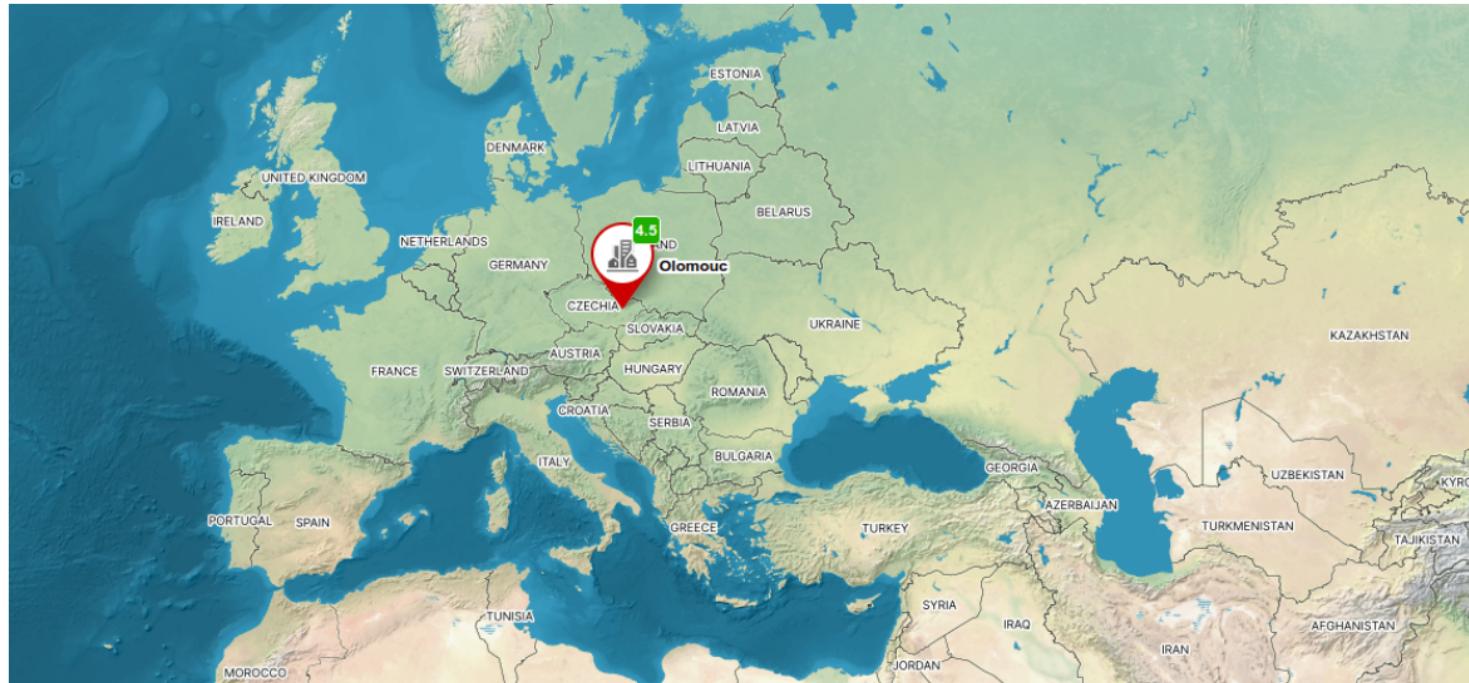


Quantum Non-Gaussian Optomechanics

Andrey A. Rakhubovsky,
Foroud Bemani, Darren W. Moore, Radim Filip

Department of Optics, Palacký University, Czech Republic

MOST 2024,
October 9, 2024,
Samarkand





The Optomechanics Group in Olomouc [within R. Filip's group]

Radim Filip



Foroud Bemani



Darren Moore



Alisa Manukhova



Najmeh Etehadi-Abari



(Now @KIT, Germany)

Surabhi Yadav



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PhD and Postdoc positions available

Introduction

Quantum Optomechanics

Quantum non-Gaussianity

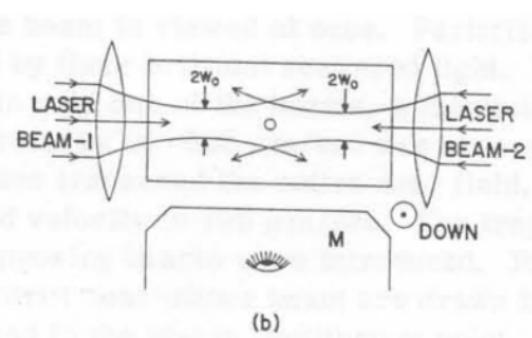
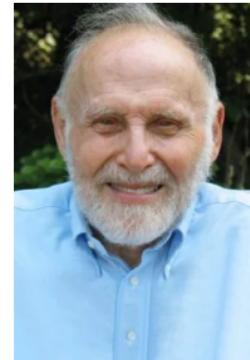
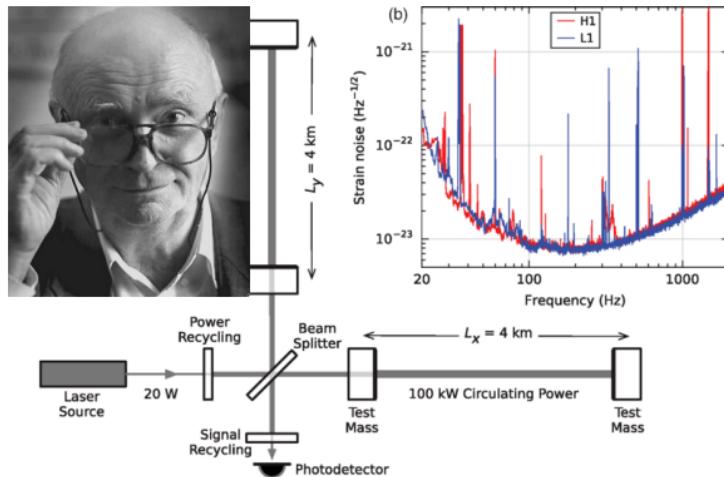
Verification of quantum non-Gaussianity

Motional Nonlinearities

Single-Phonon Addition/Subtraction



Quantum Optomechanics



Braginsky & Manukin, Soviet JETP **25**, 653 (1967)

Right img: B. P. Abbott et al., PRL **116**, 061102 (2016)

A. Ashkin, PRL **24**, 156 (1970)

$$\mathcal{H} = \omega_{\text{mech}} a_m^\dagger a_m + \omega_{\text{cav}} a_L^\dagger a_L + g(a_m + a_m^\dagger)(a_L + a_L^\dagger).$$

Gaussian vs Quantum non-Gaussian states

Gaussian states

Squeezed displaced state

$$\hat{\rho}_x = \hat{S}(r)\hat{D}(\alpha)|0\rangle.$$

Convex mixtures thereof

$$\hat{\rho}_G = \int p(x)\hat{\rho}_x dx.$$

$p(x)$ is a probability density:

- ★ $p(x) > 0$
- ★ “not more singular” than Dirac δ .

