

Quantitative approach to strain modelling

using Python, Numpy and Matplotlib

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- Concept of homogeneous deformation
- From finite to continuous deformation
- Superposition of deformations

Deformation in terms of displacement

A change in the configuration of a continuum body results in a displacement from an **initial or undeformed configuration** to a current or **deformed configuration**. The displacement of a body has two components:

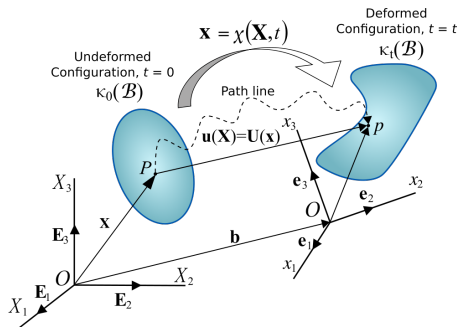
- **Rigid-body displacement**

- Translation
- Rotation

- **Deformation or strain**

- Distortion - isochoric change in shape
- Dilation - change in volume

implies the change in shape and/or size of the body from an initial or undeformed configuration



Kinematics of continuum body

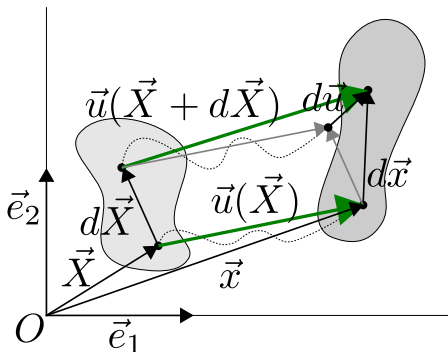
The motion of a continuum body is a **continuous** time sequence of displacements. Thus, the material body will occupy **different configurations** at different times so that a particle occupies a series of points in space which describe a **pathline**.

There is **continuity** during deformation or motion of a continuum body in the sense that:

- The material points forming a closed curve at any instant will always form a closed curve at any subsequent time.
- The material points forming a closed surface at any instant will always form a closed surface at any subsequent time and the matter within the closed surface will always remain within.

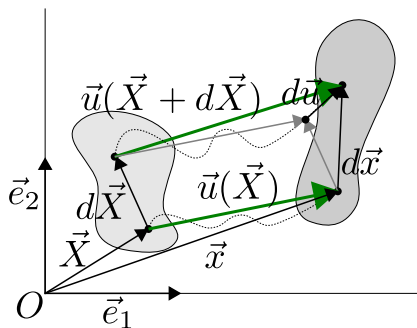
Kinematics: deformation and motion I.

It is convenient to identify a **reference configuration or initial condition** which all subsequent **deformed configurations** are referenced from. Often, the configuration at $t = 0$ is considered the reference configuration.



The components x_i of the position vector \vec{x} of a particle, taken with respect to the reference configuration, are called the **material or reference coordinates**.

Kinematics of infinitesimal deformation I.



The displacement of first point is described as:

$$\vec{x} = \vec{X} + \vec{u}(\vec{X})$$

while displacement of second surrounding point is described as:

$$\vec{x} + d\vec{x} = \vec{X} + d\vec{X} + \vec{u}(\vec{X} + d\vec{X})$$

Substituting first equation into second we got:

$$\vec{X} + \vec{u}(\vec{X}) + d\vec{x} = \vec{X} + d\vec{X} + \vec{u}(\vec{X} + d\vec{X})$$

which simplifies to:

$$d\vec{x} = d\vec{X} + \vec{u}(\vec{X} + d\vec{X}) - \vec{u}(\vec{X})$$

Note that difference between displacements is: $d\vec{u} = d\vec{x} - d\vec{X}$.

Detour on Taylor's theorem

Taylor's theorem states that any function that is infinitely differentiable may be represented by a Taylor series expansion:

$$f(X + dX) = f(X) + \frac{f'(X)}{1!}dX + \frac{f''(X)}{2!}dX^2 + \dots = \sum_{k=0}^{\infty} \frac{f^{(k)}(X)}{k!}dX^k$$

than

$$f(X + dX) - f(X) = \frac{f'(X)}{1!}dX + \frac{f''(X)}{2!}dX^2 + \dots = \sum_{k=0}^{\infty} \frac{f^{(k)}(X)}{k!}dX^k$$

neglecting higher terms as $|d\vec{X}| \ll 1$ as dX^k is very small, it is:

$$f(X + dX) - f(X) = \frac{f'(X)}{1!}dX = (\nabla f)dX$$

where ∇f is gradient of vector field.

Kinematics of infinitesimal deformation II.

Using that for infinitesimal deformation equation

$$d\vec{x} = d\vec{X} + \vec{u}(\vec{X} + d\vec{X}) - \vec{u}(\vec{X})$$

it could be written in terms of gradient as:

$$d\vec{x} = d\vec{X} + (\nabla \mathbf{u})d\vec{X}$$

where $\nabla \mathbf{u}$ is gradient of displacement field or **displacement gradient**.

