or some supersymmetries

- Four-point functions can reveal **hidden particle physics**, e.g, **helicities**
- Collider trispectra *don't* require a **linear mixing** with the inflaton
- Until recently, we *only* had constraints on
  - **Local effects** ( $g_{\text{NL}}^{\text{loc}}, \tau_{\text{NL}}^{\text{loc}}$ )
  - **Self-interactions** (from the EFT of inflation:  $g_{\text{NL}}^{\text{eq}} \times 3$ )



# How to Measure a Four-Point Function

Measuring the CMB trispectrum is a challenge!

- The trispectrum is **five-dimensional** [after symmetries] and depends on  $10^{28}$  sets of points!

$$\langle T(\theta_1, \phi_1)T(\theta_2, \phi_2)T(\theta_3, \phi_3)T(\theta_4, \phi_4) \rangle \leftrightarrow \langle a_{\ell_1 m_1}^T a_{\ell_2 m_2}^T a_{\ell_3 m_3}^T a_{\ell_4 m_4}^T \rangle$$

- We can use **compression** as for the bispectrum:

$$\widehat{g}_{\text{NL}} \sim \sum_{\ell_1 m_1 \ell_2 m_2 \ell_3 m_3 \ell_4 m_4} \langle a_{\ell_1 m_1} a_{\ell_2 m_2} a_{\ell_3 m_3} a_{\ell_4 m_4} \rangle_{\text{theory}}^\dagger \times (C^{-1}a)_{\ell_1 m_1} (C^{-1}a)_{\ell_2 m_2} (C^{-1}a)_{\ell_3 m_3} (C^{-1}a)_{\ell_4 m_4}$$

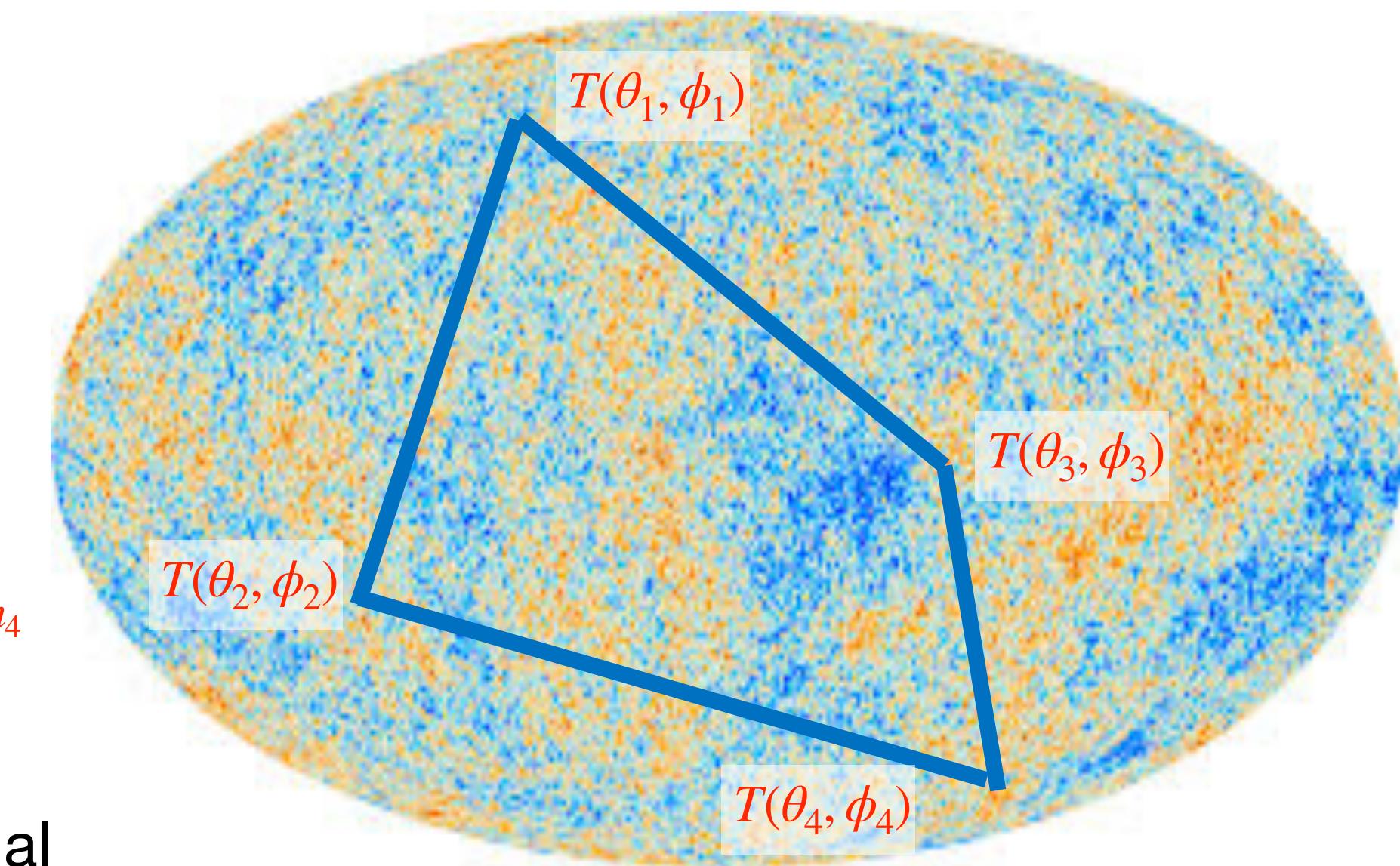
Model

Data

- To **compute** the  $\ell, m$  sum we use a variety of tricks, including low-dimensional integrals, harmonic transforms, and Monte Carlo summation

$$T_\zeta(k_1, k_2, k_3, k_4, s, t, u) \sim F(k_1)G(k_2)H(k_3)I(k_4)J(s^{1/2}) + \dots$$

- If the trispectrum can be (integral-)**factorized**, this reduces the complexity from  $\mathcal{O}(N_{\text{pix}}^4)$  to  $\mathcal{O}(N_{\text{pix}} \log N_{\text{pix}})$



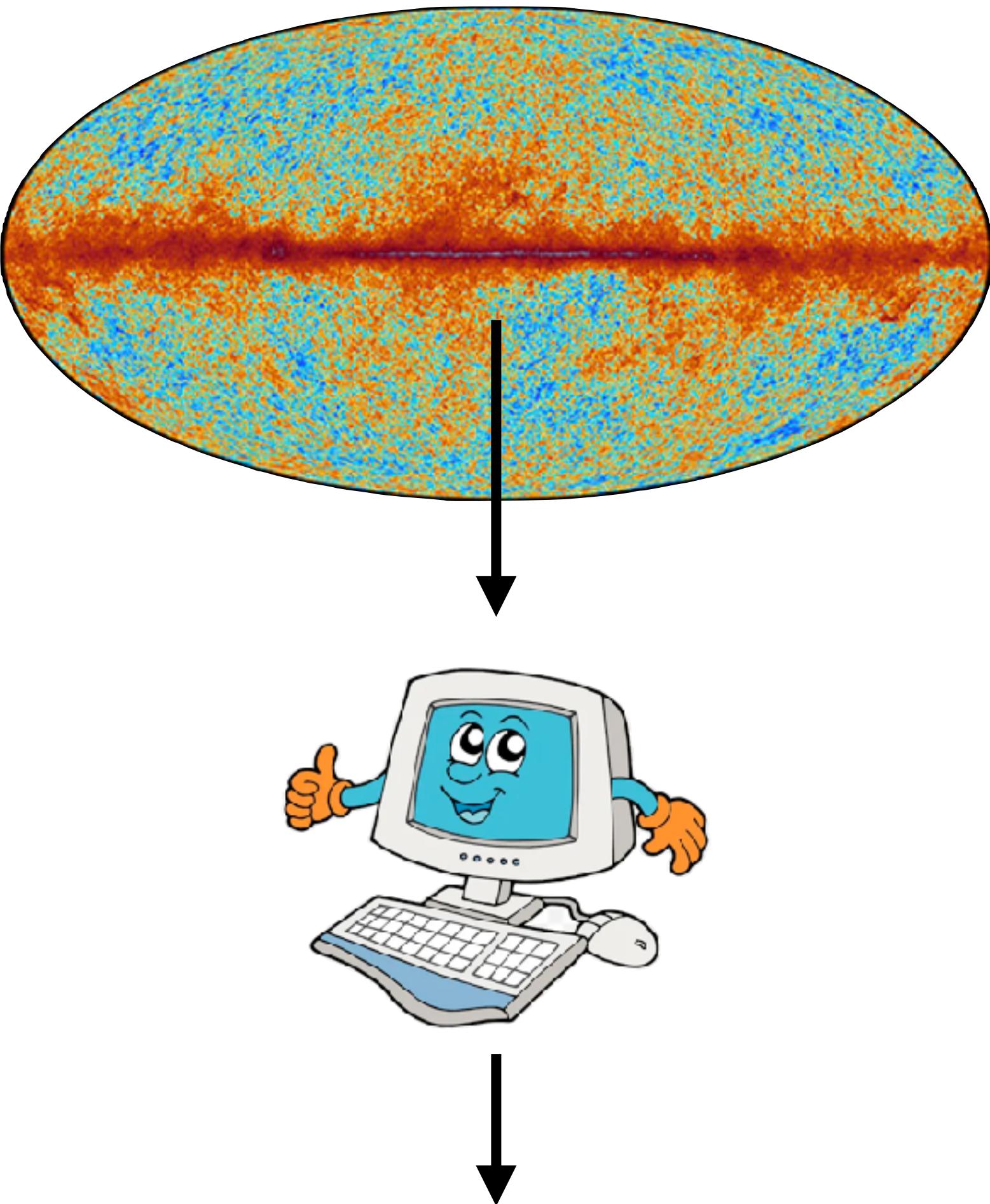
# Optimal Trispectrum Analyses



The result: **fast** estimation of four-point amplitudes!