Gaussian vs Quantum non-Gaussian states

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Quantum non-Gaussian states

Cannot be represented as convex mixtures of squeezed displaced states

$$\hat{\rho}_{NG} \neq \int p(x)\hat{\rho}_x dx.$$

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Quantum non-Gaussian states

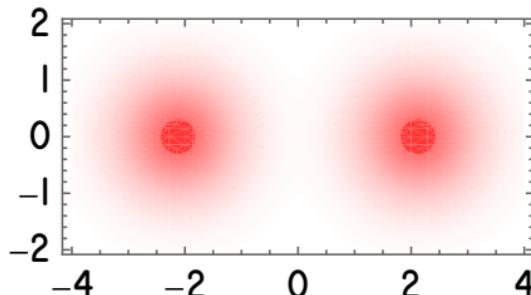
Cannot be represented as convex mixtures of squeezed displaced states

$$\hat{\rho}_{NG} \neq \int p(x)\hat{\rho}_x dx.$$

Examples

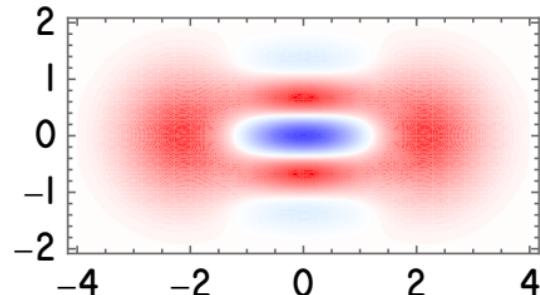
Classically NG state

$$\hat{\rho} \propto |\alpha\rangle\langle\alpha| + |-\alpha\rangle\langle-\alpha|$$



Quantum NG state

$$|\psi\rangle \propto |\alpha\rangle + |-\alpha\rangle$$



Gaussian vs Quantum non-Gaussian states

Advantages of QNG states

- ★ Universal quantum computing
- ★ Quantum sensing
- ★ Fundamental studies

QNG is a resource

F. Albarelli *et al.*, Phys. Rev. A **98**, 052350 (2018)

M. Walschaers, PRX Quantum **2**, 030204 (2021)

Quantum non-Gaussian states

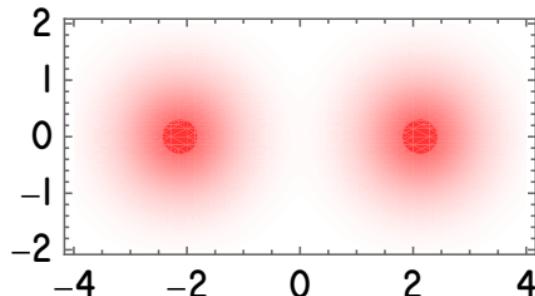
Cannot be represented as convex mixtures of squeezed displaced states

$$\hat{\rho}_{\text{NG}} \neq \int p(x) \rho_x dx.$$

Examples

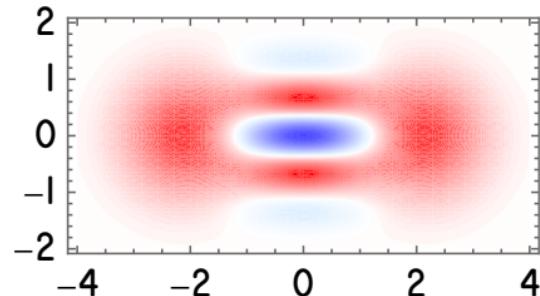
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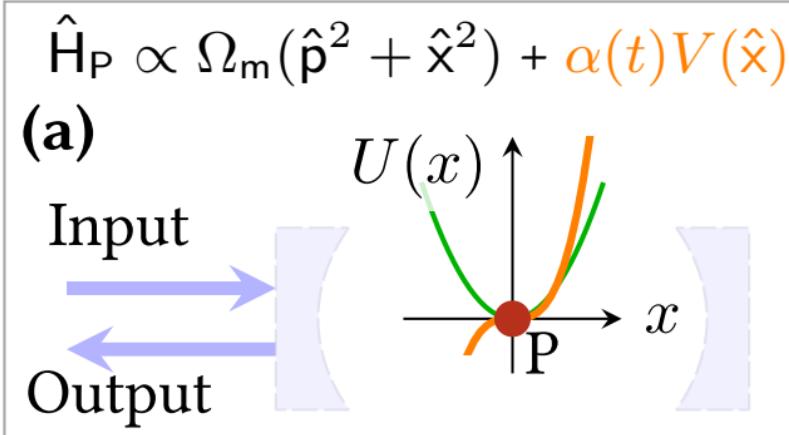
Quantum NG state

$$|\psi\rangle \propto |\alpha\rangle + |-\alpha\rangle$$



Routes to quantum non-Gaussianity in optomechanics

Add a nonlinear element



Nonlinear potential of mechanical motion

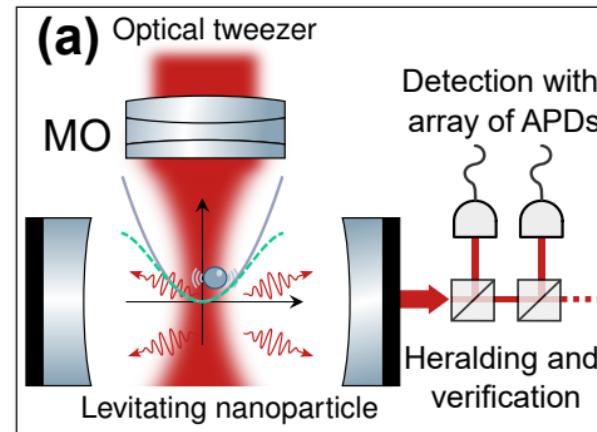
AR, R. Filip, Npj Quantum Inf 7, 120 (2021)

D.W. Moore, AR, R. Filip, NJP 21, 113050 (2019)

We don't consider here upload of QNG states

AR, R. Filip, Sci. Rep. 7, 46764 (2017)

Use non-linear detection



Counting photons

AR, R. Filip, Accepted in QST , doi:10.1088/2058-9565/ad8304 (2024)

F. Bemani, AR, R. Filip, Submitted , (2024)

Verification of quantum non-Gaussianity (QNG)