Displacement gradient

The **displacement gradient** is the derivative of each component of the linear element displacement $d\vec{u}$ with respect to each component of the reference element $d\vec{X}$:

$$\nabla \mathbf{u} = u_{i,j} = \frac{\partial u_i}{\partial X_j} = \begin{bmatrix} \frac{\partial u_1}{\partial X_1} & \frac{\partial u_1}{\partial X_2} & \frac{\partial u_1}{\partial X_3} \\ \frac{\partial u_2}{\partial X_1} & \frac{\partial u_2}{\partial X_2} & \frac{\partial u_2}{\partial X_3} \\ \frac{\partial u_3}{\partial X_1} & \frac{\partial u_3}{\partial X_2} & \frac{\partial u_3}{\partial X_3} \end{bmatrix}$$

and characterise the local change of the displacement field at a material point with position vector \vec{X} . Knowing that:

$$d\vec{u} = d\vec{x} - d\vec{X}$$

it could be also written as:

$$d\vec{u} = (\nabla \mathbf{u})d\vec{X}$$

Deformation gradient

Recalling that $d\vec{u} = d\vec{x} - d\vec{X}$

$$(\nabla \mathbf{u}) = \frac{\partial u_i}{\partial X_j} = \frac{\partial (x_i - X_i)}{\partial X_j} = \frac{\partial x_i}{\partial X_j} - \frac{\partial X_i}{\partial X_j} = \mathbf{F} - \mathbf{I}$$

where \mathbf{F} is so called **deformation gradient**, i.e the derivative of each component of the deformed linear element $d\vec{x}$ with respect to each component of the reference element $d\vec{X}$:

$$\mathbf{F} = x_{i,j} = \frac{\partial x_i}{\partial X_j} = \begin{bmatrix} \frac{\partial x_1}{\partial X_1} & \frac{\partial x_1}{\partial X_2} & \frac{\partial x_1}{\partial X_3} \\ \frac{\partial x_2}{\partial X_1} & \frac{\partial x_2}{\partial X_2} & \frac{\partial x_2}{\partial X_3} \\ \frac{\partial x_3}{\partial X_1} & \frac{\partial x_3}{\partial X_2} & \frac{\partial x_3}{\partial X_3} \end{bmatrix}$$

and characterizes the local deformation at a material point with position vector \vec{X} , assuming continuity. As

$$d\vec{u} = d\vec{x} - d\vec{X} = (\nabla \mathbf{u})d\vec{X} = (\mathbf{F} - \mathbf{I})d\vec{X} = \mathbf{F}d\vec{X} - d\vec{X}$$

$$d\vec{x} = \mathbf{F}d\vec{X}$$

Properties of deformation gradient

Deformation gradient \mathbf{F} contains all the required local information about the changes in length, volumes and angles due to the deformation as follows:

- When vector \vec{N} in the reference configuration is deformed into the vector \vec{n} , these vectors are related as: $\vec{n} = \mathbf{F}\vec{N}$
- The ratio between the local volume of the deformed configuration to the local volume in the reference configuration is equal to the determinant of the deformation gradient tensor: $J = \det \mathbf{F}$
- Two infinitesimal areas with da and dA being their magnitudes and \vec{n} and \vec{N} are unit vectors perpendicular to them, then the relationship is given by: $(da)\vec{n} = \det(\mathbf{F})(dA)\mathbf{F}^{-T}\vec{N}$
- An isochoric deformation is a deformation preserving local volume, i.e., $\det \mathbf{F} = 1$ at every point
- A deformation is called homogeneous if \mathbf{F} is constant at every point. Otherwise, the deformation is called non-homogeneous
- The physical restriction of possible deformation: $\det \mathbf{F} > 0$

Homogeneous deformation

A **homogeneous deformation** is one where the deformation gradient is uniform, i.e. independent of the coordinates, and the associated motion is termed **affine**. Every part of the material deforms as the whole does, and straight parallel lines in the reference configuration map to straight parallel lines in the deformed configuration.

$$\begin{aligned}x &= aX + bY + t_X \\ y &= cX + dY + t_Y\end{aligned}$$

or in matrix form:

$$\begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{bmatrix} X \\ Y \end{bmatrix} + \begin{bmatrix} t_X \\ t_Y \end{bmatrix}$$

Properties of homogeneous deformation are not spatially dependent.

Deformation gradient

Without translation the homogeneous deformation (rotation and strain) could be described as:

$$x = aX + bY$$

$$y = cX + dY$$

$$\begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \begin{bmatrix} X \\ Y \end{bmatrix}$$

or

$$\vec{x} = \mathbf{F} \vec{X}$$

where \mathbf{F} is so called **deformation gradient**.