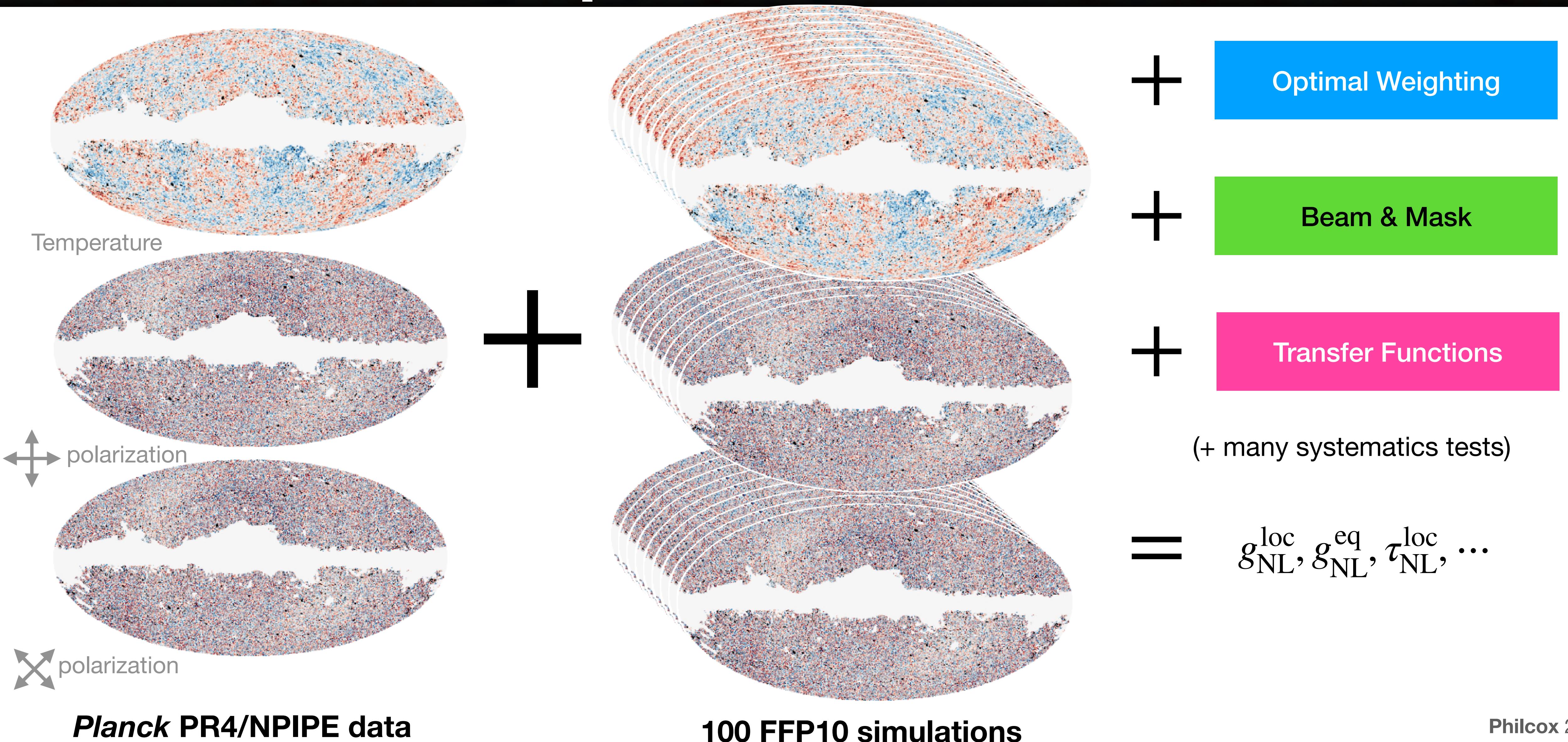
The estimators are

- **Unbiased** (by the mask, geometry, beams, lensing, ...)
- **Efficient** (limited by spherical harmonic transforms)
- **Minimum-Variance** (they saturate the Cramer-Rao bound)
- **Open-Source** (entirely written in Python/Cython)
- **General** (17 classes of **factorizable** model included so far)



Public at <https://github.com/oliverphilcox/PolySpec>

# The Planck Trispectrum



# Detecting Non-Gaussianity?



## What did we try to detect?

1. **Cubic local** shape ( $g_{\text{NL}}^{\text{loc}}$ )
2. **Quadratic<sup>2</sup> local** shape ( $\tau_{\text{NL}}^{\text{loc}}$ )
3. **Constant** shape ( $g_{\text{NL}}^{\text{con}}$ )
4. **Effective Field Theory of Inflation** shapes ( $\times 3$ )
5. **Direction-dependent** shapes
6. **Cosmological Collider** shapes [non-analytic part]
7. Weak **Gravitational Lensing**
8. Unresolved **Point-Sources**
9. ISW-lensing **Trispectra**

**All of these can be integral-factorized!**

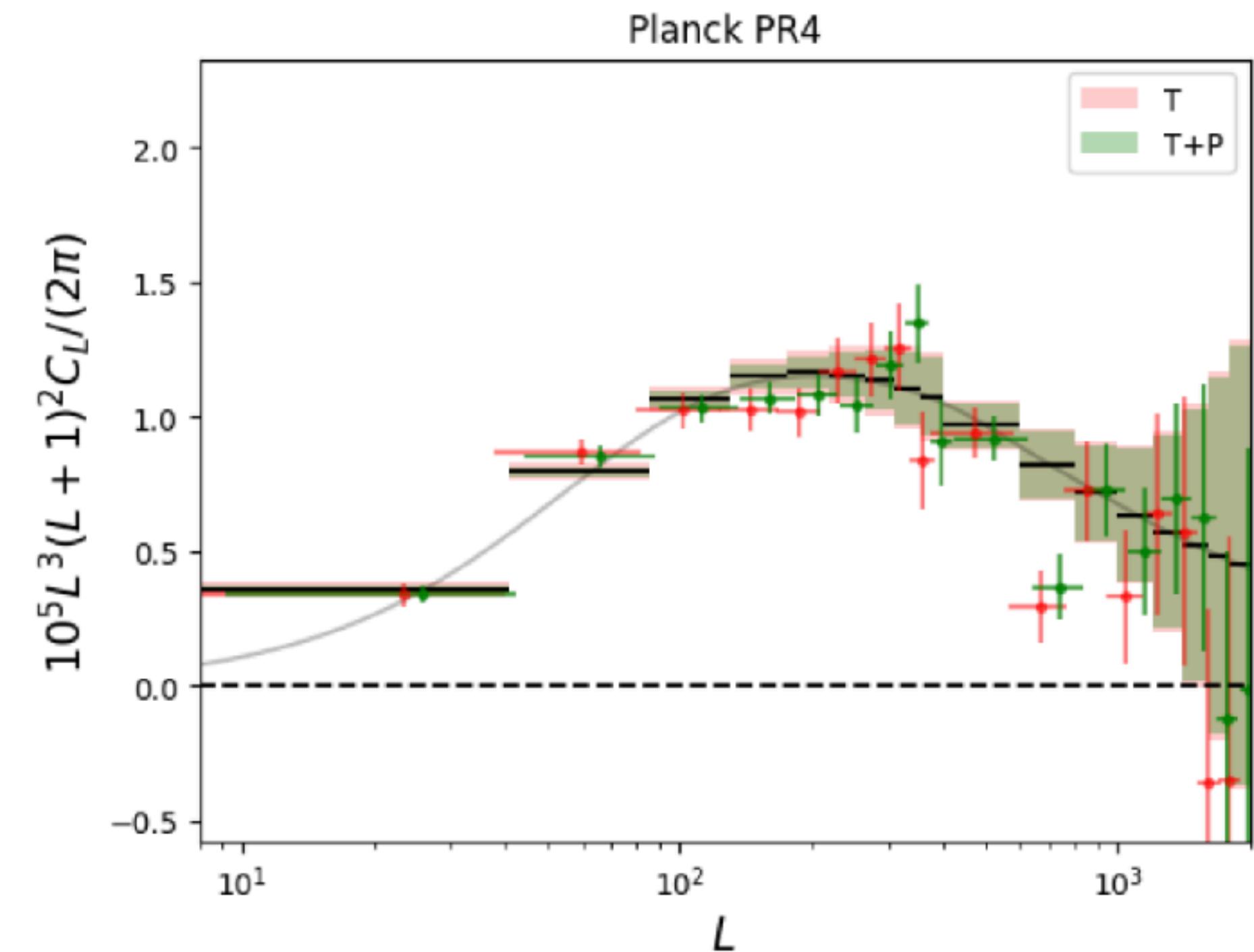
# Detecting Non-Gaussianity?