Linear Gaussian Dynamics



Universal Quantum Control



We need better figures of merit than fidelity

Introduction

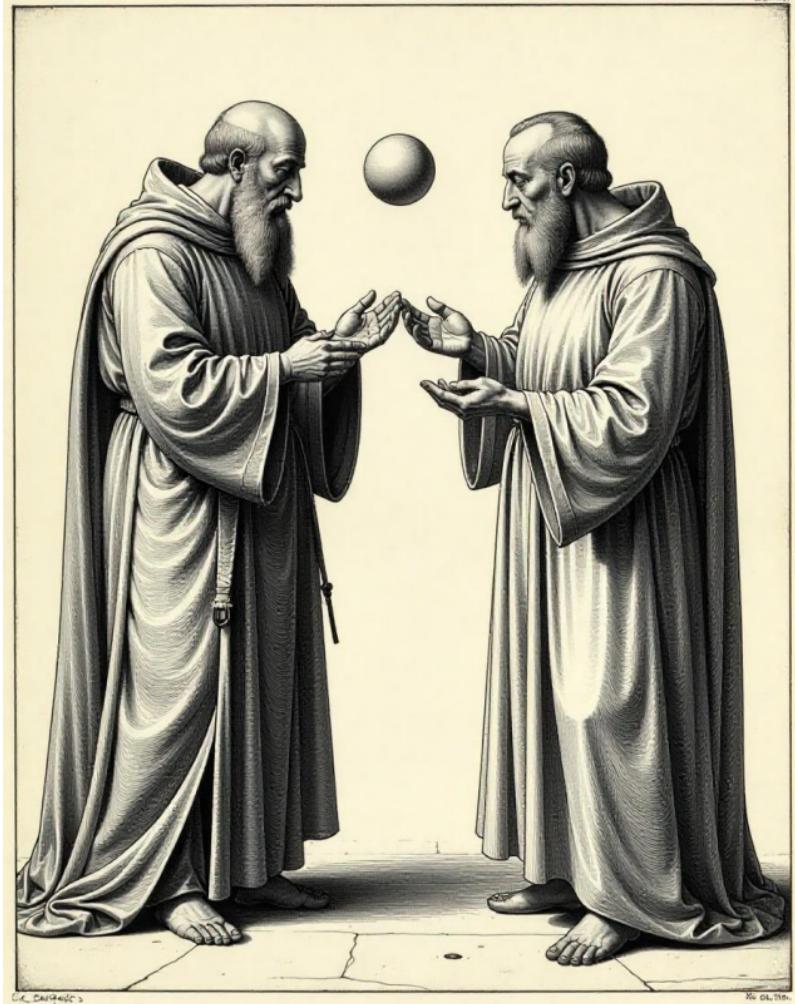
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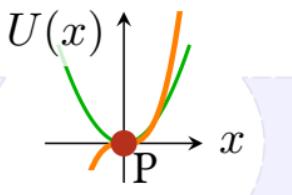
Nonlinear potential for a levitated nanoparticle

$$\hat{H}_P \propto \Omega_m(\hat{p}^2 + \hat{x}^2) + \alpha(t)V(\hat{x})$$

(a)

Input

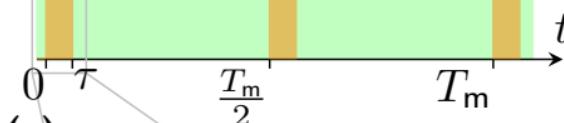
Output



Periodic Temporal Dynamics

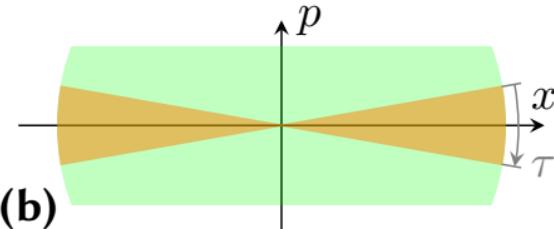
$$\alpha(t)$$

(c)



Phase Space Evolution

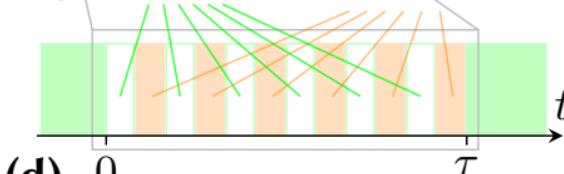
(b)



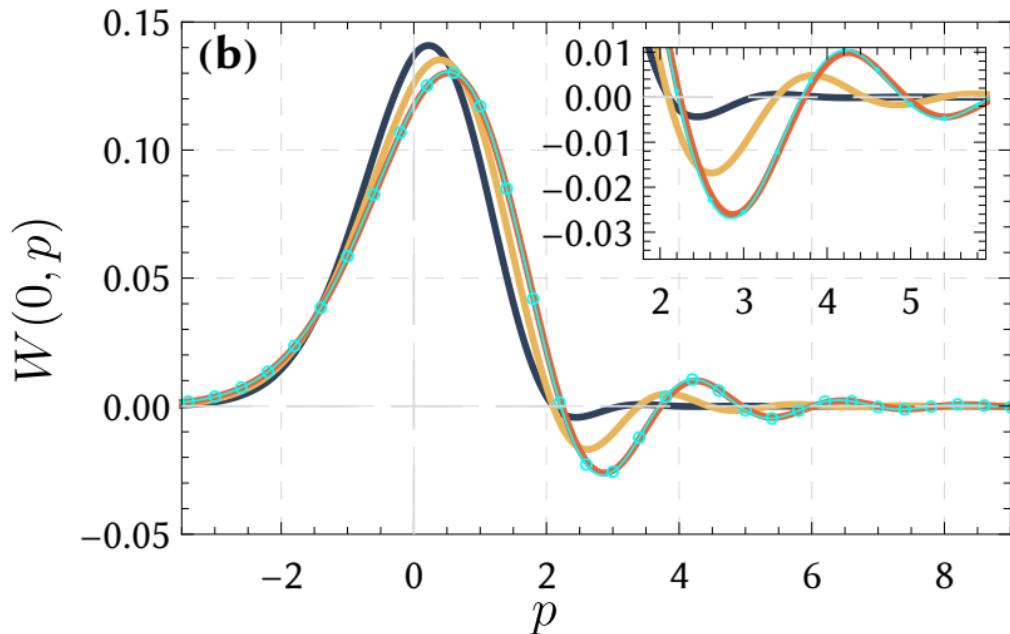
Suzuki-Trotter Simulation

$$e^{-i\hat{H}_{HO}\tau/N} e^{-iV(\hat{x})\tau/N}$$

(d)



Nonlinear potential for a levitated nanoparticle

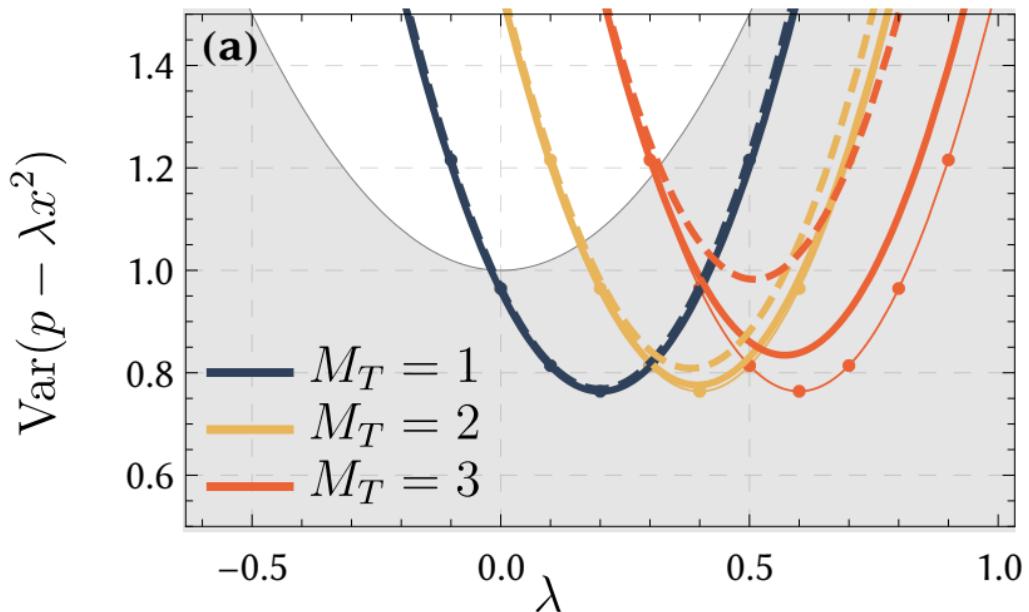


Wigner functions

Poor-man's fidelity (red & cyan)