Note, that as we excluded translation, the origin of coordinates do not change during deformation:

$$\begin{bmatrix} 0 \\ 0 \end{bmatrix} = \mathbf{F} \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Displacement gradient

Displacement of particle is vector between initial and final position, i.e:

$$u = x - X = aX + bY - X = (a - 1)X + bY$$

$$v = y - Y = cX + dY - Y = cX + (d - 1)Y$$

$$\begin{bmatrix} u \\ v \end{bmatrix} = \begin{bmatrix} a - 1 & b \\ c & d - 1 \end{bmatrix} \begin{bmatrix} X \\ Y \end{bmatrix}$$

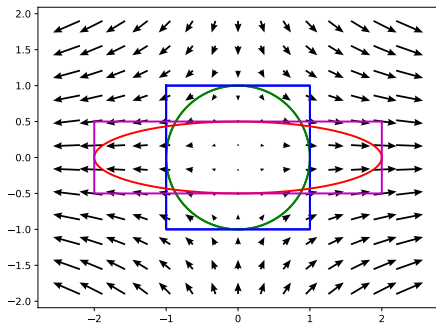
or

$$\vec{u} = (\mathbf{F} - \mathbf{I})\vec{X} = \nabla \mathbf{u} \vec{X}$$

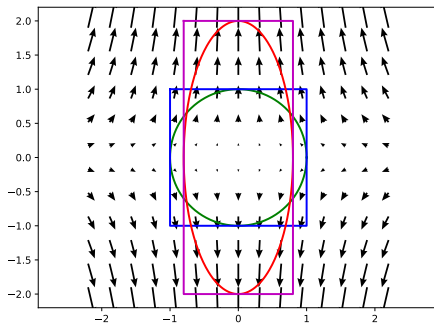
where $\nabla \mathbf{u}$ is so called **displacement gradient**.

Examples of pure shear

$$\mathbf{F} = \begin{bmatrix} 2 & 0 \\ 0 & 0.5 \end{bmatrix}$$
$$\nabla \mathbf{u} = \begin{bmatrix} 1 & 0 \\ 0 & -0.5 \end{bmatrix}$$

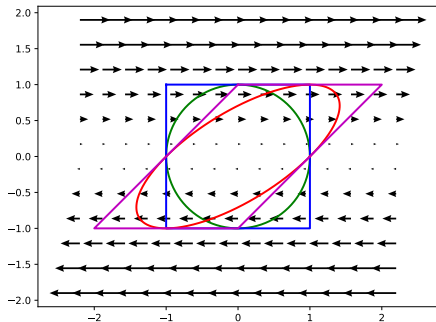


$$\mathbf{F} = \begin{bmatrix} 0.8 & 0 \\ 0 & 2 \end{bmatrix}$$
$$\nabla \mathbf{u} = \begin{bmatrix} -0.2 & 0 \\ 0 & 1 \end{bmatrix}$$

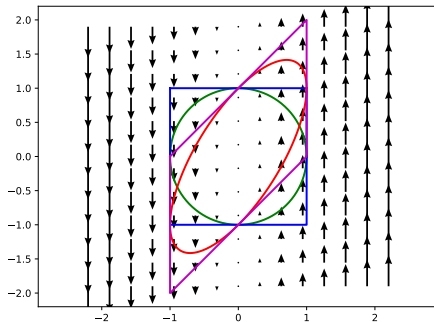


Examples of simple shear

$$\mathbf{F} = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix}$$
$$\nabla \mathbf{u} = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

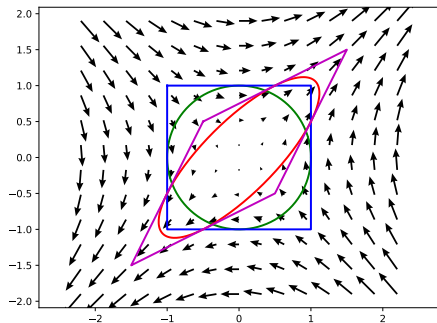


$$\mathbf{F} = \begin{bmatrix} 1 & 0 \\ 1 & 1 \end{bmatrix}$$
$$\nabla \mathbf{u} = \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}$$

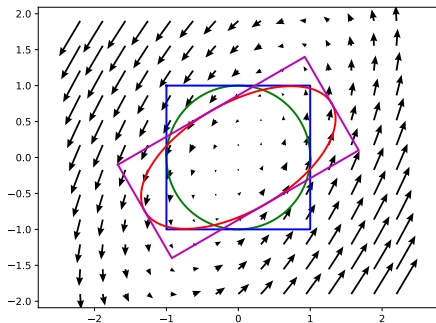


Examples of general shear

$$\mathbf{F} = \begin{bmatrix} 1 & 0.5 \\ 0.5 & 1 \end{bmatrix}$$
$$\nabla \mathbf{u} = \begin{bmatrix} 0 & 0.5 \\ 0.5 & 0 \end{bmatrix}$$



$$\mathbf{F} = \begin{bmatrix} 1.300 & -0.375 \\ 0.750 & 0.650 \end{bmatrix}$$
$$\nabla \mathbf{u} = \begin{bmatrix} 0.300 & -0.375 \\ 0.750 & -0.350 \end{bmatrix}$$



Time to think...

\boldsymbol{F} maps any undeformed vector into its deformed state. This vector can also be a position vector of a point. Therefore \boldsymbol{F} also maps any point into its new position after deformation. Considering two successive deformations \boldsymbol{F}_1 and \boldsymbol{F}_2 write transformation equation....