## What did we try to detect?

- |   |            |
|---|------------|
| 1. <b>Cubic local</b> shape ( $g_{\text{NL}}^{\text{loc}}$ )                    | No         |
| 2. <b>Quadratic<sup>2</sup> local</b> shape ( $\tau_{\text{NL}}^{\text{loc}}$ ) | No         |
| 3. <b>Constant</b> shape ( $g_{\text{NL}}^{\text{con}}$ )                       | No         |
| 4. <b>Effective Field Theory of Inflation</b> shapes ( × 3)                     | No ( × 3)  |
| 5. <b>Direction-dependent</b> shapes  | No ( × 8)  |
| 6. <b>Cosmological Collider</b> shapes [non-analytic part]                      | No ( × 17) |
| 7. Weak <b>Gravitational Lensing</b>  | Yes!!!     |
| 8. Unresolved <b>Point-Sources</b>  | No         |
| 9. ISW-lensing <b>Trispectra</b>  | No         |

## Did we detect it?



**Gravitational Lensing**  
( $43\sigma$ , but not a new detection)

**All of these can be integral-factorized!**

# Equilateral Non-Gaussianity

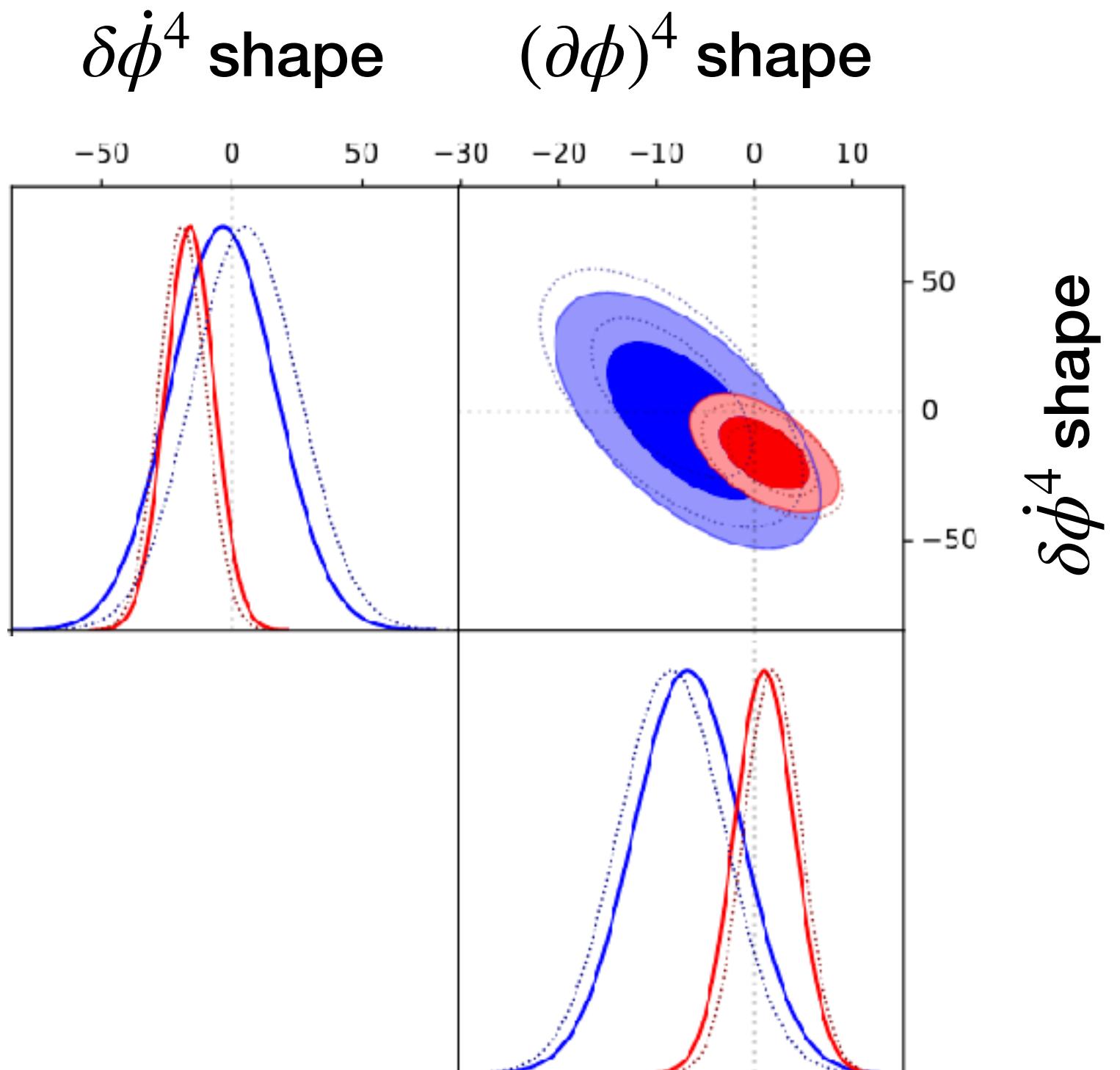
We can constrain cubic *self-interactions* in inflation

- Constrains models such as:
  - **Effective Field Theory** couplings
  - **DBI** inflation (*string theory + small sound-speed*)
  - **Generic** single-field inflation (including *Lorentz Invariant* models)
  - **Ghost** inflation,  $k$ -inflation, and beyond...

**Outcome:** Consistent with zero!

- (50 – 150%) better than any previous constraints!

**T+Pol  $\gg$  T-only**



The third shape –  $\delta\dot{\phi}^2(\partial\phi)^2$  – is very correlated, so we don't plot it [but we don't detect it]

# Cosmological Colliders

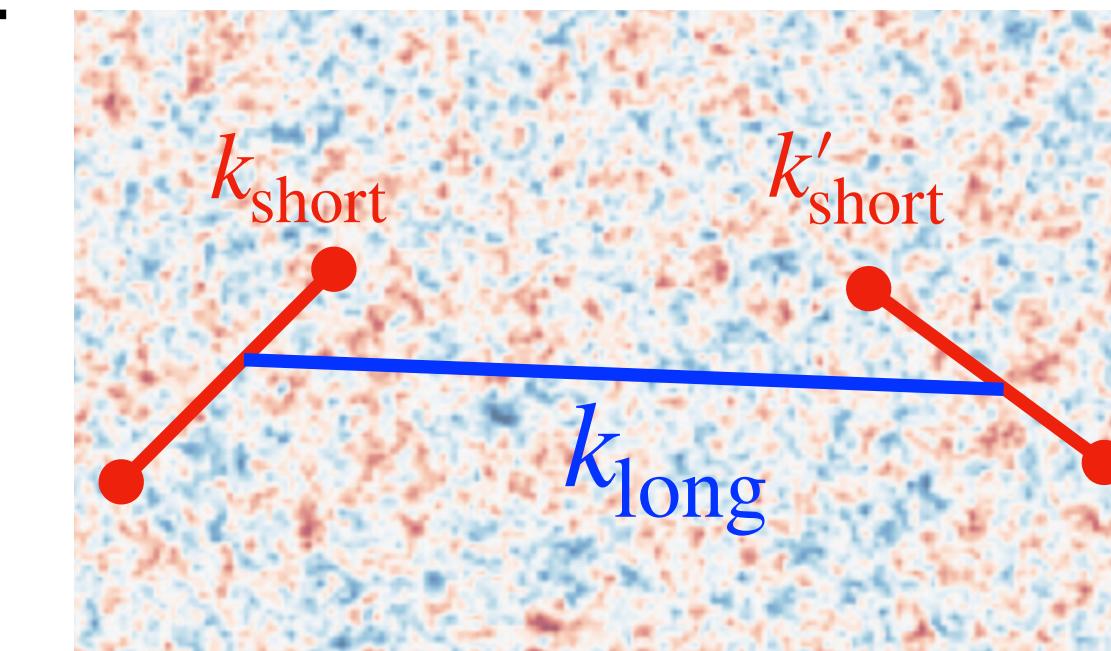
We can search for **massive** and **spinning** particles *(non-analytic piece)*

$$\langle \zeta^4 \rangle \sim P_\zeta(k_{\text{short}})P(k'_{\text{short}})P_\zeta(k_{\text{long}}) \times \left( \frac{k_{\text{long}}^2}{k_{\text{short}} k'_{\text{short}}} \right)^{3/2 \pm i \sqrt{m_\sigma^2/H^2 - 9/4}}$$