$$4\pi \int dx dy W_{\text{red}}(x, y) W_{\text{cyan}}(x, y) = 0.9877.$$

Nonlinear potential for a levitated nanoparticle



Nonlinear variance

$$v_3 \equiv \text{Var}(\hat{p} - \lambda \hat{x}^2)$$

Compare with conventional squeezing:

$$v_2 \equiv \text{Var}(\hat{x} \cos \theta + \hat{p} \sin \theta)$$

$$= \sin^2 \theta \cdot \text{Var}(\hat{p} + \lambda \hat{x}),$$

with $\lambda = \cot \theta$.

Related works about levitated NPs in nonlinear potentials (currently all theory)

Palacký University, cubic potential, multiple periods

AR, R. Filip, *Npj Quantum Inf* **7**, 120 (2021) (arxiv 2019)

University of Vienna: “Super Mario”, cubic potential, short pulse

L. Neumeier *et al.*, *PNAS* **121**, e2306953121 (2024) (arxiv 2022)

University of Innsbruck, more complex potentials

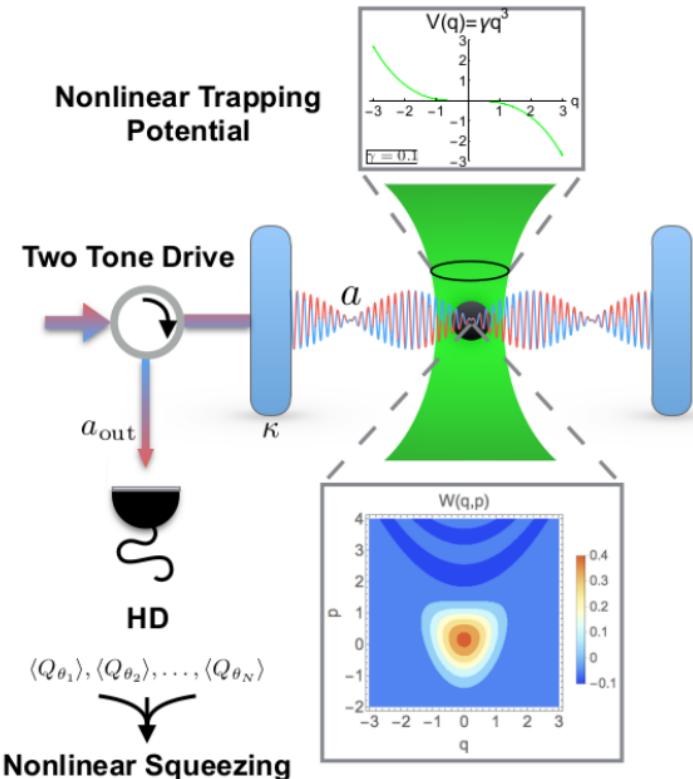
M. Roda-Llordes *et al.*, *Phys. Rev. Res.* **6**, 013262 (2024),

