

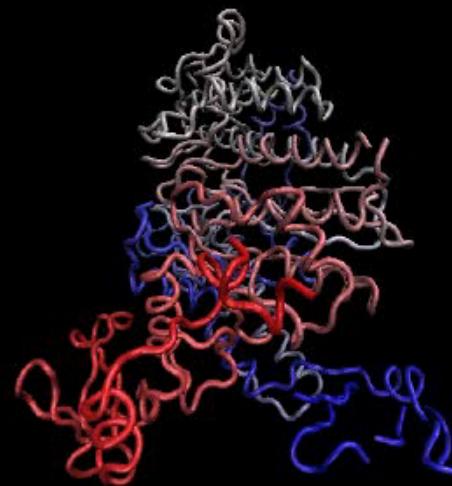
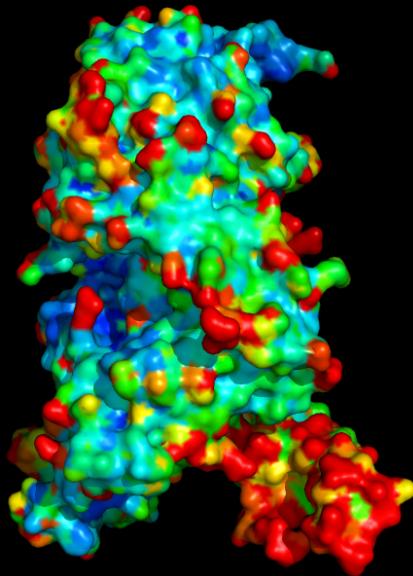
# **Modal Analysis of Myosin II and Identification of Functionally Important Sites**

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# Jiggling and wiggling of atoms



"Certainly no subject or field is making more progress on so many fronts at the present moment, than biology, and if we were to name the most powerful assumption of all, which leads one on and on in an attempt to understand life, it is that all things are made of atoms, and that everything that living things do can be understood in terms of the jiggling and wiggling of atoms."

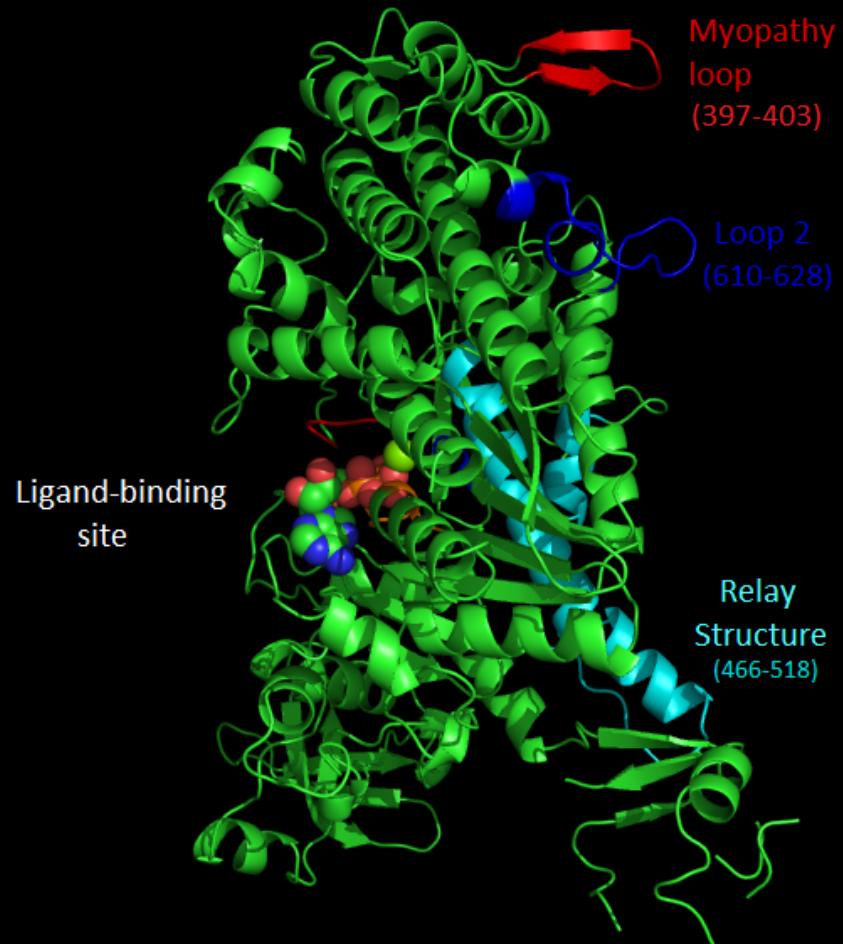
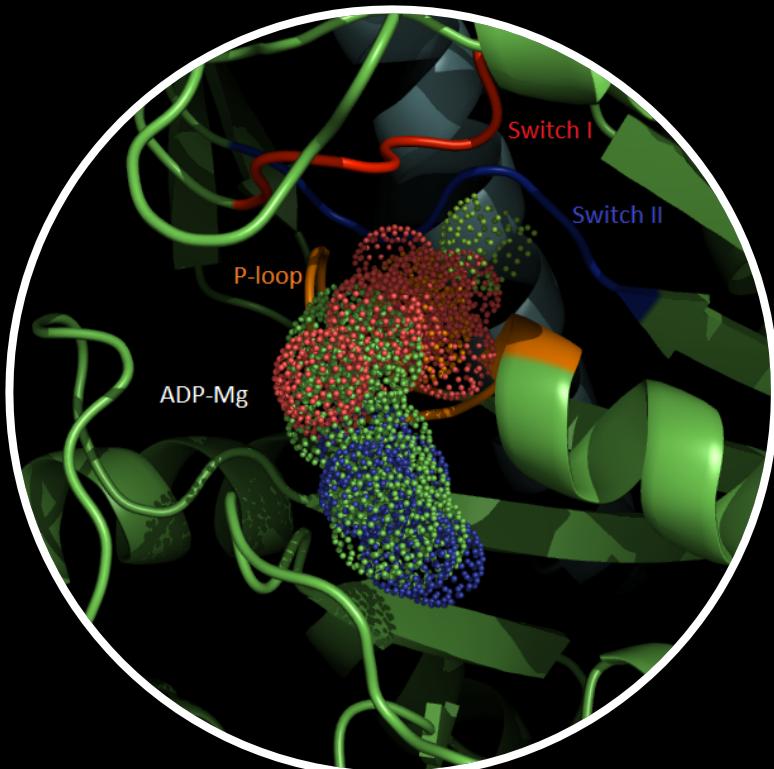
Richard Philips Feynman