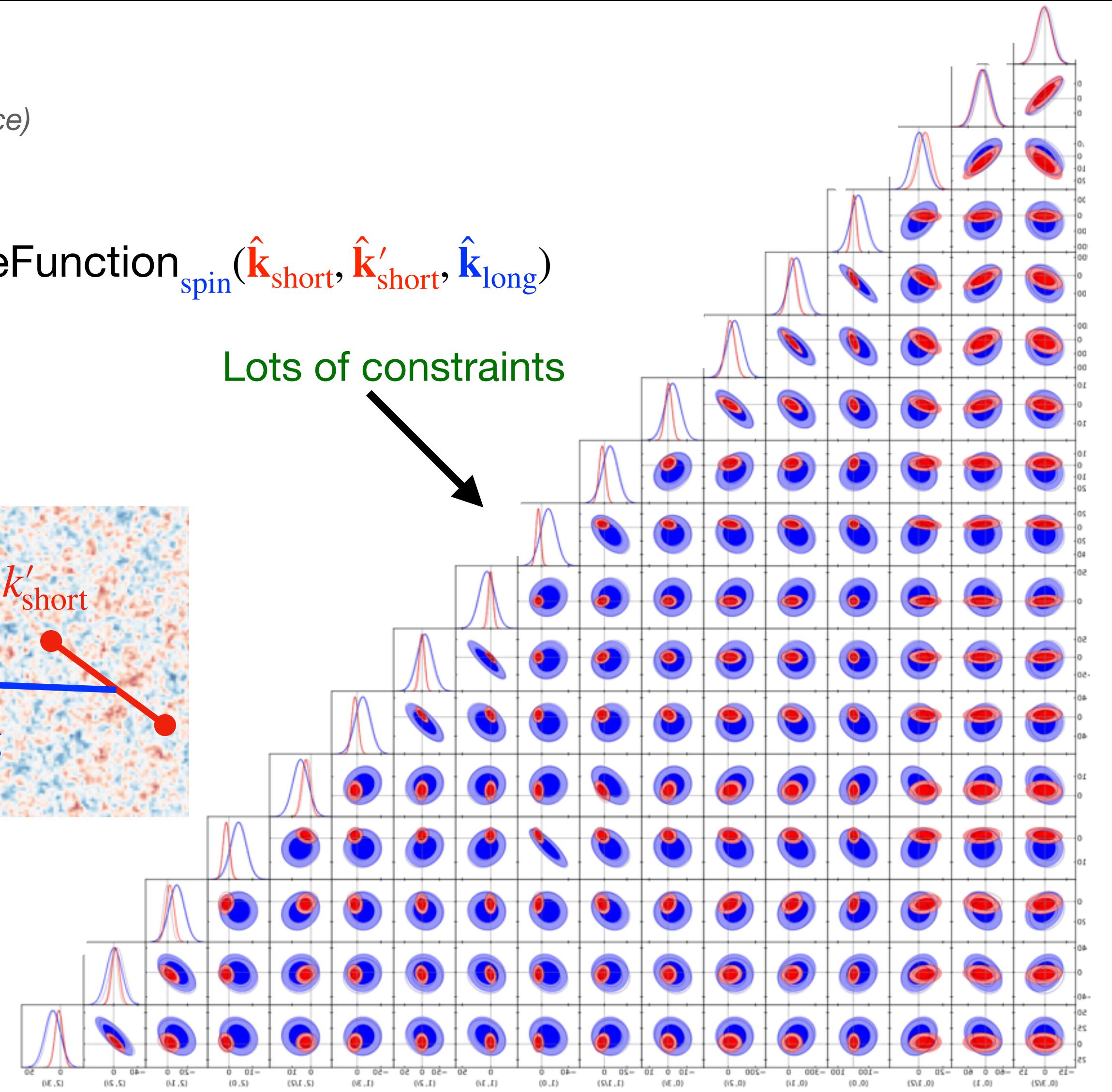
AngleFunction<sub>spin</sub>( $\hat{\mathbf{k}}_{\text{short}}, \hat{\mathbf{k}}'_{\text{short}}, \hat{\mathbf{k}}_{\text{long}}$ )

- Several regimes, including:
  - **Light Fields** (Complementary Series):  
 $m_\sigma \lesssim 3H/2$
  - **Conformally Coupled Fields**:  
 $m_\sigma = \sqrt{2}H$
  - **Heavy Fields** (Principal Series):  
 $m_\sigma \gtrsim 3H/2$

No detections!



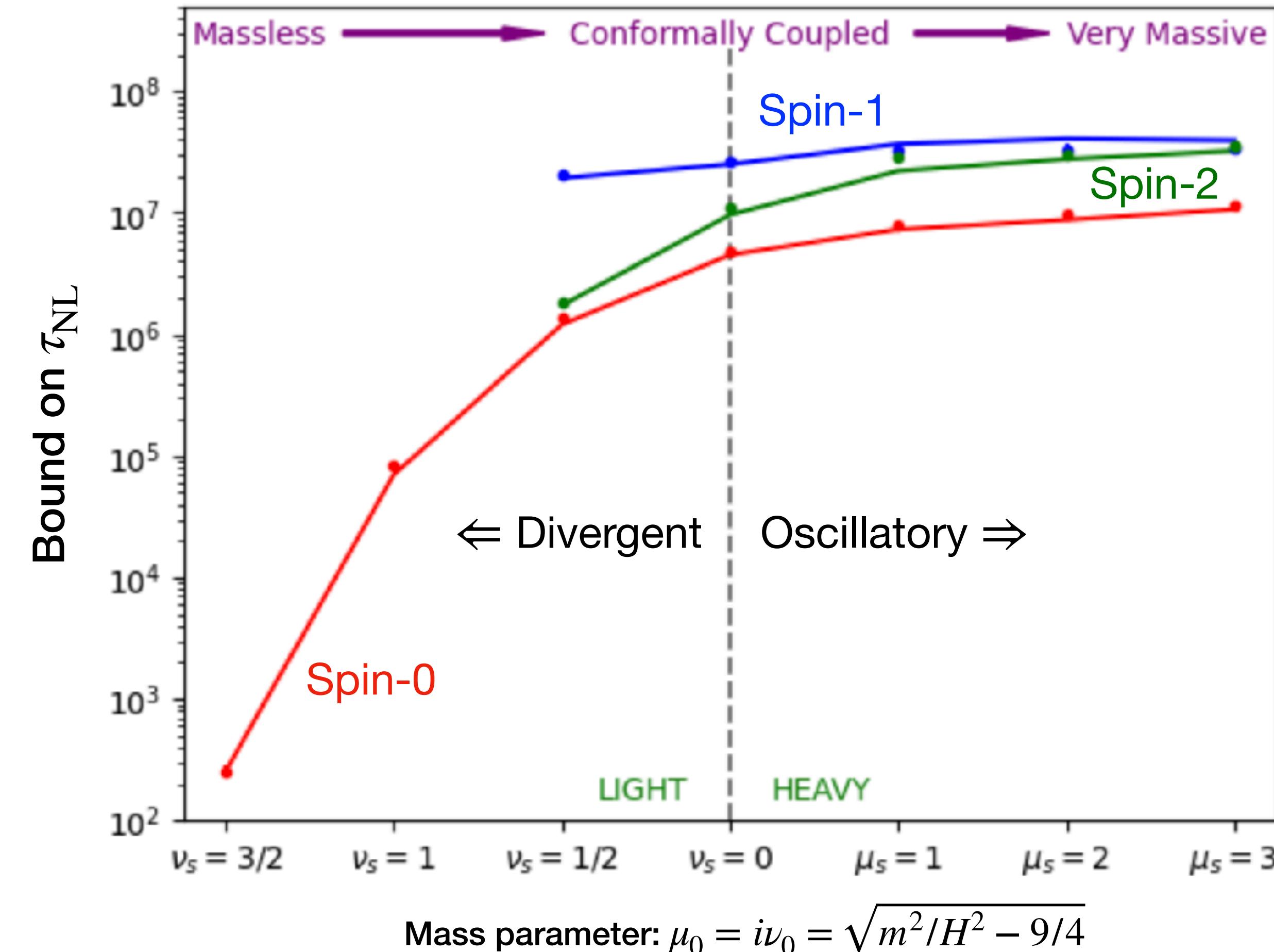
Lots of constraints



# Cosmological Colliders

We can search for **massive** and **spinning** particles *(non-analytic piece)*

- Several regimes, including:
  - **Light Fields** (Complementary Series):  $m_\sigma \lesssim 3H/2$
  - **Conformally Coupled Fields**:  $m_\sigma = 3H/2$
  - **Heavy Fields** (Principal Series):  $m_\sigma \gtrsim 3H/2$
- As expected, **light fields** are easiest to constrain since their trispectrum **diverges**
- Odd-spins are **hard** to constrain due to cancellations!
- **Note:** many of the collider signals are **orthogonal** to the standard templates! [Suman+25, Sohn+24]



# What's Next For the Trispectrum?

There are *many* ways to extend.

## 1. More Data

$$\sigma(\tau_{\text{NL}}^{\text{loc}}) \sim \ell_{\text{max}}^{-2}$$

- ACT, SPT, Simons Observatory, LiteBird, CMB-HD, ... will provide data down to **much** smaller scales!
- **Polarization** will be particularly useful and could benefit from **delensing**

## 2. More Models

- Lighter particles? Heavier particles? Unparticles?
- Tensor non-Gaussianity?
- Collider physics beyond the collapsed limit?
- Thermal baths? Higher-spin particles? Modified sound speeds? Loops? Fermions?
- Scale-dependence? Isocurvature? Primordial magnetic fields?

# Separable Inflationary Correlators

- Efficient **bispectrum** and **trispectrum** analyses require **factorizable** primordial signals.

- Separable  $N$ -point function  $\rightarrow \mathcal{O}(N_{\text{pix}} \log N_{\text{pix}})$  algorithm

- Non-separable  $N$ -point function  $\rightarrow \mathcal{O}(N_{\text{pix}}^N)$  algorithm

$$B_\zeta(k_1, k_2, k_3) \sim F(k_1)G(k_2)H(k_3)$$

$$T_\zeta(k_1, k_2, k_3, k_4, s, t, u) \sim F(k_1)G(k_2)H(k_3)I(k_4)J(s^{1/2}) + \dots$$

- Many models of interest are **not separable**

- Some require complex oscillatory integrals (via *in-in*)

- Others cannot be expressed analytically (e.g., *numerical methods*)

- To analyze these models, we have two options:

1. **Bin the statistic** [lossy, and expensive to compute theory predictions!]