M. Roda-Llordes *et al.*, *Phys. Rev. Lett.* **132**, 023601 (2024),

A. Riera-Campeny *et al.*, arXiv:2307.14106 (2023)



Detecting Nonlinear squeezing



Approximate cubic phase state in mechanics

$$e^{i\gamma q^3} \hat{S} |0\rangle .$$

Pulsed QND interaction

$$H_{\text{int}} \propto x_{\text{light}}(q \cos \phi + p \sin \phi).$$

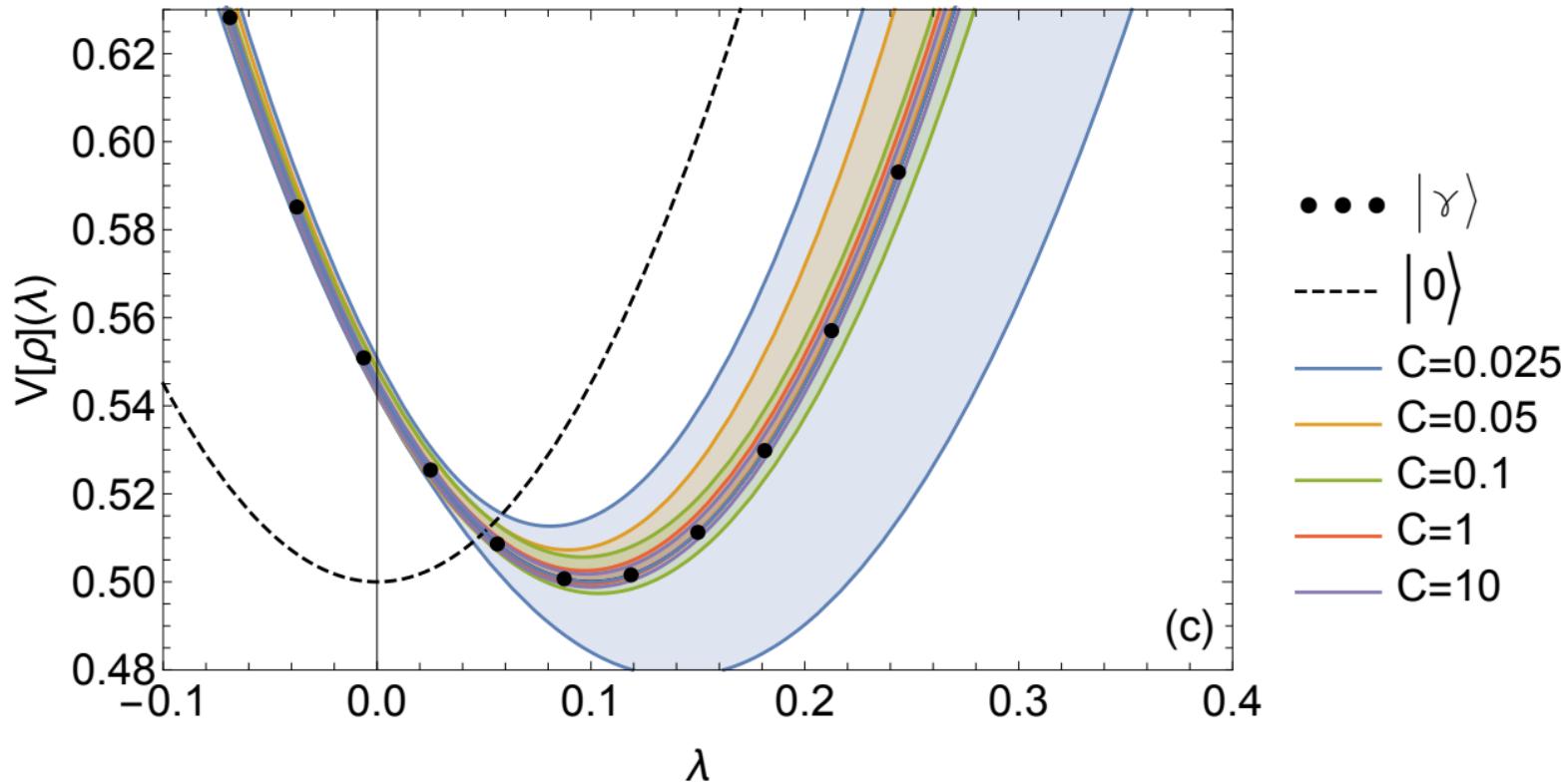
Detect leaking light

Estimate nonlinear variance

$$q_{\text{NL}} = p - \lambda x^2 \Rightarrow \text{Var}[q_{\text{NL}}]$$

D.W. Moore, AR, R. Filip, NJP **21**, 113050 (2019)

$$\text{Evaluation } V[\rho](\lambda) = \text{Tr}(\rho[\Delta(p - \lambda x^2)]^2)$$



Introduction

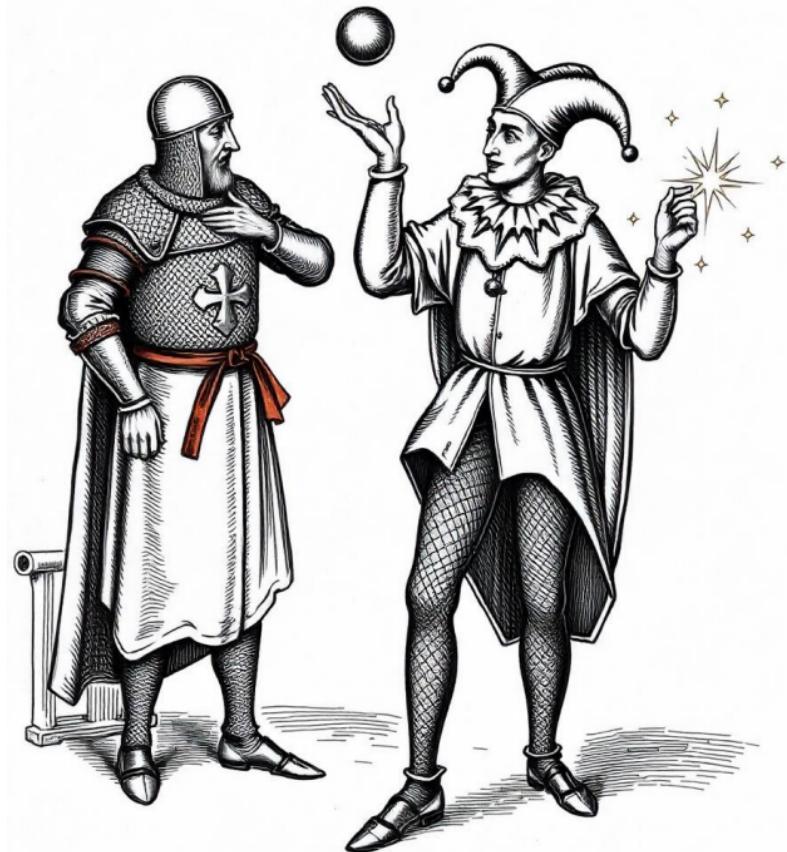
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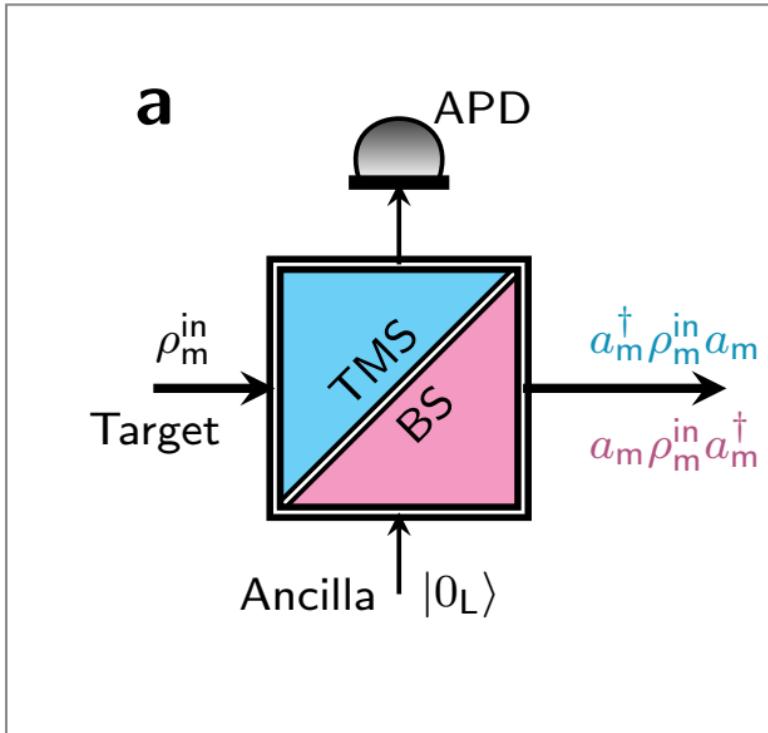
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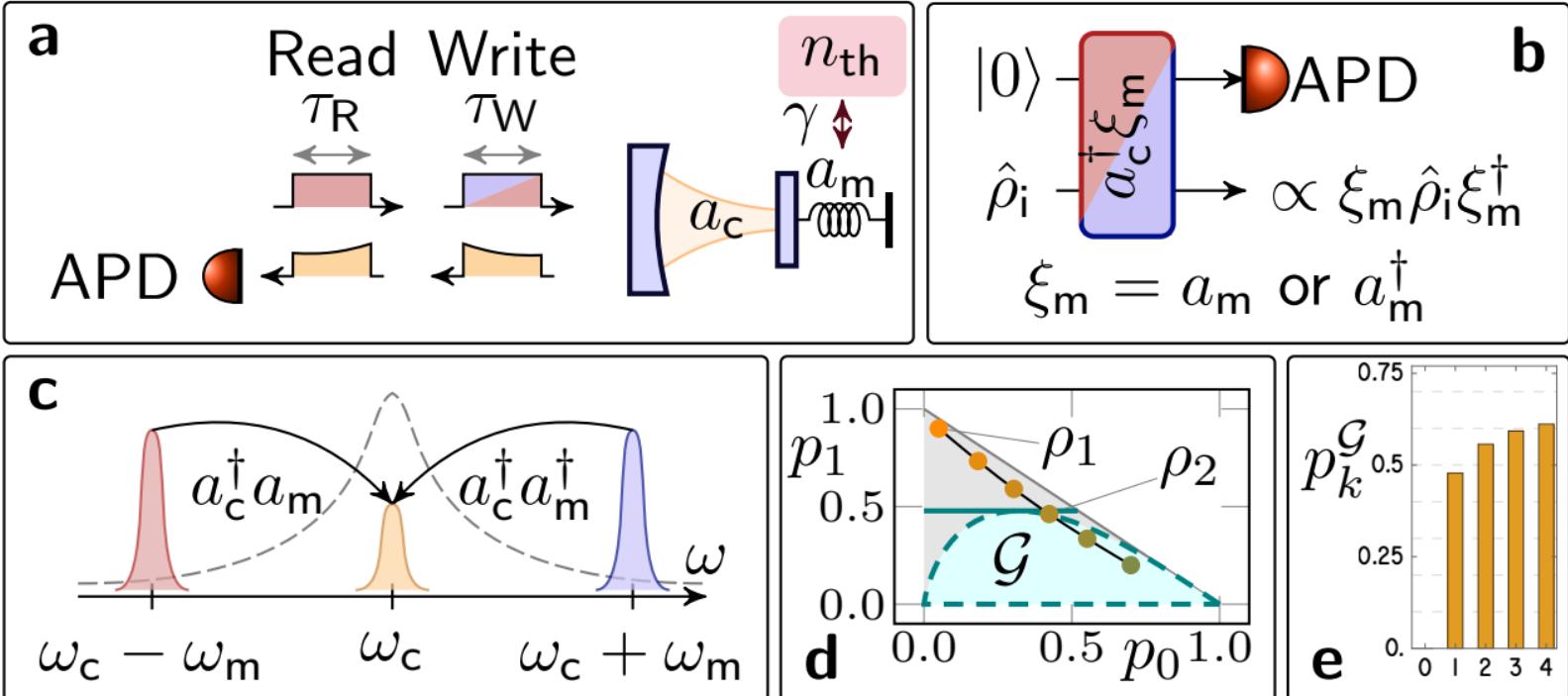
Single-Phonon Addition/Subtraction



