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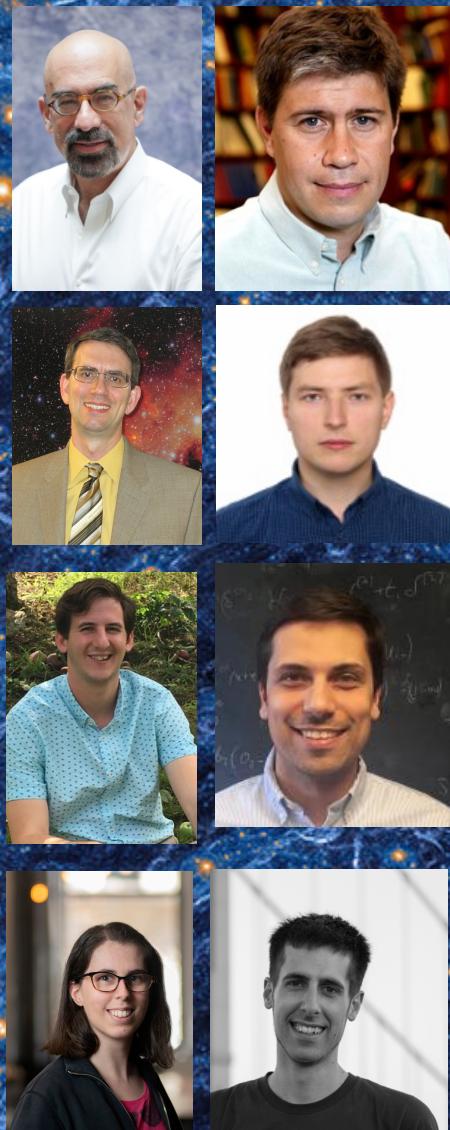
INSTITUTE FOR  
ADVANCED STUDY



# Have We Exhausted the Galaxy Two-Point Function?

**Oliver Philcox (Princeton / IAS)**

Gravity Group, Princeton University, 4/2/21



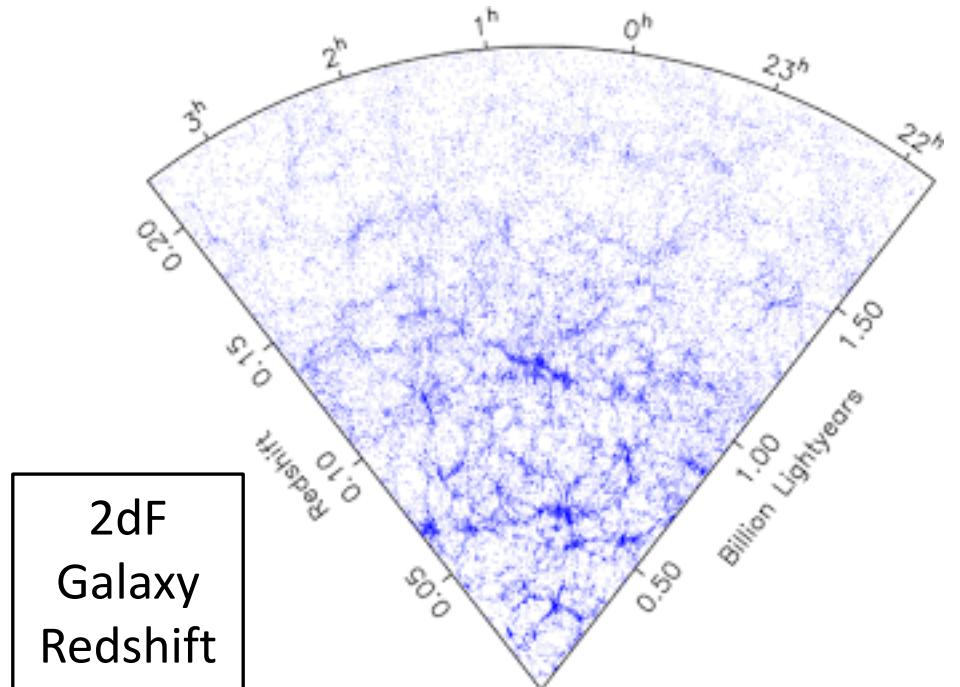
Based on: [1912.01010](https://arxiv.org/abs/1912.01010), [2005.01739](https://arxiv.org/abs/2005.01739) [2009.03311](https://arxiv.org/abs/2009.03311), [2006.10055](https://arxiv.org/abs/2006.10055), [2012.09389](https://arxiv.org/abs/2012.09389), [2102.08384](https://arxiv.org/abs/2102.08384)

# Extracting Information from Galaxy Surveys

- Fundamental observable: the galaxy **overdensity** field

$$\delta(\mathbf{r}) = \frac{1}{\bar{n}} [n_g(\mathbf{r}) - n_r(\mathbf{r})]$$

Mean density      Galaxy positions      Random particle positions



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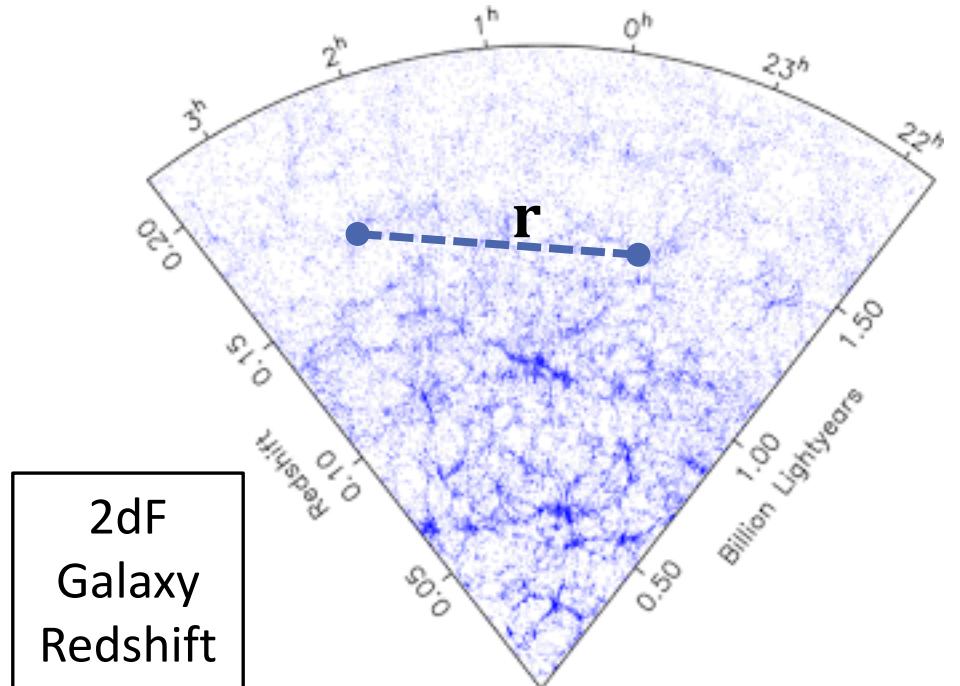
- Analyze with **summary statistics**:

- Two-point correlation function (2PCF),  $\xi(\mathbf{r})$

$$\xi(r) = \langle \delta(\mathbf{x})\delta(\mathbf{x} + \mathbf{r}) \rangle$$

- Power spectrum,  $P(\mathbf{k})$

$$(2\pi)^3 \delta_D(\mathbf{k} + \mathbf{k}') P(k) = \langle \delta(\mathbf{k})\delta(\mathbf{k}') \rangle$$



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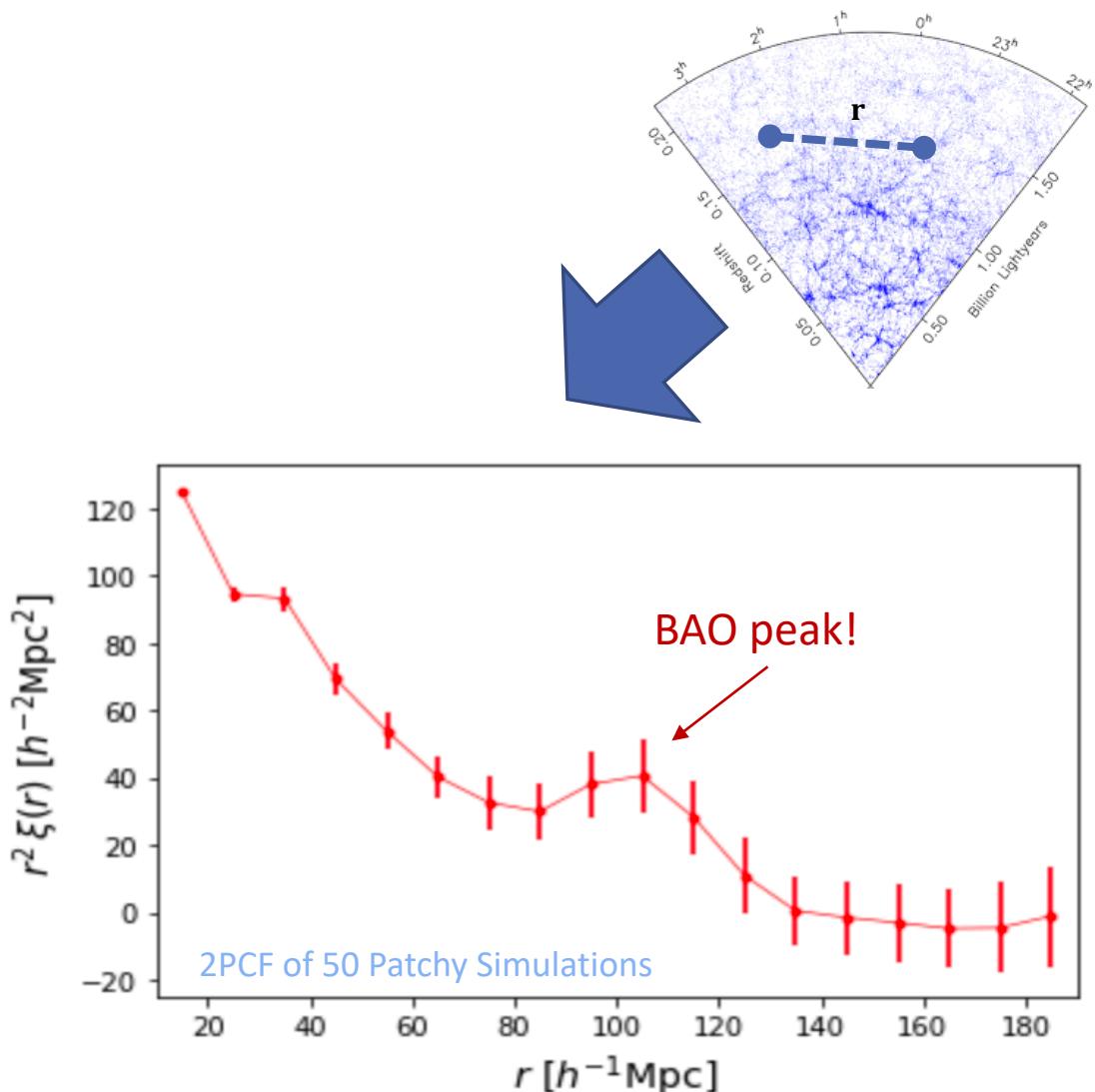
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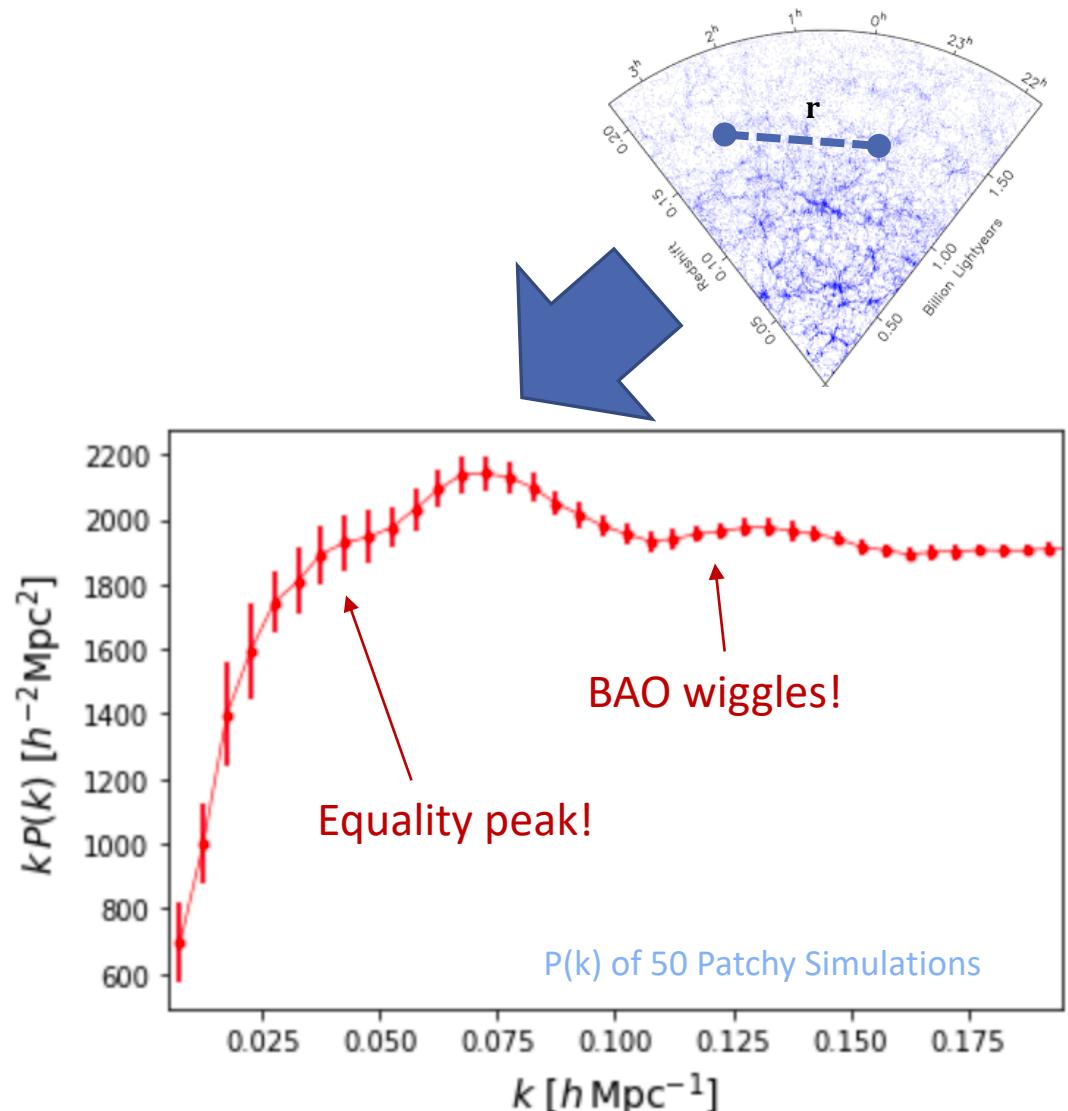
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# Understanding Anisotropy

- Redshift-space distortions lead to anisotropy

- Parametrize by **galaxy separation** and angle to **line-of-sight**,  $\hat{\mathbf{n}}$

$$\xi(\mathbf{r}) = \sum_{\ell} \xi_{\ell}(r) L_{\ell}(\hat{\mathbf{r}} \cdot \hat{\mathbf{n}}) \quad P(\mathbf{k}) = \sum_{\ell} P_{\ell}(k) L_{\ell}(\hat{\mathbf{k}} \cdot \hat{\mathbf{n}})$$

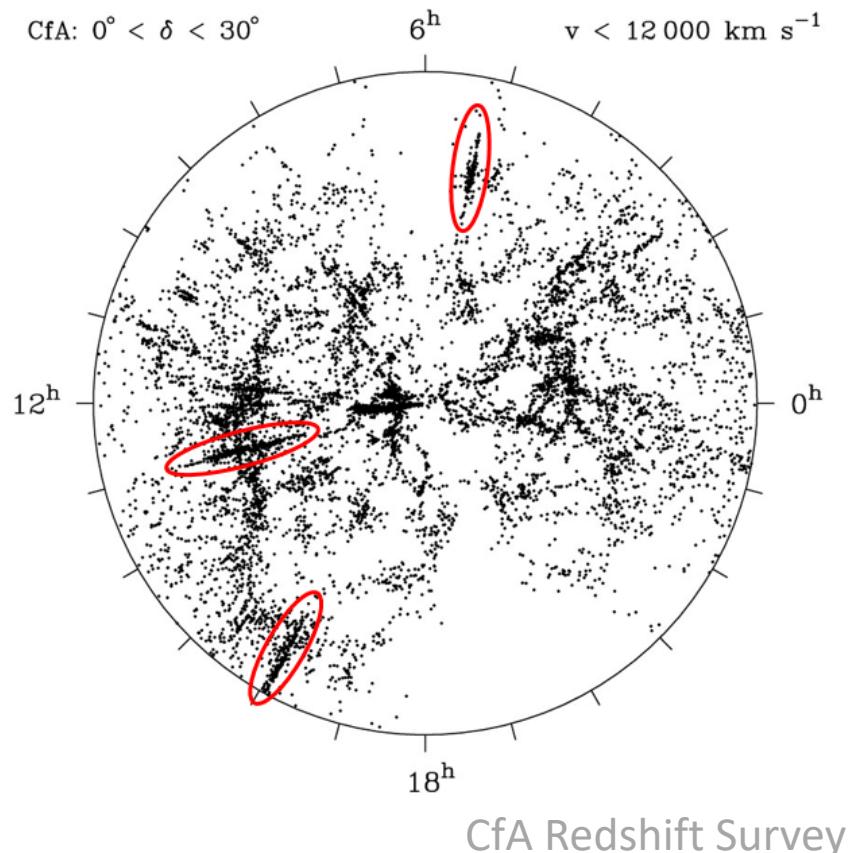
- Define the **multipoles**:

$$\hat{\xi}_{\ell}(r) = (2\ell + 1) \int \frac{d\Omega_r}{4\pi} \int d\mathbf{x} \delta(\mathbf{x}) \delta(\mathbf{x} + \mathbf{r}) L_{\ell}(\hat{\mathbf{r}} \cdot \hat{\mathbf{n}})$$

Legendre  
Polynomials

$$\hat{P}_{\ell}(k) = \frac{(2\ell + 1)}{V} \int \frac{d\Omega_k}{4\pi} \int d\mathbf{r}_1 d\mathbf{r}_2 e^{-i\mathbf{k} \cdot (\mathbf{r}_1 - \mathbf{r}_2)} \delta(\mathbf{r}_1) \delta(\mathbf{r}_2) L_{\ell}(\hat{\mathbf{k}} \cdot \hat{\mathbf{n}})$$

$= \delta(\mathbf{k}) \delta^*(\mathbf{k})$



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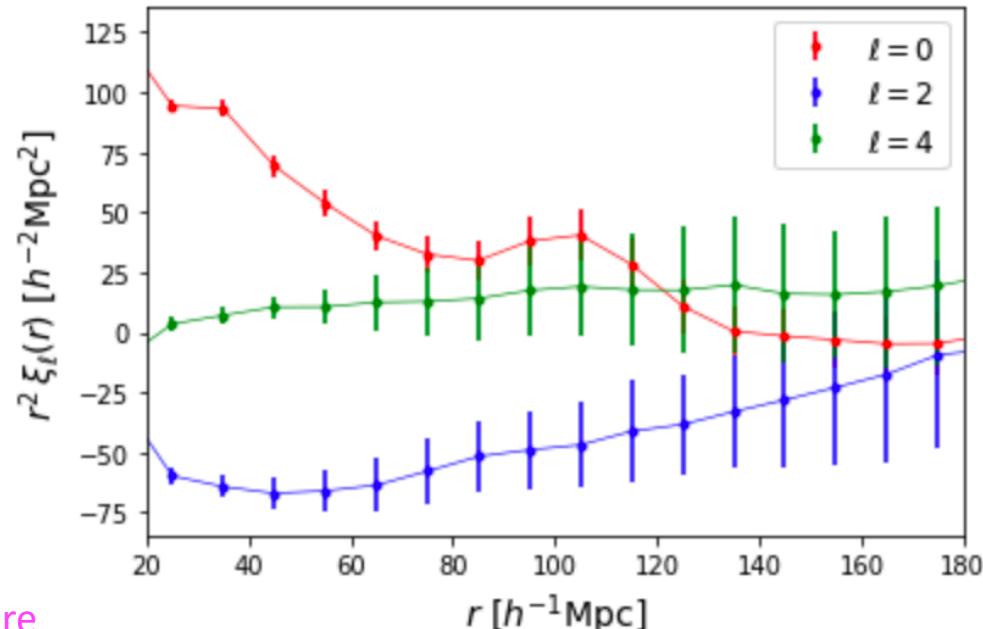
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Legendre  
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# Understanding Anisotropy

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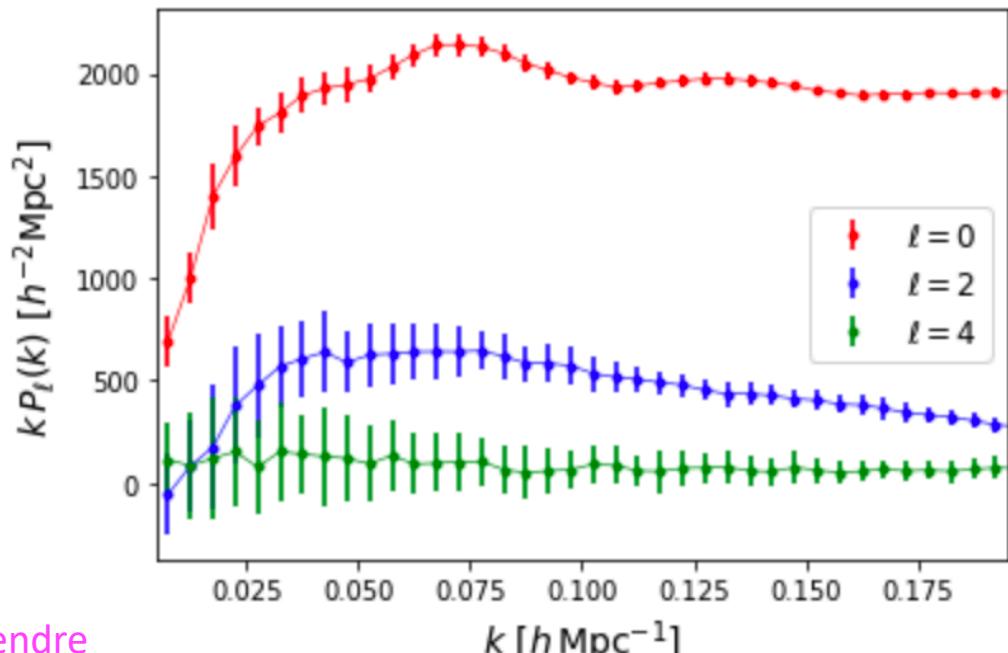
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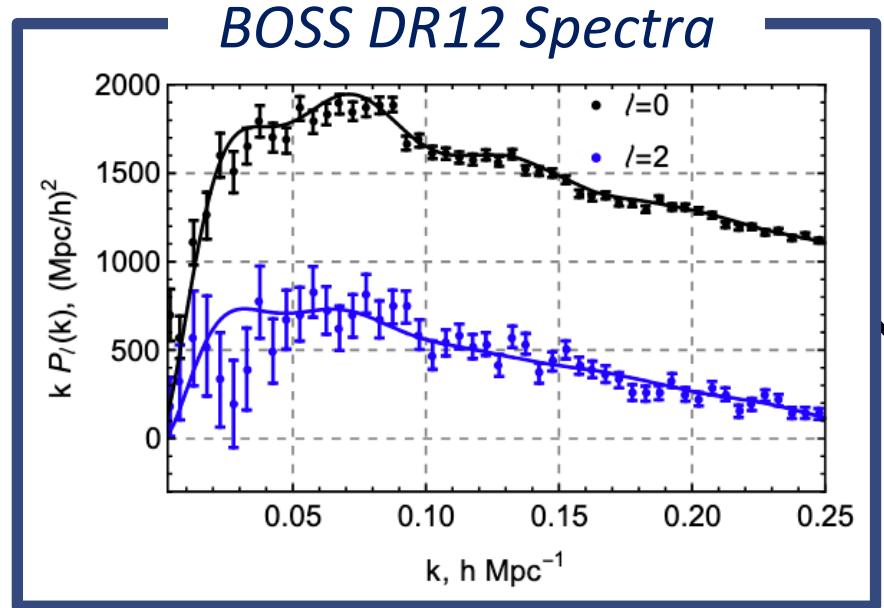
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$$\begin{aligned} \hat{P}_{\ell}(k) &= \frac{(2\ell + 1)}{V} \int \frac{d\Omega_k}{4\pi} \int d\mathbf{r}_1 d\mathbf{r}_2 e^{-i\mathbf{k} \cdot (\mathbf{r}_1 - \mathbf{r}_2)} \delta(\mathbf{r}_1) \delta(\mathbf{r}_2) L_{\ell}(\hat{\mathbf{k}} \cdot \hat{\mathbf{n}}) \\ &= \delta(\mathbf{k}) \delta^*(\mathbf{k}) \end{aligned}$$



# Parameter Inference

CMB-Strength  
Parameter Constraints,  
including 1.6% on  $H_0$ !



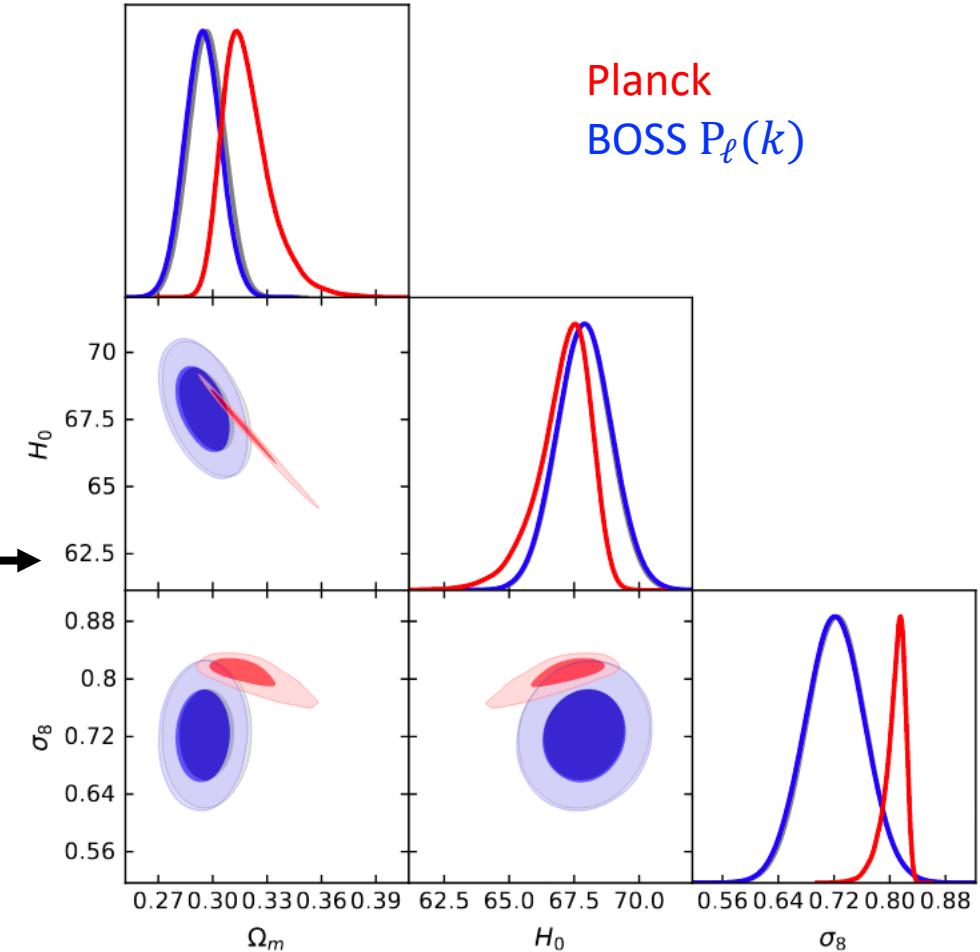
$P_{g,\ell}(k) = P_{g,\ell}^{\text{tree}}(k) + P_{g,\ell}^{\text{1-loop}}(k) + P_{g,\ell}^{\text{noise}}(k) + P_{g,\ell}^{\text{ctr}}(k)$

Linear Theory    1-loop PT    Shot-noise    Counterterms

*Theory Model*  
(Effective Field Theory)

MCMC

Mock Datasets



# Beyond 2-Point Statistics

The Universe is **non-Gaussian**

Information in **higher-point** functions, e.g.