# Outline

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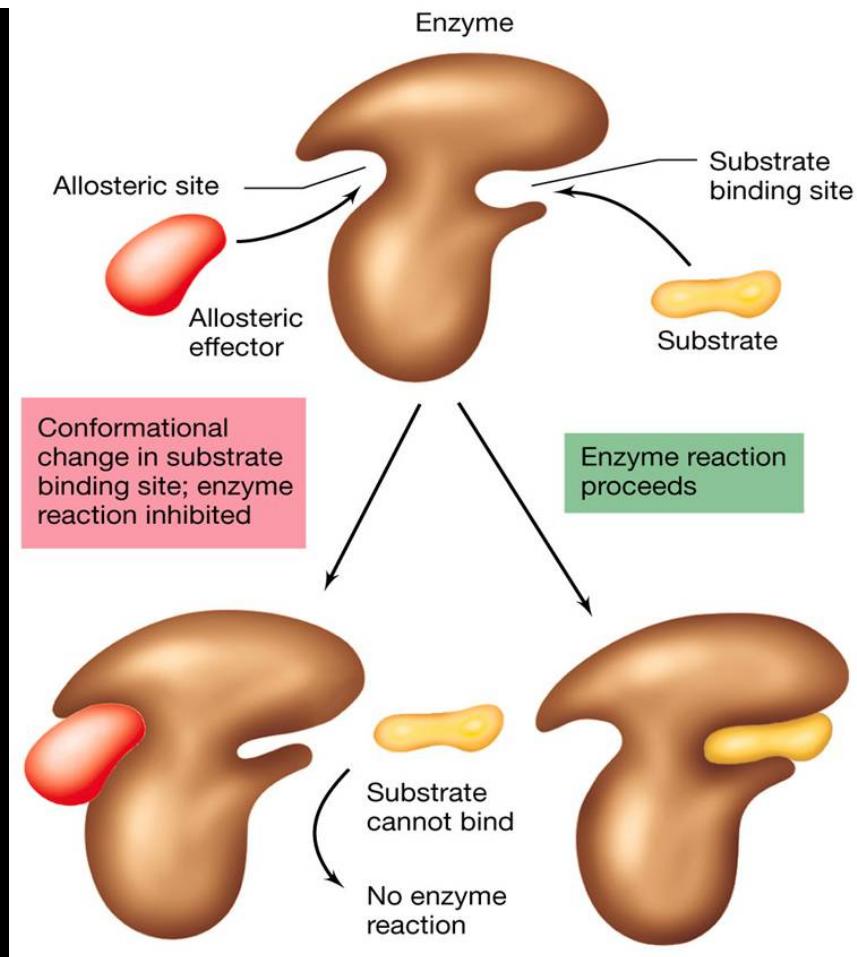
- Myosin structure
- Allostery and myosin kinetic cycle
- Our approach
- Results

# Myosin II

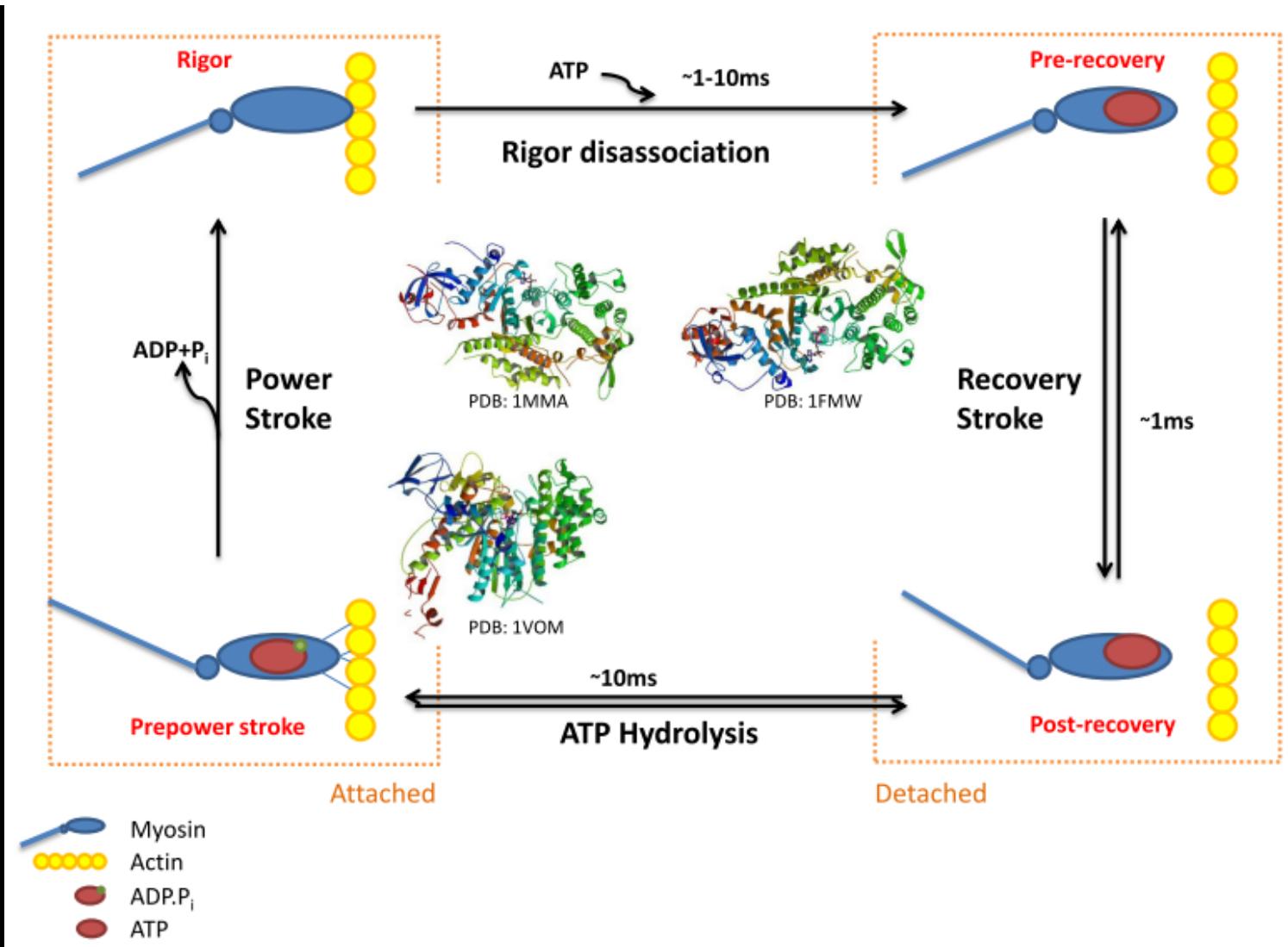


# Allostery

- Allostery may cause conformational change on distant site
- Affinity of substrate binding changes
- Energy transfer from allosteric site to catalytic domain

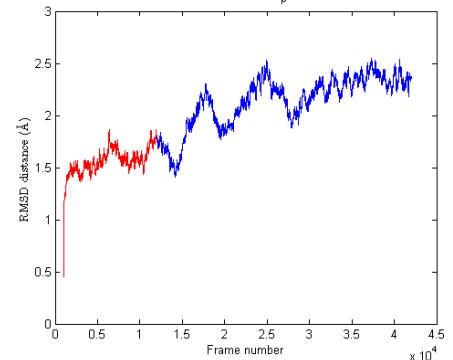
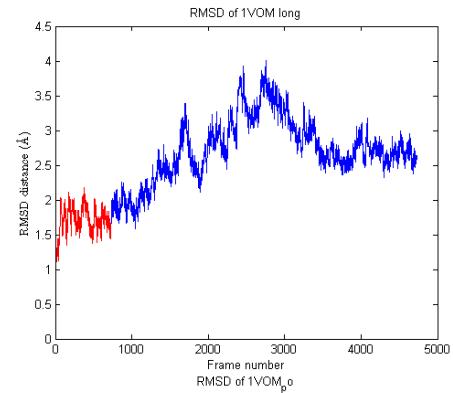
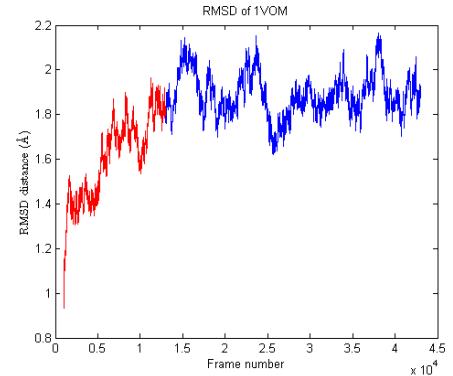


# Myosin kinetic cycle



# MD Simulation

- NAMD
- Charmm 27
- Langevin dynamics
- NPT ensemble
- Explicit solvent (15Å cushion)
- 1fs time steps
  - 50fs frame rate 1VOM & 1VOM\_po (~2ns)
  - 5000fs for 1VOM long (~20ns)



