

## Practical work - To be returned before January 27

### Finite Differences for the Wave Equation

## 1 1D Model

We consider the following 1D wave equation with initial and boundary conditions:

$$(E_1) \quad \begin{cases} \frac{\partial^2 u}{\partial t^2} - c^2 \frac{\partial^2 u}{\partial x^2} = 0 & \text{for } t > 0, -L < x < L \\ u(x, 0) = f(x) & \text{for } -L \leq x \leq L \\ \frac{\partial u}{\partial t}(x, 0) = 0 & \text{for } -L \leq x \leq L \\ u(-L, t) = u(L, t) = 0 & \text{for } t > 0 \end{cases}$$

$u(x, t)$  represents the planar wave and  $c$  is the speed of the wave (here we take  $c > 0$ ).

1. Show that the solution of equation  $(E_1)$  can be written as the sum of two progressive waves  $F$  and  $G$  propagating respectively with speeds  $c$  and  $-c$ :

$$u(x, t) = F(x - ct) + G(x + ct)$$

We propose to approximate the solution of  $(E_1)$  on  $[-L, L] \times [0, T]$ .

We subdivide the interval  $[-L, L]$  into a regular mesh formed of  $N$  nodes  $x_j$ . We denote by  $u_j^n$  the approximate solution at node  $x_j = j\Delta x$  and at time  $t^n = n\Delta t$ , and we set  $c \frac{\Delta t}{\Delta x} = \lambda$ .

A finite difference scheme centered in space and time can be written as follows:

$$\frac{u_j^{n+1} - 2u_j^n + u_j^{n-1}}{\Delta t^2} - c^2 \frac{u_{j+1}^n - 2u_j^n + u_{j-1}^n}{\Delta x^2} = 0 \quad (1)$$

2. Study the truncation error, the order and the consistency of the numerical scheme (1).

3. Using Fourier Von-Neumann analysis, show that the scheme (1) is stable under the condition

$$c \frac{\Delta t}{\Delta x} \leq 1$$

4. Implement the numerical algorithm and perform the simulations on the following two test cases, using a mesh with  $N = 100$  nodes.

#### Test case 1:

- $L = 10 \text{ m}$  ;  $c = 1 \text{ m/s}$  ;  $f(x) = e^{-x^2}$ .
- Represent the results at physical times  $t_0 = 0 \text{ s}$ ,  $t_1 = 1 \text{ s}$ ,  $t_2 = 4 \text{ s}$ ,  $t_3 = 7 \text{ s}$ .

### Test case 2:

- $L = 1\text{ m}$  ;  $c = 1\text{ m/s}$  ;  $f(x) = \cos\left(\frac{(2k+1)\pi}{2} \frac{x}{L}\right)$ , with  $k = 3$ .
- Represent the results at times  $t_0 = 0\text{ s}$ ,  $t_1 = \frac{T}{5}$ ,  $t_2 = \frac{2T}{3}$ , with  $T = \frac{4L}{(2k+1)c}$ .

## 2 2D Model

The 2D wave equation can be written as follows:

$$(E_2) \quad \begin{cases} \frac{\partial^2 u}{\partial t^2} - c^2 \Delta u = 0 & \text{for } t > 0, (x, y) \in \overset{\circ}{\Omega} \\ u(x, y, 0) = f(x, y) & \text{for } (x, y) \in \Omega \\ \frac{\partial u}{\partial t}(x, y, 0) = 0 & \text{for } (x, y) \in \Omega \\ u(x, y, t) = 1 & \text{for } t > 0, (x, y) \in \partial\Omega \end{cases}$$

$\Omega = [-L_x, L_x] \times [-L_y, L_y]$  being the computational domain, which is assumed to be rectangular, and  $\partial\Omega$  is the boundary of  $\Omega$ .

The computational domain is discretized by a rectangular mesh with  $N_x$  nodes in  $x$  direction and  $N_y$  nodes in  $y$  direction. The space steps are  $\Delta x = \frac{L_x}{N_x - 1}$  and  $\Delta y = \frac{L_y}{N_y - 1}$ .

We denote by  $X_{i,j}$  and  $Y_{i,j}$  the coordinates of a node of index  $(i, j)$  and by  $u_{i,j}^n$  the approximate solution at the node  $(i, j)$  and at the time  $t^n = n\Delta t$ .

5. Based on the scheme (1) for the 1D wave equation ( $E_1$ ), propose a finite difference scheme for the 2D wave equation ( $E_2$ ).

### 6. Test case:

- $L_x = L_y = 10\text{ m}$  ;  $c = 1\text{ m/s}$  ;  $f(x, y) = e^{-(x^2+y^2)}$  ;  $N_x = N_y = 30$ .
- Represent the contour plots of  $u$  at times  $t_0 = 0\text{ s}$ ,  $t_1 = 1\text{ s}$ ,  $t_2 = 4\text{ s}$ ,  $t_3 = 7\text{ s}$ .