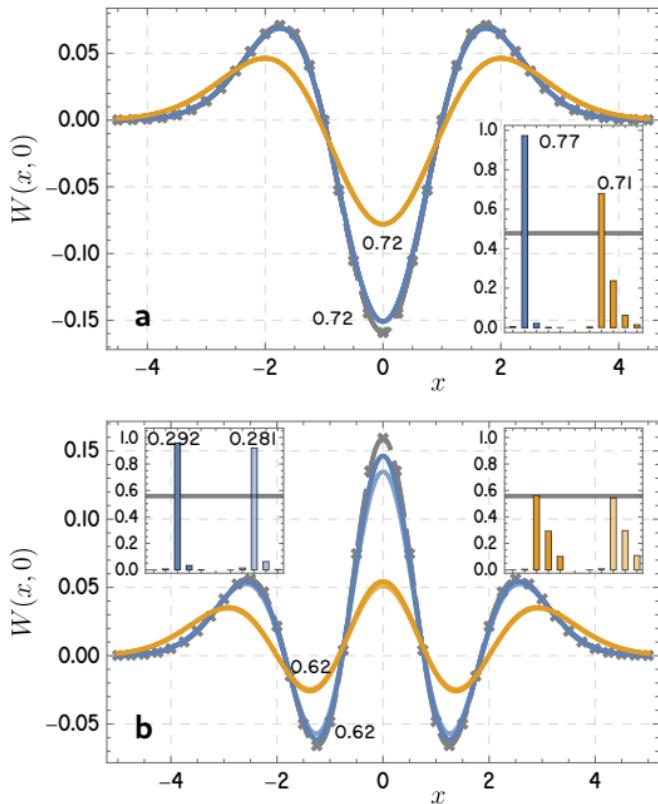
Single-phonon addition or subtraction in optomechanics



Single-phonon addition or subtraction in optomechanics



Evaluation of multiphonon quantum non-Gaussianity (superfluid He)



Multiphonon probabilities

$$p_k = \langle k | \rho | k \rangle$$

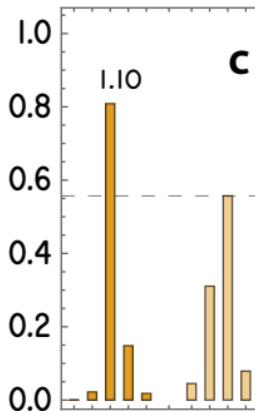
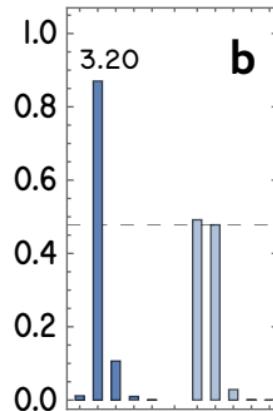
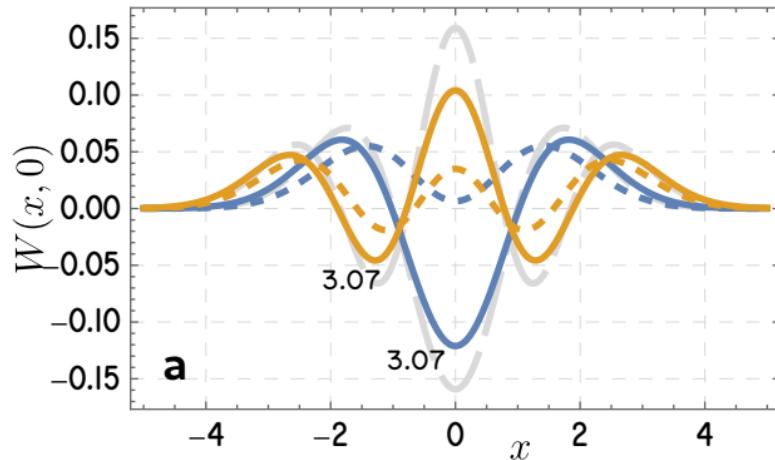
Criteria of absolute n -phonon quantum non-Gaussianity

$$p_k^G = \max_{\alpha, r, \{c_i\}} \left| \left\langle k \left| \hat{D}(\alpha) \hat{S}(r) \sum_{i=0}^{k-1} c_i |i\rangle \right| \right\rangle \right|^2.$$

$$p_2^G = \max_{\alpha, r, c_0, c_1} \left| \left\langle 2 \left| \hat{D}(\alpha) \hat{S}(r) \left(c_0 |0\rangle + c_1 |1\rangle \right) \right| \right\rangle \right|^2.$$