# Our approach


$$\Delta \mathbf{R}_{3N_c,f} \xrightarrow[\text{decomposition}]{\text{modal}} \Delta \mathbf{r}_{3N_c-6,f}$$

$H_\nu$   
 $C_\nu$  } Hermite polynomial for rank  $\nu$

$f_o$  distributions from normal distribution

$f_1$  distributions using normal distribution and Hermite polynomials and constants

PDF (probability density function) estimation for **MD**,  $f_o$  and  $f_1$  samples using **KDE** (Kernel Density Estimation) after transform back into **coordinate space**

Differences between distributions **MD**,  $F_o$  and  $F_1$  measured by KL metric

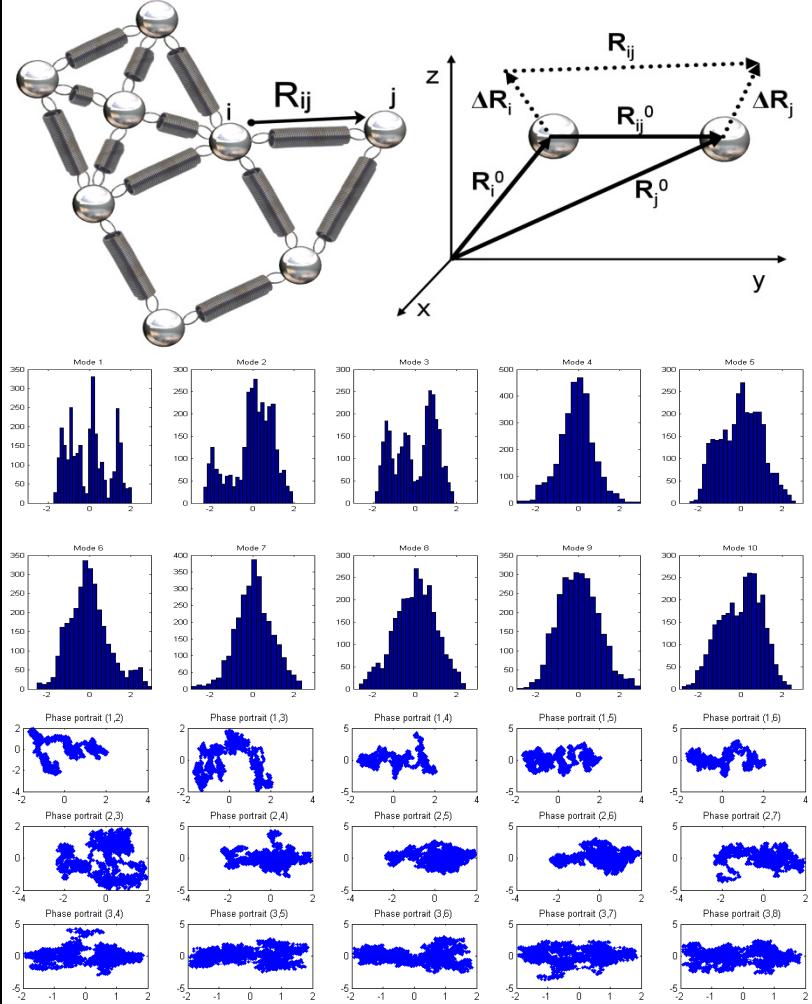
# Modal decomposition

- Transform fluctuations into modal space
- Fluctuations can be seen as multivariate probability distribution function.

$$C = \langle \Delta \mathbf{R} \Delta \mathbf{R}^T \rangle$$

$$\langle \Delta \mathbf{R} \Delta \mathbf{R}^T \rangle^{-\frac{1}{2}} = diag \lambda^{-1/2} \mathbf{e}^T$$

$$\Delta \mathbf{r} = C^{-\frac{1}{2}} \Delta \mathbf{R}$$

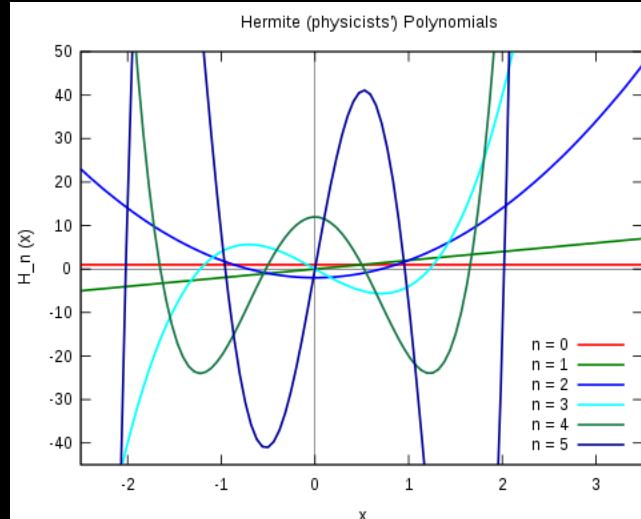


# Hermite expansion

$$f(\Delta \mathbf{r}) = \frac{1}{\sqrt{(2\pi)^N}} e^{-\sum_i \Delta r_i^2/2} [1 + \sum_i \sum_{\nu=3}^{\infty} \frac{1}{\nu!} \langle H_{\nu}(\Delta r_i) \rangle H_{\nu}(\Delta r_i) \\ + \sum_{i \neq j} \sum_{\nu=3}^{\infty} \frac{1}{\nu!} \sum_{p=1}^{\nu-1} \binom{\nu}{p} \langle H_p(\Delta r_i) H_{\nu-p}(\Delta r_j) \rangle H_p(\Delta r_i) H_{\nu-p}(\Delta r_j) + \\ \sum_{i \neq j \neq k} \dots]$$

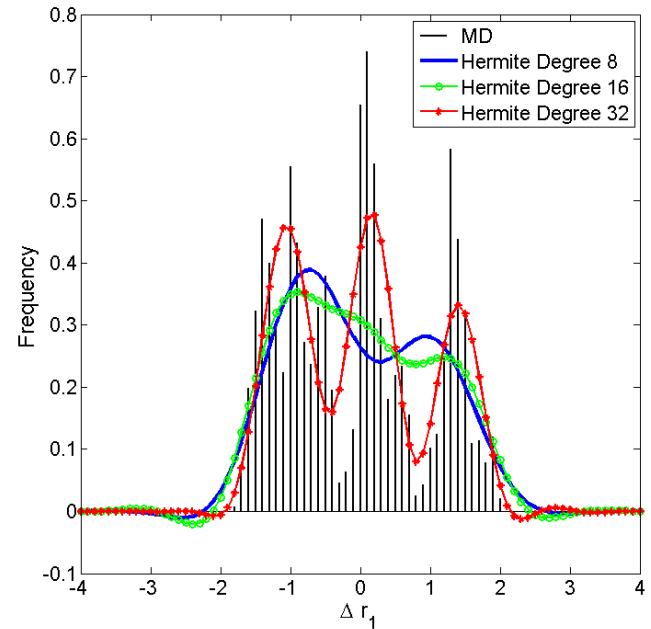
- Hermite polynomials are orthogonal w.r.t. weight function
- $H_i$  hermite polynomial rank  $i$
- Recursion equation

$$H_{n+1}(x) = xH_n(x) - H_n'(x)$$



# Harmonic model & Anharmonic corrections

- Samples from  $f_0$  (harmonic)
- Samples  $f_1$  (anharmonic model) obtained using  $f_0$  and first degree corrections
- Random sampling of  $f_1$  by rejection sampling



$$f_0(\Delta \mathbf{r}) = \frac{1}{\sqrt{(2\pi)^N}} e^{-\sum_i \Delta r_i^2/2} \quad f_1(\Delta \mathbf{r}) = \frac{1}{\sqrt{(2\pi)^N}} e^{-\sum_i \Delta r_i^2/2}$$
$$\times \prod_i [1 + \sum_i \sum_{\nu=3}^{\infty} \frac{1}{\nu!} \langle H_\nu(\Delta r_i) \rangle H_\nu(\Delta r_i)]$$