- Bispectrum / 3PCF [Gil-Marín+16, Slepian+15, d'Amico+19]

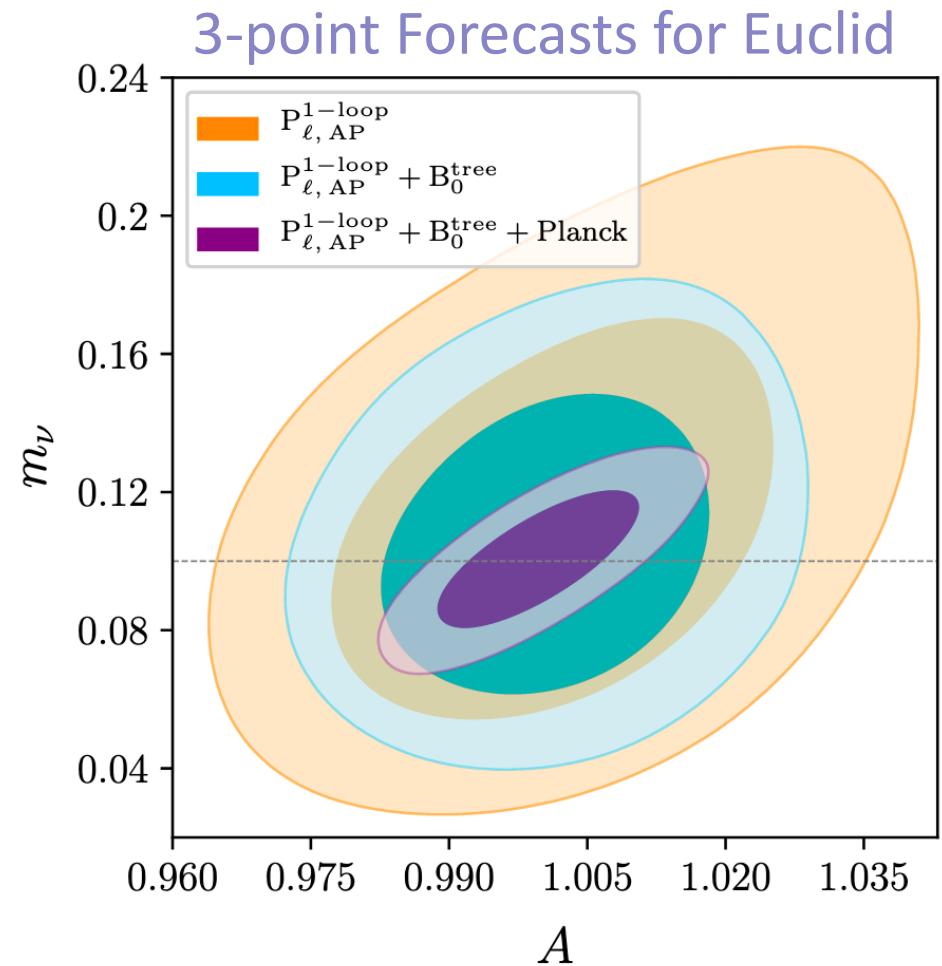
$$(2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B(\mathbf{k}_1, \mathbf{k}_2) = \langle \delta(\mathbf{k}_1) \delta(\mathbf{k}_2) \delta(\mathbf{k}_3) \rangle$$

- Trispectrum / 4PCF [Gualdi+20]

$$(2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4) T(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) = \langle \delta(\mathbf{k}_1) \delta(\mathbf{k}_2) \delta(\mathbf{k}_3) \delta(\mathbf{k}_4) \rangle$$

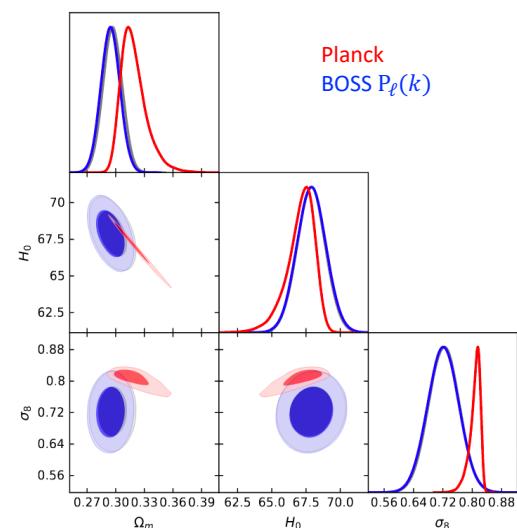
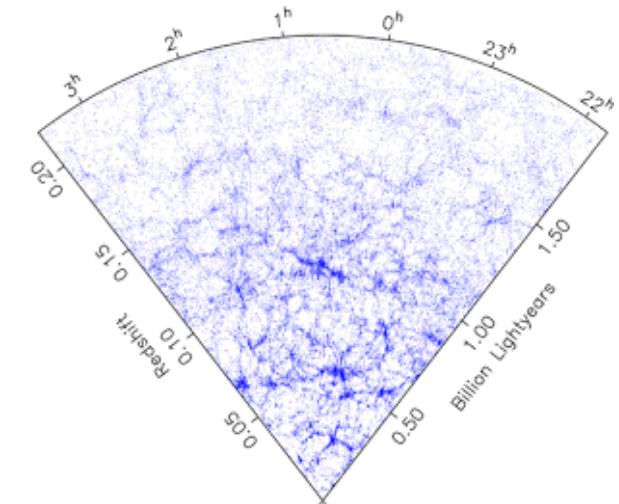
These get steadily larger and harder to measure.

- Not used in many cosmological analyses yet!



# Cosmology from $P_\ell(k)$ : A Summary

- Fundamental observable: the galaxy **overdensity** field
- $P_\ell(k)$  parametrized by pair separation and **line-of-sight angle**
- Power spectrum estimators measure  $|\delta(\mathbf{k})|^2 L_\ell(\hat{\mathbf{k}} \cdot \hat{\mathbf{n}})$
- Computed using **Fast Fourier Transforms** (FFTs)
- Compare data and theory with MCMC



# Cosmology from $P_\ell(k)$ : A Summary

- Fundamental observable: the galaxy **overdensity** field *Is this the best field to use?*
- $P_\ell(k)$  parametrized by pair separation and **line-of-sight angle** *How do we define this angle?*
- Power spectrum estimators measure  $|\delta(\mathbf{k})|^2 L_\ell(\hat{\mathbf{k}} \cdot \hat{\mathbf{n}})$  *Can we estimate it more optimally?*
- Computed using **Fast Fourier Transforms** (FFTs) *Are FFTs always the most efficient?*
- Compare data and theory with MCMC *Can data-compression help?*

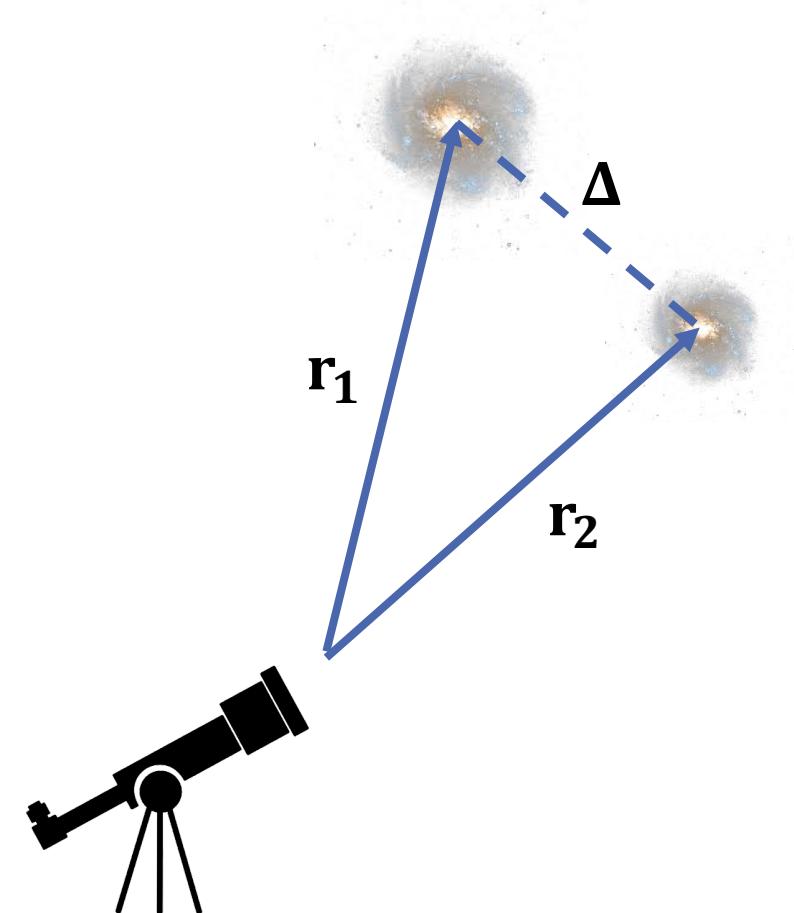


# 1. Parametrizing Anisotropy

# Choosing the Line-of-Sight

- Galaxy correlation function depends on the angle between the separation vector  $\Delta$  and the line-of-sight  $\hat{\mathbf{n}}$ :

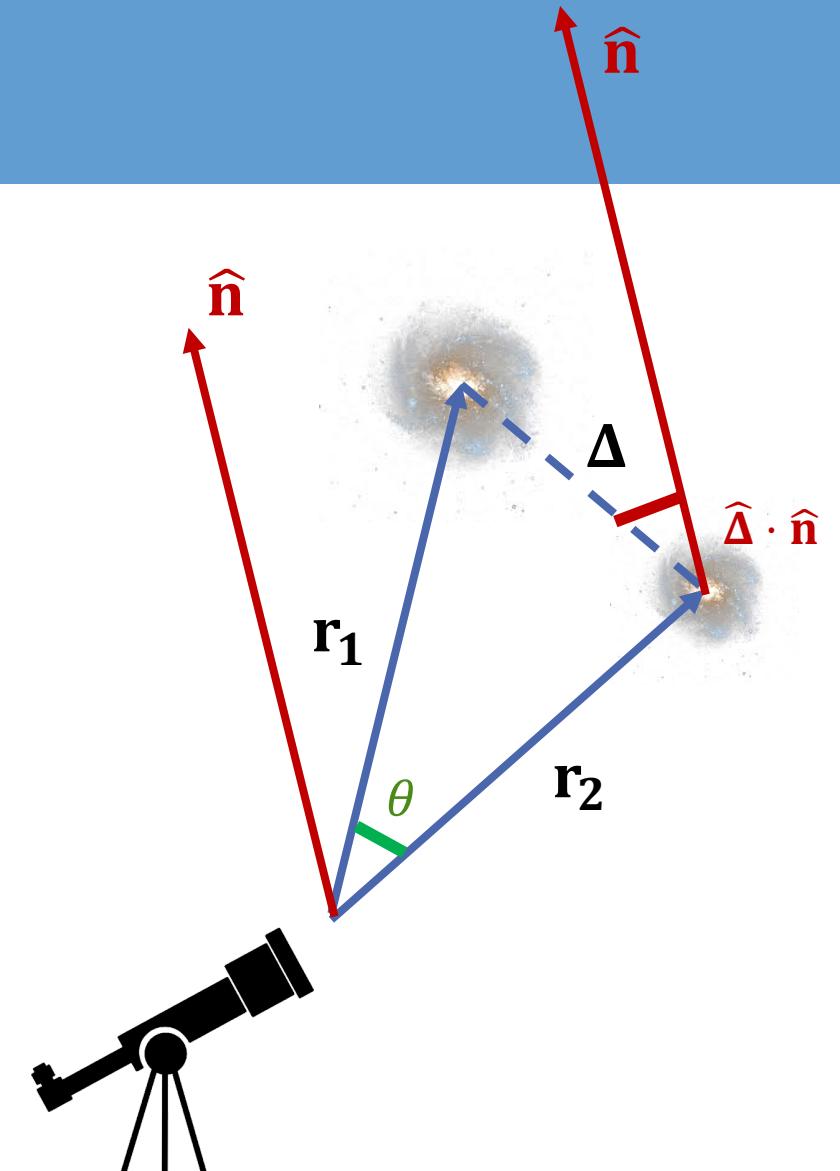
$$\hat{\xi}_\ell(r) = \frac{2\ell+1}{V} \int d\mathbf{r}_1 d\mathbf{r}_2 \underbrace{\delta(\mathbf{r}_1)\delta(\mathbf{r}_2)}_{\text{Density Fields}} \underbrace{L_\ell(\hat{\Delta} \cdot \hat{\mathbf{n}})}_{\text{Angular Dependence}} \left[ \frac{\delta_D(r - \Delta)}{4\pi r^2} \right] \underbrace{\frac{\delta_D(r - \Delta)}{4\pi r^2}}_{\text{Binning}}$$



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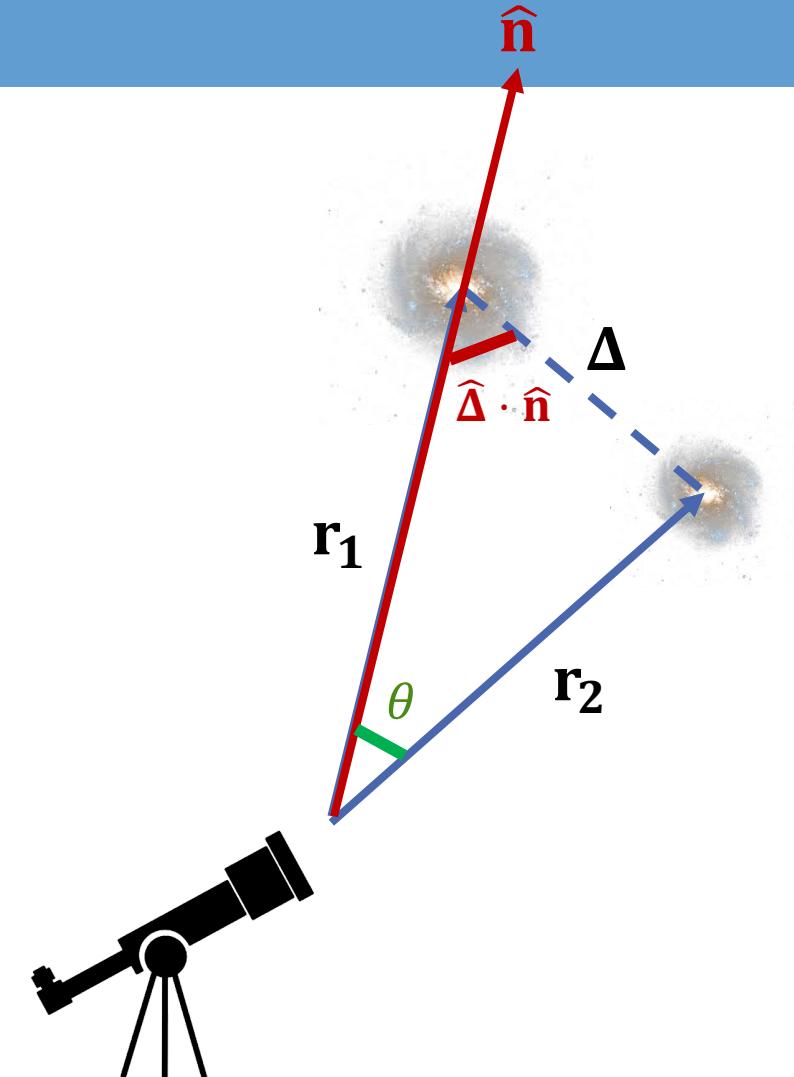
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  - Yamamoto approximation:  $\hat{\mathbf{n}} = \hat{\mathbf{r}}_1$ ,  $\mathcal{O}(\theta^2)$  error

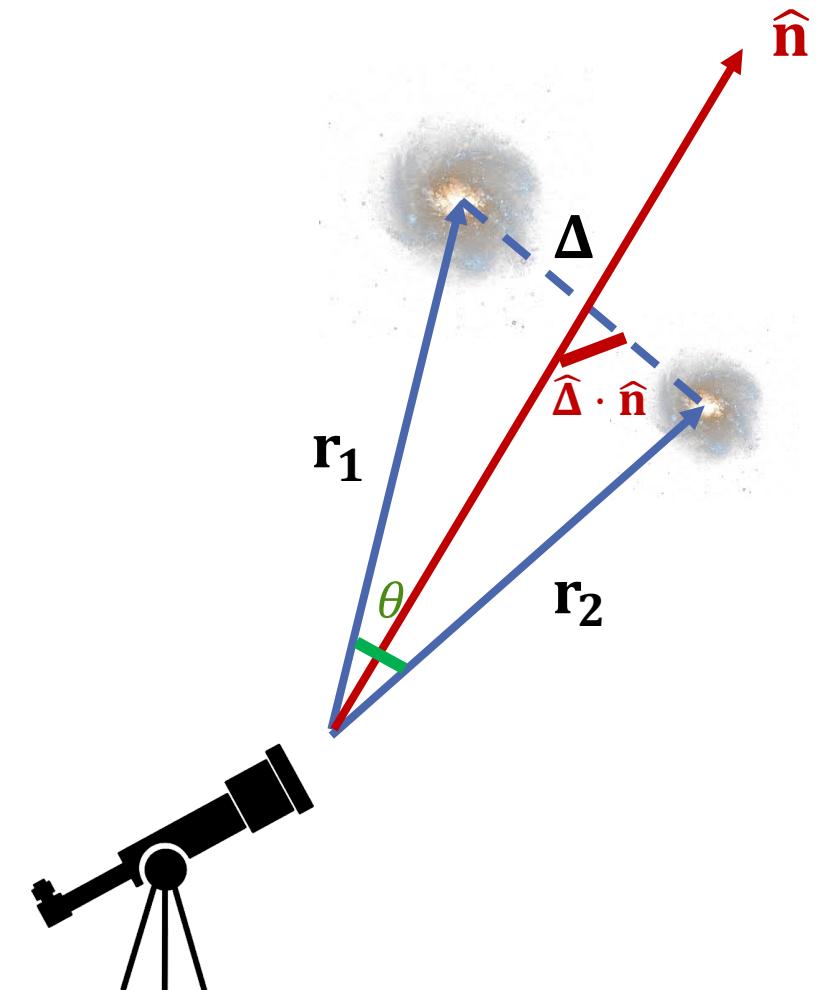


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  - Midpoint method:  $\hat{\mathbf{n}} = \widehat{\mathbf{r}_1 + \mathbf{r}_2}$ ,  $\mathcal{O}(\theta^{4+})$  error



# Lines-of-Sight in the Power Spectrum

- Same for the **power spectrum**:

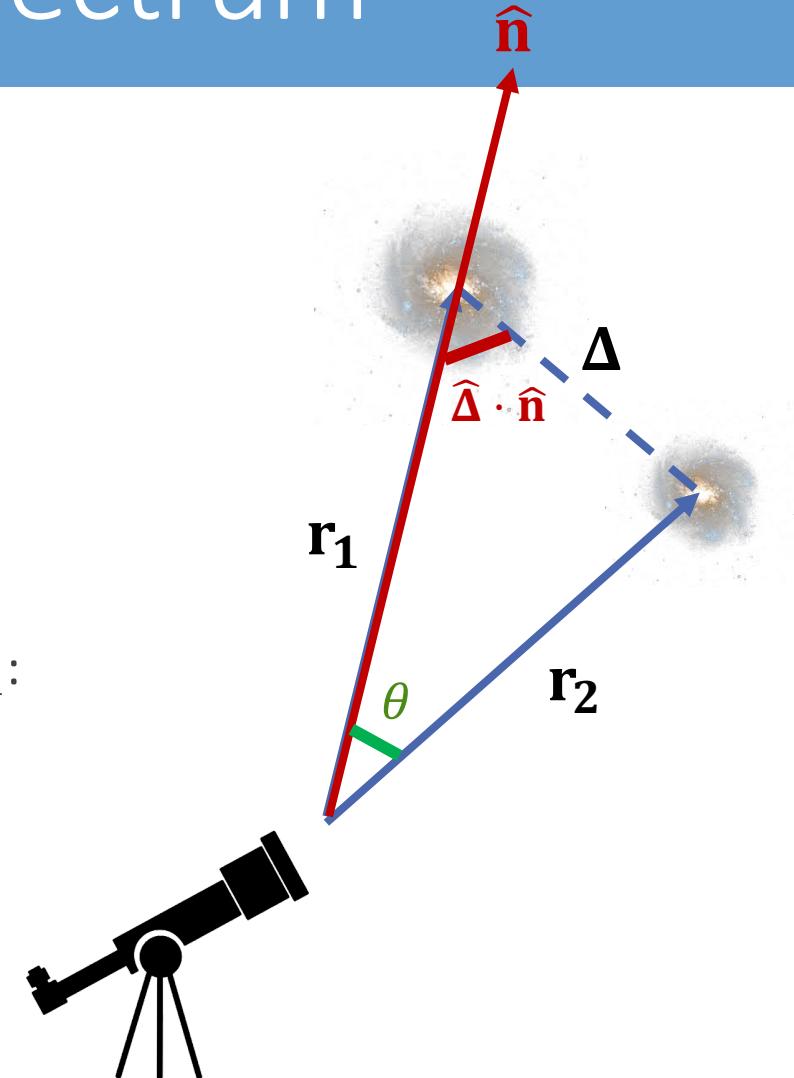
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Fourier Transform      Angular Dependence  
Density Fields

- This is **easy to implement** for the Yamamoto approximation,  $\hat{\mathbf{n}} = \hat{\mathbf{r}}_1$ :

$$\hat{P}_\ell^{\text{Yama}}(k) = \frac{4\pi}{V} \int_{\Omega_k} \left[ \sum_{m=-\ell}^{\ell} Y_\ell^{m*}(\hat{\mathbf{k}}) \mathcal{F}[Y_\ell^m \delta](\mathbf{k}) \right] \delta^*(\mathbf{k})$$

Spherical Harmonics



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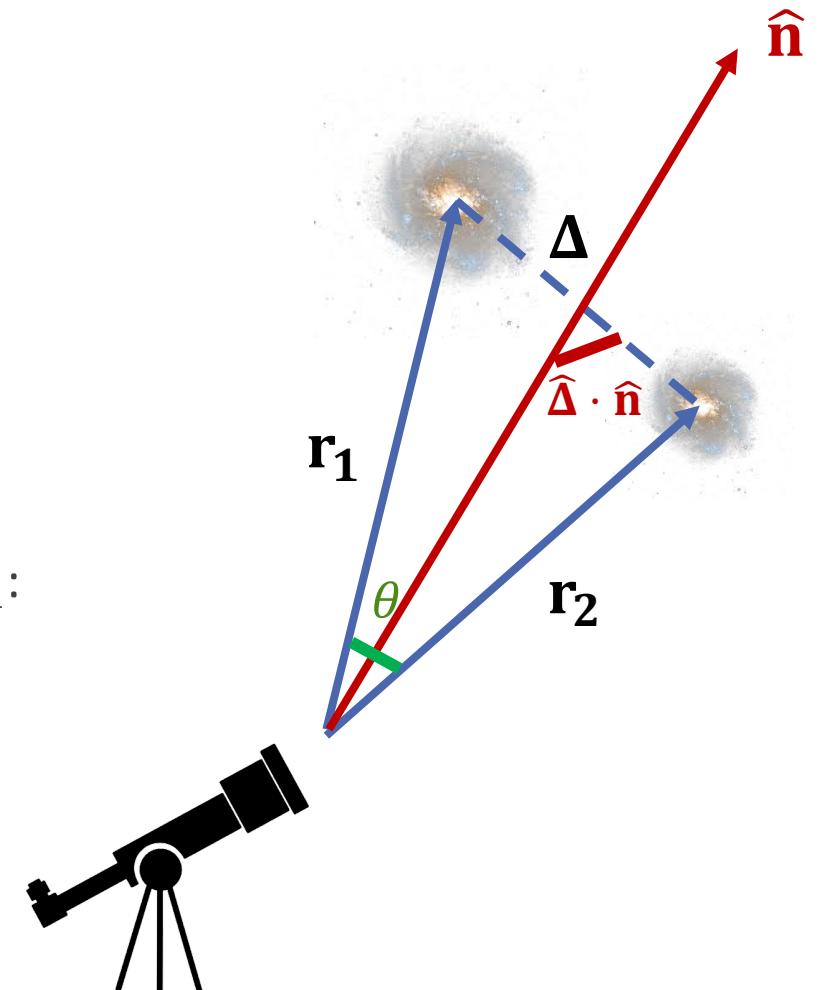
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Spherical Harmonics

- But **not separable** for the midpoint method!