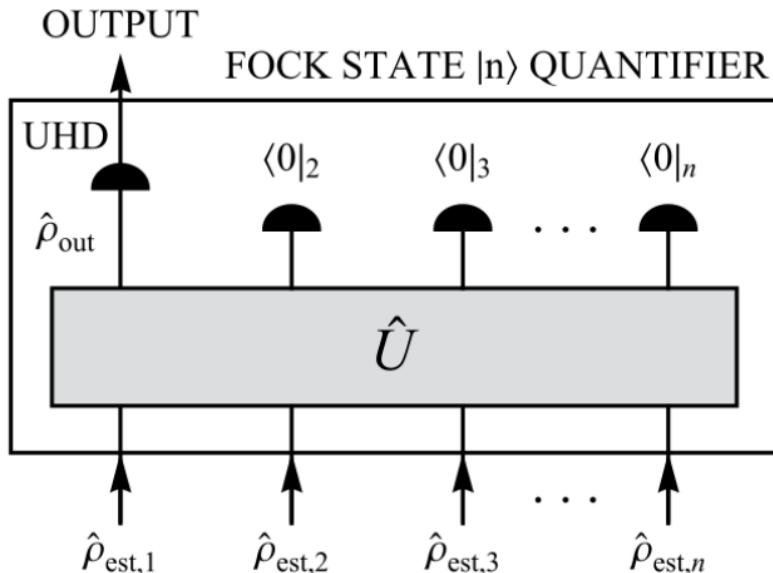
k	1	2	3
p_k^G	0.478	0.557	0.593

Readout and verification



Inset numbers show QNG depth: loss (in dB) to lose QNG.

Bunching capability

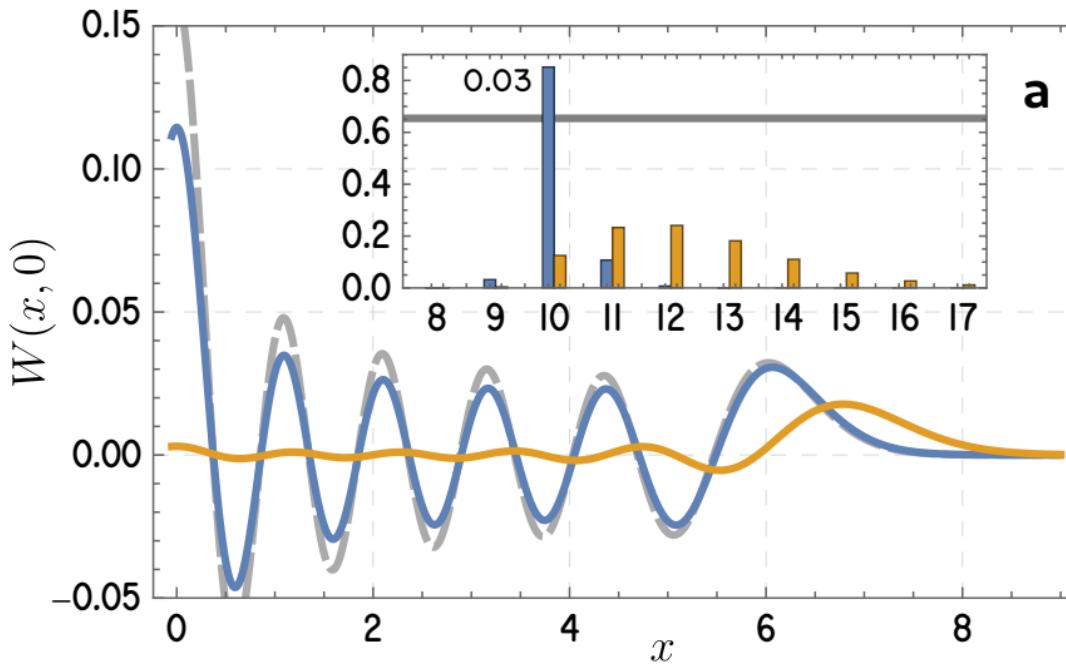


The recipe

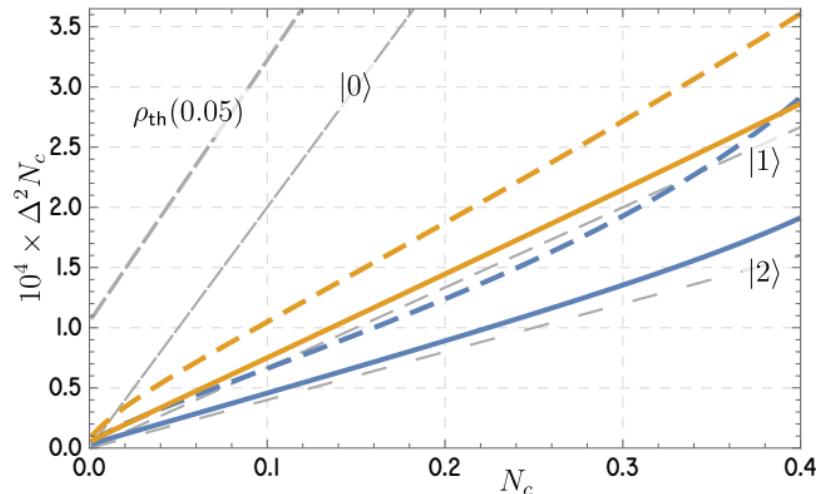
- ★ measure the statistics $\langle k|\hat{\rho}_{\text{est}}|k\rangle$
- ★ compute hypothetical bunching state

Original proposal P. Zapletal and R. Filip, Sci Rep 7, 1 (2017)
Implementations with OPA: P. Zapletal *et al.*, OPTICA 8, 743 (2021).

Bunching capability



Application: detection of phase-randomized displacement



Phase-randomized displacement

$$\rho_{\text{in}} \mapsto \int_0^{2\pi} \frac{d\phi}{2\pi} \hat{D}(\sqrt{N_c} e^{i\phi}) \rho_{\text{in}} \hat{D}^\dagger(\sqrt{N_c} e^{i\phi})$$