2. Create a **separable approximation** [e.g., modal decompositions]

$$B_\zeta(k_1, k_2, k_3) \sim \frac{k_1 k_2 k_3}{(k_1 + k_2 + k_3)^3}$$

# Separable Inflationary Bispectra

- **Modal approach:** represent the bispectrum as a **sum of polynomials**:

$$(k_1 k_2 k_3)^2 B_\zeta(k_1, k_2, k_3) \sim \sum_{p+q+r=0} \alpha_{pqr} k_1^p k_2^q k_3^r \quad (\text{or Legendre polynomials})$$

- Given a **target** bispectrum, the coefficients  $\alpha_{pqr}$  are computed with **linear algebra**
  - Each term is **factorizable**, so we can build an efficient CMB estimator!
  - However, this basis is **big** ( $\approx 5000$  terms used in the *Planck* collider analysis) and does not represent all shapes of interest.
- 
- **Alternative approach:** learn the basis from the **theory** itself

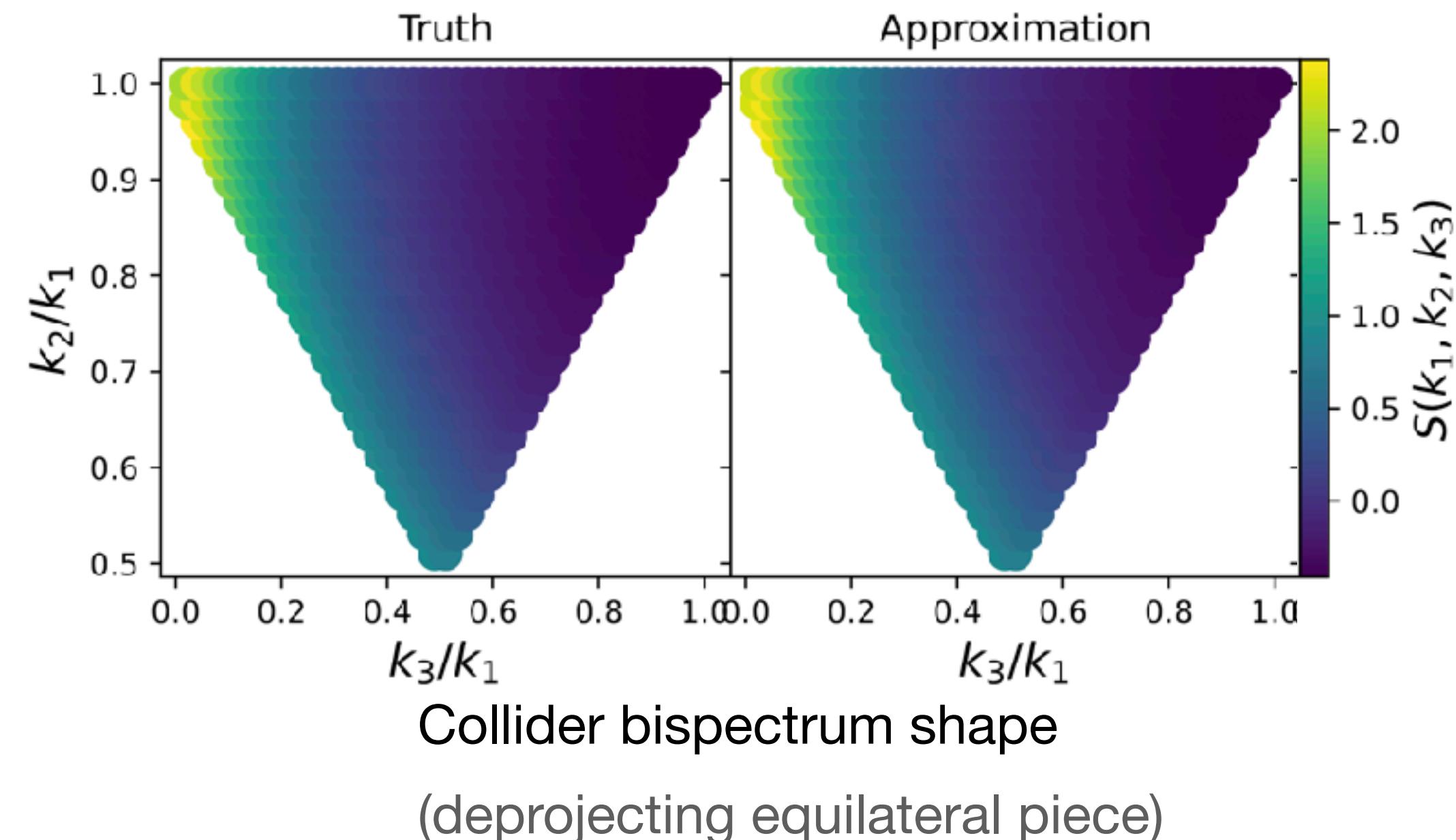
$$(k_1 k_2 k_3)^2 B_\zeta(k_1, k_2, k_3) \sim \sum_n w_n \alpha_n(k_1) \beta_n(k_2) \gamma_n(k_3) + \text{perms.}$$

- Given a **target** bispectrum, the **functions**  $\alpha_n, \beta_n, \gamma_n$  and **weights**  $w_n$  are computed using **machine learning**
- By carefully choosing the loss function, we can **optimize** the decomposition for the task of interest, e.g., *Planck* CMB analysis
- This typically requires **far fewer terms** ( $N \lesssim 3$ ) to compute the bispectra!

# Separable Bispectra in Practice

- This is implemented in the `separable_bk` code, which includes:
  - Simple **neural network** architecture, supplemented with permutation symmetries
  - Training with **stochastic gradient descent**
  - Fast pytorch implementation, giving basis functions in  $\mathcal{O}(\text{minutes})$
- We test `separable_bk` using **numerical** bispectra obtained with the **CosmoFlow** code
  - We use a **strongly-mixed** collider shape that **cannot** be computed analytically
  - With just **three** terms, we find approximations with > 99.9 % accuracy!

$$(k_1 k_2 k_3)^2 B_\zeta(k_1, k_2, k_3) \sim \sum_n w_n \alpha_n(k_1) \beta_n(k_2) \gamma_n(k_3) + \text{perms.}$$

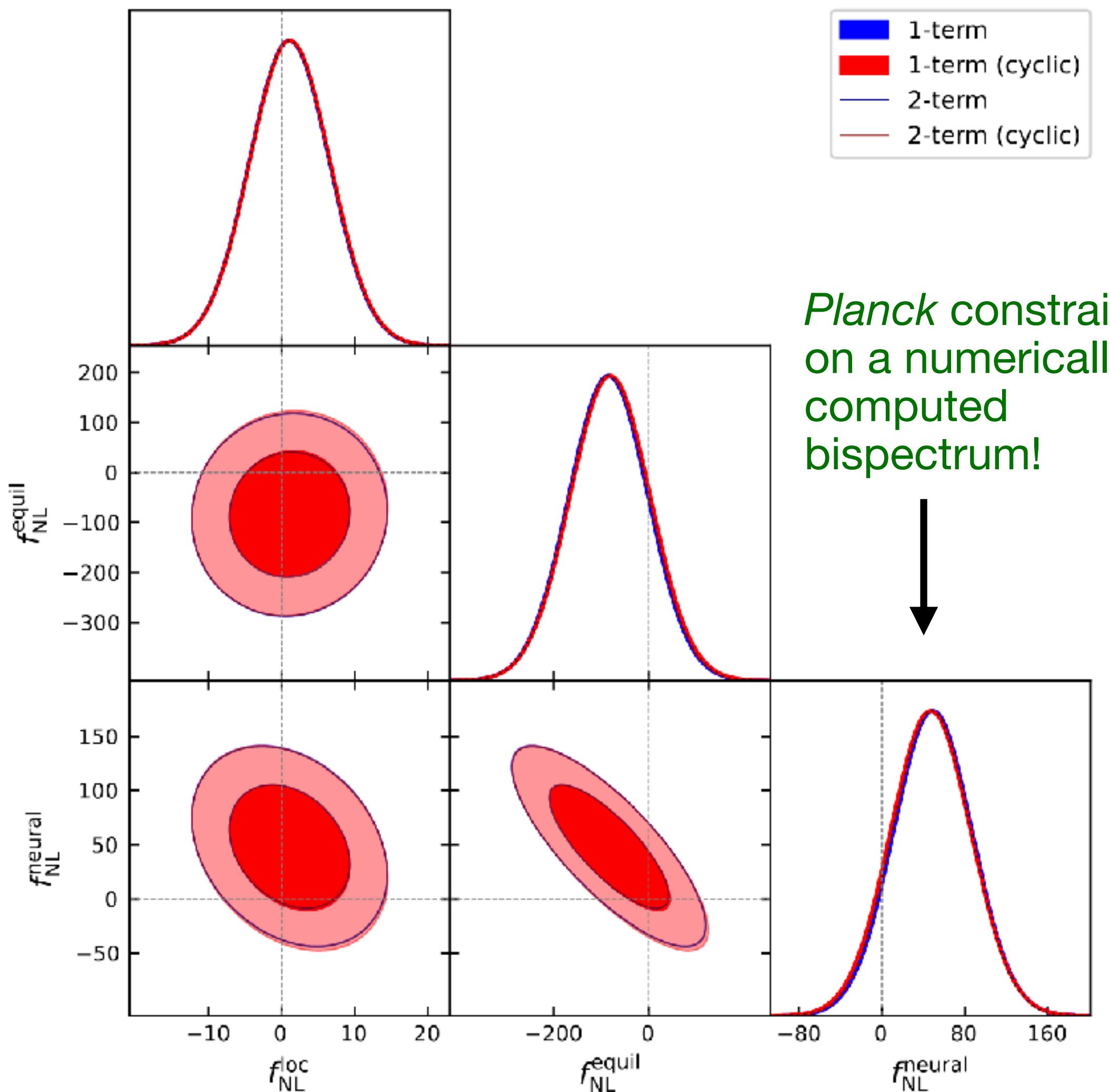


Available at [https://github.com/KunhaoZhong/separable\\_bk](https://github.com/KunhaoZhong/separable_bk)