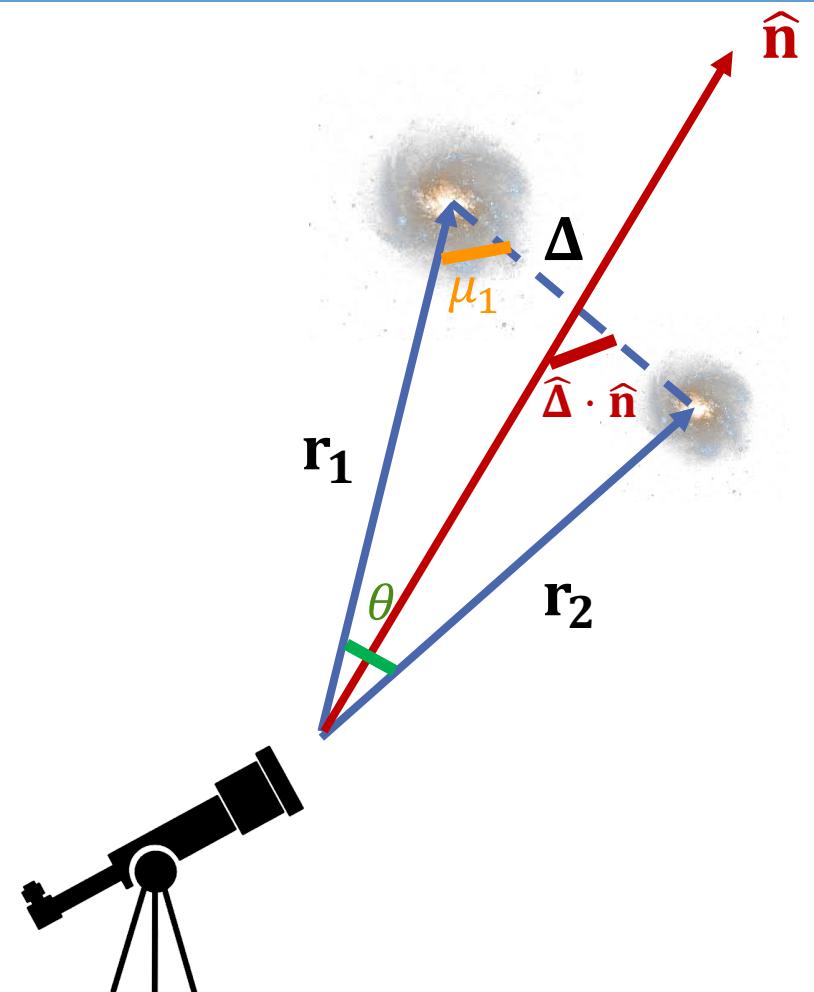
# Implementing the Midpoint Method

- Use a trick to make the integrals separable:

- Expand in powers of  $\theta \sim \Delta/r_1$ :

$$L_\ell(\hat{\Delta} \cdot \widehat{\mathbf{r}_1 + \mathbf{r}_2}) = \sum_{\alpha=0}^{\infty} \sum_{J=0}^{\ell+\alpha} \underbrace{f_J^{\alpha,\ell}}_{\text{Coefficients}} \underbrace{\left(\frac{\Delta}{2r_1}\right)^\alpha}_{\text{Survey Angle, } \ll 1} \underbrace{L_J(\hat{\Delta} \cdot \mathbf{r}_1)}_{\text{Yamamoto Piece}}$$

$$\begin{aligned} L_2(\hat{\Delta} \cdot \widehat{\mathbf{r}_1 + \mathbf{r}_2}) &= L_2(\mu_1) + \frac{6}{5} \left( \frac{\Delta}{2r_1} \right) [L_1(\mu_1) - L_3(\mu_1)] \\ &\quad + \frac{1}{35} \left( \frac{\Delta}{2r_1} \right)^2 [7L_0(\mu_1) - 55L_2(\mu_1) + 48L_4(\mu_1)] \\ &\quad - \frac{4}{105} \left( \frac{\Delta}{2r_1} \right)^3 [9L_1(\mu_1) - 49L_3(\mu_1) + 40L_5(\mu_1)] \\ &\quad + \frac{1}{385} \left( \frac{\Delta}{2r_1} \right)^4 [11L_0(\mu_1) + 165L_2(\mu_1) - 816L_4(\mu_1) + 640L_6(\mu_1)] \\ &\quad + \dots \end{aligned}$$



# Implementing the Midpoint Method

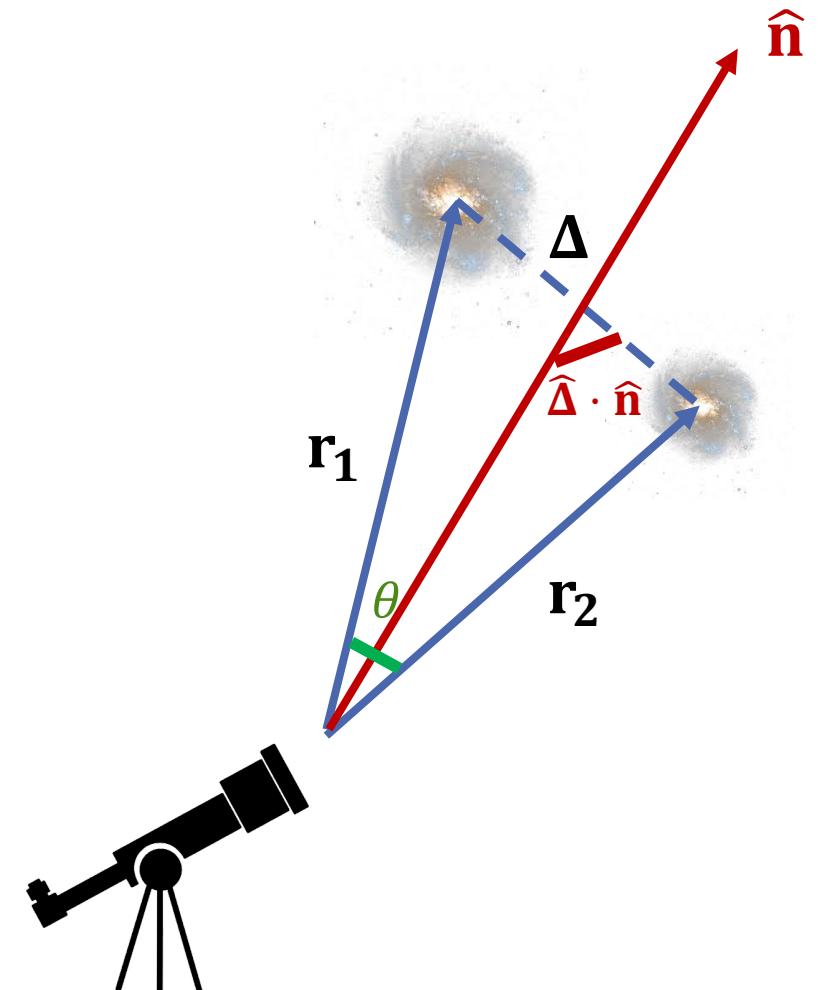
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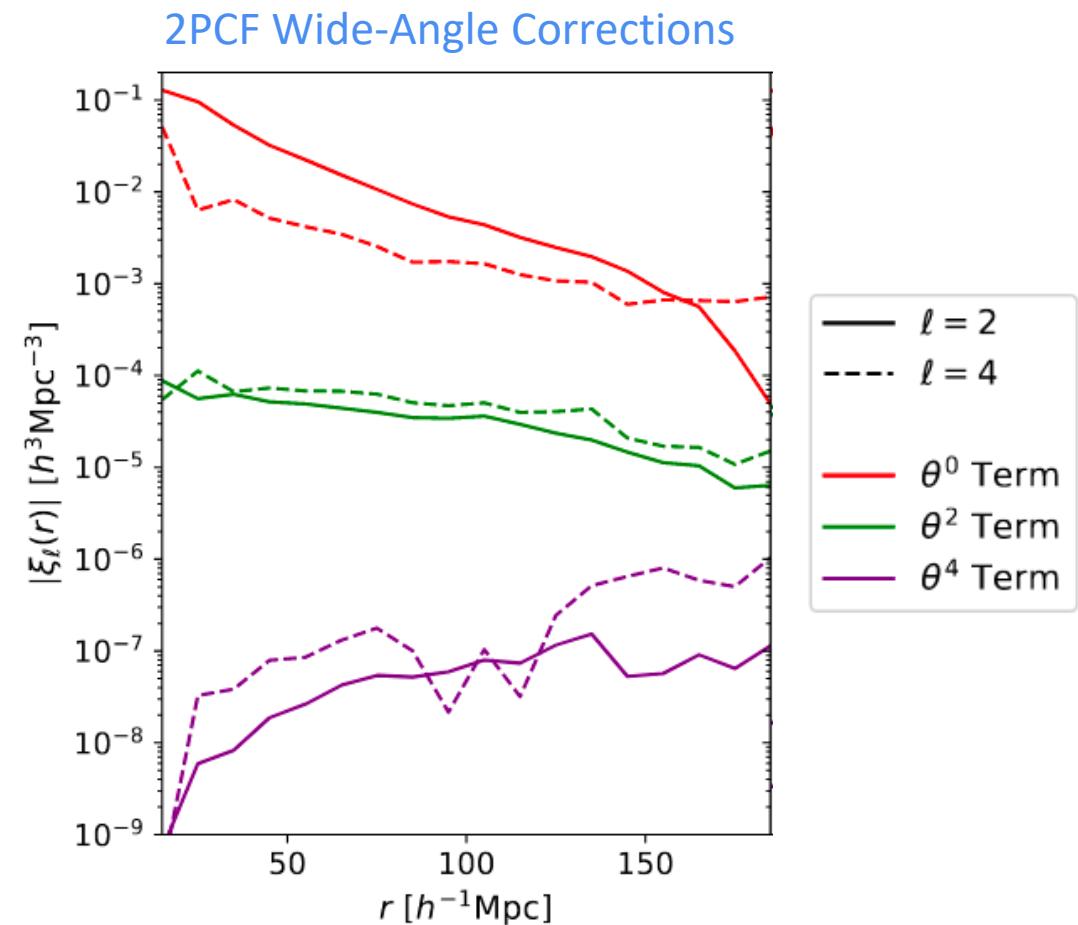
Coefficients                                  Yamamoto Piece

- Can now compute the 2PCF using Fourier transforms!
- Also applies to the **power spectrum**
- Same computational complexity as Yamamoto approximation



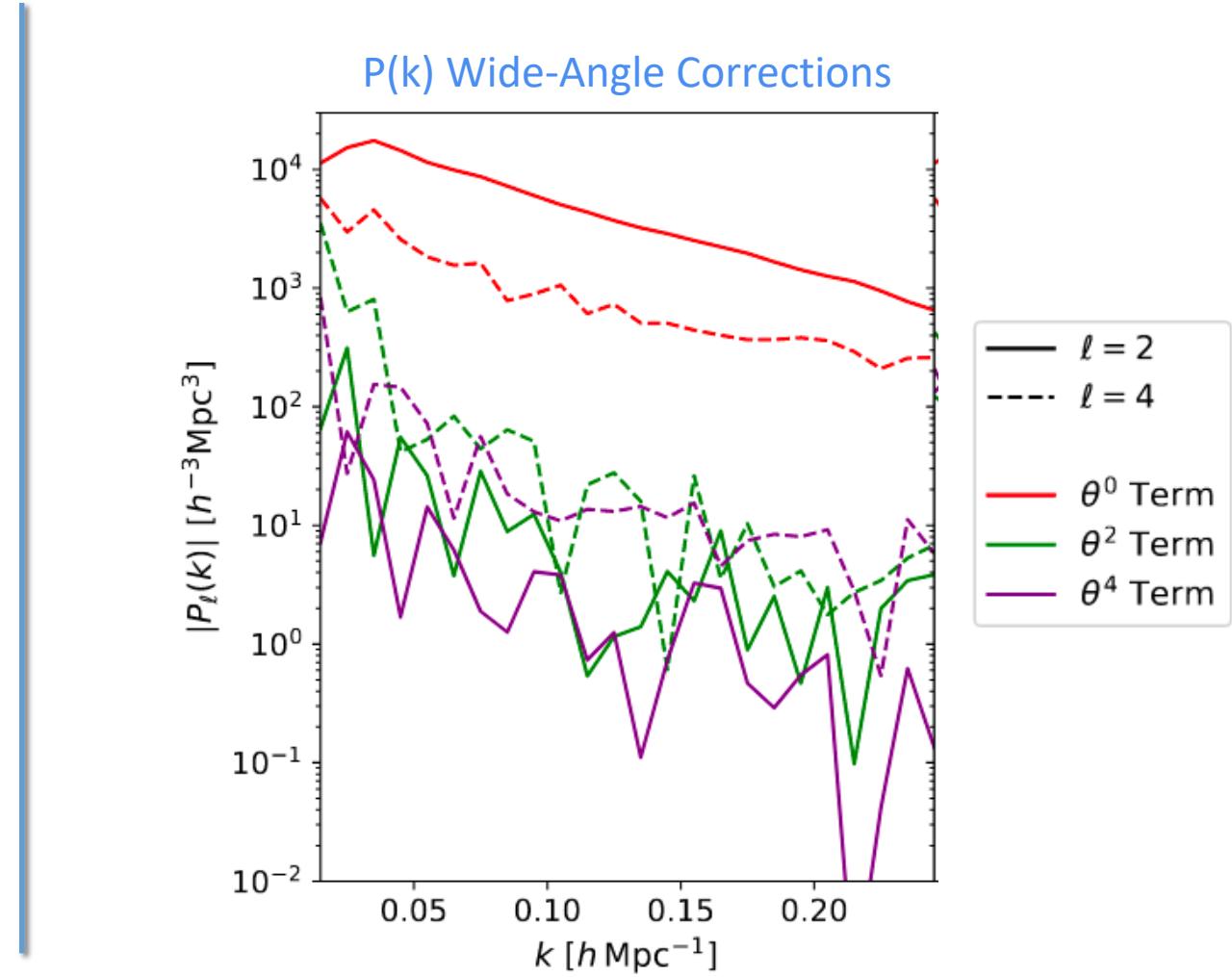
# The Midpoint Method in Practice

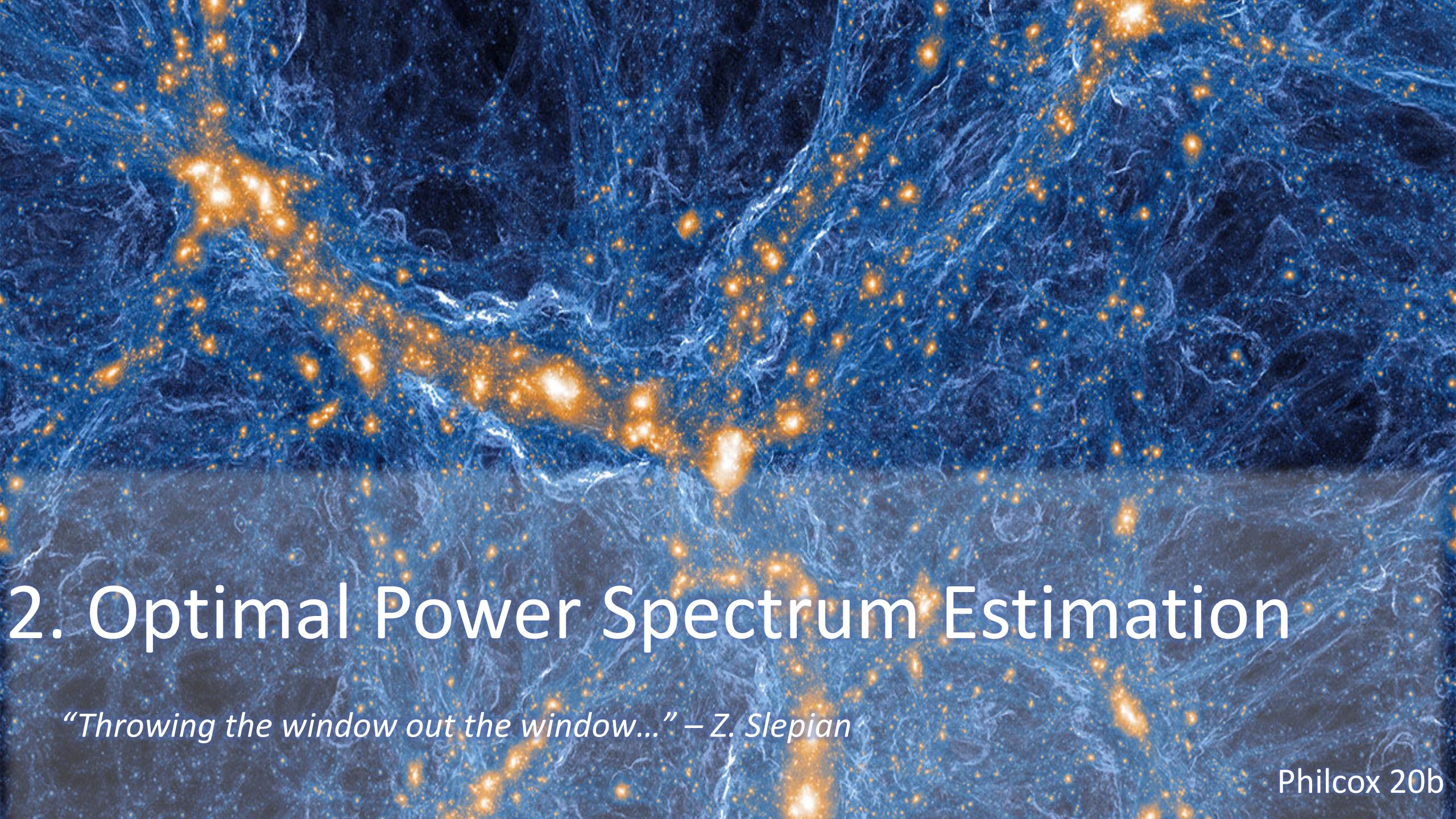
- BOSS correlation function:
  - $\theta \sim 0.1 - 0.2$  at the BAO scale
  - Larger  $r \Rightarrow$  Larger corrections
  - Still  $\ll 1\sigma$  for BOSS



# The Midpoint Method in Practice

- BOSS correlation function:
  - $\theta \sim 0.1 - 0.2$  at the BAO scale
  - Larger  $r \Rightarrow$  Larger corrections
  - Still  $\ll 1\sigma$  for BOSS
- BOSS  $P(k)$ :
  - Spectrum is an integral over **all**  $r$  in survey
  - $\theta \sim 1$  for the largest-modes
  - Corrections **are** marginally important at **all**  $k$
- Most important for **wide surveys at low redshifts**





## 2. Optimal Power Spectrum Estimation

*“Throwing the window out the window...” – Z. Slepian*

# The FKP Estimator

Power spectrum **isn't** just  $|\delta(\mathbf{k})|^2$ .

- Neglects **inhomogeneous noise** and **survey window functions**
- 1. Define  $\delta(\mathbf{r})$  as the difference between **galaxy** and **random** densities
- 2. Add an **FKP weight** to incorporate **Poisson noise** densities (and systematics)

This is the **optimal solution** on small-scales with Poisson noise

**But:**

- **Not** optimal on large scales
- Measures the **window-convolved** power spectrum

$$\hat{P}(k) = \int \frac{d\Omega_k}{4\pi} \int d\mathbf{r}_1 d\mathbf{r}_2 e^{-i\mathbf{k}\cdot(\mathbf{r}_1-\mathbf{r}_2)} \delta(\mathbf{r}_1)\delta(\mathbf{r}_2)$$

$$\delta(\mathbf{r}) \rightarrow \frac{w(\mathbf{r})[n_g(\mathbf{r}) - \alpha_r n_r(\mathbf{r})]}{I^{1/2}}, \quad I \equiv \int d\mathbf{r} w^2(\mathbf{r}) \bar{n}^2(\mathbf{r})$$

$$w(\mathbf{r}) = \frac{w_{\text{sys}}(\mathbf{r})}{1 + P_{\text{FKP}} n(\mathbf{r})}$$

*Poisson Noise Correction,  $P_{\text{FKP}} \sim 10^4$*

Feldman+94  
Tegmark+98

# Optimal Estimators

Maximize the **likelihood** for data,  $\mathbf{d}$ , with band-powers  $\mathbf{p}$  and pixel covariance  $\mathbf{C}(\mathbf{p})$

$$-2 \log L[\mathbf{d}](\mathbf{p}) = \mathbf{d}^T \mathbf{C}^{-1}(\mathbf{p}) \mathbf{d} + \text{Tr} \log \mathbf{C}(\mathbf{p}) + \text{const.} \quad \xleftarrow{\boxed{\text{Gaussian likelihood}}}$$

Gives a **maximum-likelihood** estimator for the **unwindowed** power spectrum:

$$\hat{p}_\alpha^{\text{ML}} = p_\alpha^{\text{fid}} + \sum_\beta F_{\alpha\beta}^{-1} (\hat{q}_\beta - \bar{q}_\beta)$$

Fisher matrix (from sims)

Known  
Fiducial Model

Bias term  
(from sims)

Quadratic  
estimator  
(from data)

Estimator is a **quadratic** function of the data,  $\hat{q}_\beta$

# Implementing the ML Estimator

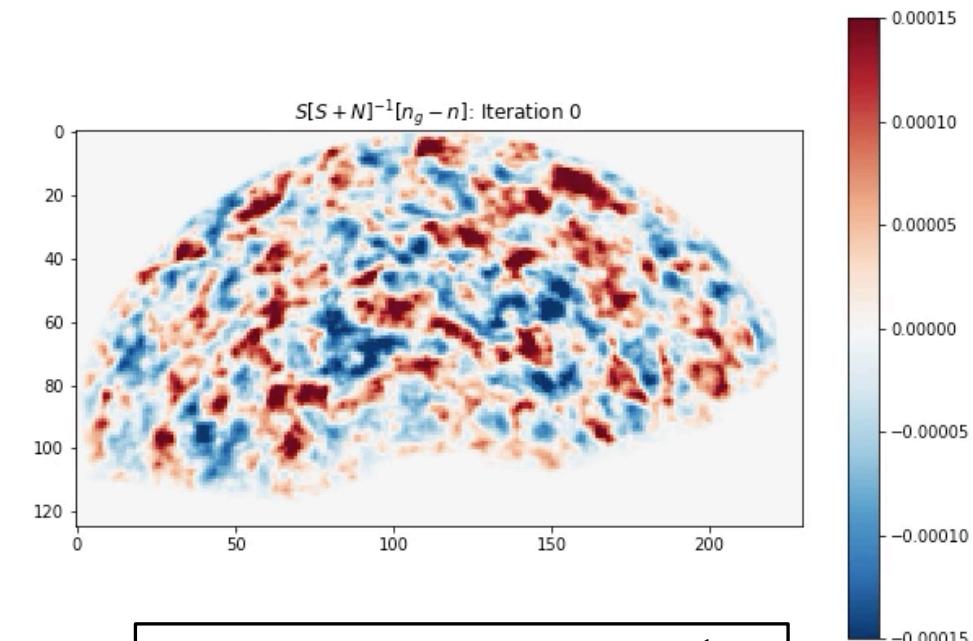
$$\hat{p}_\alpha^{\text{ML}} = p_\alpha^{\text{fid}} + \sum_\beta F_{\alpha\beta}^{-1} (\hat{q}_\beta - \bar{q}_\beta)$$

- Need the **quadratic estimator**  $\hat{q}_\beta$ :

$$\hat{q}(k) = \int \frac{d\Omega_k}{4\pi} \int d\mathbf{r} d\mathbf{r}' e^{-i\mathbf{k}\cdot(\mathbf{r}-\mathbf{r}')} [\mathbf{C}^{-1}\mathbf{d}](\mathbf{r}) [\mathbf{C}^{-1}\mathbf{d}](\mathbf{r}')$$

- Just a power spectrum of the **inverse-covariance weighted data**
- Need the covariance for each pair of **pixels**:

$$C(\mathbf{r}, \mathbf{r}') = n(\mathbf{r})n(\mathbf{r}') \underbrace{\int_{\mathbf{k}} e^{i\mathbf{k}\cdot(\mathbf{r}-\mathbf{r}')} \sum_\ell P_\ell(k) L_\ell(\hat{\mathbf{k}} \cdot \hat{\mathbf{r}}')}_{\text{Signal}} + (1 + \alpha)n(\mathbf{r})\delta_D(\mathbf{r} - \mathbf{r}') \underbrace{+ (1 + \alpha)n(\mathbf{r})\delta_D(\mathbf{r} - \mathbf{r}')}_{\text{Poisson Noise}}$$



Iterative computation of  $C^{-1}\mathbf{d}$   
from an initial guess

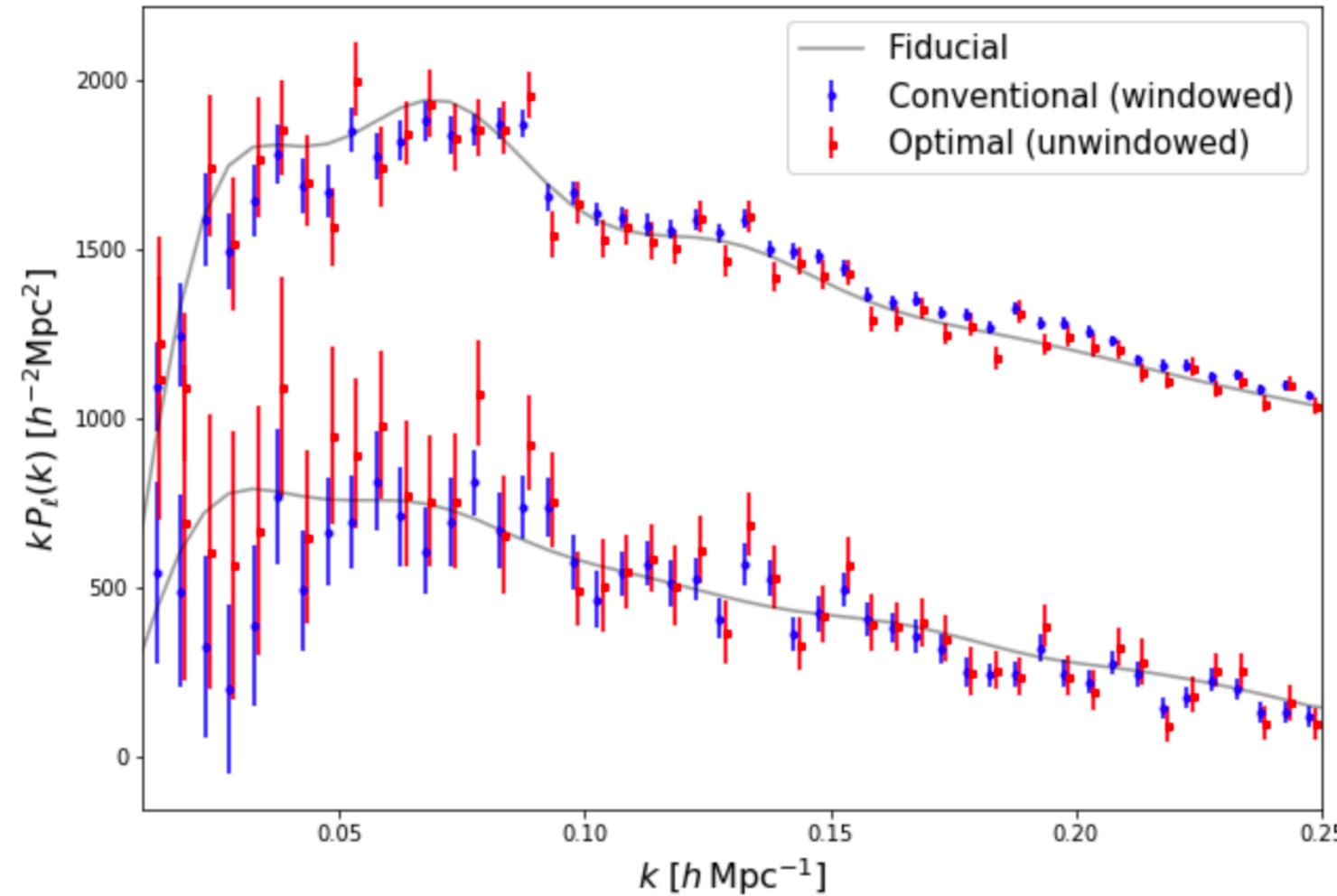
- This covariance is **gigantic** ( $N_{\text{pix}} \times N_{\text{pix}}$ )
- Never store directly
- Invert using **conjugate gradient descent** methods

# Implementing the ML Estimator

$$\hat{p}_\alpha^{\text{ML}} = p_\alpha^{\text{fid}} + \sum_\beta F_{\alpha\beta}^{-1} (\hat{q}_\beta - \bar{q}_\beta)$$

## Pipeline:

1. Choose a **fiducial cosmology**
2. Compute the **quadratic estimator** on the data,  $\hat{q}_\beta$
3. Repeat on **simulations** to get bias,  $\bar{q}_\beta$  and Fisher matrix,  $F_{\alpha\beta}$
4. Combine to get the **power spectrum**
5. **Optional:** Repeat with new cosmology