# Kernel Density Estimation

- Samples in modal space transform back into real space
- KDE used in order to estimate pdf of the distributions

$$\hat{f}_h(x) = \frac{1}{nh} \sum_{i=1}^n K\left(\frac{x - x_i}{h}\right)$$

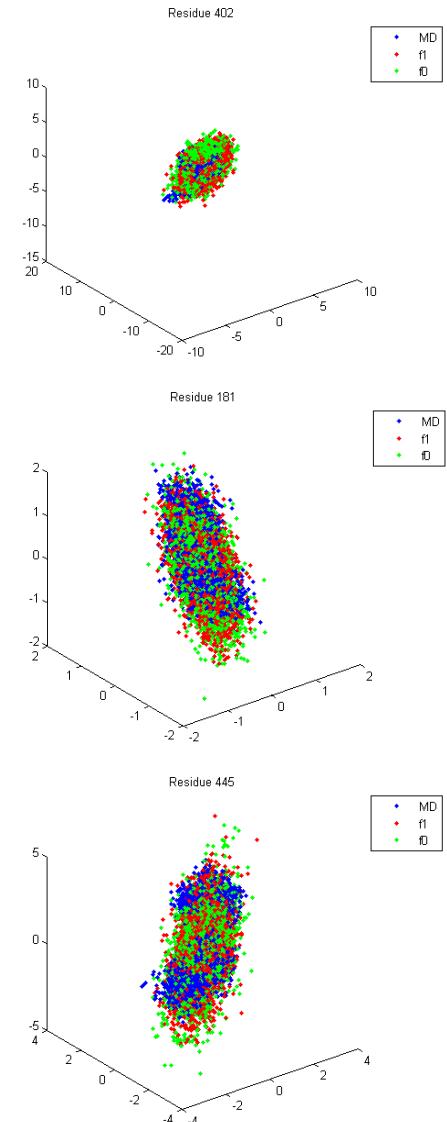
$$h = \left(\frac{4\hat{\sigma}^5}{3n}\right)^{1/5} \approx 1.06\hat{\sigma}n^{-1/5}$$

# KL Divergence

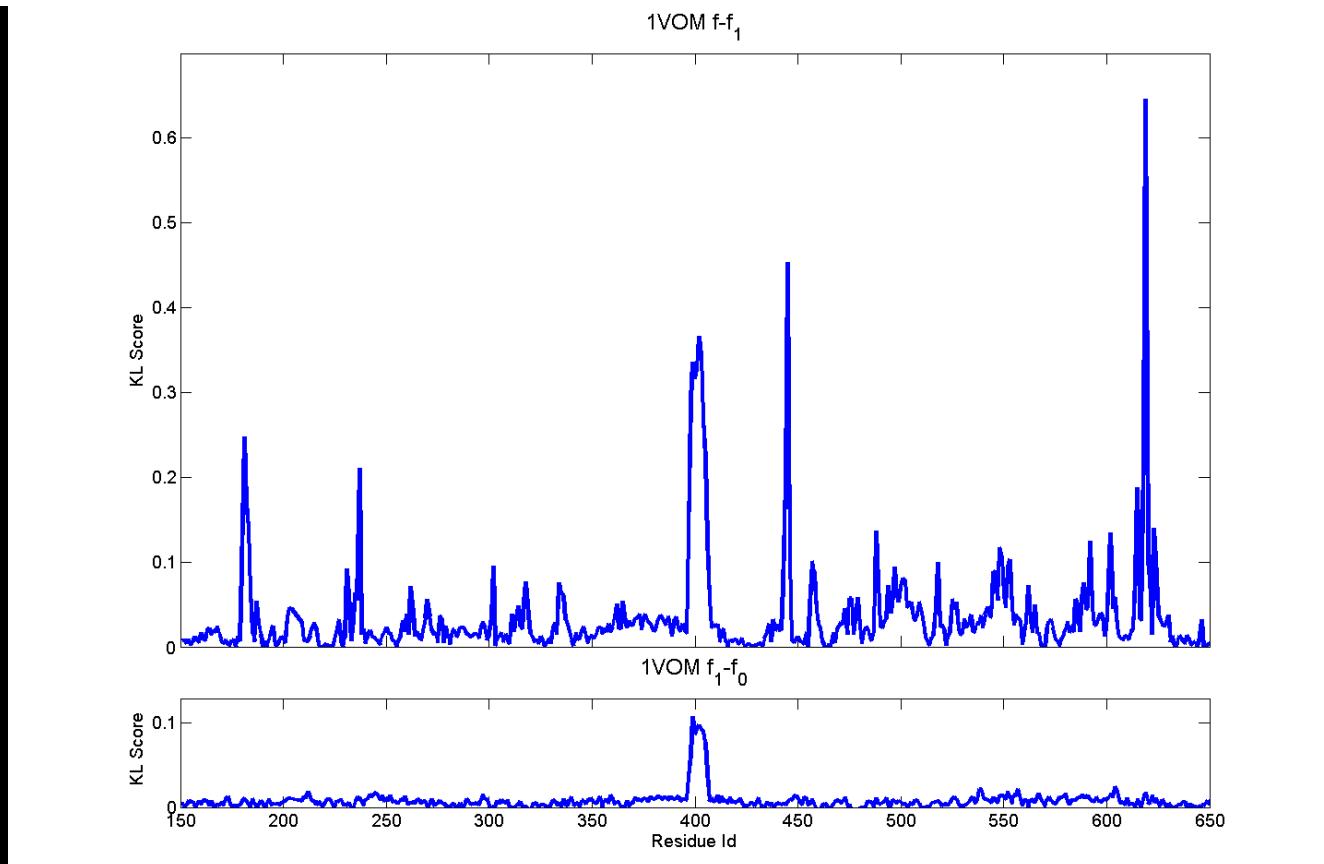
$$D_{KL}(P \parallel Q) = \sum_i P(i) \ln \frac{P(i)}{Q(i)}$$

- Kullback-Leibler divergence measure of two pdf
- P is true data (MD), Q is model distributions