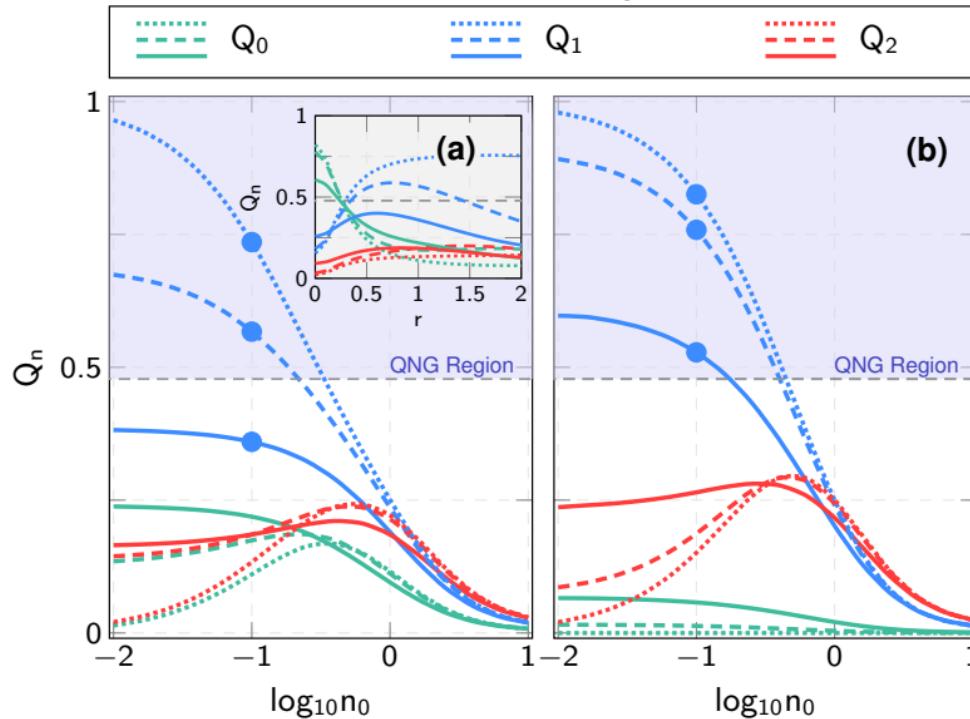
Cramér-Rao bound

$$\Delta^2 N_c \geq \frac{1}{M \cdot F(N_c)},$$

M – number of copies, F – quantum Fisher information

In levitated optomechanics

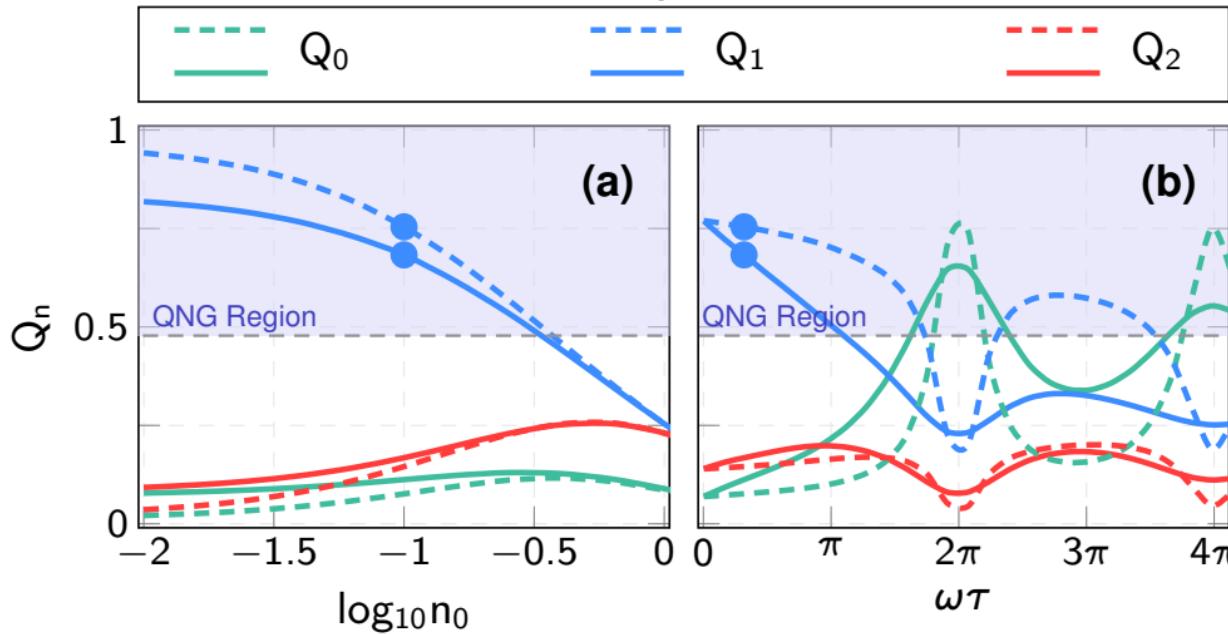
Inside a cavity



Parameters U. Delić, Science 367, 892 (2020)

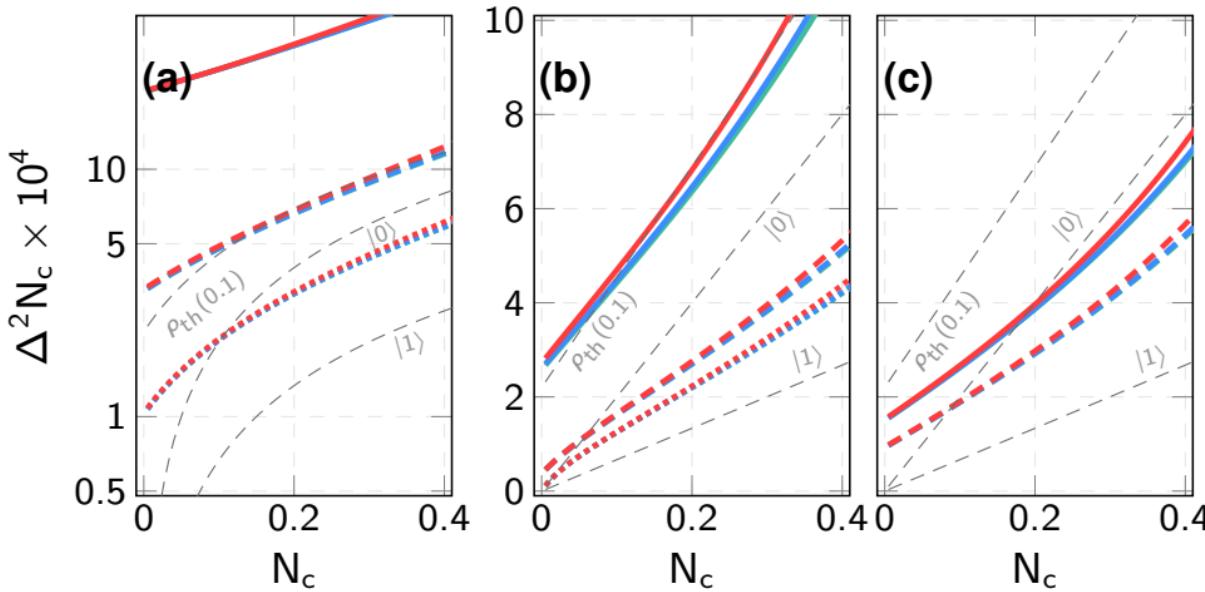
In levitated optomechanics

In free space

Parameters: L. Magrini, Phys. Rev. Lett **129**, 053601 (2022)

In levitated optomechanics

Phase-randomized displacement sensitivity



Thank You!

Ph.D. and postdoc positions available



These slides
<https://bit.ly/andrey-most>



Beware of the appendix slide!

Effective classical simulation

Consider the setup:

- ★ n quantum subsystems
- ★ t operations
- ★ local measurements performed on individual subsystems

The measurement produces one list of outcomes:

$$\mathbf{k} \equiv \{k_1, k_2, \dots, k_n\},$$

one outcome per system.

Probability of the outcomes

$$P(k_1, \dots, k_n) \quad (1)$$

A classical algorithm is said to capable of effective sampling, if it

- ★ is efficient (polynomial) in t and n
- ★ provides outcomes \mathbf{k} draws from the same probability as (1)

The very last frame which is empty