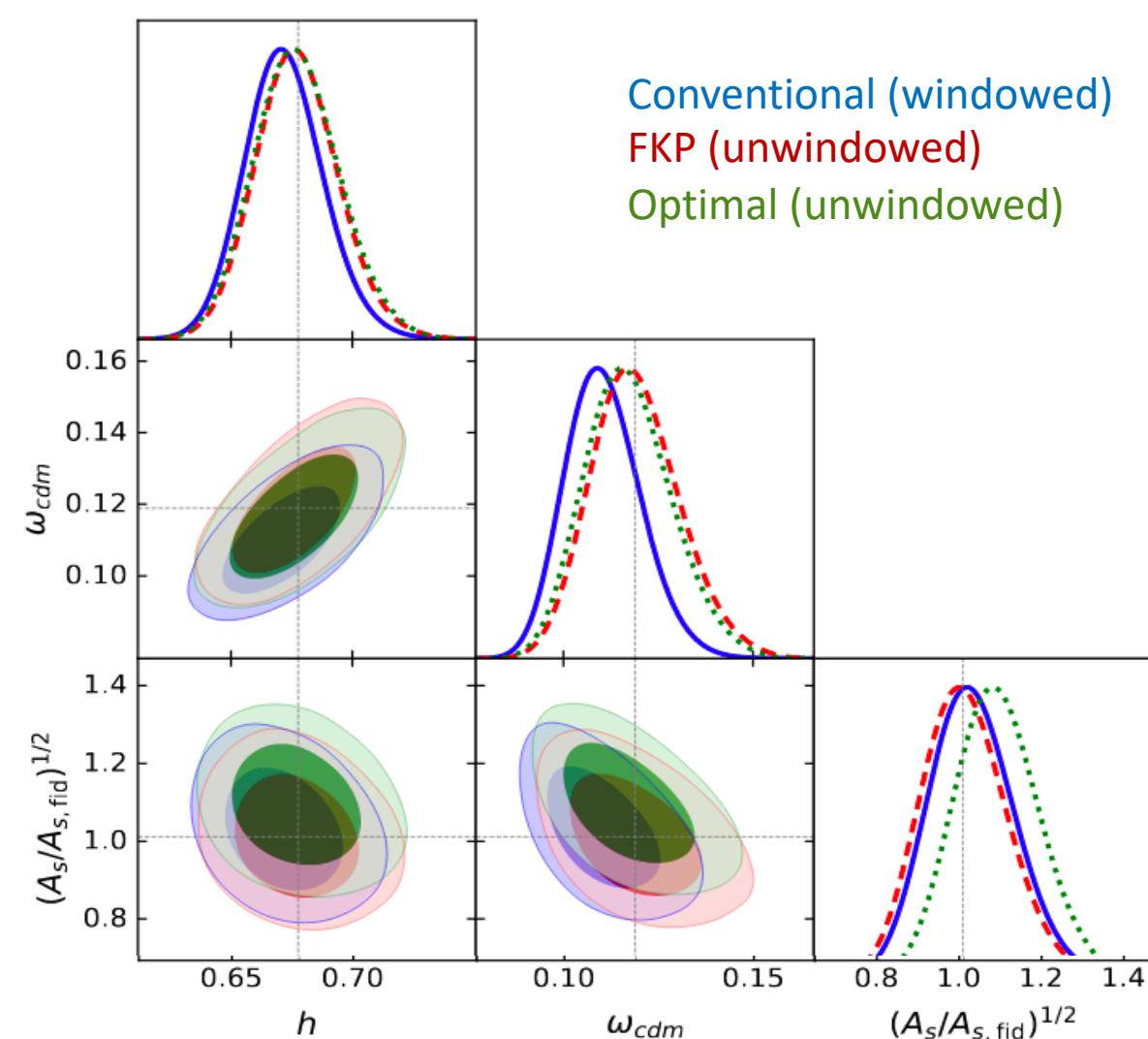


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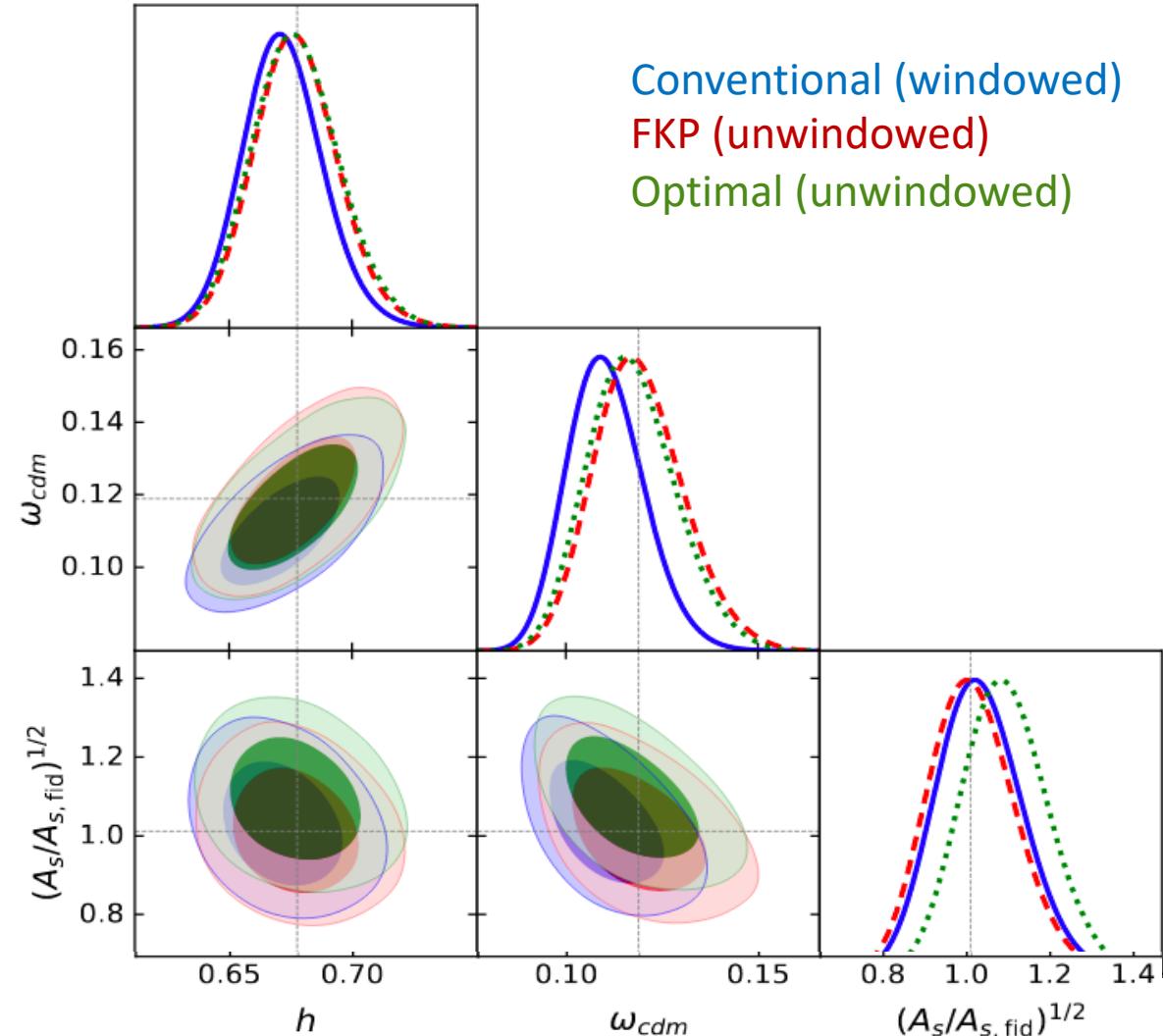
# Is this useful?

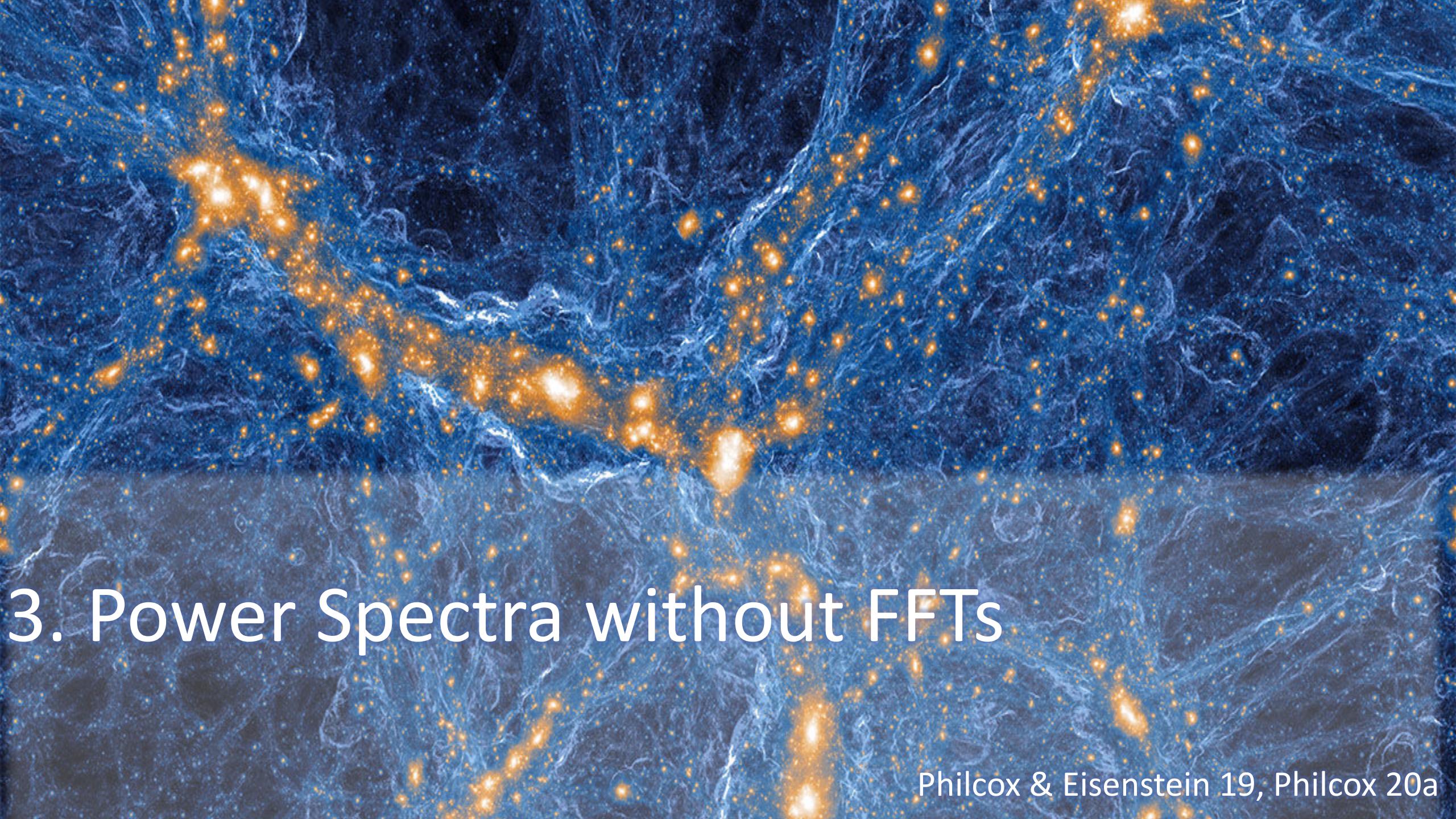
- Benefits:

- No window-convolution
- **Optimal** error-bars if Gaussian
- Less gridding
- Less shot-noise

Best for **small, dense, anisotropic surveys, and large-scale modes**

- Especially useful for  $f_{NL}$  and the **bispectrum**





### 3. Power Spectra without FFTs

Philcox & Eisenstein 19, Philcox 20a

# Configuration-Space $P(k)$ Estimators

- $P(k)$  usually estimated using Fast Fourier Transforms

$$\hat{P}(k) = \int \frac{d\Omega_k}{4\pi} |\text{FFT}[\delta](\mathbf{k})|^2$$

- **Complexity:**  $\mathcal{O}(N_g \log N_g)$  for  $N_g$  grid points

- Small scales need **large  $N_g$**  ⇒ **slow computation and high memory usage!**

$$\text{Time} \propto k_{\max} \log k_{\max}$$

- 2PCF estimated by **counting pairs of particles** with  $\mathcal{O}(N^2)$  complexity

$$\xi^a = \int d\mathbf{r}_1 d\mathbf{r}_2 \delta(\mathbf{r}_1)\delta(\mathbf{r}_2)\Theta^a(|\mathbf{r}_1 - \mathbf{r}_2|) = \sum_{i \neq j} w_i w_j \Theta^a(|\mathbf{r}_i - \mathbf{r}_j|)$$

Weights  
Binning function  
Sum over galaxies

- This is **fast** on small scales!

$$\text{Time} \propto N n R_{\max}^3$$

# Configuration-Space $P(k)$ Estimators

- We can do the same for  $P(k)$ :

$$P(k) \propto \int \frac{d\Omega_k}{4\pi} \sum_{i \neq j} w_i w_j e^{-i\mathbf{k} \cdot (\mathbf{r}_i - \mathbf{r}_j)} = \sum_{i \neq j} w_i w_j j_0(k|\mathbf{r}_i - \mathbf{r}_j|)$$

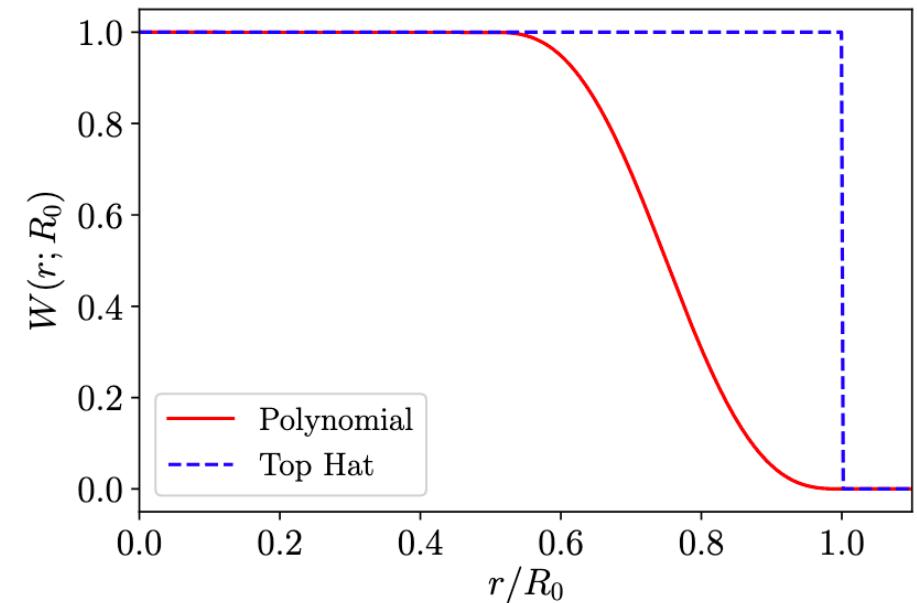
Weights      0<sup>th</sup> order Bessel function  
Sum over galaxies

- But we need to sum over all  $N^2$  pairs of galaxies in the survey!

- Only sum up to some maximum radius  $R_0$ , via a smooth function  $W(r; R_0)$

$$P(k; R_0) \propto \sum_{i \neq j} w_i w_j j_0(k|\mathbf{r}_i - \mathbf{r}_j|) W(|\mathbf{r}_i - \mathbf{r}_j|; R_0)$$

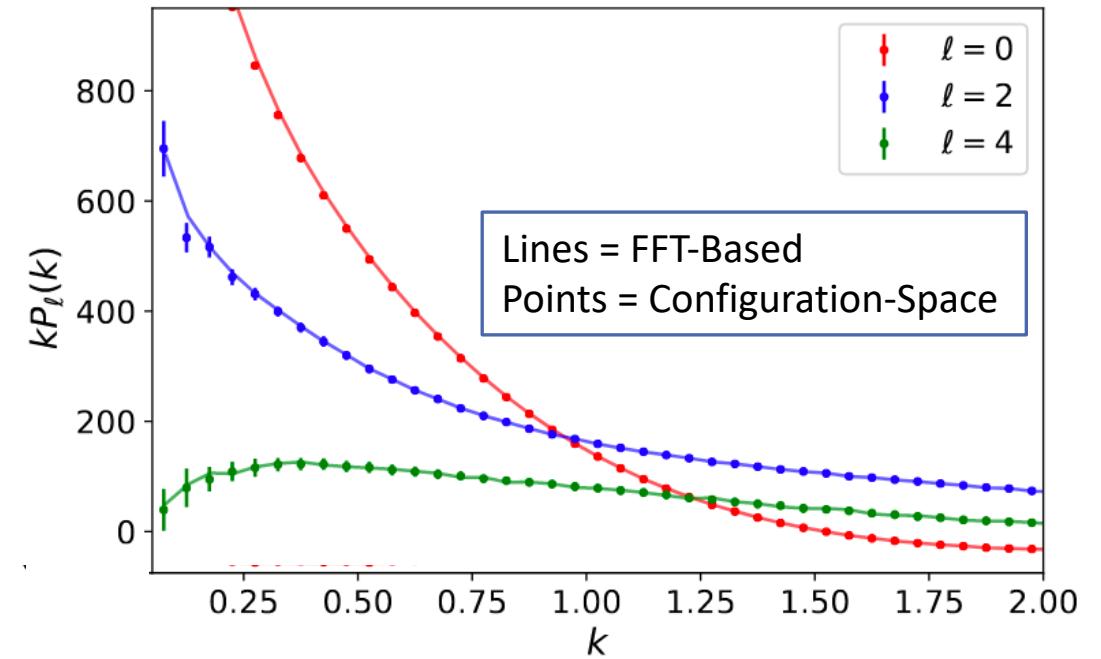
Time  $\propto N n R_0^3 \propto k_{\min}^{-3}$



# Configuration-Space $P_\ell(k)$ Estimators

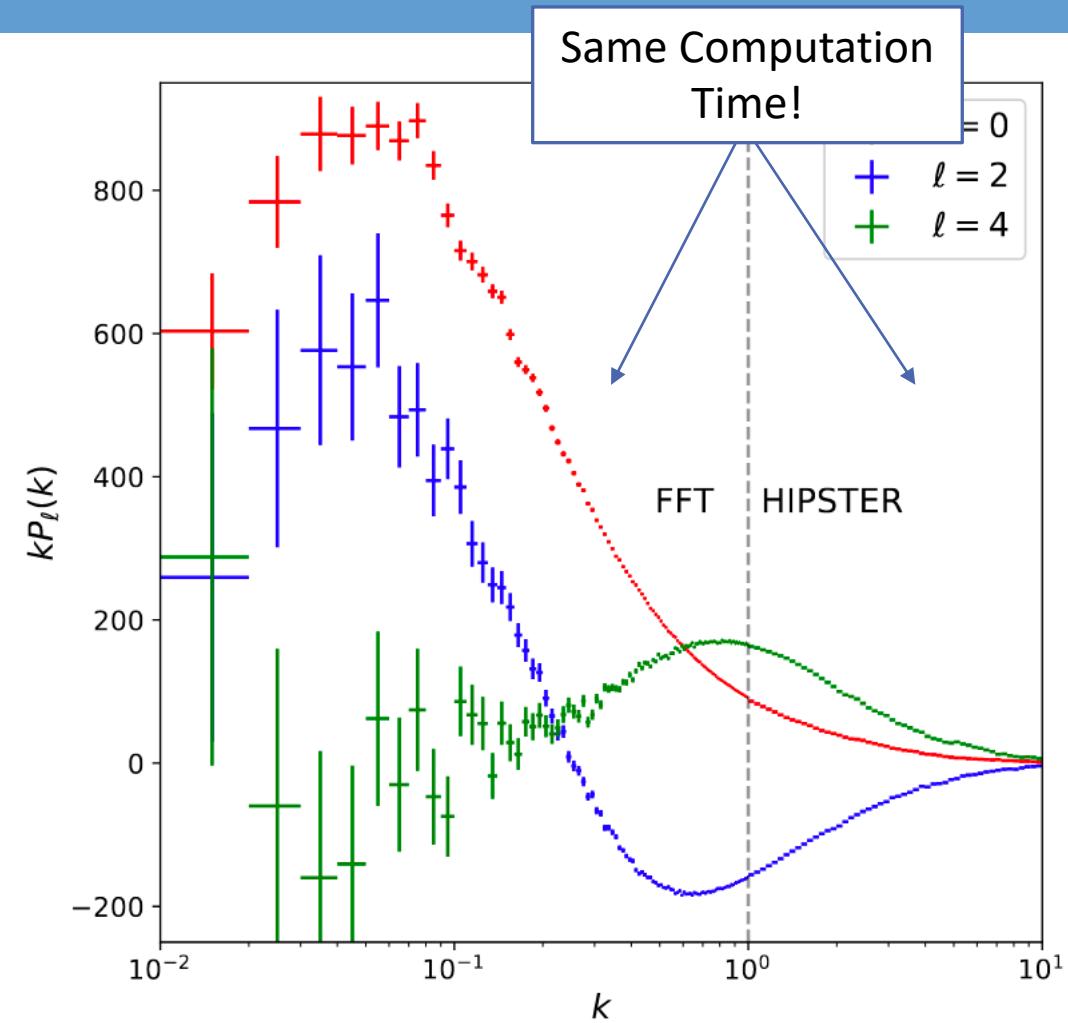
## Benefits

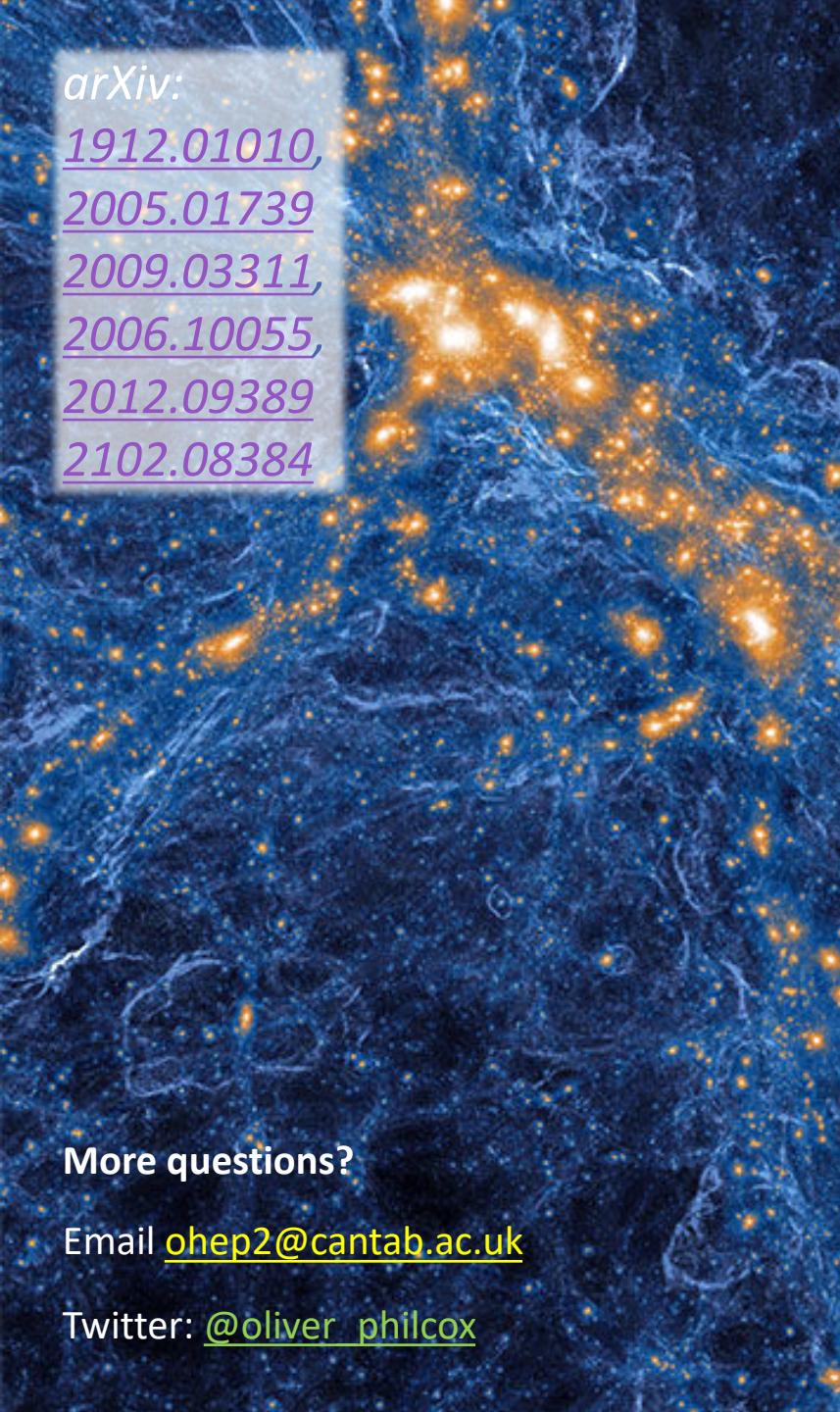
- **Speed**
  - Time scales as  $k_{\min}^{-3}$
- **Memory**
  - No storage of large FFT grids
- **Aliasing**
  - No gridding!
- **Shot-noise**
  - Removes self-counts -> Poissonian shot-noise!
- **Window function**
  - Can remove survey window, just as for 2PCF



# Configuration-Space $P_\ell(k)$ Estimators

- Implemented in the **HIPSTER** code
- Combine with FFT-based treatments:
  - FFTs are fastest on **large** scales (time  $\sim k_{\max} \log k_{\max}$ )
  - HIPSTER is fastest on **small** scales (time  $\sim k_{\min}^{-3}$ )
- Can be similarly applied to **bispectra**
  - Time  $\propto Nn^2R_0^6 \propto k_{\min}^{-6}$
  - Same scaling with number density as for  $P(k)$ !





arXiv:

[1912.01010](https://arxiv.org/abs/1912.01010),  
[2005.01739](https://arxiv.org/abs/2005.01739),  
[2009.03311](https://arxiv.org/abs/2009.03311),  
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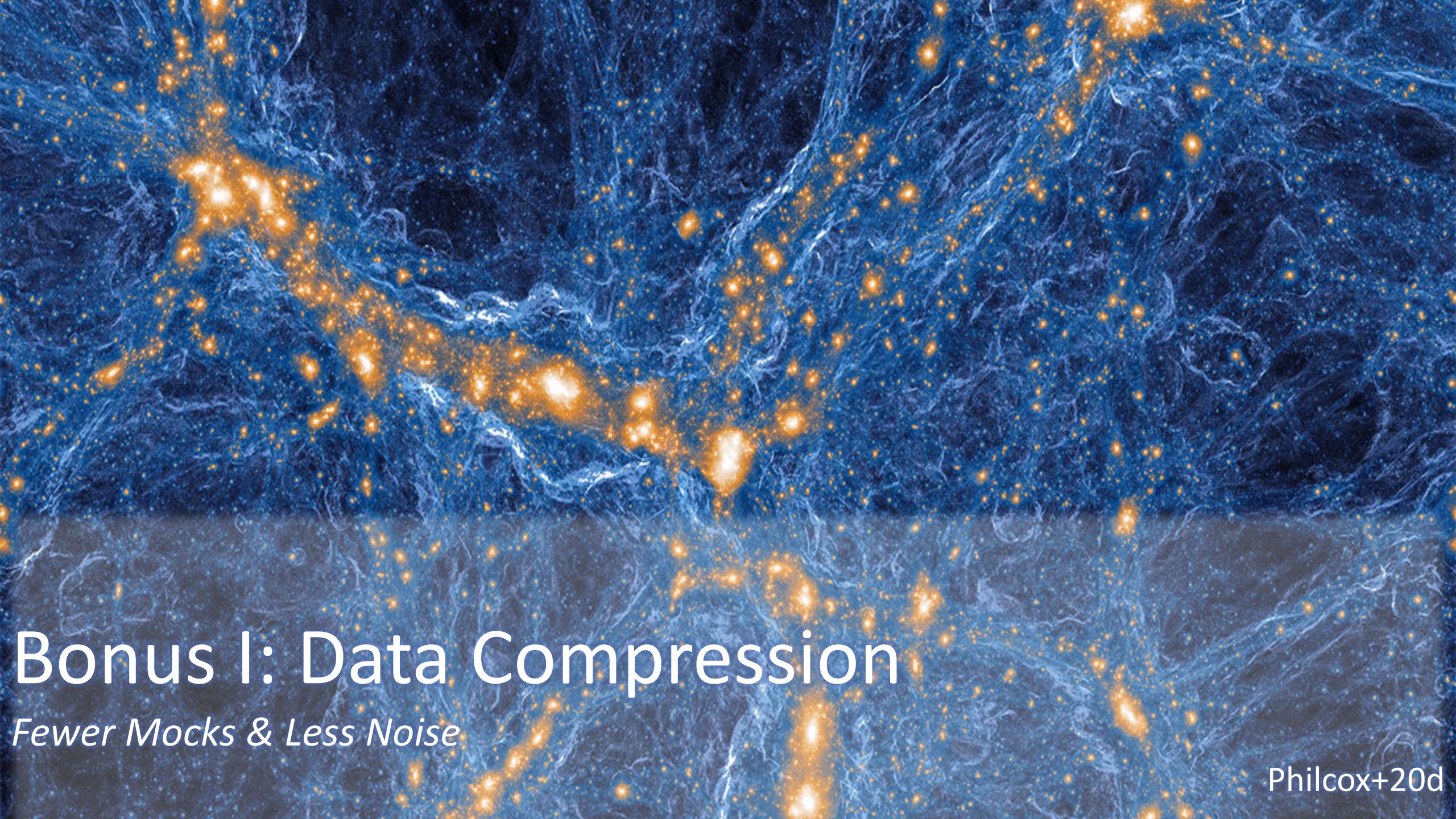
Twitter: [@oliver\\_philcox](https://twitter.com/@oliver_philcox)

# Conclusions

- We're not finished with the galaxy power spectrum yet!
  
