

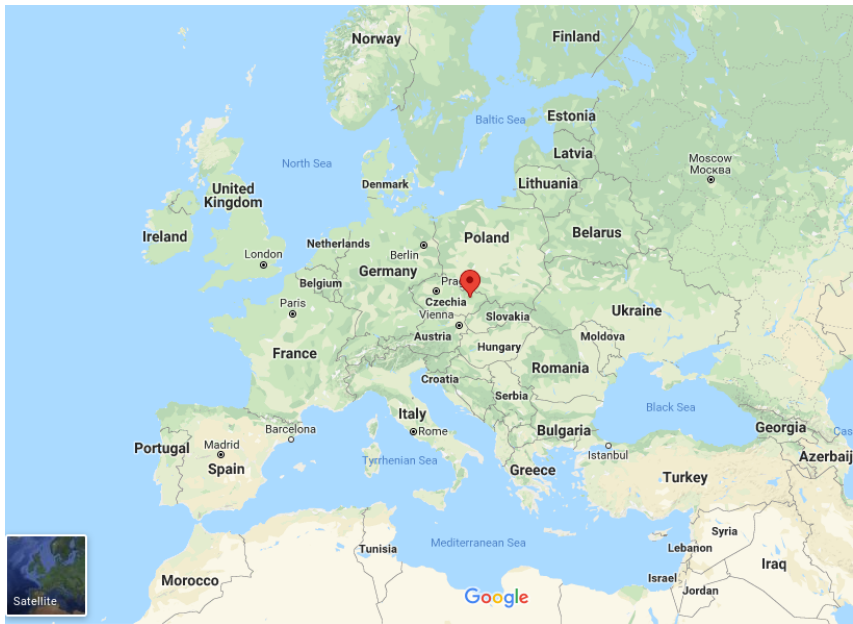
Stroboscopic High-Order Nonlinearity in Optomechanics

Andrey A. Rakhubovsky, Darren W. Moore, Radim Filip

Department of Optics, Palacký University, Czech Republic

New J. Phys. **21** 113050 and npj Quantum Inf. **7**, 120

Quantum Engineering of Levitated Systems
Benasque 2022



Main Results

Levitated Nanoparticle

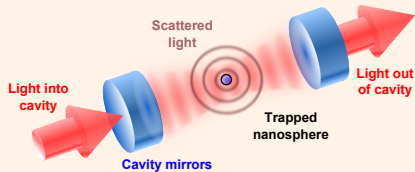
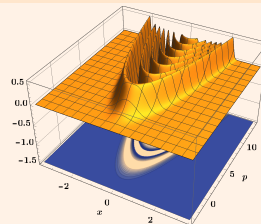


Fig. from Genoni et al. NJP **17**, 073019 (2015).

Approximate Cubic Phase State



$$|\gamma\rangle = e^{i\gamma x^3} |p\rangle \approx e^{i\gamma x^3} \hat{S} |0\rangle$$

Introduction

Prerequisites: Nonlinear Potential
Cubic States

Cubic Phase State in Levitated Optomechanics

CPS Preparation
CPS Evaluation

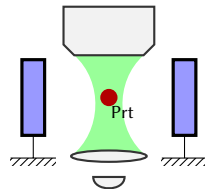


Levitated Nanoparticles in Engineered Potentials

$$H_{\text{trap}} = -\frac{1}{2} \int_{\text{Vol}} d\mathbf{r} \, \mathbf{P}(\mathbf{r}) \mathbf{E}(\mathbf{r}) \propto - \int d\mathbf{r} \, \mathbf{E}^2(\mathbf{r}),$$

$$\mathbf{P} \propto \mathbf{E},$$

$$\text{Equiv. potential: } V(\mathbf{r}; t) \propto -I(\mathbf{r}; t),$$



Levitated Nanoparticles in Engineered Potentials

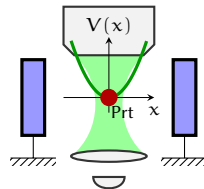
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Gaussian intensity profile

$$I(x) \propto \exp \left[-\frac{x^2}{2x_0^2} \right] \approx 1 - \frac{x^2}{2x_0^2}.$$



$$V(x) \propto \omega_m^2 x^2$$

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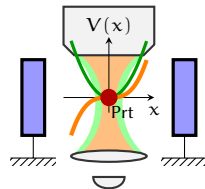
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Gaussian intensity profile

$$I(x) \propto \exp \left[-\frac{x^2}{2x_0^2} \right] \approx 1 - \frac{x^2}{2x_0^2}.$$

Can engineer complicated $I(x)$, particularly cubic $I \propto x^3$



$$V(x) \propto \omega_m^2 x^2 + k_3 x^3$$

Requirements

Ability to switch the nonlinear contribution faster than ω_m .