# Separable Bispectra in Practice

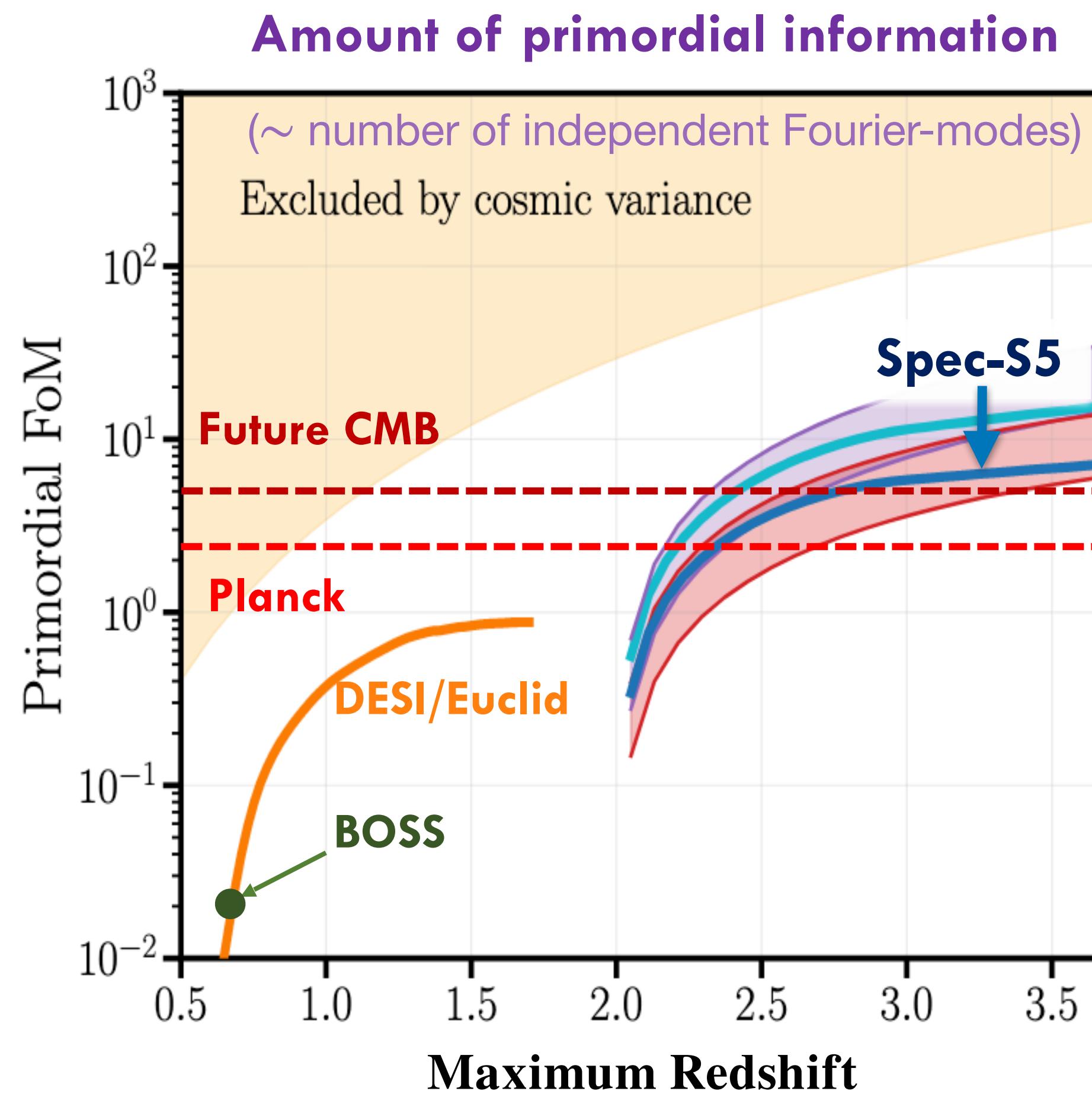
- The output of `separable_bk` is a set of **basis functions** describing a particular primordial bispectrum.
- These can be interfaced with the **PolySpec** code to compute  $f_{NL}$  bounds, e.g., from *Planck*
- Since the number of separable terms is **small**, we can analyze **collider** bispectra in similar computation time to standard shapes, such as local and equilateral!
- This will allow us to constrain **arbitrary** primordial bispectra, including those that can only be computed numerically.
- There is **lots** to explore, e.g., analysis of strongly-mixed colliders, and extension to trispectra

Available at [https://github.com/KunhaoZhong/separable\\_bk](https://github.com/KunhaoZhong/separable_bk)



# The Future of Non-Gaussianity

- Future **CMB** experiments will improve bounds on PNG by  $\lesssim 3 \times$ 
  - This is a **two-dimensional** field
  - We're running out of modes to look at
    - Large-scales are **cosmic-variance-limited**
    - Small-scales are limited by **secondaries** and **Silk damping**
- What about **galaxy surveys**?
  - The data precision is **rapidly increasing**
    - Legacy surveys map **a million** galaxies [BOSS]
    - New surveys map  $\sim 100 \times$  more! [DESI, Euclid, Rubin, Roman, SphereX,...]
  - This is a **three-dimensional** field
    - We aren't limited by **projection effects**
  - There are new observables e.g., galaxy **shapes**, kSZ cross-correlations, ...



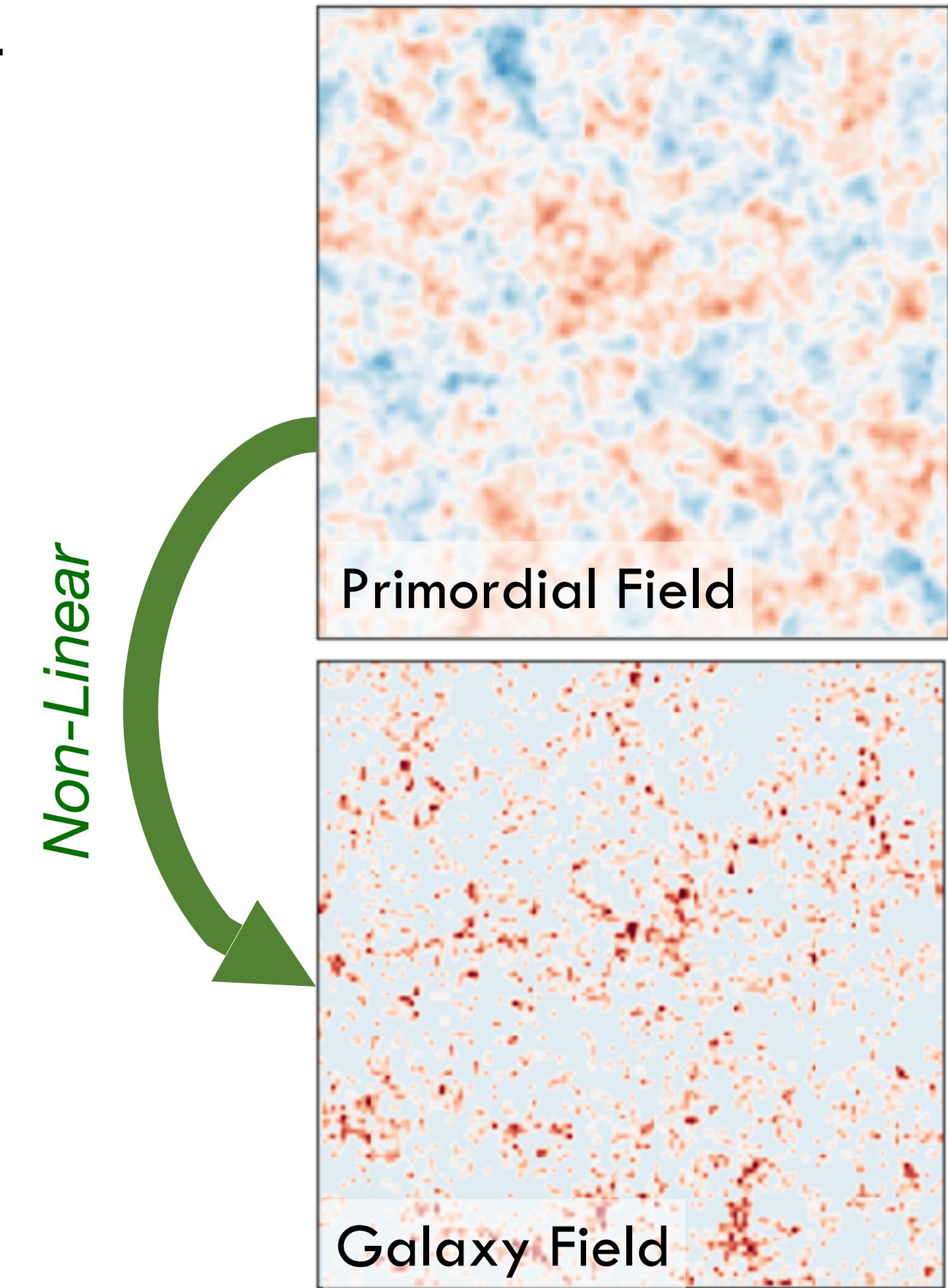
# Inflation from Galaxy Surveys

- Modern galaxy surveys map of the distribution of galaxies in three-dimensions:  $\delta_g(\mathbf{x}, z)$
- This **traces dark matter evolution and the initial conditions**
- To extract **inflationary information**, we need a **joint** model of all effects:

$$\langle \delta_g \delta_g \delta_g \rangle \sim \text{Primordial Physics} + \text{Gravity} + \text{cross-terms}$$