$$KL_{score} = KL_x + KL_y + KL_z$$

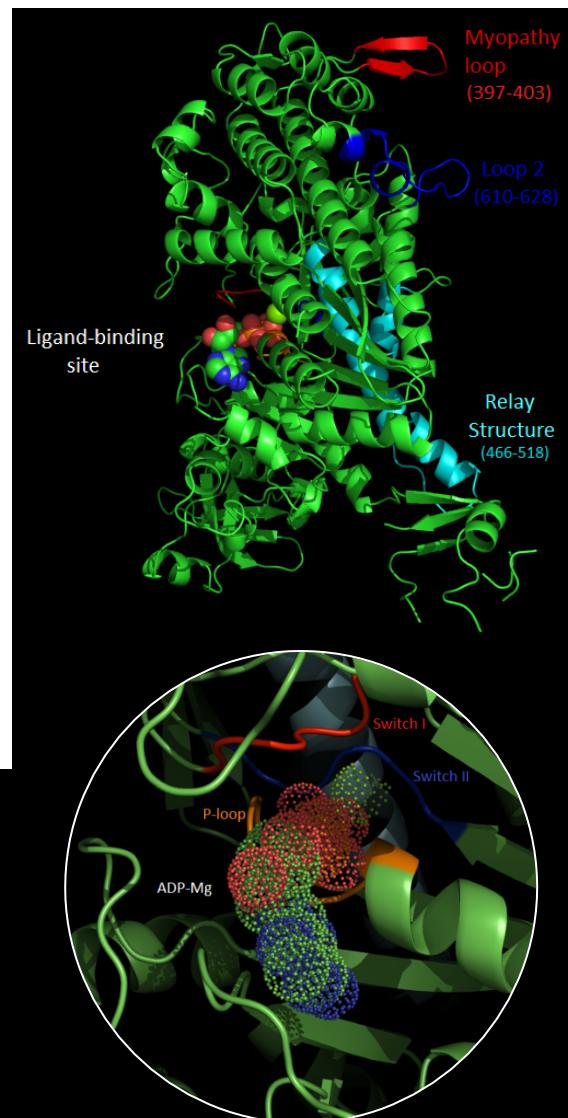
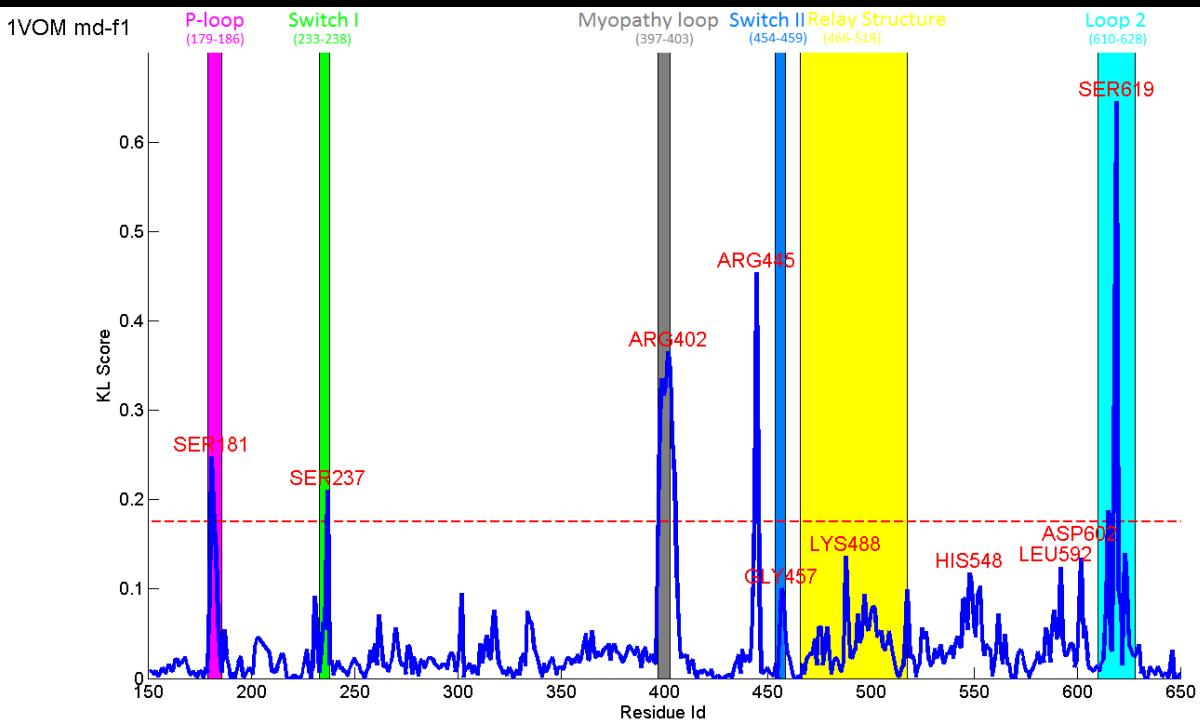


# Anharmonicity vs. Mode-coupling



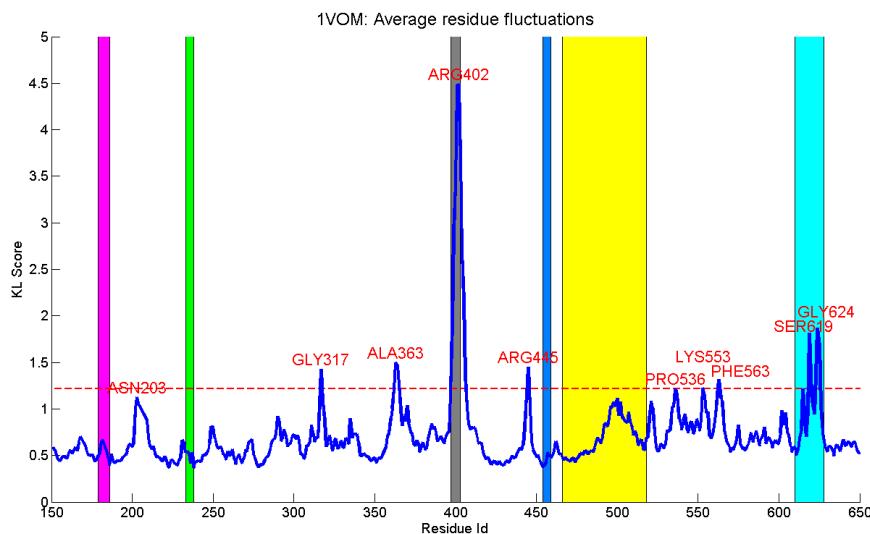
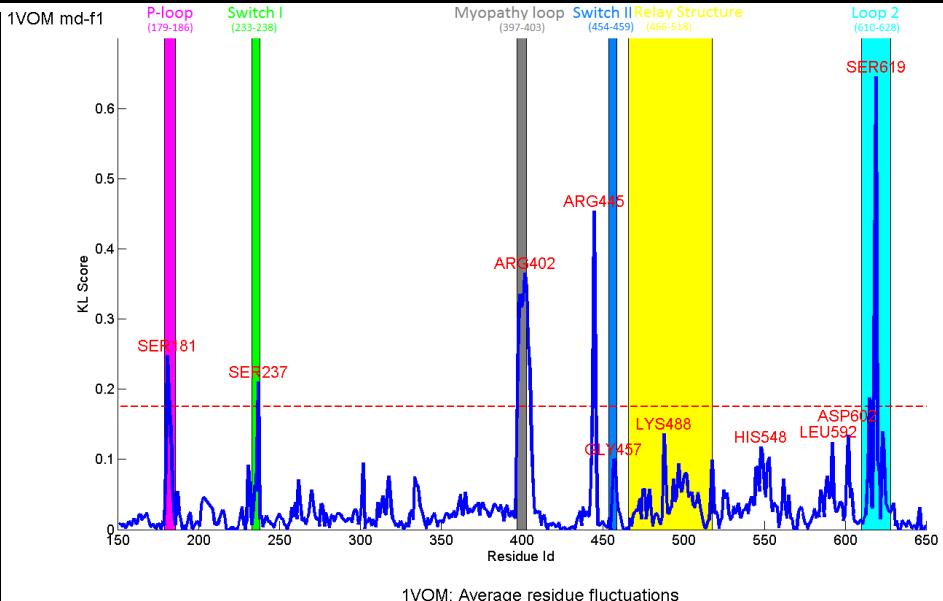
- Contribution of the mode-coupling is greater than anharmonicity which is also observed before in entropy change