Cubic phase state

Devised by Gottesman, Kitaev and Preskill ¹

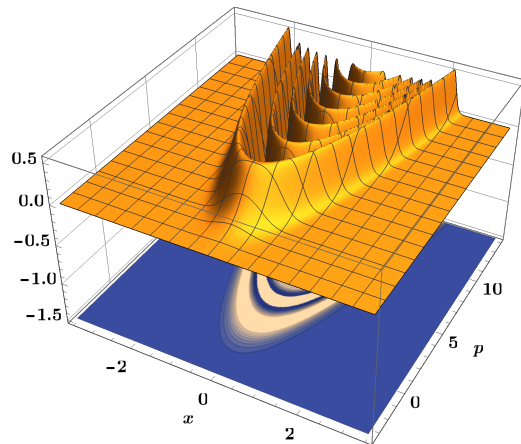
$$|\gamma_{\text{GKP}}\rangle \propto e^{-i\gamma x^3} |p=0\rangle,$$

Wigner Function

$$W_{\text{GKP}}(x, p) \propto \text{Ai} \left[\left(\frac{4}{3\gamma} \right)^{1/3} (3\gamma x^2 - p) \right].$$

More physical is an approximation

$$|\gamma\rangle = e^{-i\gamma x^3} \hat{S} |0\rangle.$$



¹Gottesman, Kitaev, Preskill, PRA **64**, 012310 (2001)

Motivation: Measurement-based computing

PHYSICAL REVIEW A **97**, 022329 (2018)

General implementation of arbitrary nonlinear quadrature phase gates

Petr Marek,^{1,*} Radim Filip,¹ Hisashi Ogawa,² Atsushi Sakaguchi,² Shuntaro Takeda,² Jun-ichi Yoshikawa,^{2,3}
and Akira Furusawa^{2,†}

¹*Department of Optics, Palacký University, 17. listopadu 1192/12, 77146 Olomouc, Czech Republic*

²*Department of Applied Physics, School of Engineering, The University of Tokyo, 7-3-1 Hongo, Bunkyo-ku, Tokyo 113-8656, Japan*

³*Quantum-Phase Electronics Center, School of Engineering, The University of Tokyo, 7-3-1 Hongo, Bunkyo-ku, Tokyo 113-8656, Japan*



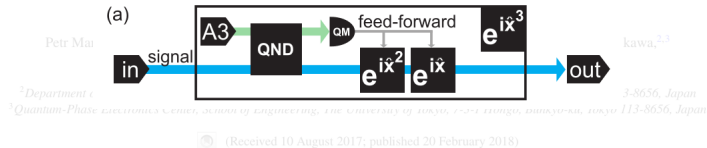
(Received 10 August 2017; published 20 February 2018)

We propose general methodology of deterministic single-mode quantum interaction nonlinearly modifying single quadrature variable of a continuous-variable system. The methodology is based on linear coupling of the system to ancillary systems subsequently measured by quadrature detectors. The nonlinear interaction is obtained by using the data from the quadrature detection for dynamical manipulation of the coupling parameters. This measurement-induced methodology enables direct realization of arbitrary nonlinear quadrature interactions without the need to construct them from the lowest-order gates. Such nonlinear interactions are crucial for more practical and efficient manipulation of continuous quadrature variables as well as qubits encoded in continuous-variable systems.