- Recent updates include:
  - More accurate lines-of-sight
  - Closer to optimal large-scale  $P_\ell(k)$  estimation
  - Faster small-scale computation *without* FFTs
  - (*Powerful analysis-specific data compression*)

*Coming soon:*

- Estimating the bispectrum and beyond!



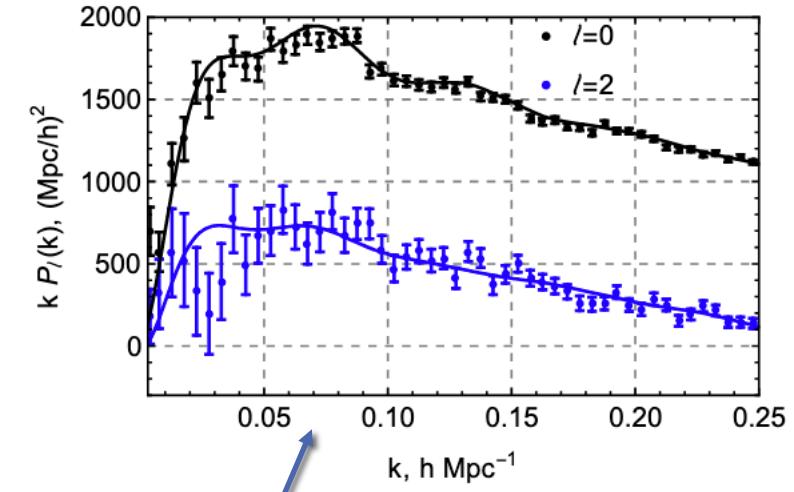
# Bonus I: Data Compression

*Fewer Mocks & Less Noise*

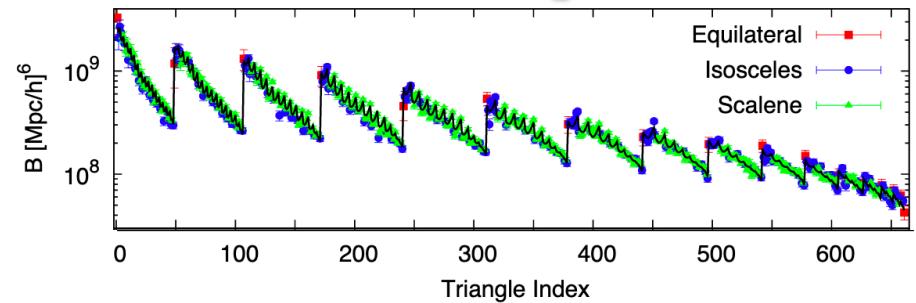
Philcox+20d

# The Curse of Dimensionality

- $P_\ell(k)$  is **high-dimensional**, e.g.;
  - BOSS has  $\sim 100$  bins
  - Only use these to measure  $\sim 10$  parameters
- Conventional likelihoods use a **sample covariance**
  - Need  $N_{mocks} > N_{bins}$  to invert
  - Too few mocks  $\Rightarrow$  **parameter shifts or error inflation**
- We should **compress** our data!



Power Spectrum      Bispectrum

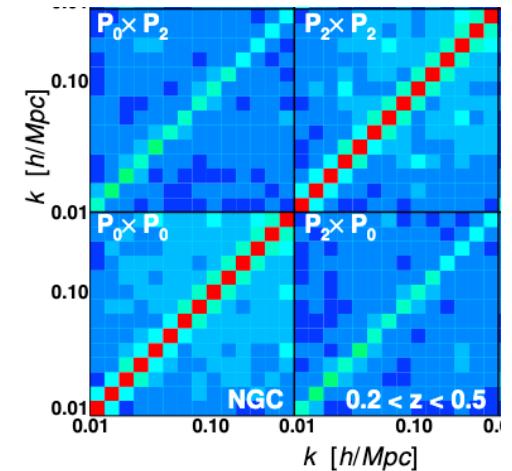


# Data Compression via PCA

Beutler+16

- A canonical approach: [e.g. Scoccimarro 2000]
  - Compute the theoretical **covariance** matrix
  - Perform a **Principal Component Analysis**
  - Project the data onto the first few components
- This chooses the basis vectors that contribute most to the signal-to-noise
- Signal-to-noise isn't everything!

Power Spectrum Covariance



PCA

$$P(k) \approx \sum_i a_i W_i(k)$$

Basis Vectors

Coefficients

See also MOPED: Heavens+00, Alsing+18 , KL: Tegmark+97

# Data Compression via Subspace Projection

## New\* approach

- Draw sets of parameters from the **priors**
- Compute the **theory model** at each point
- Perform a **Singular Value Decomposition** on the **noise-weighted** samples
- Use these **basis vectors** to perform the compression

*Parameters used in the analysis*

$$\theta = \{\omega_{\text{cdm}}, A_s/A_{s,\text{fid}}, h, \dots\} \times \{b_1, b_2, b_{G_2}, b_4, c_{s,0}, c_{s,2}, P_{\text{shot}}\}$$

Picks out directions contributing most to the **log-likelihood**

\*somewhat inspired by gravitational wave analyses [e.g. Roulet+19]

$$X_a(\theta) \equiv \sum_{ab} C_{ab}^{-1/2} [P_b(\theta) - \bar{P}_b]$$

*Covariance Estimate*      *Theory Model*

$$X_{ia} = \sum_{\alpha} U_{i\alpha} D_{\alpha} V_{\alpha a}$$

*Basis Vectors*

$$X_a^{(i)} \approx \sum_{\alpha=1}^{N_{\text{SV}}} c_{\alpha}^{(i)} V_{\alpha a}$$

*Subspace Coefficients*

# Data Compression via Subspace Projection

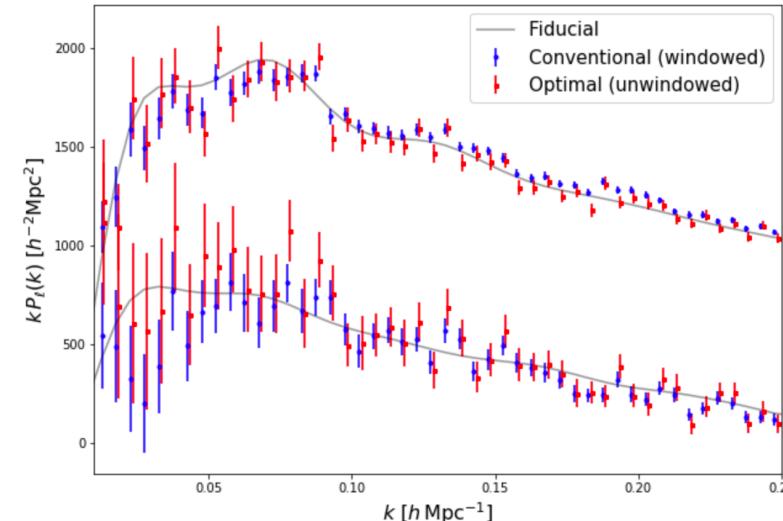
- This is the **best** linear compression for a **specific** analysis
- Set the **number** of basis vectors **robustly**
- Estimate coefficients **optimally**

For BOSS 10-parameter analysis:

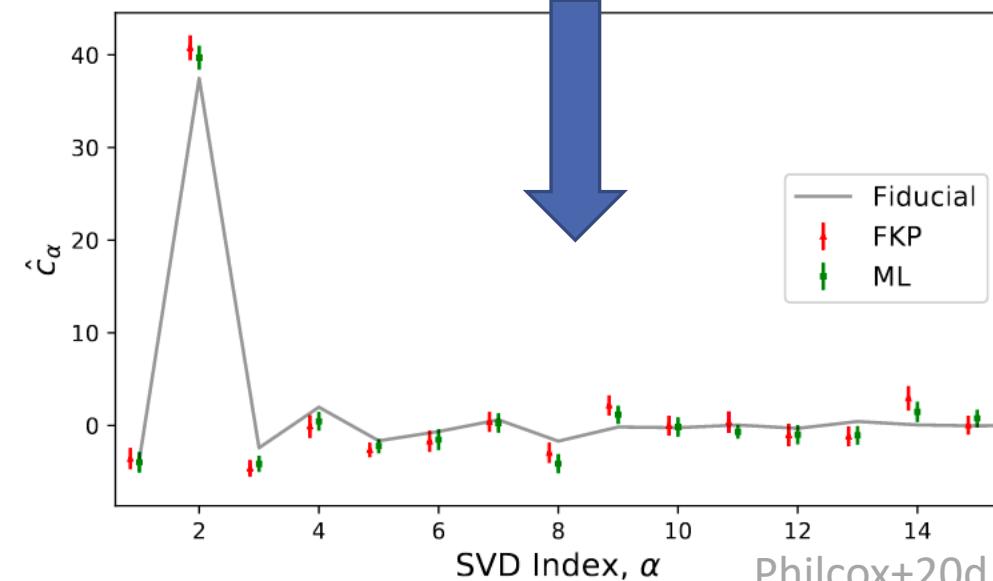
- 100-bin  $P(k)$  ----> 12 subspace coefficients
- 2135-bin  $B(k_1, k_2)$  ----> 8 subspace coefficients

Applicable to **any** analysis given:

1. Theory Model
2. Parameter Priors
3. Rough Covariance Estimate



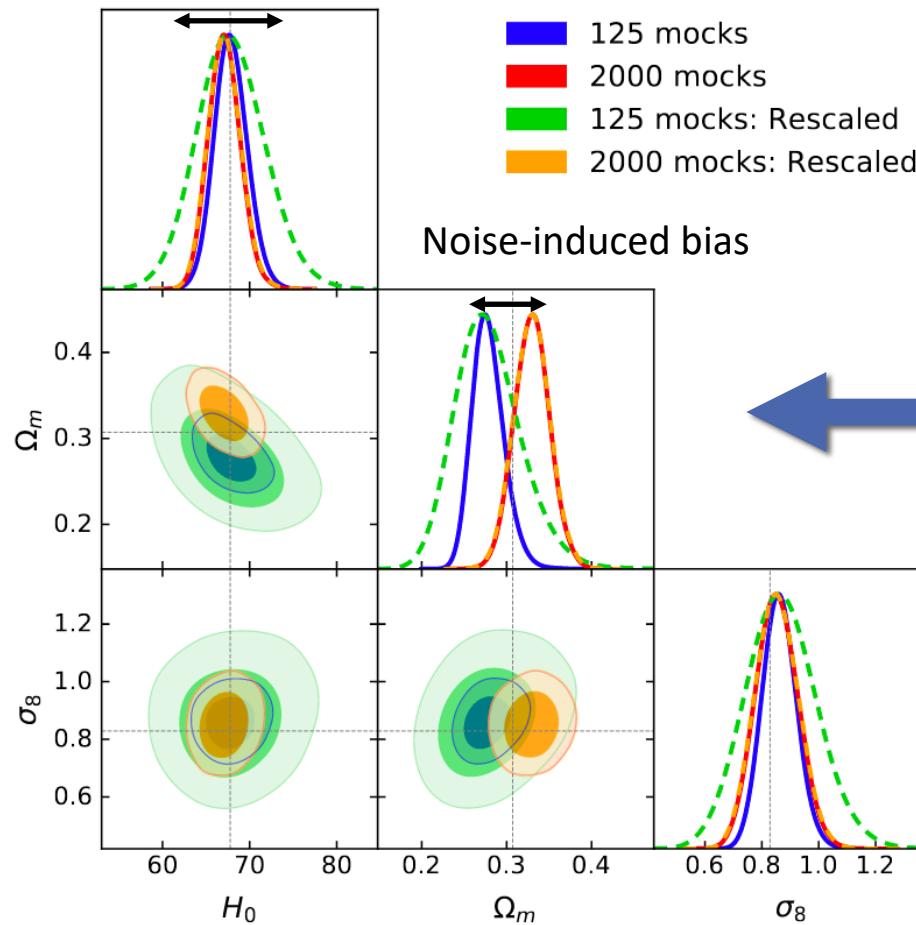
Power Spectrum



Subspace Coefficients

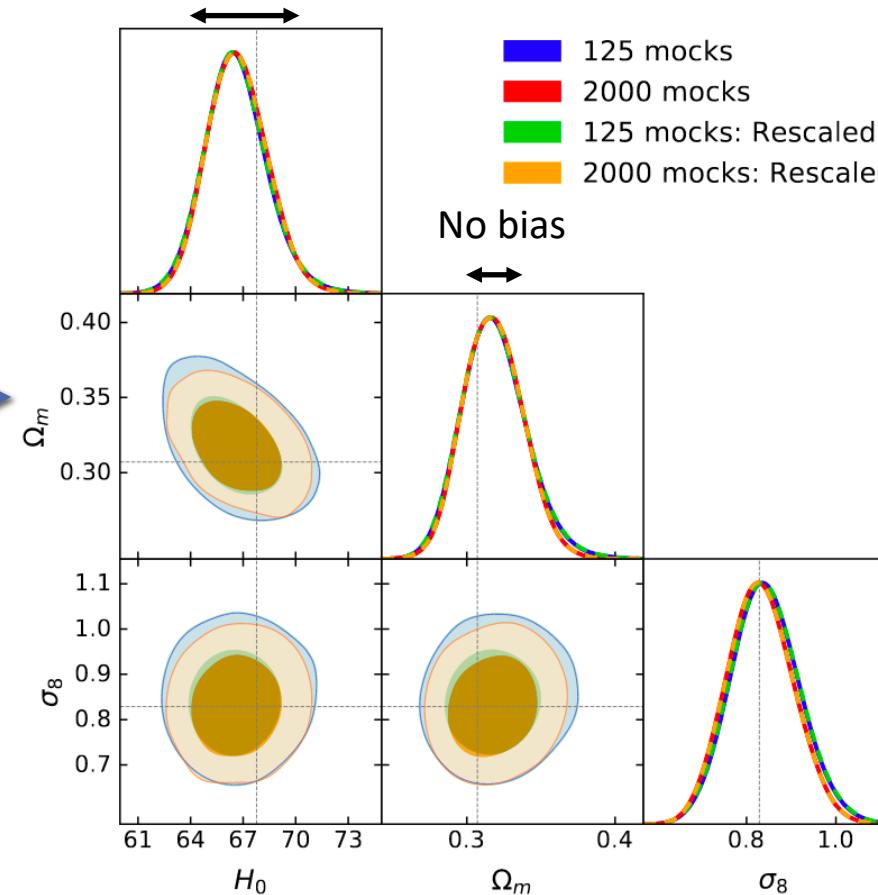
# Too Few Mocks $\rightarrow$ Parameter Biases

Error Inflation

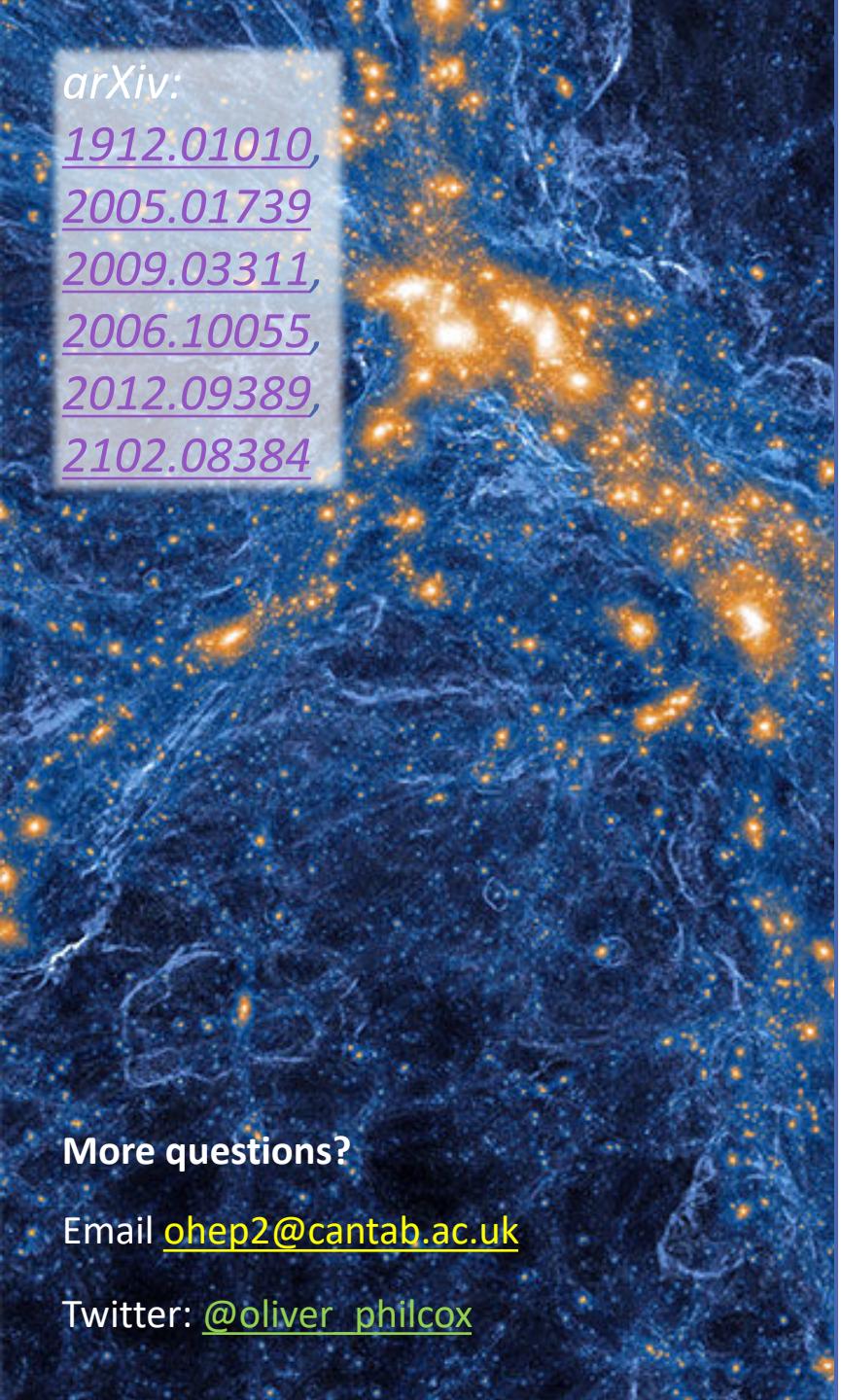


(a) 96-bin Power Spectrum

No Inflation



(c) 12 Subspace Coefficients



arXiv:

[1912.01010](https://arxiv.org/abs/1912.01010),  
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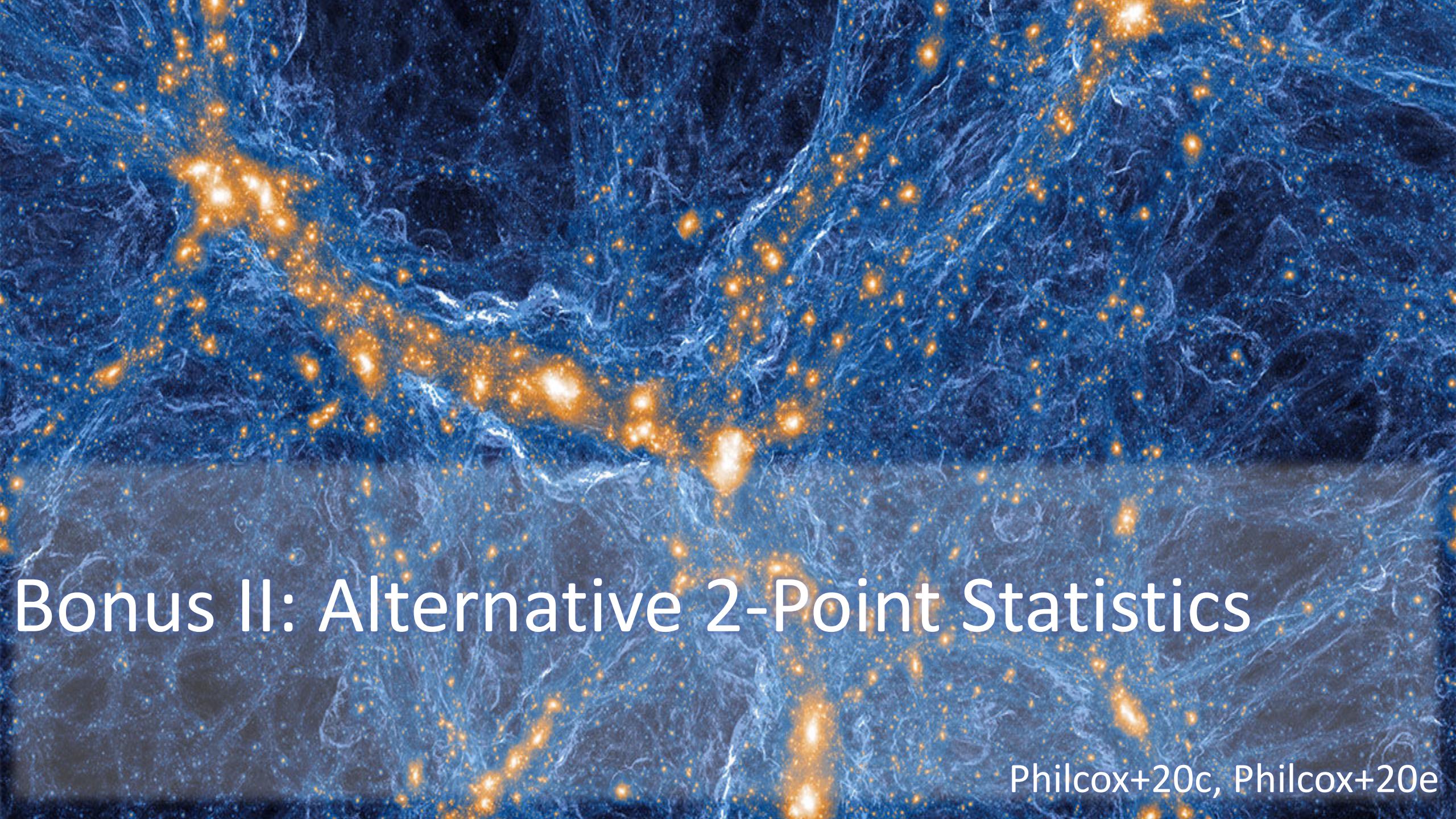
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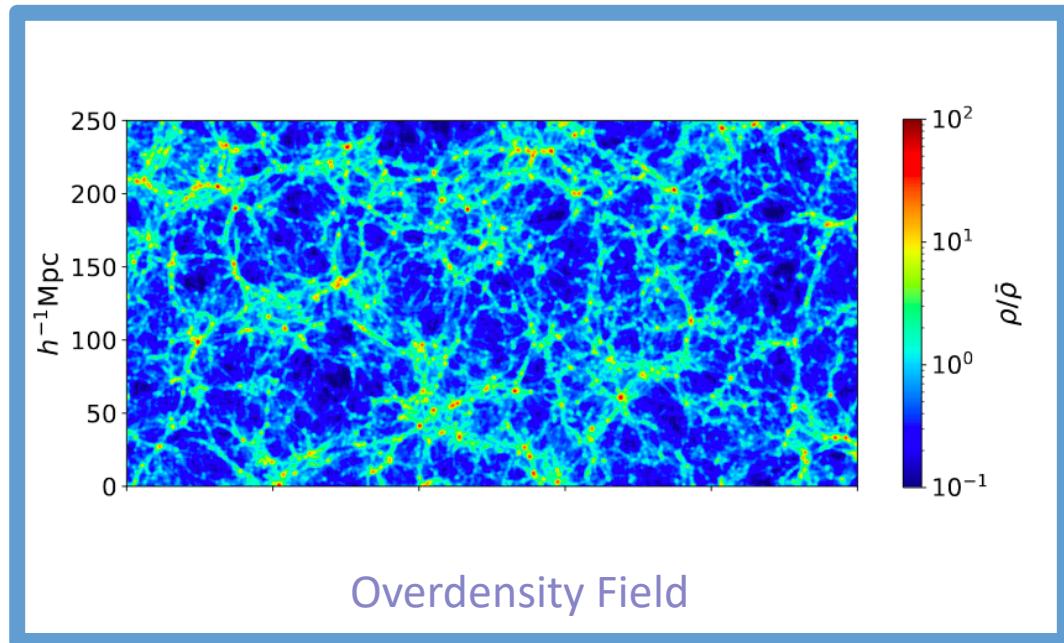