# Functionally Important Sites



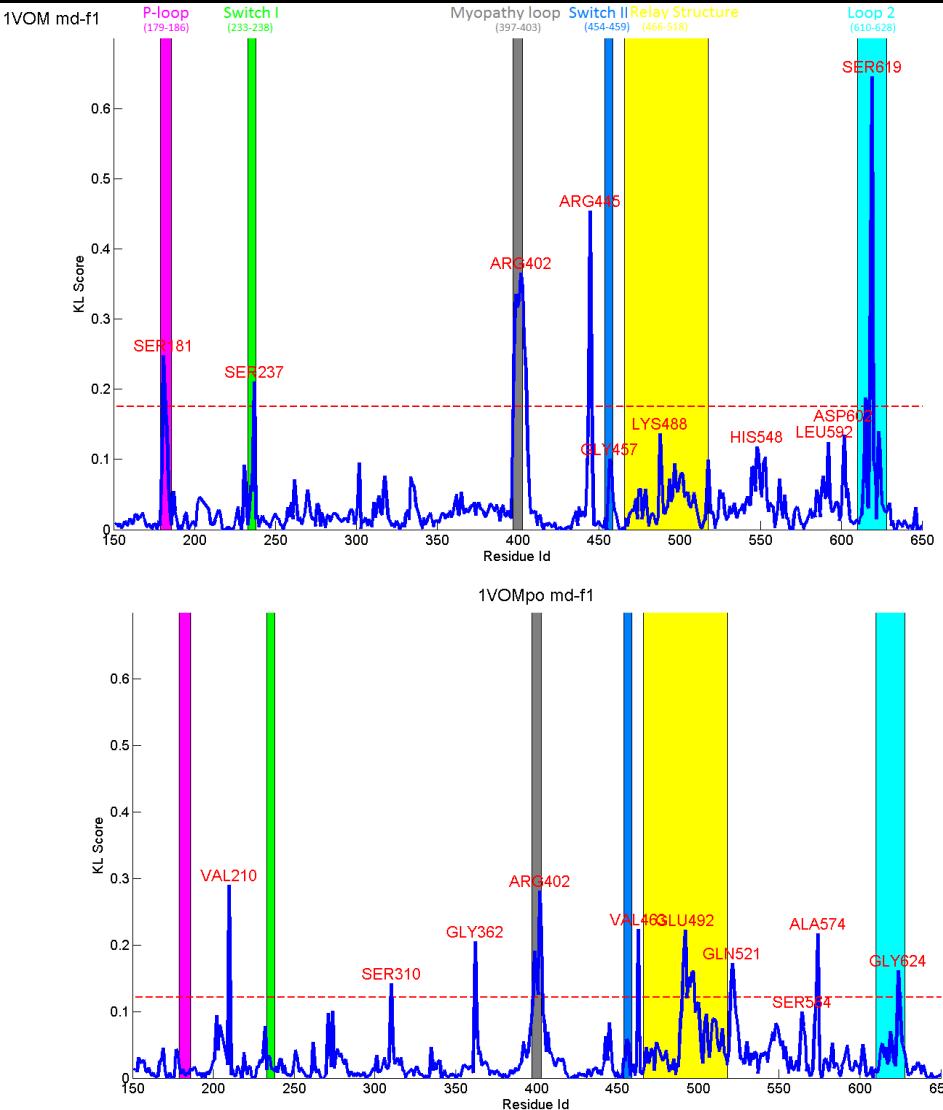
# Compare with fluctuations

- Fluctuations inform about residue displacement
- Mostly on surface loop regions



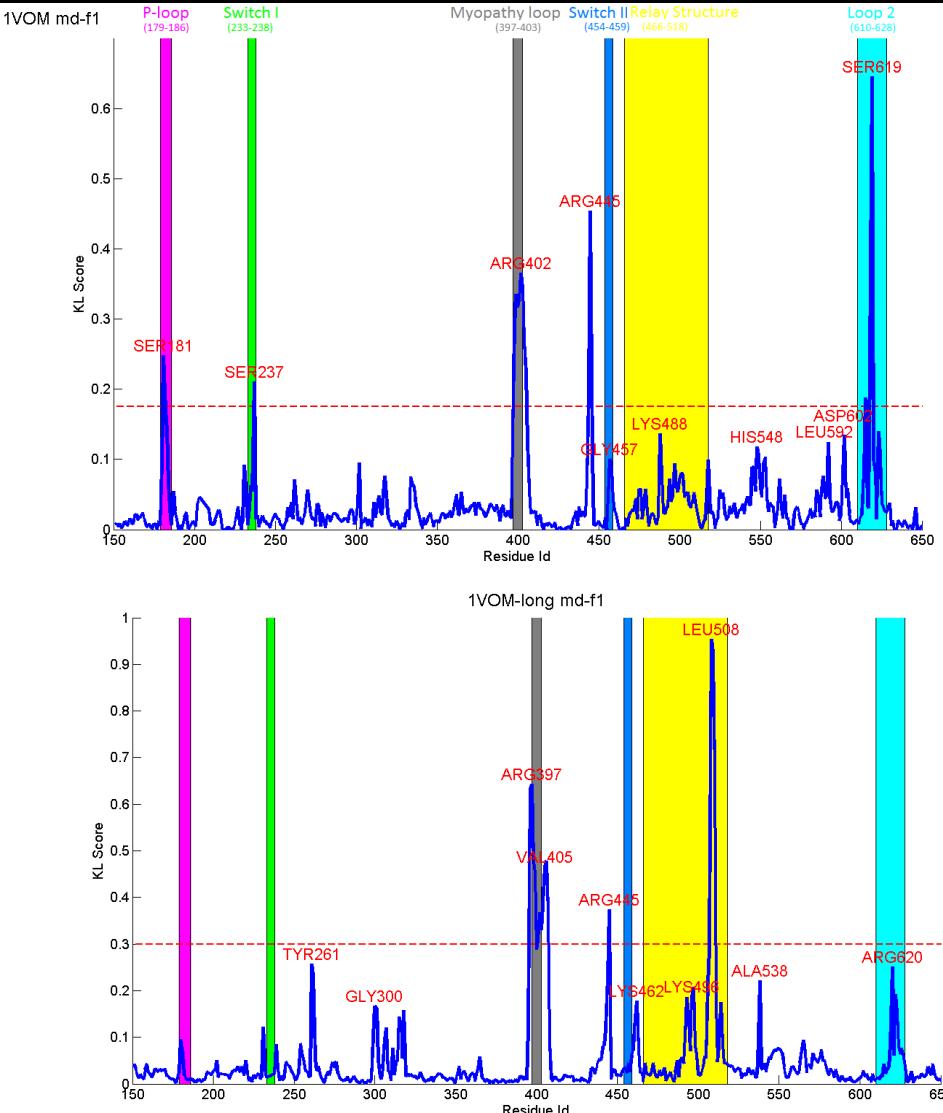
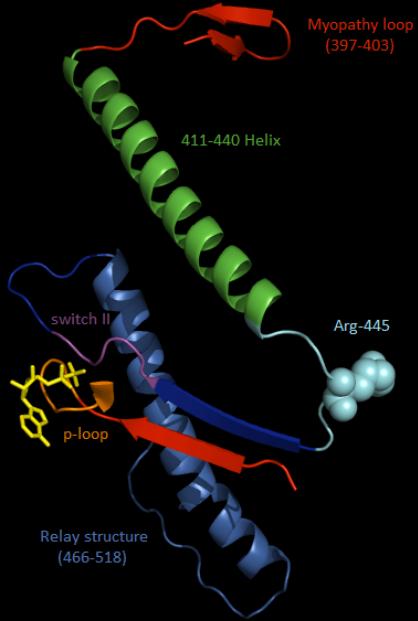
# Ligand bound vs. unbound

- Ligand bonding compress activation of other sites
- Ligand release reduces activation on actin-binding sites



# Longer simulation

- Long simulation shows characteristics of near-rigor state
- Central beta-sheets and relay structure drive transition