DOI: [10.1103/PhysRevA.97.022329](https://doi.org/10.1103/PhysRevA.97.022329)

Motivation: Measurement-based computing

PHYSICAL REVIEW A 97, 022329 (2018)



We propose general methodology of deterministic single-mode quantum interaction nonlinearly modifying single quadrature variable of a continuous-variable system. The methodology is based on linear coupling of the i th quadrature of the system with an ancillary state of the k th order squeezed in the i th quadrature. This is obtained by the interaction with the ancillary state. The practical variational method is proposed. DOI: 10.1103/PhysRevA.97.022329

FIG. 1. Schematic circuits for various implementations of nonlinear gates. QND: quantum nondemolition interaction; QM: quadrature measurement; A_k : ancillary state of the k th order squeezed in the i th quadrature. $e^{i\hat{x}^k}$: unitary realization of k th-order nonlinear gate with arbitrary strength. (a) Cubic gate with $N = 3$; (b) $(N + 1)$ th-

Required $\text{Var}(p - \gamma x^2) \rightarrow 0$ for ancilla

Motivation: Measurement-based computing



$$\text{Noise} \propto \text{Var}_{|\text{Ancilla}\rangle}(\mathbf{p} - \gamma \mathbf{x}^2)$$

For a CPS $|\text{Ancilla}\rangle = |\gamma_{\text{GKP}}\rangle := e^{i\gamma x^3} |\mathbf{p} = 0\rangle$,

the variance vanishes $\text{Var}_{|\gamma_{\text{GKP}}\rangle}(\mathbf{p} - \gamma \mathbf{x}^2)$.

Figure of Merit: Nonlinear Variance

The Nonlinear Variance for the implementation of $\exp[-i\gamma x^k]$ is

$$\sigma_k(\lambda) = \text{Var}(p - \lambda x^{k-1}).$$

For a cubic gate $\exp[-i\gamma x^3]$,

$$\sigma_3(\lambda) = \text{Var}(p - \lambda x^2)$$

Evaluated on vacuum state

$$\sigma_3^{\text{vac}} = 1 + 2\lambda^2$$

Analogy with quadratic squeezing

A quantum state is squeezed when for some θ

$$\text{Var}(p \cos \theta + x \sin \theta) < \sigma_{\text{vac}}$$

Equivalent to

$$\sigma_2(\lambda) = \text{Var}(p + \lambda x) < \sigma_{\text{vac}}(1 + \lambda^2), \text{ with } \lambda = \tan \theta.$$



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Cubic Phase State in Levitated Optomechanics

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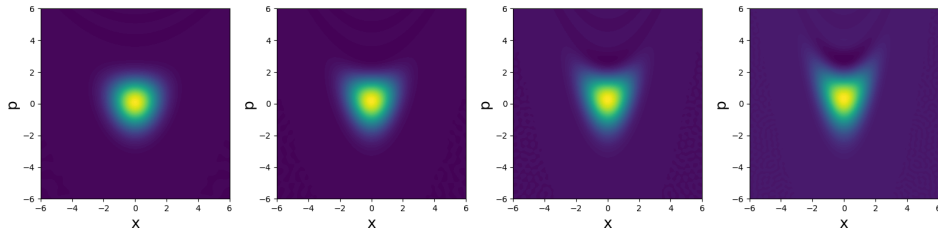
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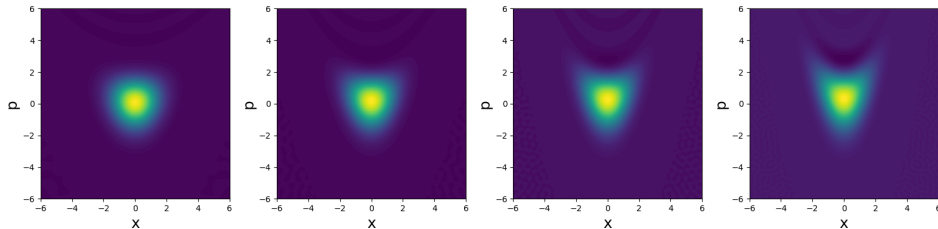
Why bother with protocols

Wigner functions from $\exp[-i(\gamma\hat{x}^3)\tau]$:

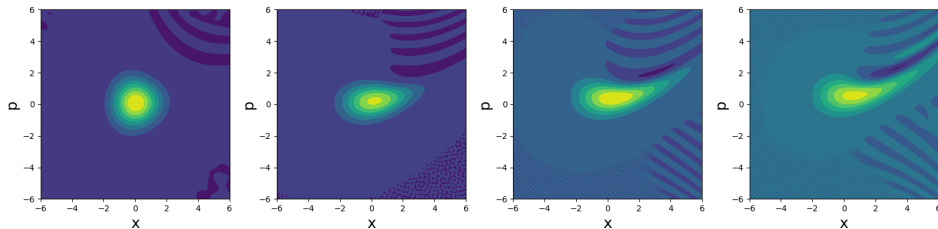


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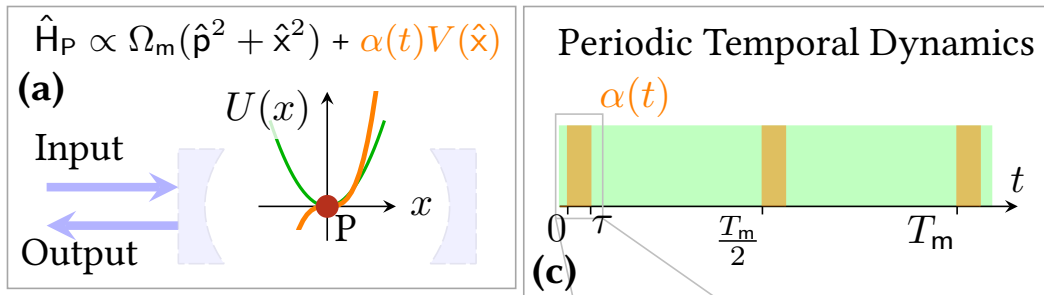


Wigner functions from $\exp[-i(H_{\text{HO}}(\hat{p}) + \gamma\hat{x}^3)\tau]$:

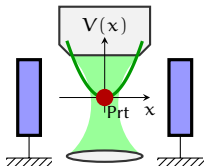


The Model

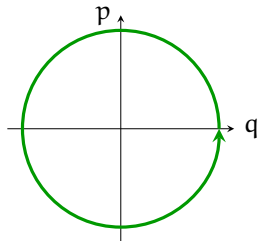
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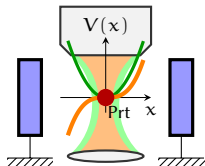
Stroboscopic Application of Nonlinear Potential



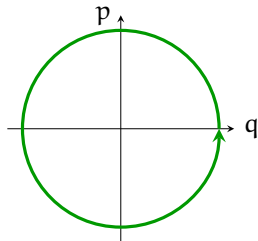
Phase space picture



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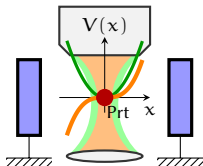


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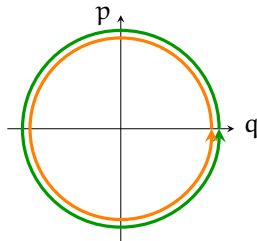


Stroboscopic Application of Nonlinear Potential

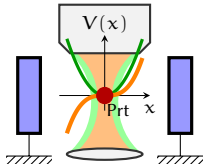
- ★ Continuous application of cubic is smeared out by harmonic evolution



Phase space picture

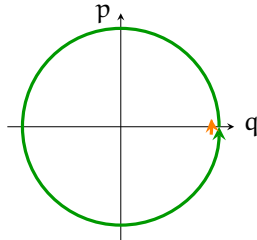


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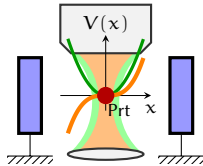


- ★ Continuous application of cubic is smeared out by harmonic evolution
- ★ Apply instantaneous nonlinearity at certain phases of oscillations $t = 0, T_m, 2T_m \dots$

Phase space picture

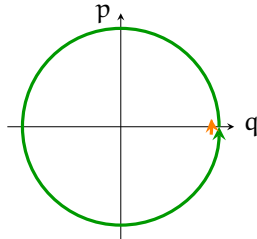


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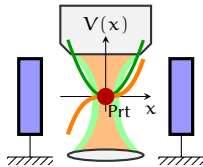


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Phase space picture

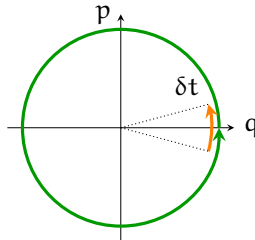


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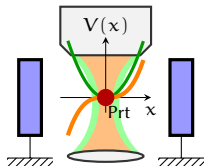


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- ★ Apply long-lasting nonlinearity for a fraction of mechanical period δt .