## Bonus II: Alternative 2-Point Statistics

Philcox+20c, Philcox+20e

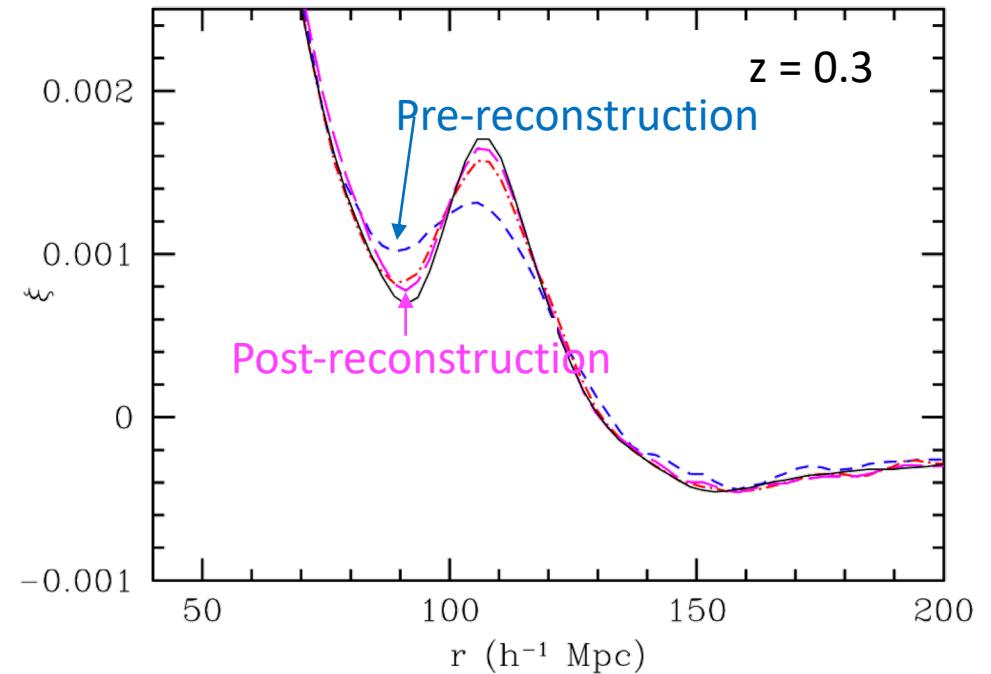
# Beyond the Density Field

- What should we compute the two-point function of?
  - For a **Gaussian** universe, the power-spectrum of galaxy overdensity contains **all** the information
- The Universe is **not** Gaussian:
  - Information **cascades** to the higher-point functions
  - **Low-density regions** carry a lot of cosmological information, and contribute little to  $\delta$  [e.g. Pisani+19]
- Can use a **transformed** field, e.g.:
  - Reconstructed Density Fields [e.g. Eisenstein+07]
  - Log-normal Transforms [Neyrinck+09, Wang+11]
  - Gaussianized Density Fields [Weinberg 92, Neyrinck+17]
  - Marked Density Fields [Stoyan 84, White 16, Massara+20]



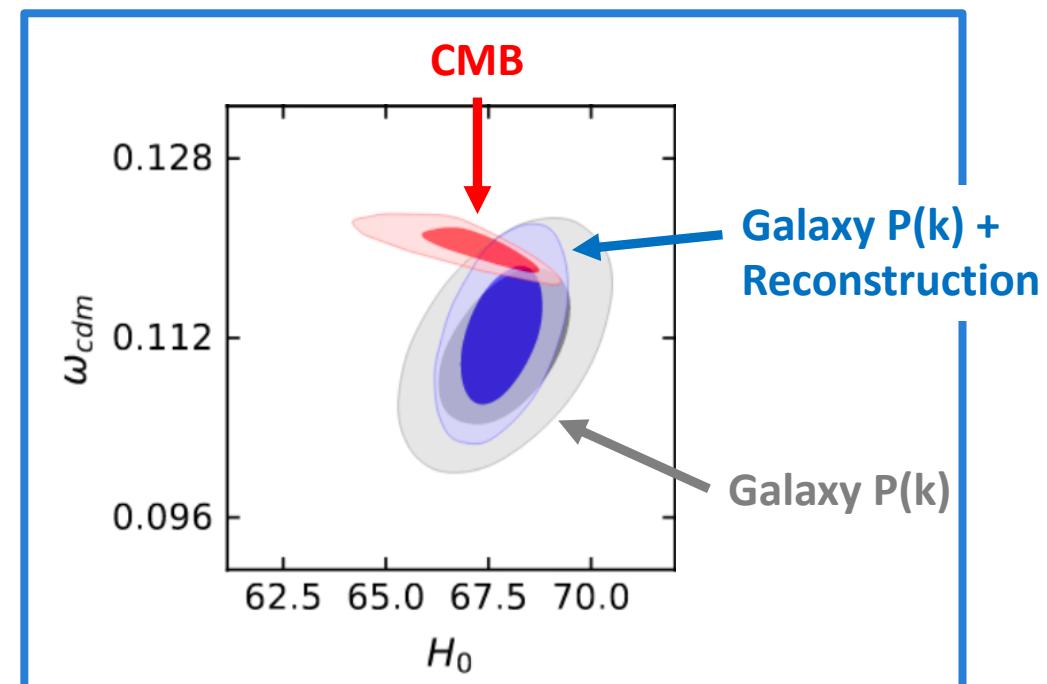
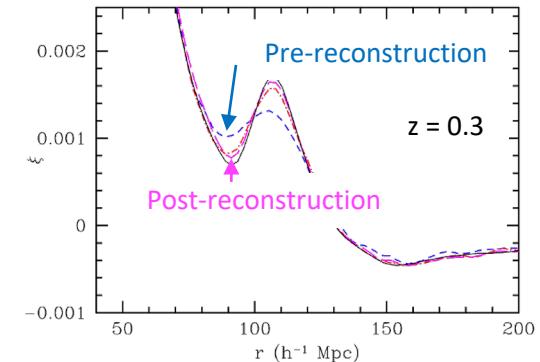
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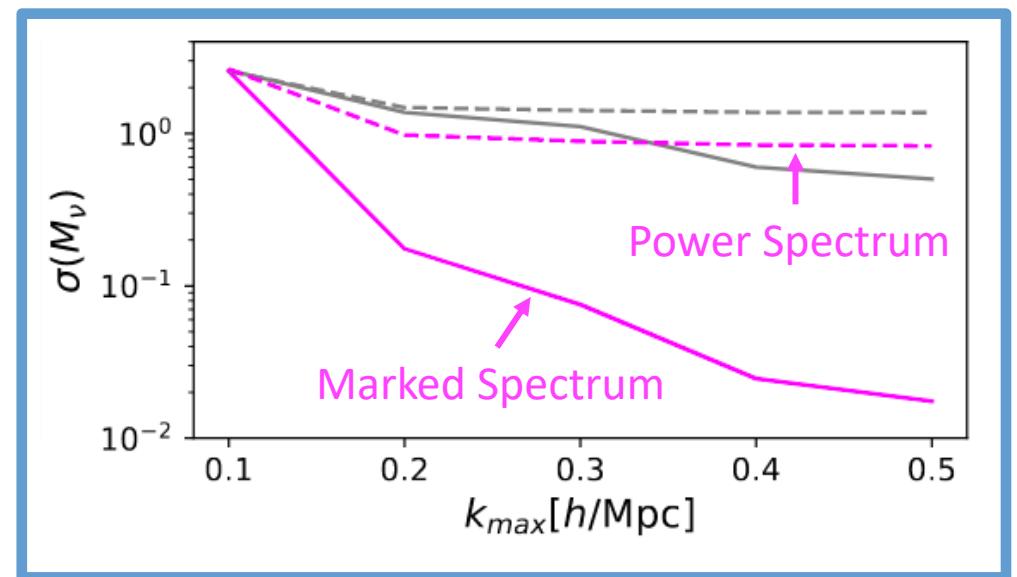
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Fisher Matrix Constraints on Neutrino Mass

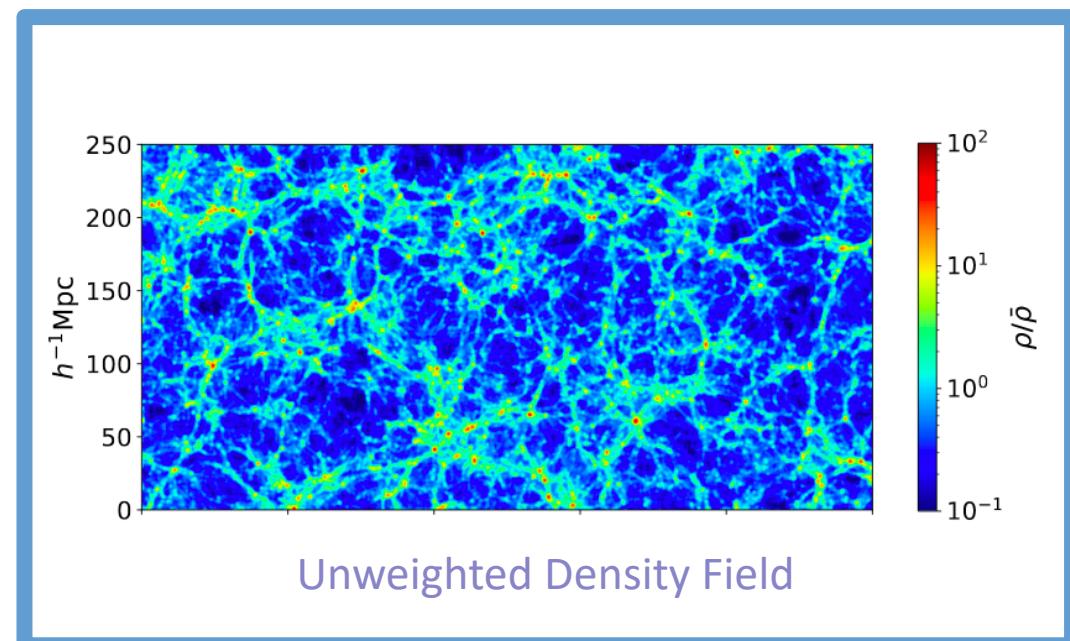
# The Marked Density Field

- Define a new density field by weighting by the **mark**

$$m(\mathbf{x}) = \left( \frac{1 + \delta_s}{1 + \delta_s + \delta_R(\mathbf{x})} \right)^p$$

$$\rho_M(\mathbf{x}) = m(\mathbf{x})n(\mathbf{x}) = m(\mathbf{x})\bar{n}[1 + \delta(\mathbf{x})]$$

depending on **smoothed** overdensity  $\delta_R(\mathbf{x})$



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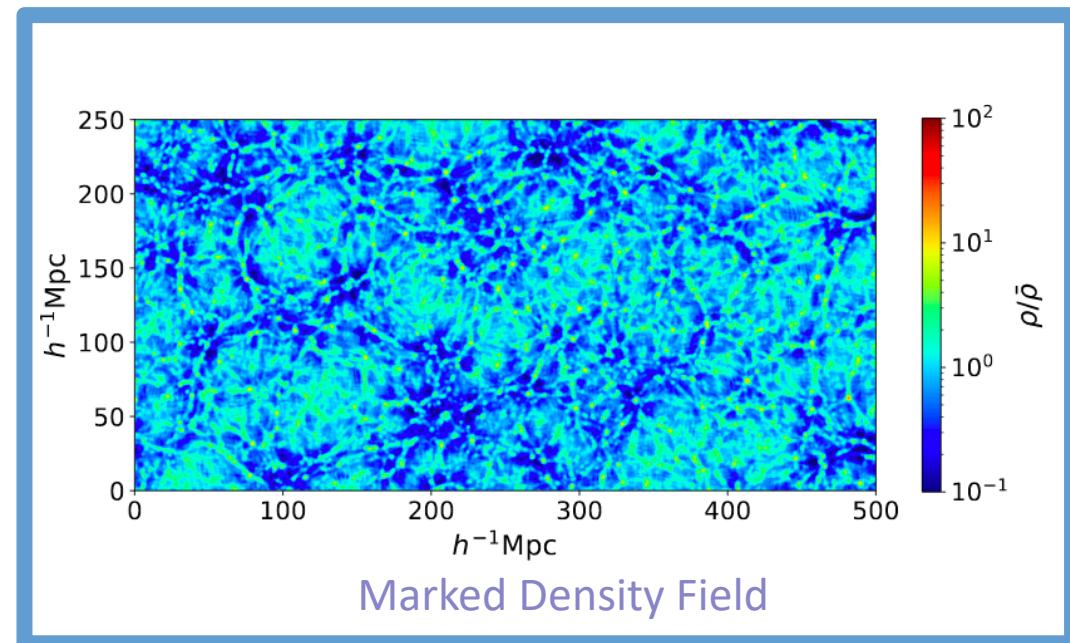
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depending on **smoothed** overdensity  $\delta_R(\mathbf{x})$

- Significantly enhances constraints on:

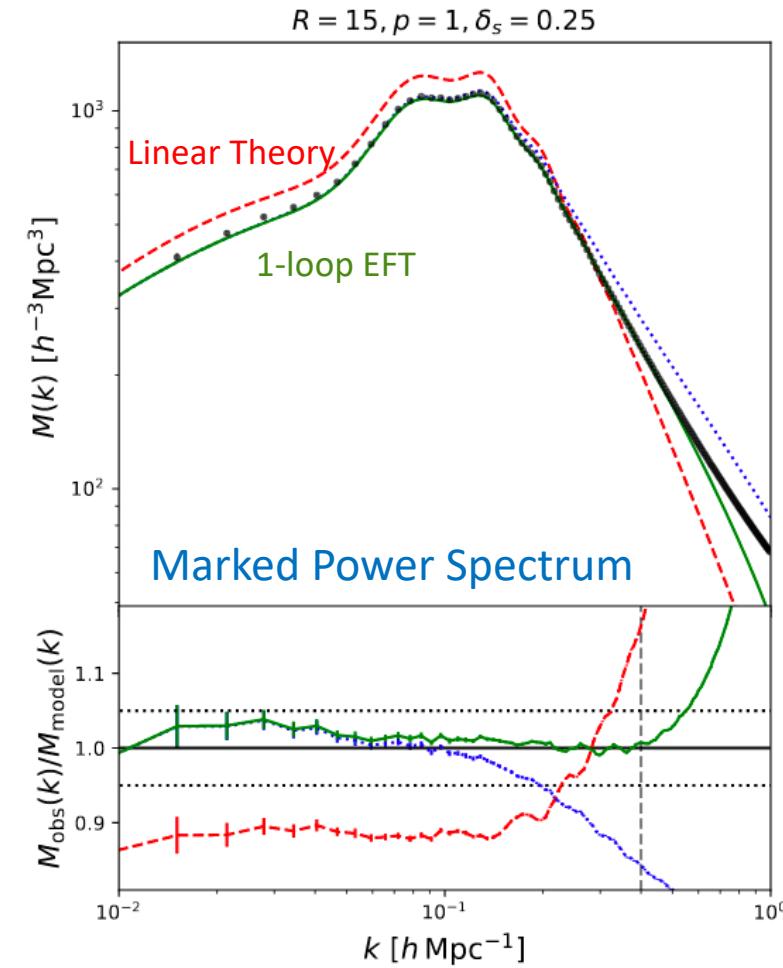
- **Neutrino masses** [Massara+20]

- **Modified gravity** [White 16]



# The Marked Density Field

- Can we **model** the marked spectrum?
  - Yes! Using **Effective Field Theory**
- Can we **understand** the impressive information content?
  - The mark couples **small-scale** non-Gaussianities to **large-scale** modes
  - So we find more neutrino information at low- $k$ !
- But:
  - Modelling is **difficult** at low- $z$
  - Is it still useful for galaxies?

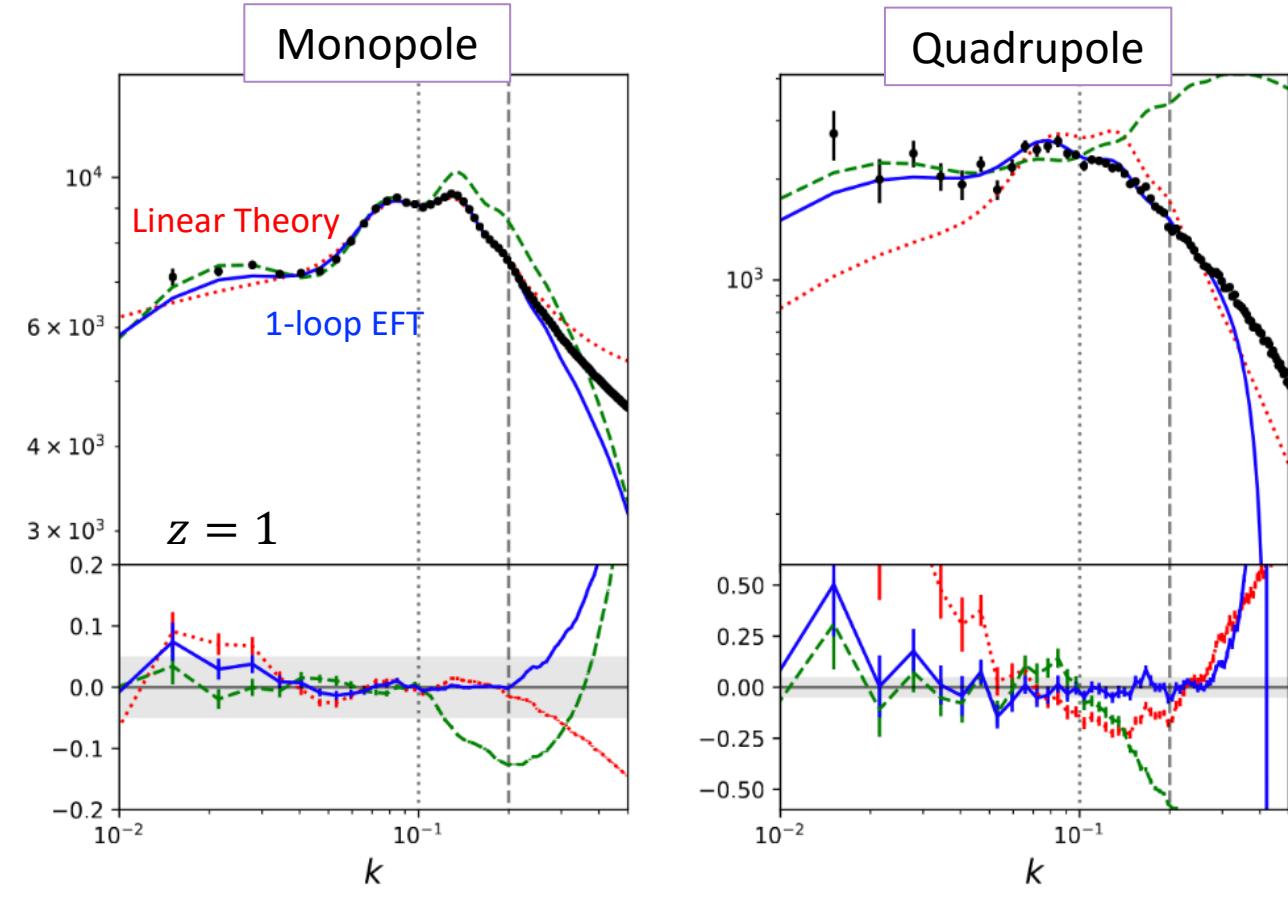


Matter at  $z = 1$

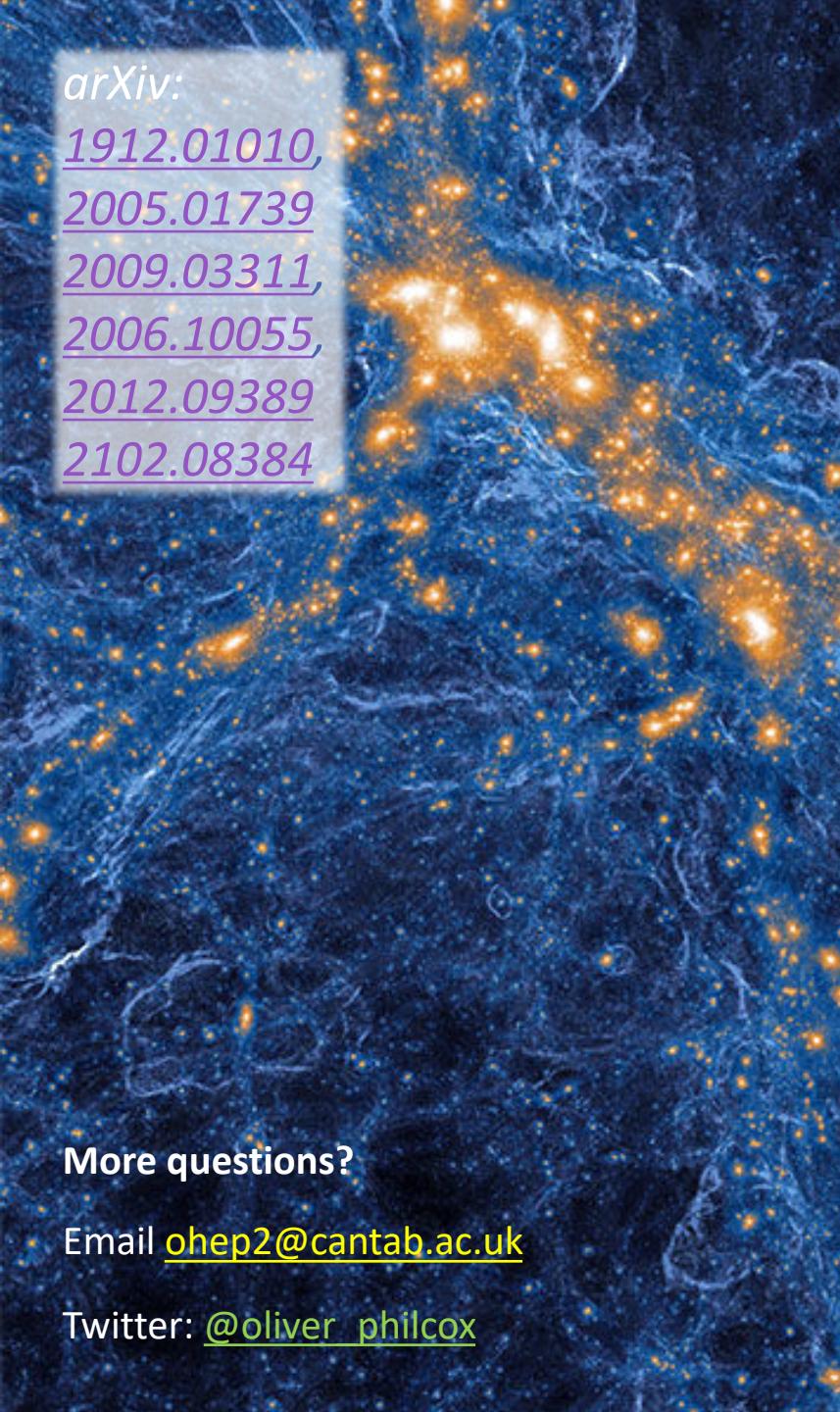
Massara+20, Philcox+20ce

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Galaxies at  $z = 1$



arXiv:

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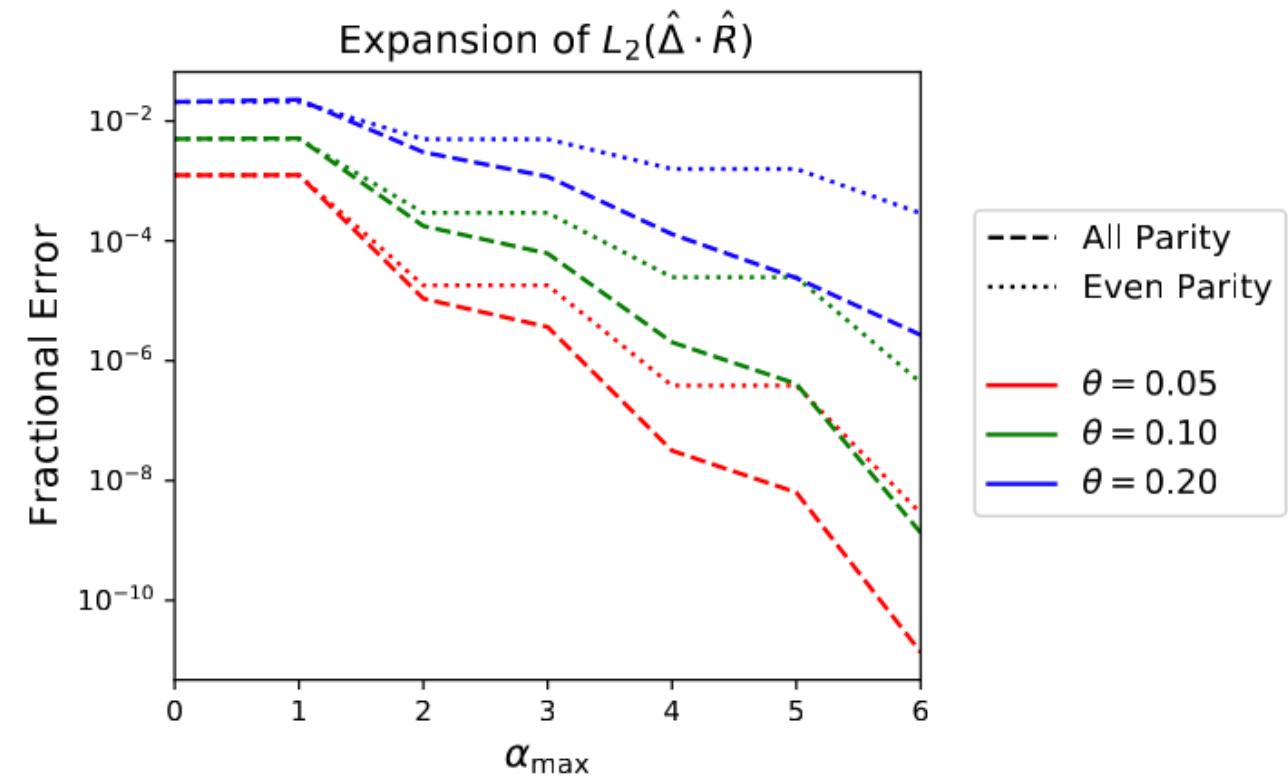
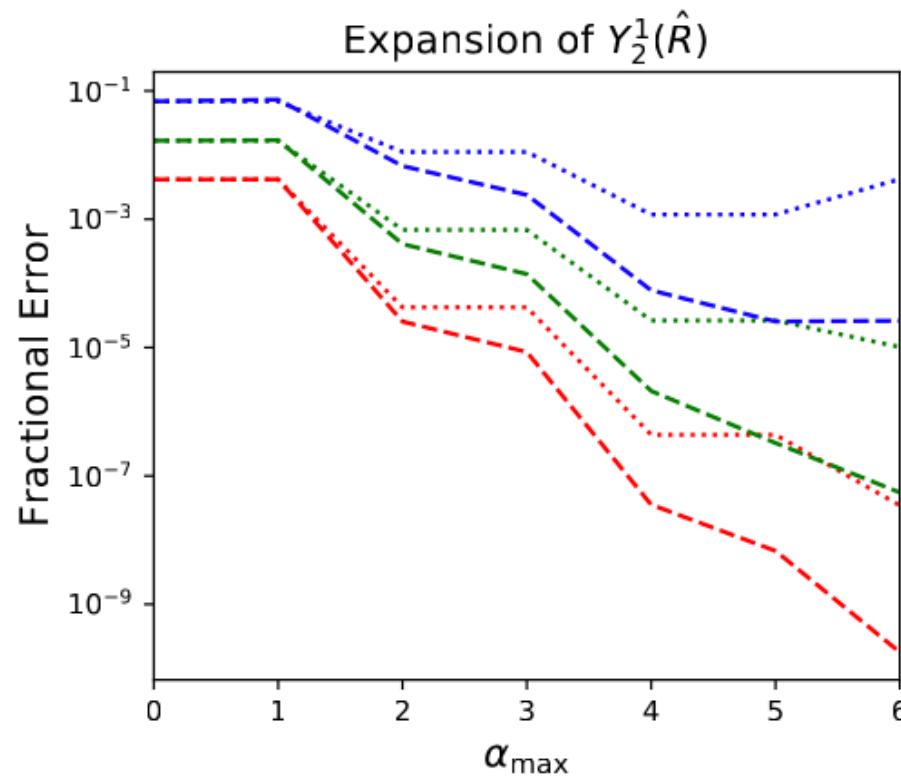
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  - Statistics **beyond** the density field