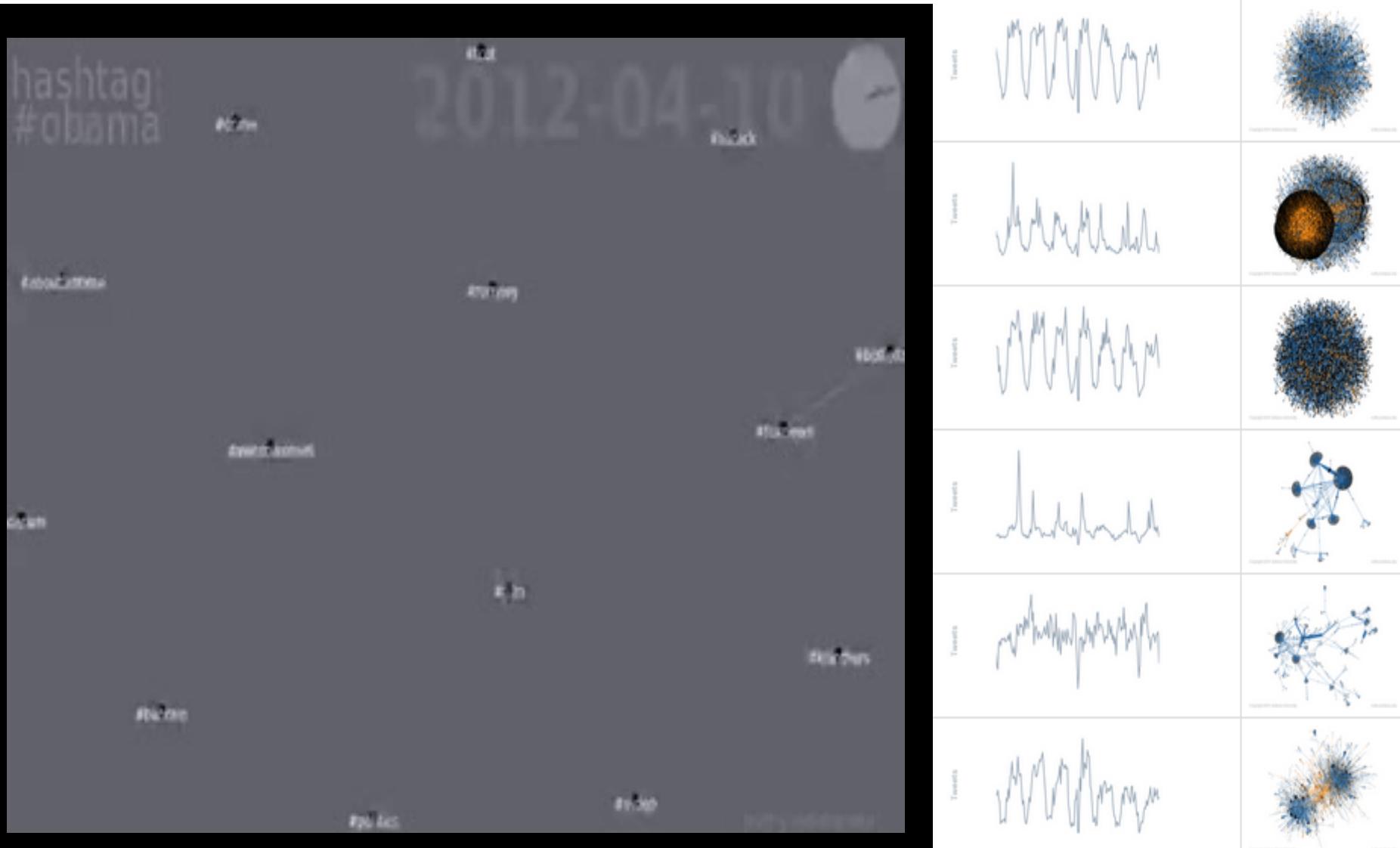


# Conclusion

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- Mode-coupling contribution highlight functionally important sites
- Ligand binding effect suppress nonspecific residues and points binding sites
- Importance of fluctuation analysis shown

# Interacting systems



# Thank you !!

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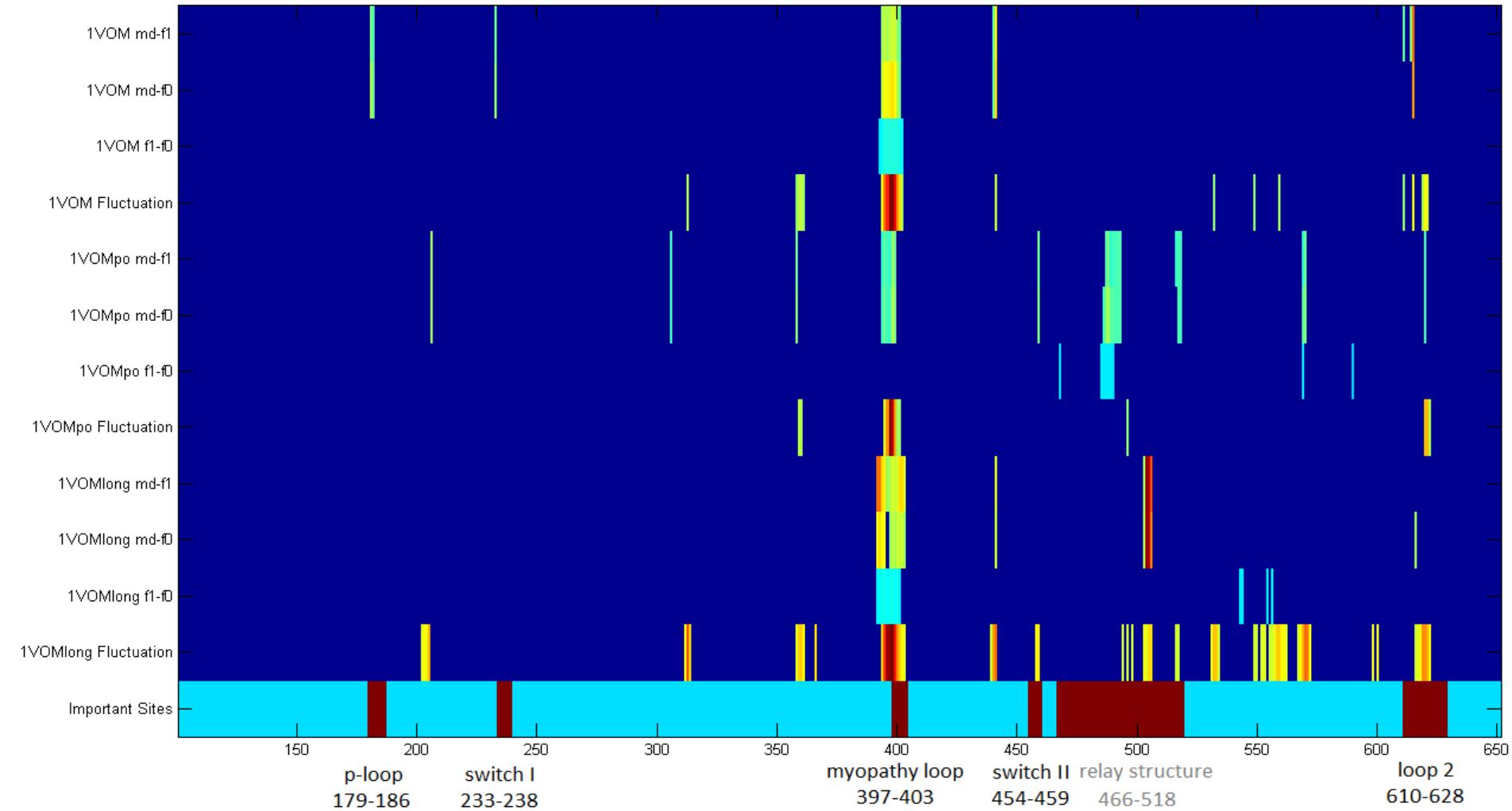
Questions?