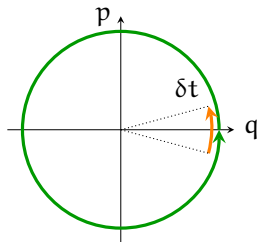
Phase space picture



Stroboscopic Application of Nonlinear Potential



Phase space picture



- ★ Continuous application of cubic is smeared out by harmonic evolution
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Tradeoff between the number of periods $M_T < \Gamma_m^{-1}/T_m$ and duration of application within a certain period δt .

M_T vs δt .

Stroboscopic QND measurement of mechanical motion

Quantum Non-Demolition measurement of mechanical position

$$H_{OM} = \Delta a_L^\dagger a_L + \Omega_m a_M^\dagger a_M + g X_L X_M$$

In order to realize a true QND coupling, get rid of the first two terms:
Tune on resonance $\Delta = 0$ and

Modulate the coupling rate

$$g \mapsto g(t) = g_0 \cos 2\Omega_m t$$

Hamiltonian in rotating frame

$$H_{OM} \mapsto \propto g_0 \tilde{X}_L \tilde{X}_M$$

Use short pulses

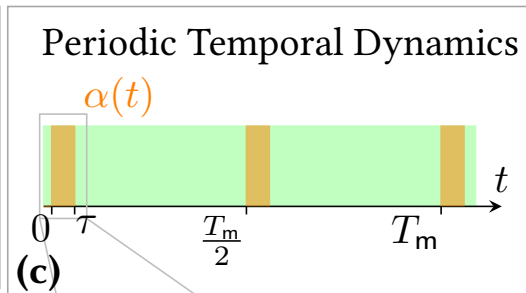
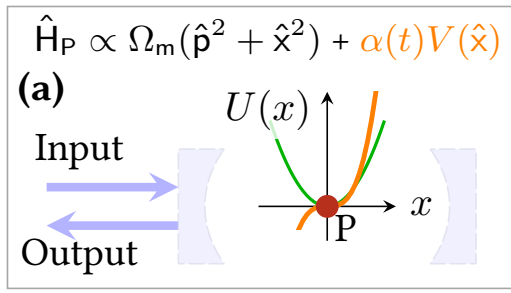
$$\Omega_m \tau \ll 1$$

Effective Hamiltonian

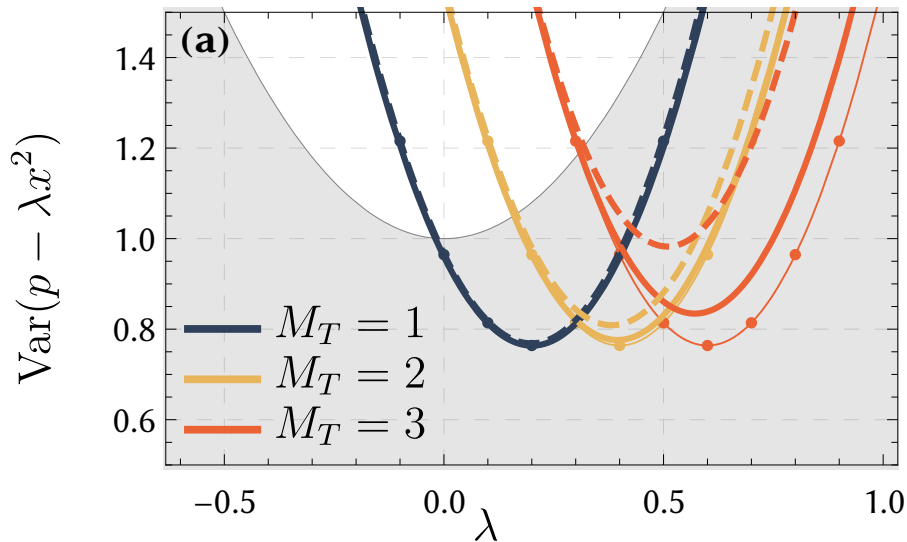
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The Model