State-of-the-art method:

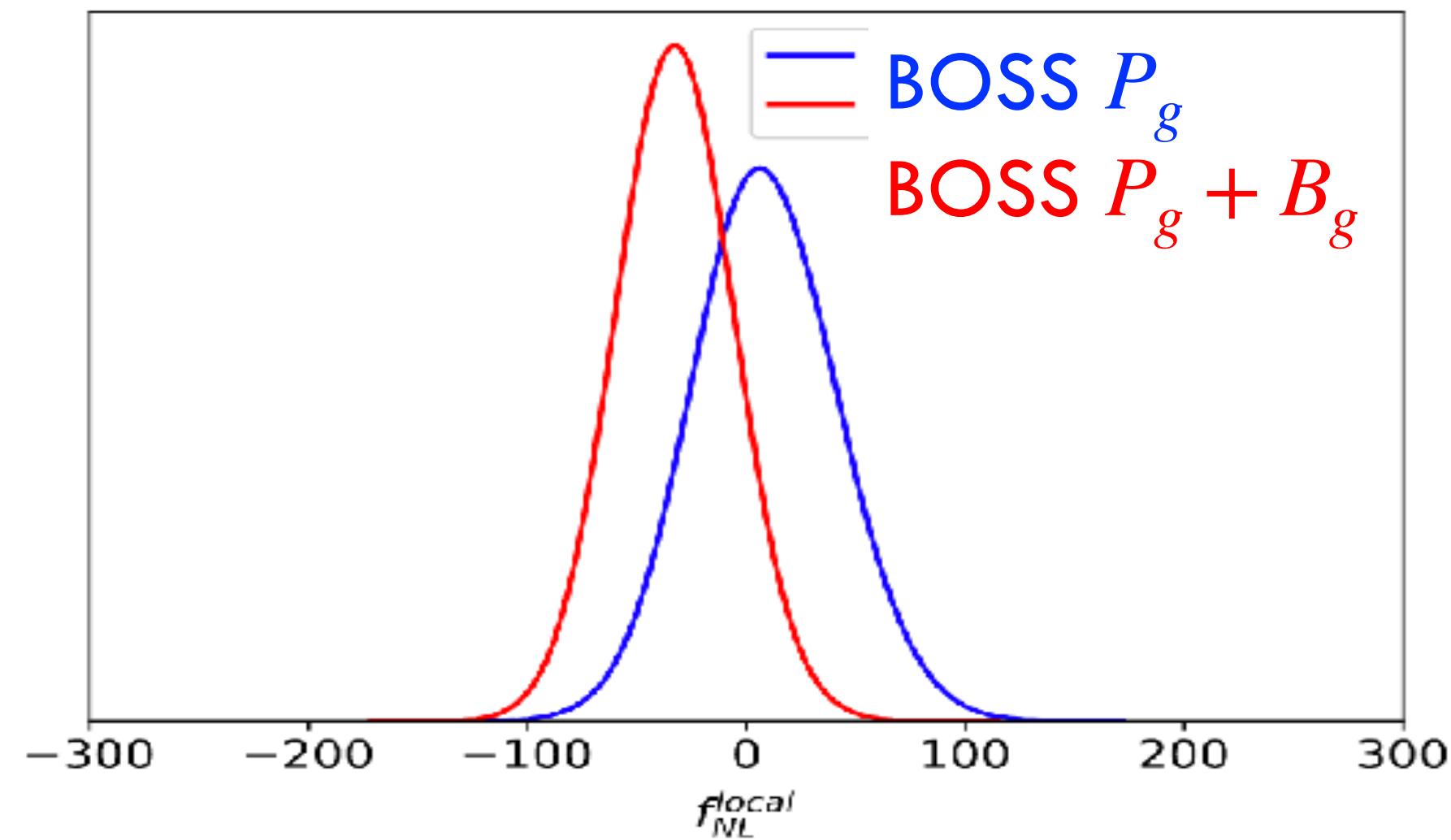
***Effective Field Theory of Large Scale Structure (EFTofLSS)***



# Inflation from Galaxy Surveys

- Recent works have constrained inflationary **bispectra** with **legacy galaxy survey** data (SDSS-BOSS):
  - $f_{NL}^{\text{loc}}$ : **Local** three-point functions from additional **light fields**
  - $f_{NL}^{\text{eq,orth}}$ : **Equilateral** three-point functions from cubic interactions in single-field inflation
  - $f_{NL}^{\text{coll}}(m_\sigma, c_\sigma)$ : **Collider** three-point functions from the exchange of massive scalar fields
- For now, the constraints are **much** worse than the CMB (5 – 20×)
- Much better data is coming soon!***

## *Light Field Constraints*

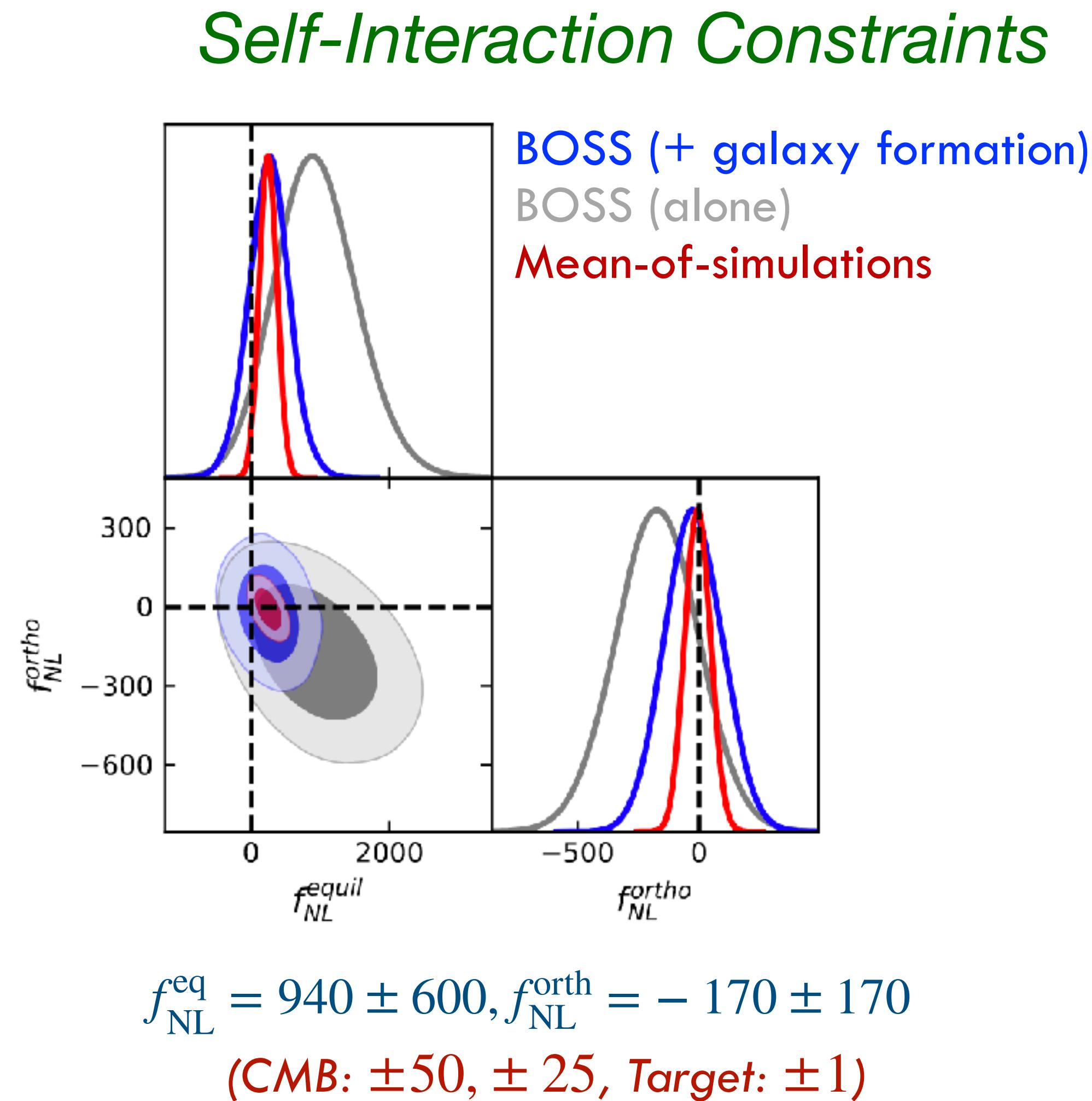


$$f_{NL}^{\text{loc}} = -33 \pm 28 \quad (9 \pm 34 \text{ w/o bispectra})$$

*(CMB:  $\pm 5$ , Target:  $\pm 1$ )*

# Inflation from Galaxy Surveys

- Recent works have constrained inflationary **bispectra** with **legacy galaxy survey** data (SDSS-BOSS):
  - $f_{NL}^{\text{loc}}$ : **Local** three-point functions from additional **light fields**
  - $f_{NL}^{\text{eq,orth}}$ : **Equilateral** three-point functions from cubic interactions in single-field inflation
  - $f_{NL}^{\text{coll}}(m_\sigma, c_\sigma)$ : **Collider** three-point functions from the exchange of massive scalar fields
- For now, the constraints are **much** worse than the CMB (5 – 20×)
- Much better data is coming soon!***



# Inflation from DESI

The first year of DESI data is now public!