## Coming soon:

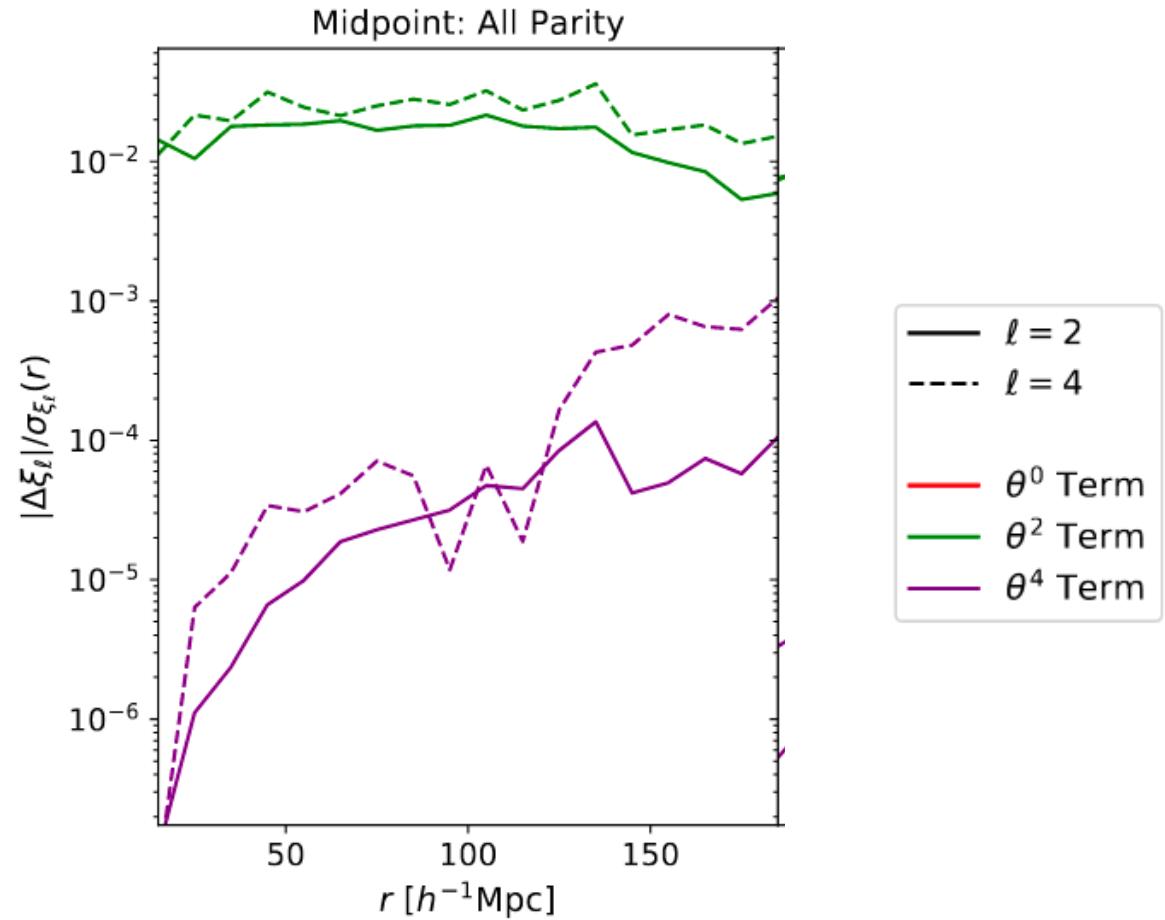
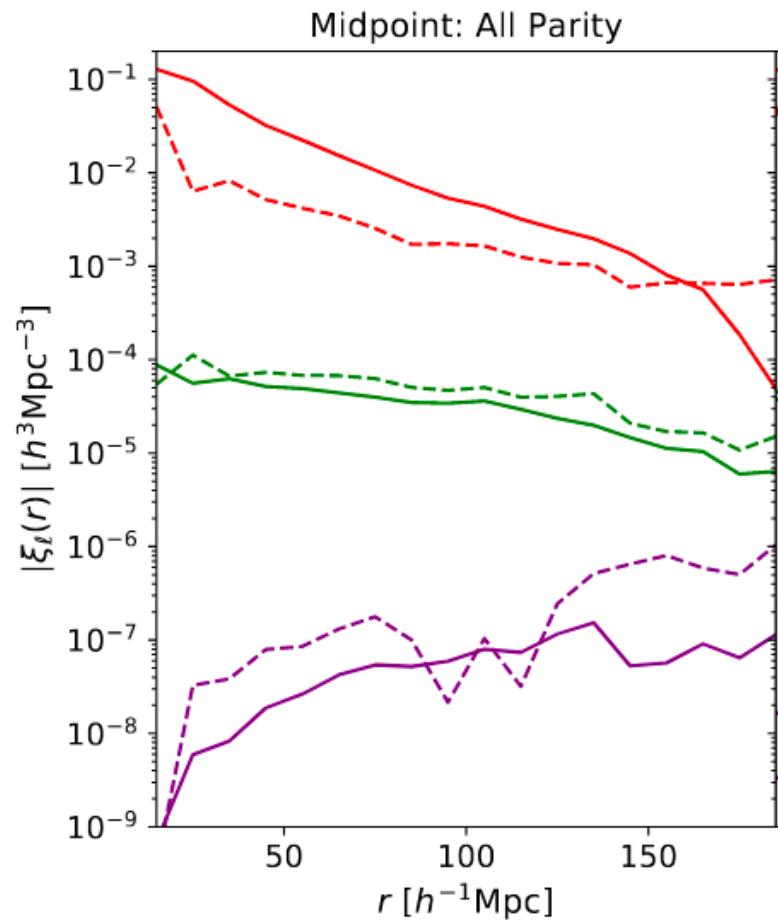
- Estimating the bispectrum and beyond!

# Shift Theorem Convergence

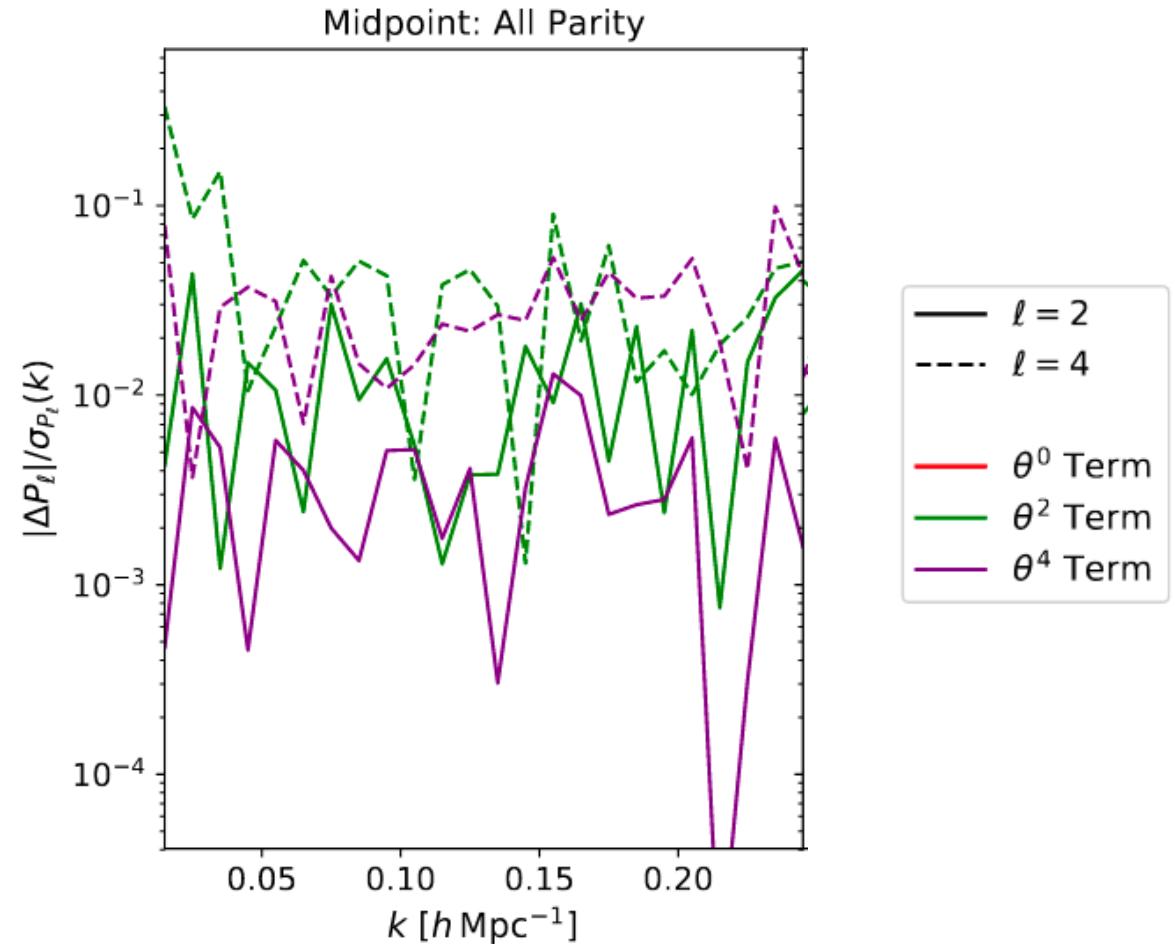
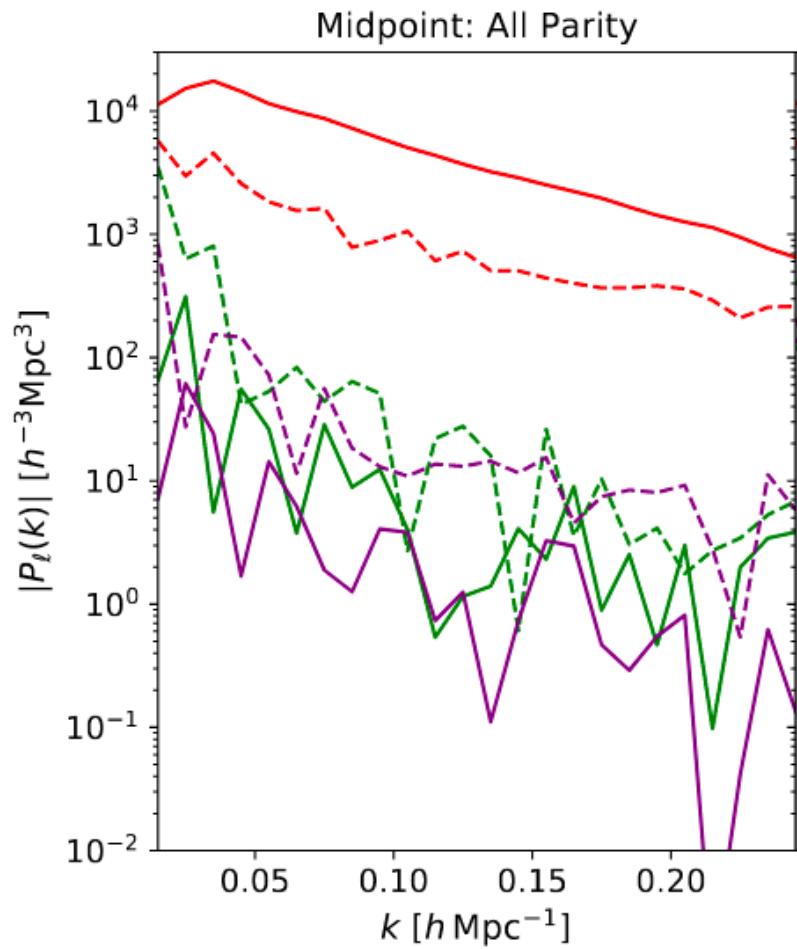


- - - All Parity
- · - Even Parity
- $\theta = 0.05$
- $\theta = 0.10$
- $\theta = 0.20$