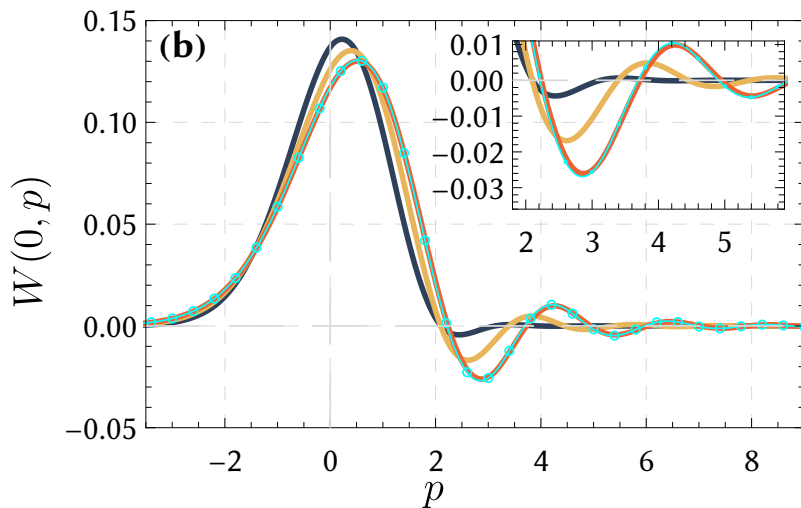
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Nonlinear Variance



Wigner Function Cuts



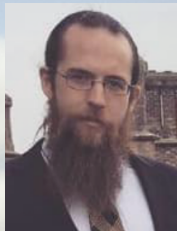
For red and cyan $\text{Tr}[\rho_{\text{red}}\rho_{\text{cyan}}]/\text{Tr}[\rho_{\text{cyan}}^2] = 0.9877$

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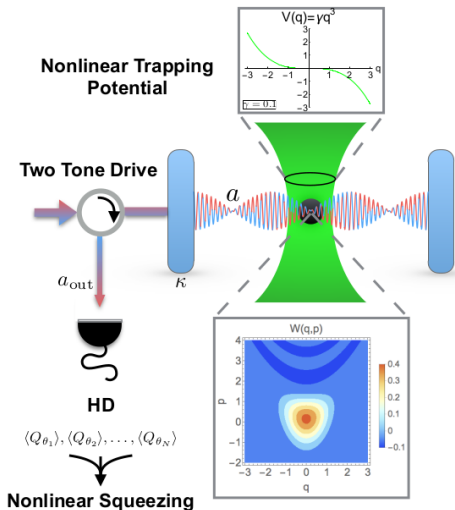
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Darren Moore
NJP 21 113050

The Model



Approximate cubic phase state in mechanics

$$e^{i\gamma q^3} \hat{S} |0\rangle.$$

Pulsed QND interaction

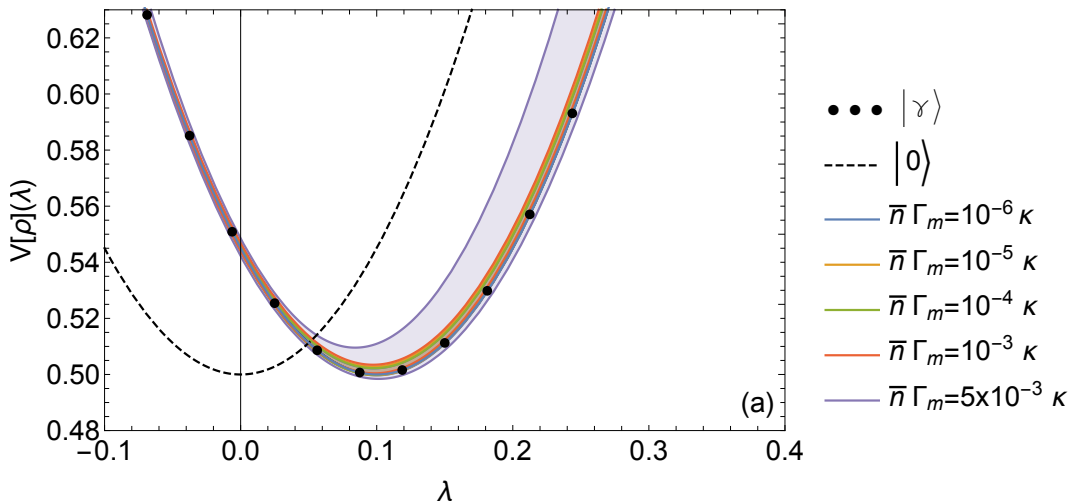
$$H_{\text{int}} \propto x_{\text{light}}(q \cos \phi + p \sin \phi).$$