- We have developed an **independent** pipeline for analyzing the **power spectrum** and **bispectrum**
- This has been used to constrain:  $\Lambda$ CDM ( $\Omega_m, H_0, \sigma_8$ ), dark energy ( $w_0 w_a$ ), curvature ( $\Omega_k$ ), neutrino masses ( $\sum m_\nu$ )

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New constraints on inflation!

• Multi-field:  $f_{\text{NL}}^{\text{loc}} = -0.1 \pm 7.4$

Compare to official DESI results:

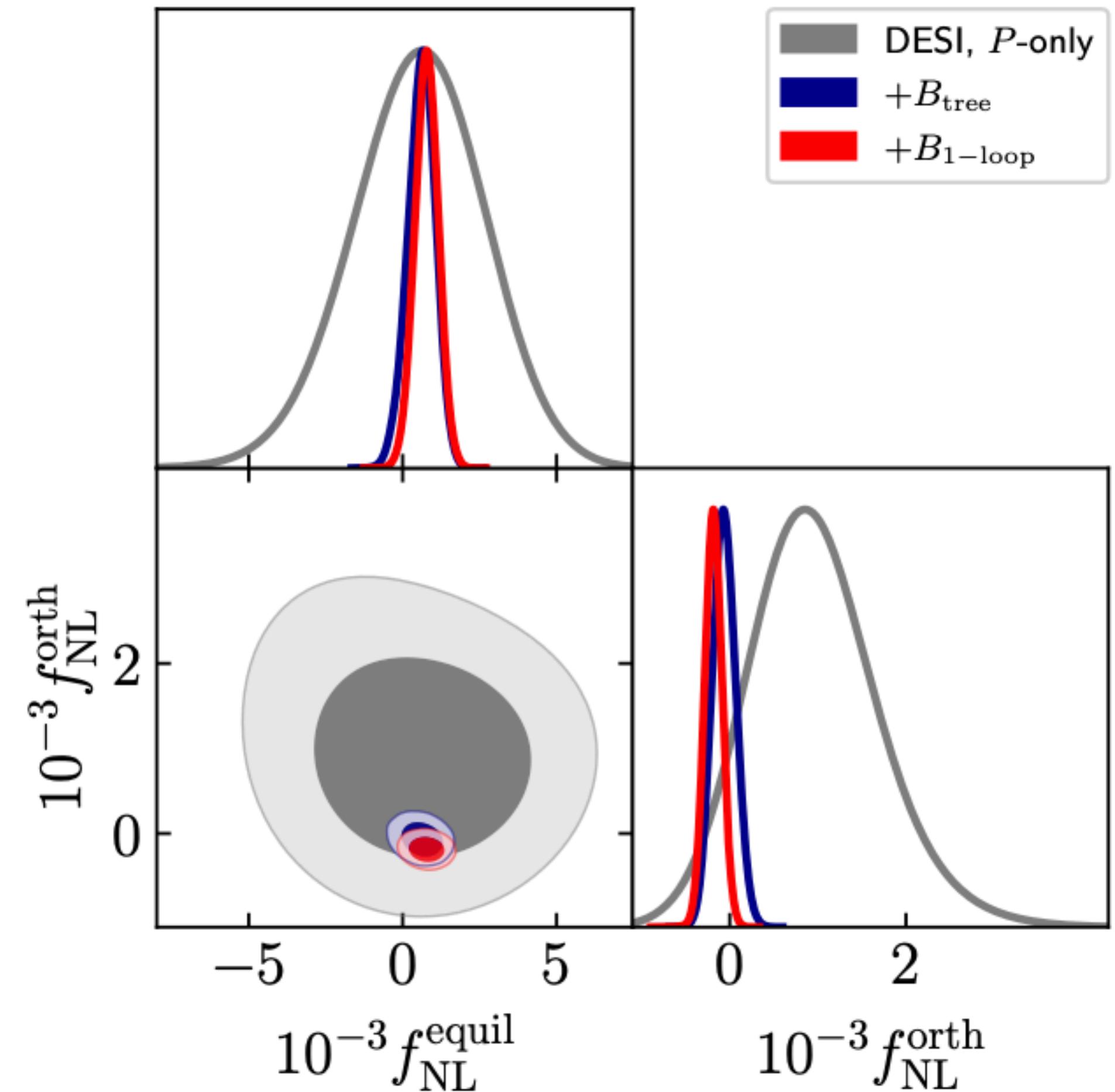
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• Single-Field:  $f_{\text{NL}}^{\text{eq}} = 200 \pm 230$ ,  $f_{\text{NL}}^{\text{forth}} = -24 \pm 86$

(Using the DESI DR1 one-loop power spectrum and bispectrum, plus the high- $z$  quasar sample)

- Adding **Planck**, we obtain the **tightest** constraint on local PNG yet!!

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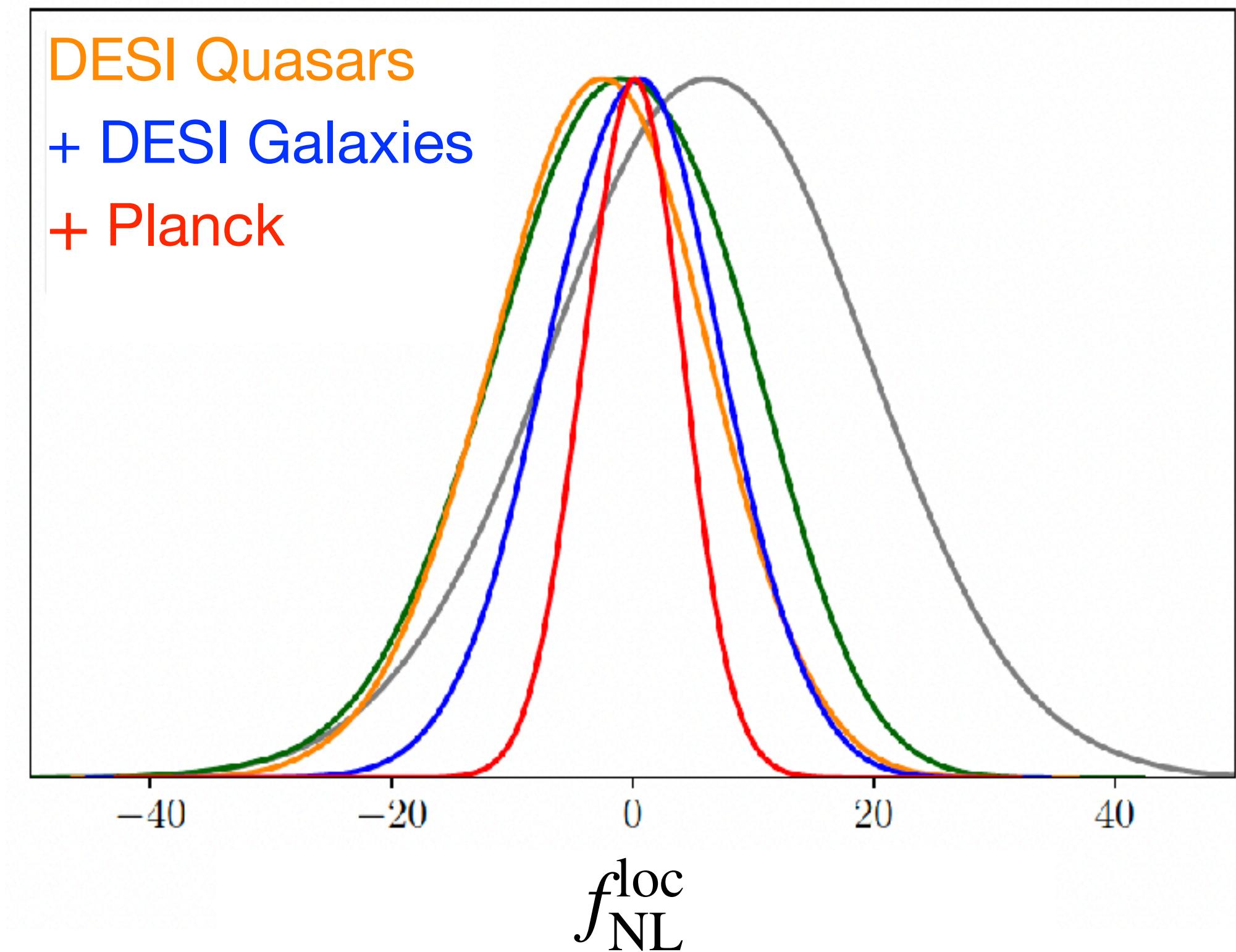
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# Summary

PNG analysis is a very active field!

- We can now constrain inflationary **four-point** functions in the CMB, including the **cosmological collider**!
- We can probe **arbitrary** non-separable bispectrum models with CMB data and **machine-learning**
- **Galaxy surveys** are providing exciting new insights into inflation and are starting to rival the CMB