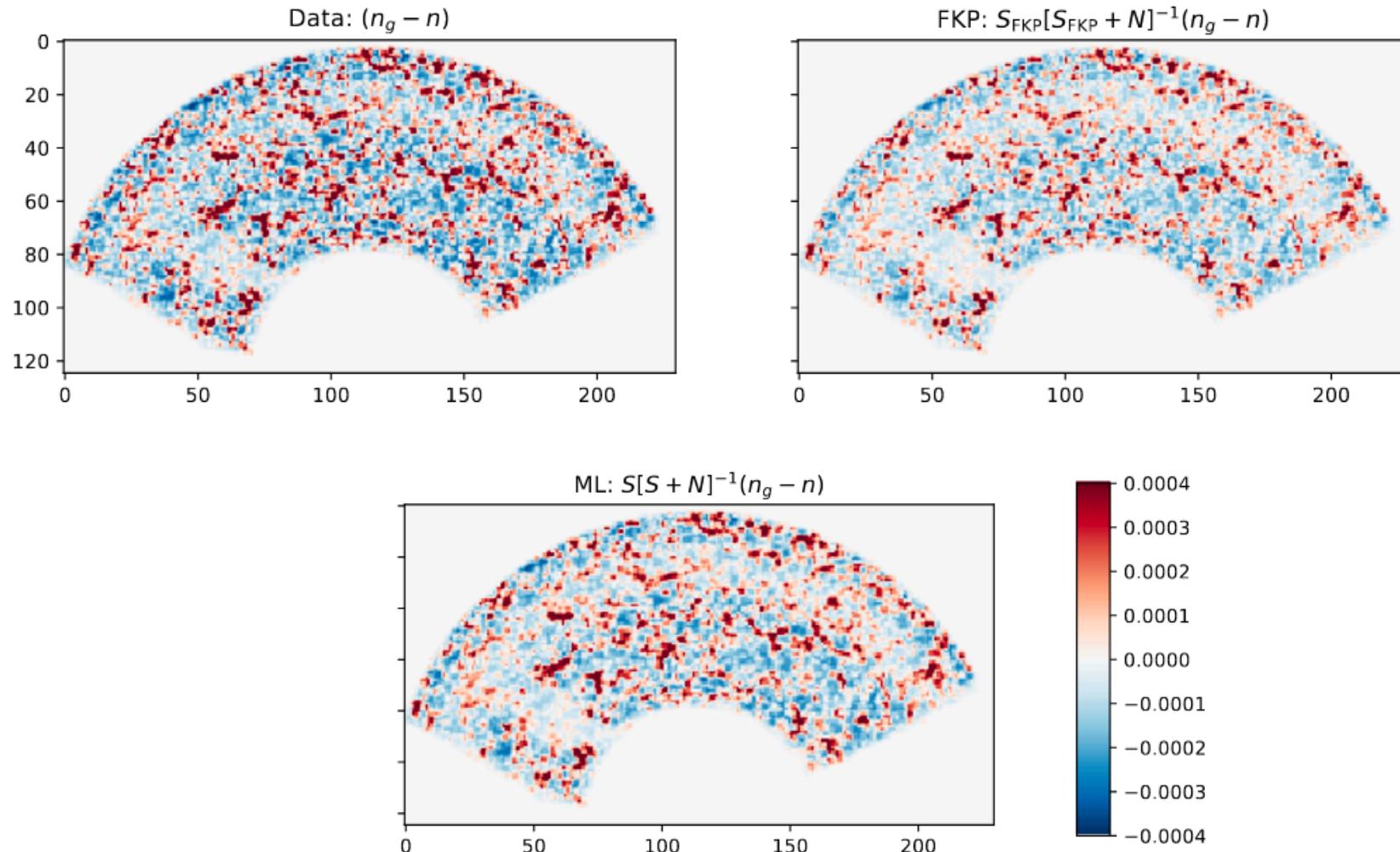
# 2PCF Wide-Angle Effects



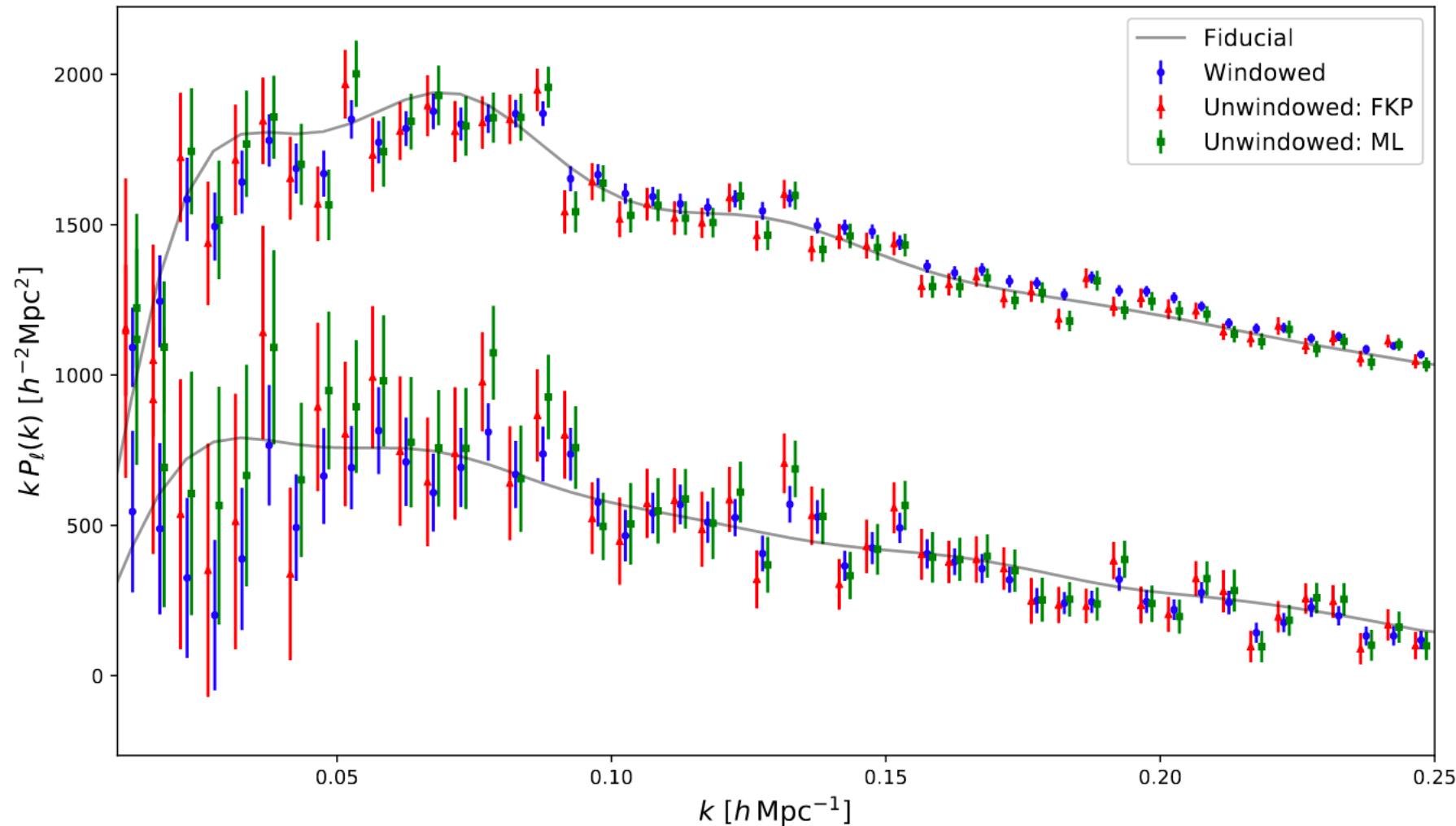
# $P(k)$ Wide-Angle Effects



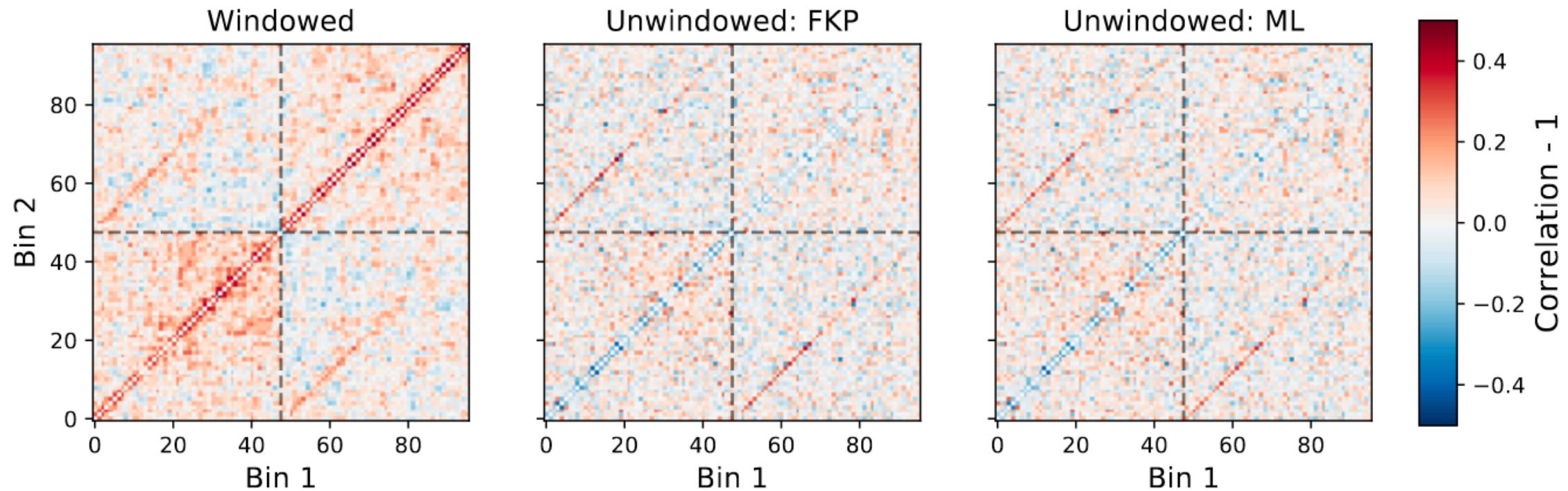
# Optimal Estimators: Filtering



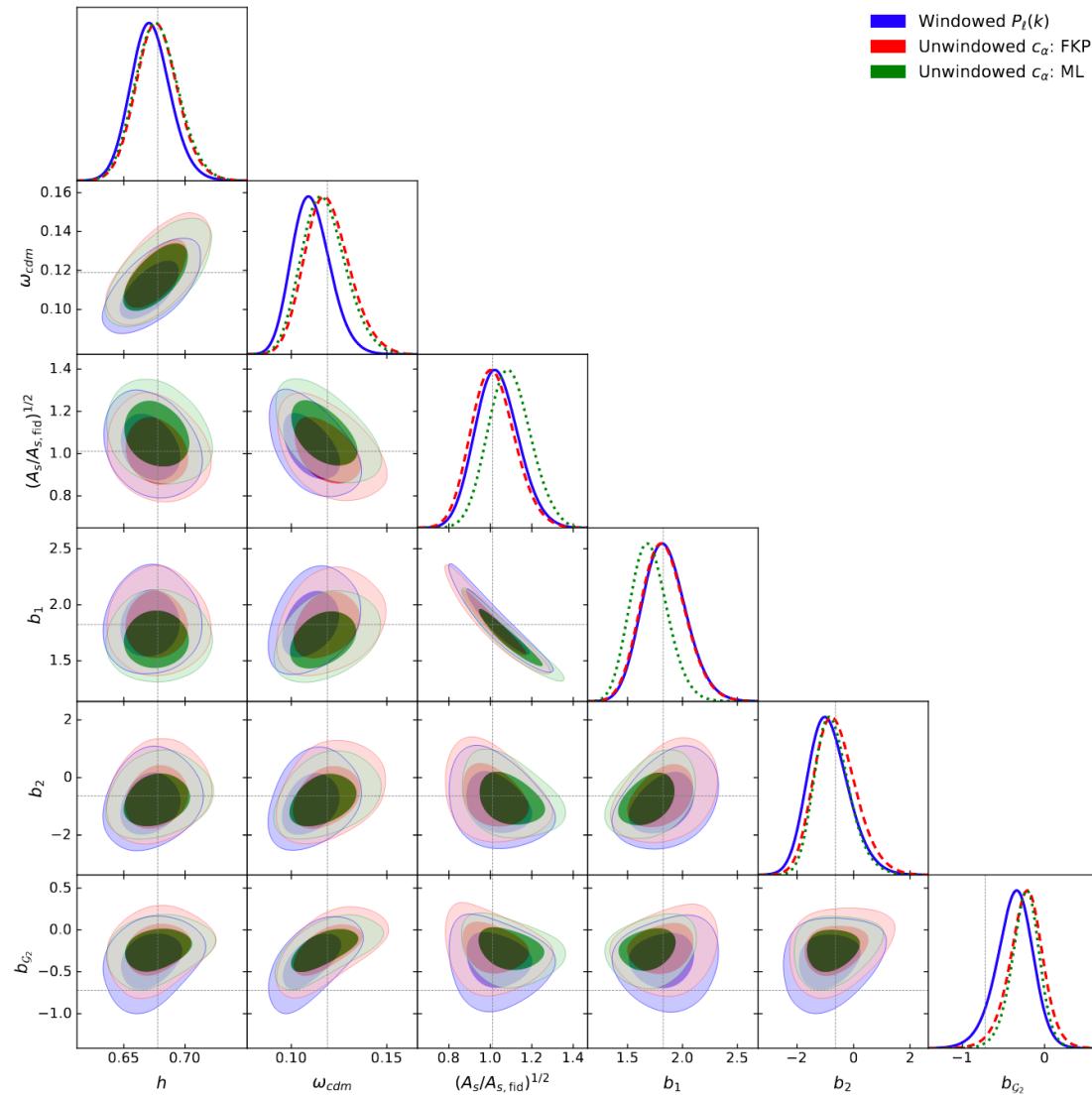
# Optimal Estimators: Spectra



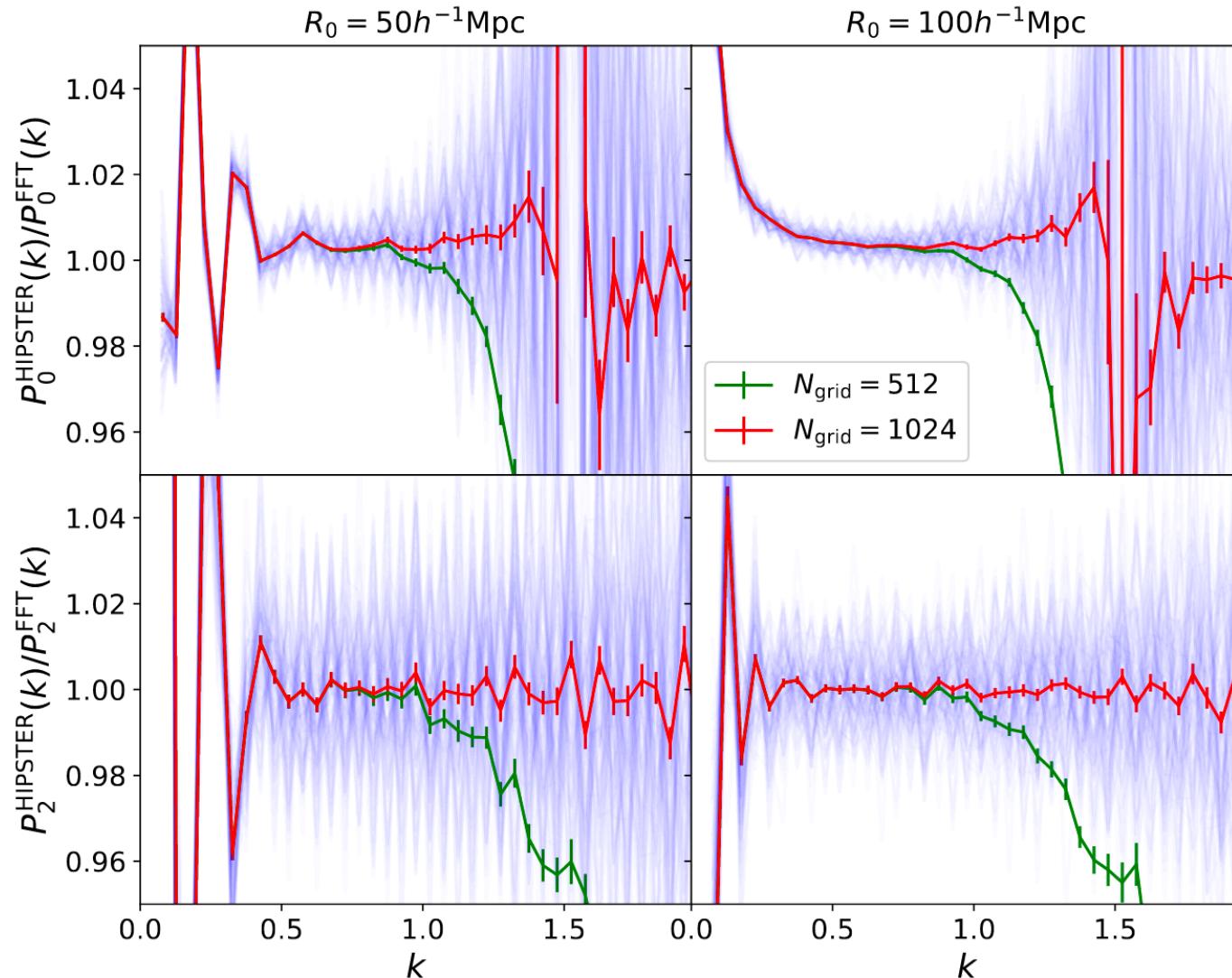
# Optimal Estimators: Covariance



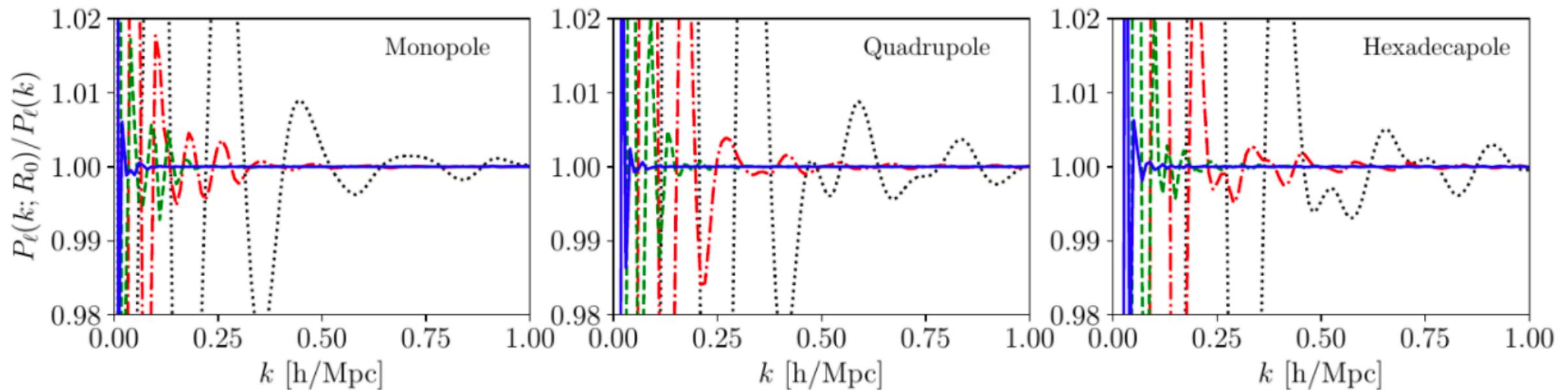
# Optimal Estimators: Results



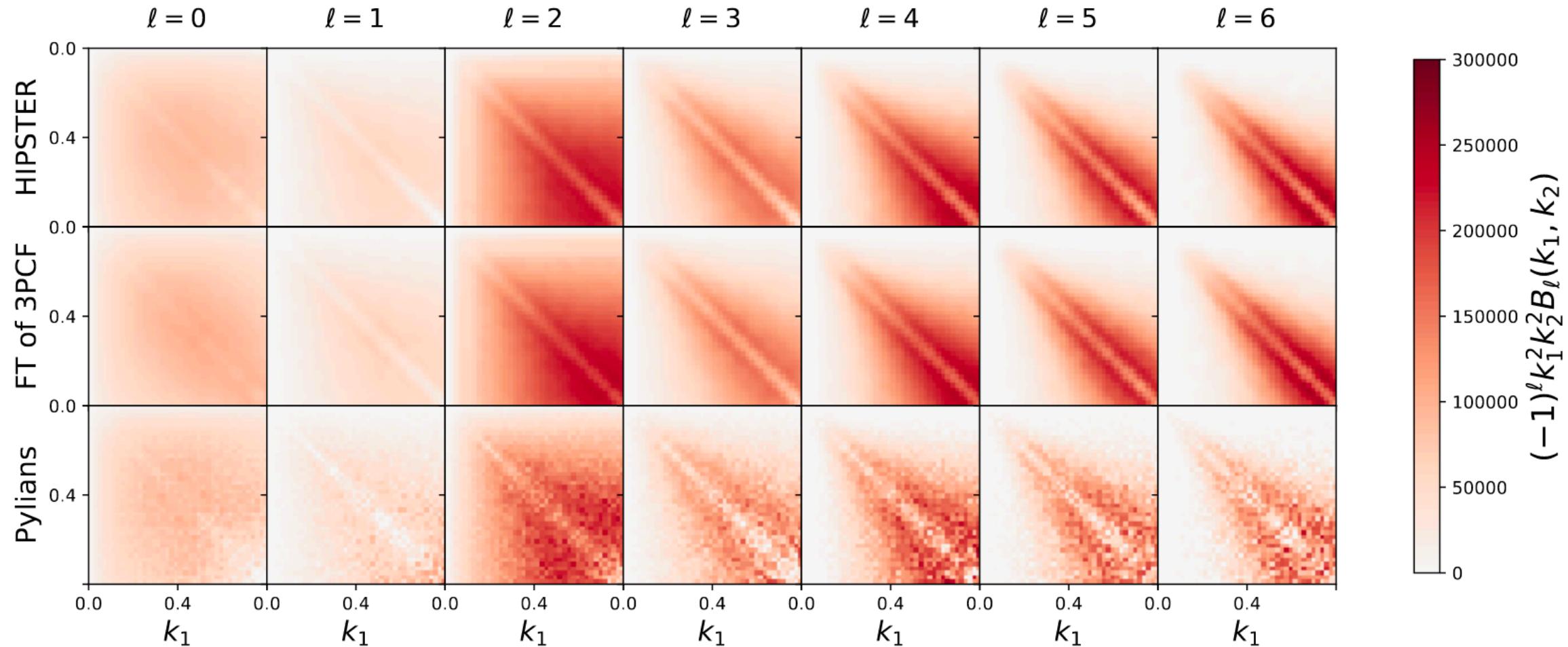
# HIPSTER: Accuracy



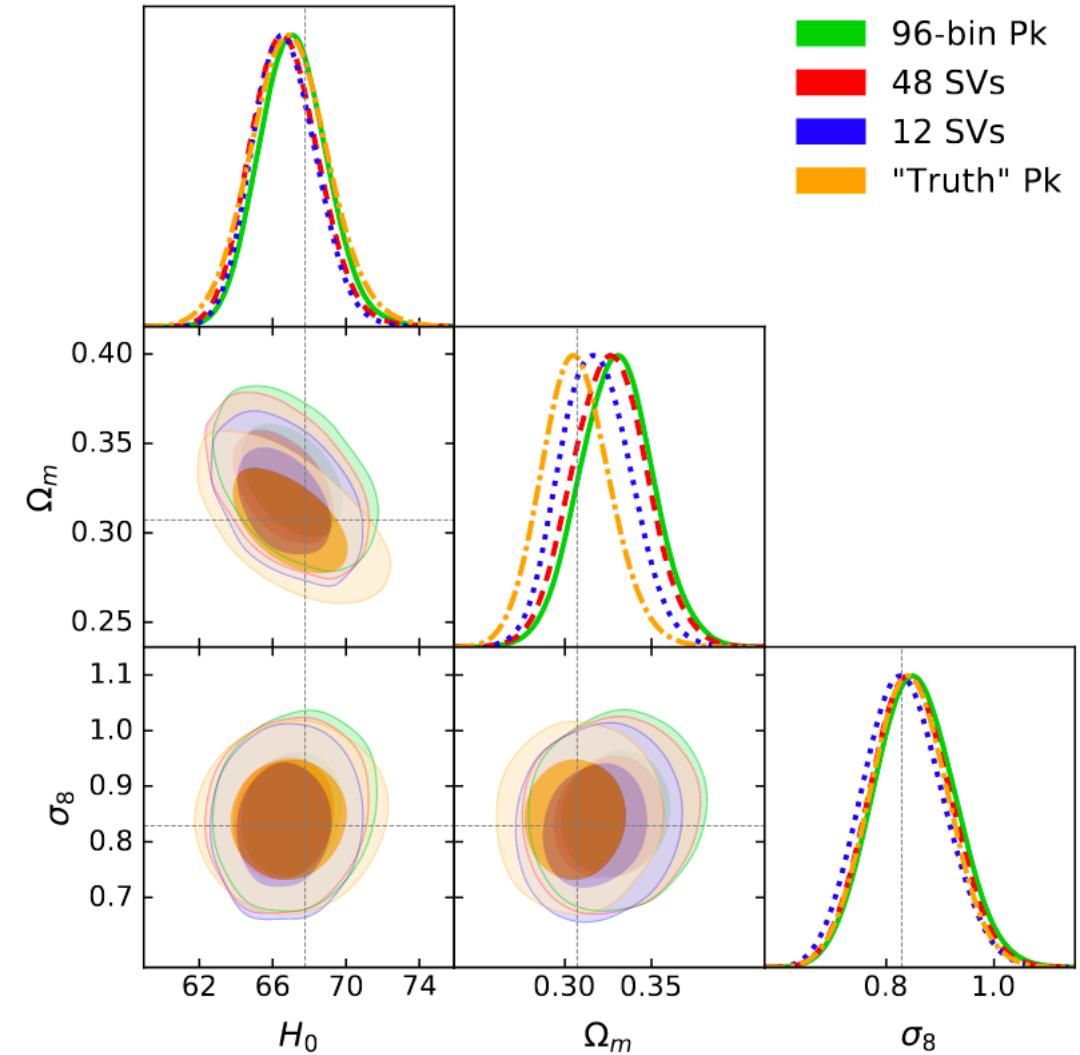
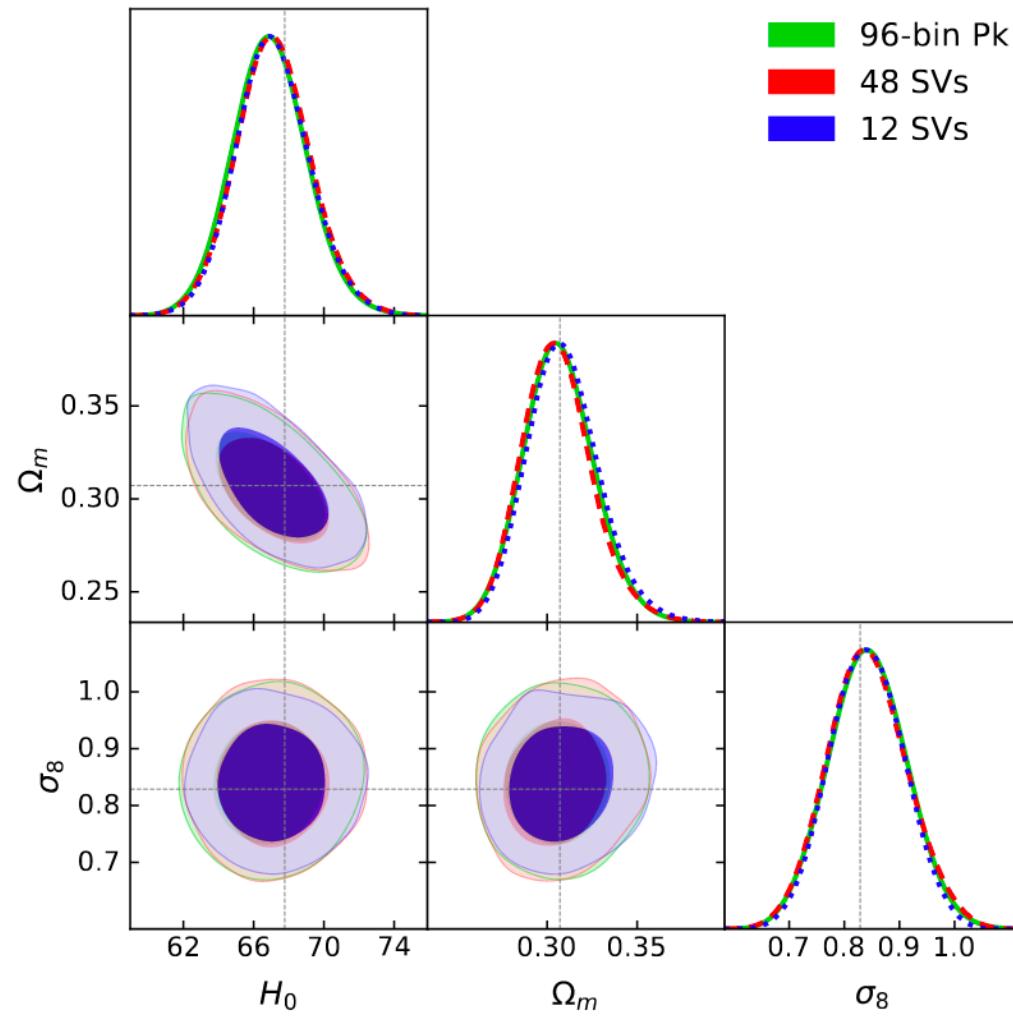
# HIPSTER: Effects of Windowing



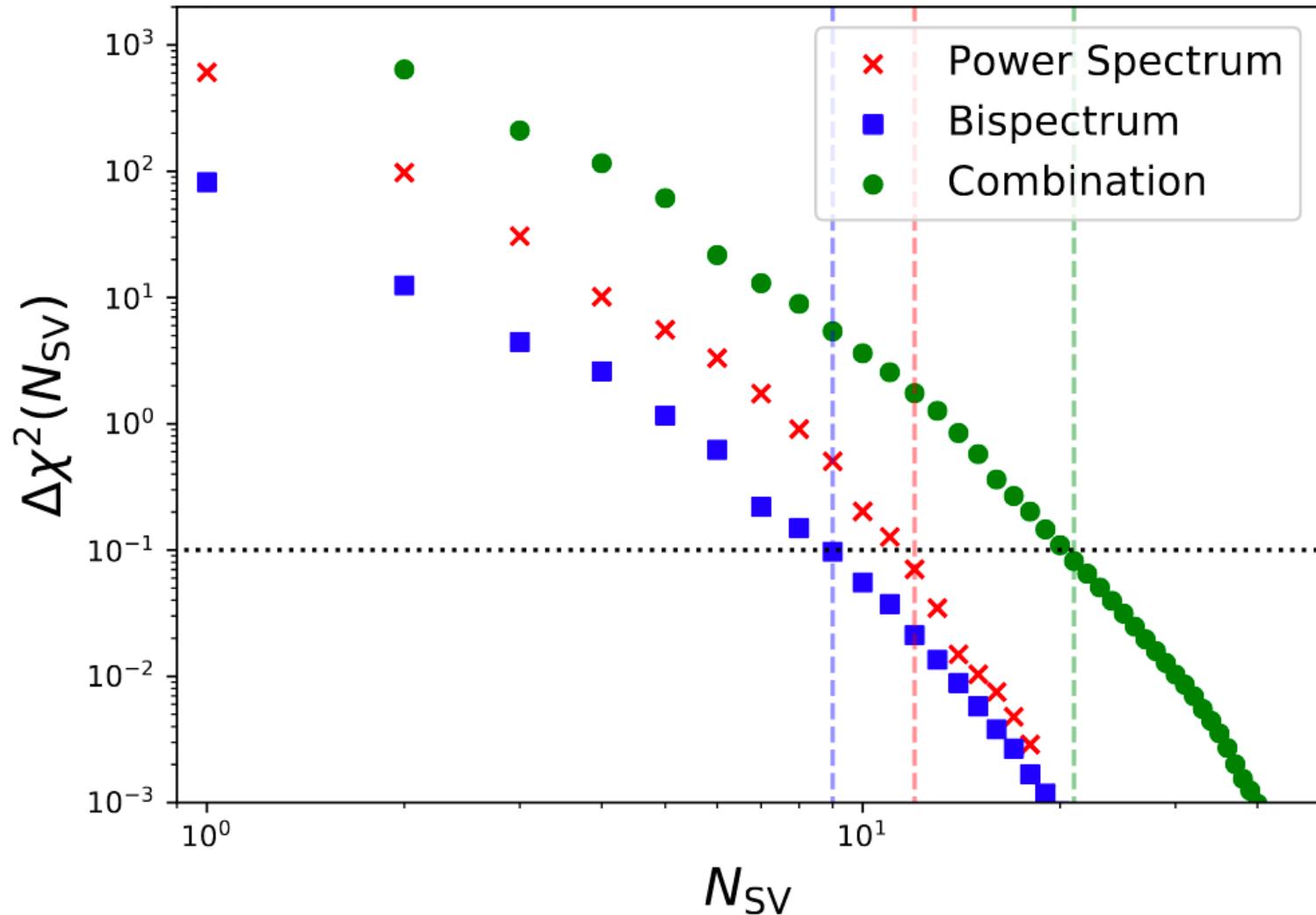
# HIPSTER: Bispectra



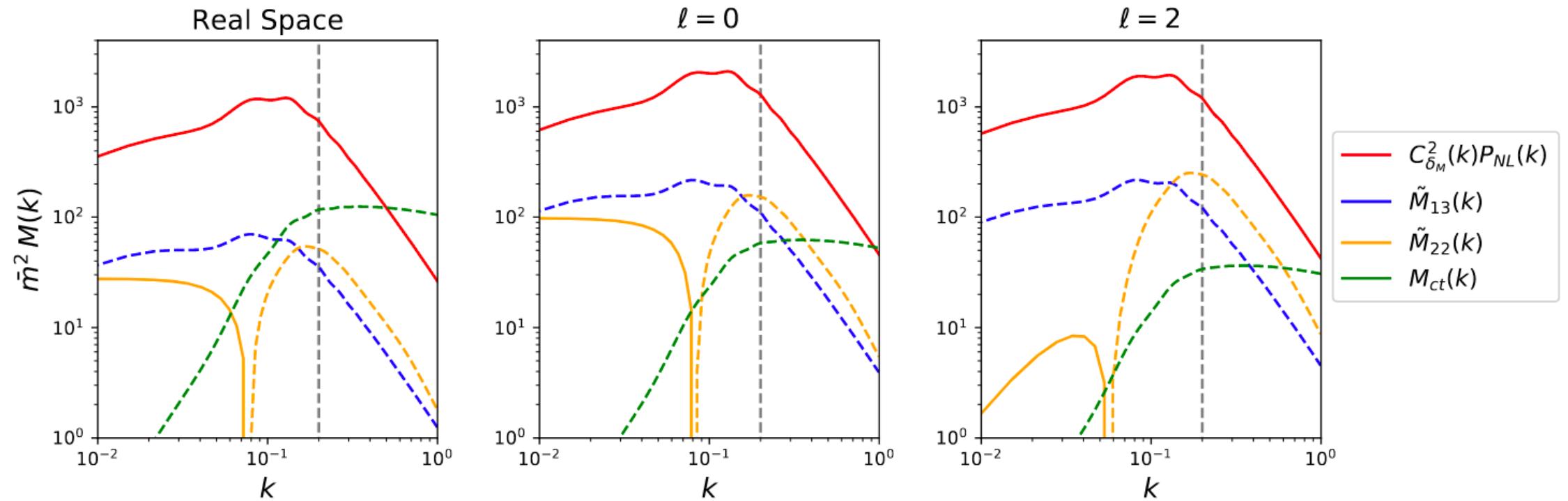
# Compression: Mean of Mocks & Single Mock



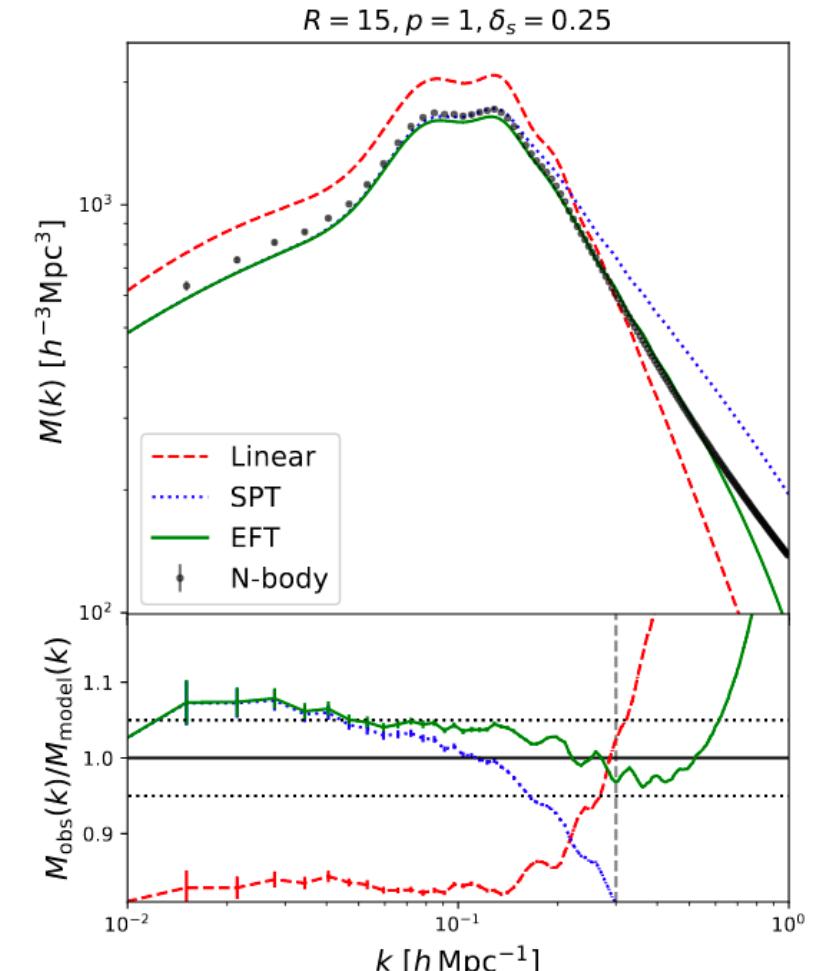
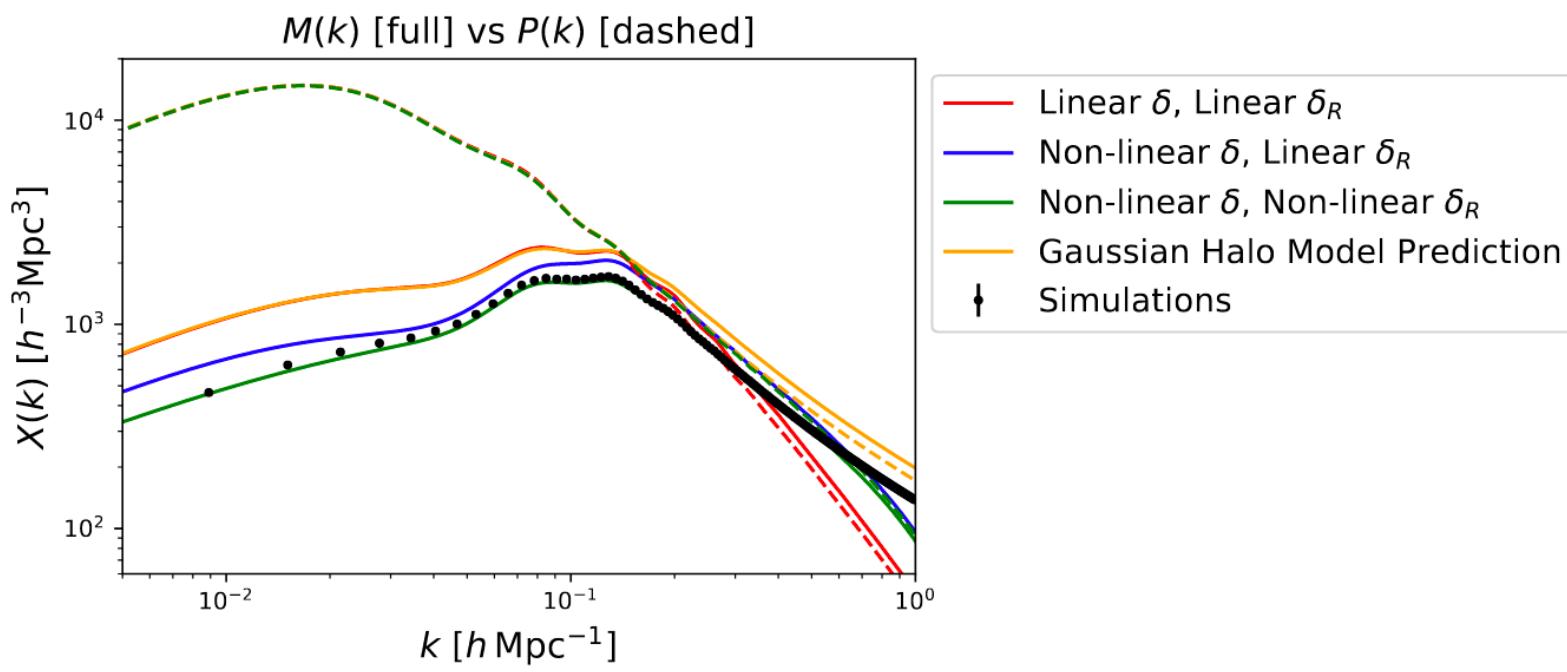
# Compression: Number of Basis Vectors



# Marked Spectra: Matter Contributions



# Marked Spectra: Information Content & Low-z



(a)  $z = 0.5$