Detect leaking light

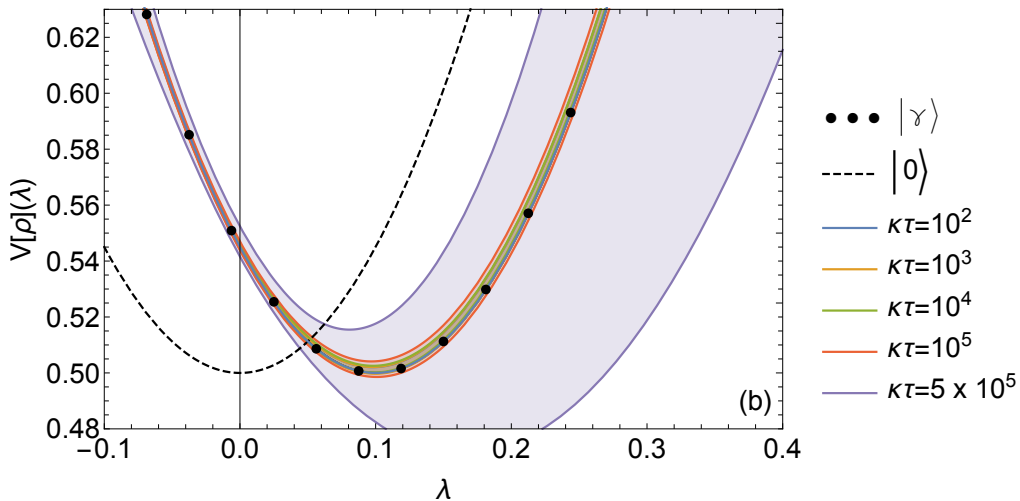
Estimate nonlinear variance

$$q_{\text{NL}} = p - \lambda x^2 \Rightarrow \text{Var}[q_{\text{NL}}]$$

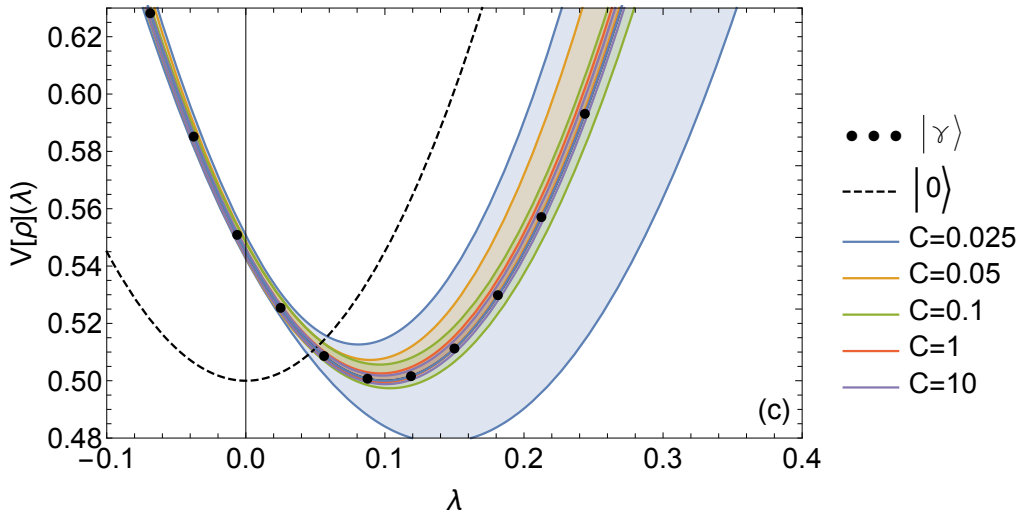
$$\text{Evaluation } V[\rho](\lambda) = \text{Tr}(\rho[\Delta(p - \lambda x^2)]^2)$$



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Conclusion

- ★ Optomechanics provides full linear control over a mechanical oscillator
- ★ Levitated nanoparticles combine advantages of linear optomechanics with possibilities to engineer nontrivial nonlinear potentials
- ★ Stroboscopic application of a cubic potential allows creation of approximate Cubic Phase States
- ★ With the toolbox of optomechanics these states can be read out, verified and used for quantum computation

Thank You!