# Backup slides

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# Precision table



# Hermite

$$f(\Delta \mathbf{r}) = \frac{1}{\sqrt{(2\pi)^{3N}}} e^{-\frac{1}{2} \sum_{i=1}^{3N} \Delta r_i^2} \left[ 1 + \sum_{\nu=3}^{\infty} \mathbf{C}_{\nu} \cdot \mathbf{H}_{\nu}(\Delta \mathbf{r}) \right]$$

$$H_{\nu}^{ij..k}(\Delta \mathbf{r}) = \frac{(-1)^{\nu}}{g(\Delta \mathbf{r})} \nabla^{ij..k} g(\Delta \mathbf{r})$$

Tensorial hermite polynomials can be obtained by successive differentiation using Rodrigues' formula

$$\mathbf{C}_{\nu} = \frac{1}{\nu!} \int_{-\infty}^{\infty} \mathbf{H}_{\nu}(\mathbf{x}) f(\Delta \mathbf{r}) d\Delta \mathbf{r} = \langle \mathbf{H}_{\nu}(\Delta \mathbf{r}) \rangle / \nu! \quad \text{Orthogonality relation}$$

Value of tensor element does not depend on the order of indices due to commutativity of the gradient operator

$$\nabla_k \nabla_l - \nabla_l \nabla_k = 0$$

$$\mathbf{H}_{\nu}^{i_1 i_2 \dots i_{\nu}}(\Delta \mathbf{r}) = \mathbf{H}_{\nu}^p(\Delta r_k, \Delta r_l)$$

where  $p$  is the number of indices equal to  $k$  (the remaining  $\nu-p$  indices equal to  $l$ )

# Hermite-2

Covariance matrix in the normal basis is diagonal

$$\mathbf{H}_\nu^p(\Delta \mathbf{r}) = H_p(\Delta r_1) \times H_{\nu-p}(\Delta r_2)$$

$$f(\Delta \mathbf{r}) = \frac{1}{\sqrt{(2\pi)^N}} e^{-\sum_i \Delta r_i^2/2} [1 + \sum_i \sum_{\nu=3}^{\infty} \frac{1}{\nu!} \langle H_\nu(\Delta r_i) \rangle H_\nu(\Delta r_i) \\ + \sum_{i \neq j} \sum_{\nu=3}^{\infty} \frac{1}{\nu!} \sum_{p=1}^{\nu-1} \binom{\nu}{p} \langle H_p(\Delta r_i) H_{\nu-p}(\Delta r_j) \rangle H_p(\Delta r_i) H_{\nu-p}(\Delta r_j) + \\ \sum_{i \neq j \neq k} \dots]$$

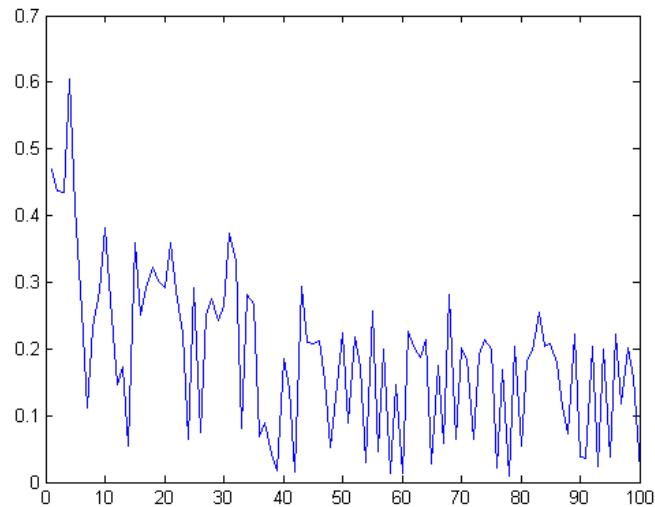
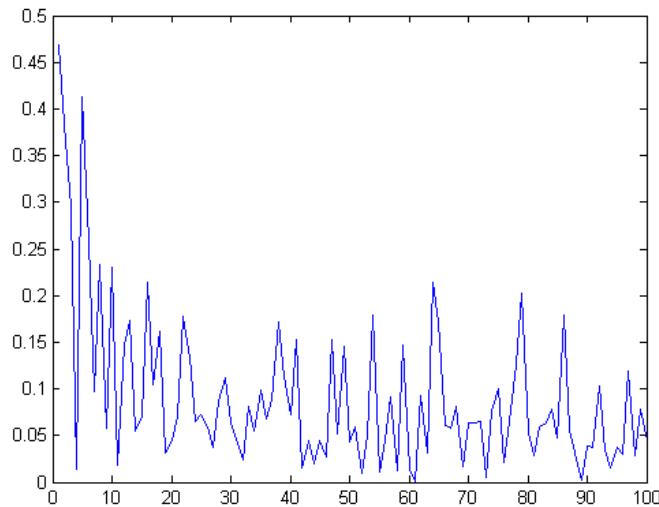
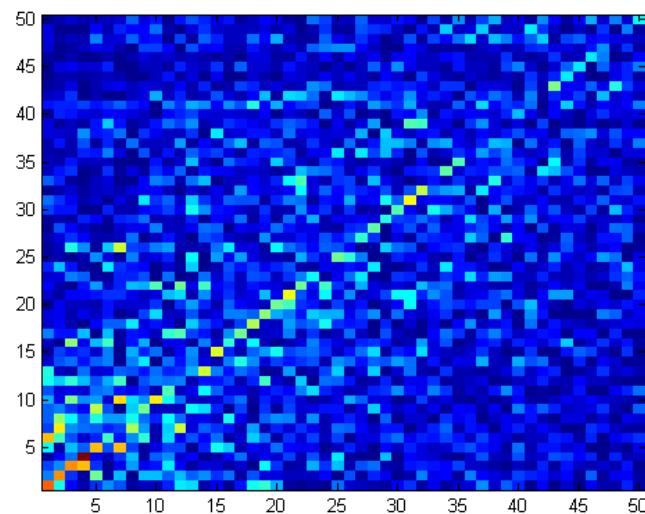
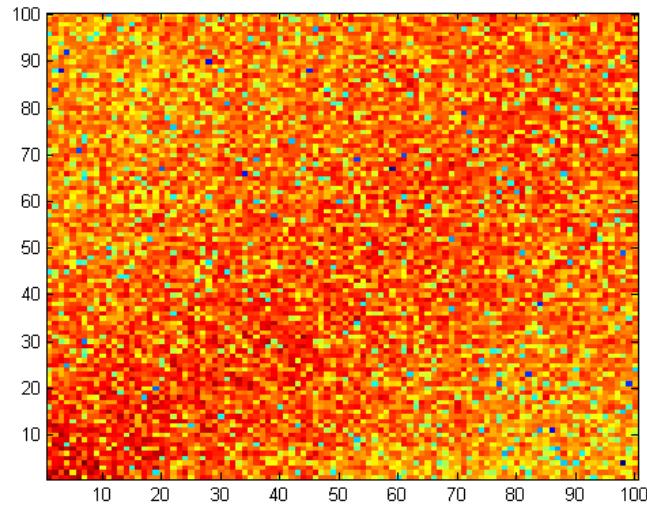
$$f_1(\Delta \mathbf{r}) = \frac{1}{\sqrt{(2\pi)^N}} e^{-\sum_i \Delta r_i^2/2} \\ \times \prod_i [1 + \sum_i \sum_{\nu=3}^{\infty} \frac{1}{\nu!} \langle H_\nu(\Delta r_i) \rangle H_\nu(\Delta r_i)]$$

The difference between full pdf and the approximation  $f_1$  is the mode-coupling correction such as  $\langle H_p(\Delta r_i) H_{\nu-p}(\Delta r_j) \rangle - \langle H_p(\Delta r_i) \rangle \langle H_{\nu-p}(\Delta r_j) \rangle \neq 0$

$$f(\Delta r_i) \equiv \int_0^\infty \prod_{j \neq i} d\Delta r_j f(\Delta \mathbf{r})$$

Marginal distributions are transparent to such Corrections as a merit of the orthogonality relation.

# Long-short simulation consistency



# Experimentally verified residues



233 236 237 238 403 405 454 456 457 458 459 464 465 467 470 472 473 475 481  
482 487 494 499 501 506 531 536 538 548 562 586 624 664 680 691 692 740 746