

# <sup>1</sup> NGSBEM: potential operators for FEM-BEM couplings in NETGEN/NGSolve

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## Software

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## <sup>6</sup> Summary

<sup>7</sup> NGSBEM is the dedicated boundary element method (BEM) module within NETGEN/NGSolve.  
<sup>8</sup> It provides the full set of tools needed to formulate and solve boundary integral equations,  
<sup>9</sup> including potential operators and their Galerkin discretizations. Seamlessly integrated with the  
<sup>10</sup> NGSolve core, NGSBEM enables efficient numerical solutions of linear interior and exterior  
<sup>11</sup> boundary value problems based on the BEM.

<sup>12</sup> Key features of NGSBEM, in conjunction with NGSolve, include:

- Use of finite element spaces already available in NGSolve, including high-order elements on curved surfaces.
- A user-friendly Python interface.
- Declarative programming for problem formulation and solution.
- Standard potential operators for electrostatics, acoustics, and electromagnetics, allowing both scalar and vector-valued densities to be discretized.
- Support for multi-domain problems through FEM–BEM coupling.
- Acceleration via a Multi-Level Fast Multipole Method (MLFMM) for approximating the arising Galerkin matrices.
- Kernel-driven implementations of potential operators, enabling users to add custom operators.
- Extensive documentation and demonstrations supporting users as they start working with.

## <sup>26</sup> Statement of Need

<sup>27</sup> Computational engineering frequently involves solving boundary value problems (BVPs),  
<sup>28</sup> including linear and nonlinear partial differential equations, as well as interior, exterior, and  
<sup>29</sup> transmission problems. While the finite element method (FEM) excels at handling interior  
<sup>30</sup> problems with nonlinearities, the BEM is the method of choice for solving linear second-order  
<sup>31</sup> problems in unbounded domains.

<sup>32</sup> A unified framework that naturally combines both approaches is rare. NETGEN/NGSolve  
<sup>33</sup> addresses this gap by providing a complete environment for Galerkin BEM and its coupling  
<sup>34</sup> with FEM.

## <sup>35</sup> Usability, Accuracy, and Efficiency

<sup>36</sup> The BEM is attractive due to its dimensional reduction and high accuracy. However, it also  
<sup>37</sup> introduces challenges, such as singular integral operators and dense matrices. NGSBEM is

<sup>38</sup> designed to resolve these difficulties internally, providing robust and efficient implementations  
<sup>39</sup> while sparing users from the underlying complexity.

#### <sup>40</sup> Usability

<sup>41</sup> NGSolve follows the philosophy of allowing users to “type formulas as they appear in the book.”  
<sup>42</sup> NGSBEM extends this philosophy to the BEM: users express boundary integral operators  
<sup>43</sup> declaratively in Python, in a syntax that mirrors their mathematical formulation. This clarity  
<sup>44</sup> and directness in representing operators is a central design goal.

<sup>45</sup> Consider, for example, the Dirichlet problem

$$\begin{cases} \Delta\phi = 0, & \text{in } \Omega, \\ \gamma_0\phi = m, & \text{on } \Gamma. \end{cases}$$

<sup>46</sup> The representation formula yields:

$$\phi(\mathbf{x}) = \text{LaplaceSL}(j)(\mathbf{x}) - \text{LaplaceDL}(m)(\mathbf{x}). \quad (1)$$

<sup>47</sup> To compute the unknown Neumann data  $j$ , one solves the Galerkin discretization of the single  
<sup>48</sup> layer operator. With trial functions  $u_j$  and test functions  $v_i$ , the matrix entries are

$$V_{ij} = \int_{\Gamma} [\text{LaplaceSL}(u_j)] v_i \, ds.$$

<sup>49</sup> In NGSBEM, the matrix is assembled as:

```
fesL2 = SurfaceL2(mesh, order=3, dual_mapping=True)
u, v = fesL2.TnT()
V = LaplaceSL(u * ds) * v * ds
```

<sup>50</sup> Given  $V$ , the linear system

$$Vj = b$$

<sup>51</sup> is solved using a standard iterative solver.

<sup>52</sup> A second example illustrates the hypersingular operator arising for the Neumann problem. Its  
<sup>53</sup> Galerkin matrix is given by

$$D_{ij} = \int_{\Gamma} [\text{LaplaceSL}(\text{curl}_{\Gamma} u_j)] \text{curl}_{\Gamma} v_i(\mathbf{x}) \, ds_x.$$

<sup>54</sup> implemented simply as

```
D = LaplaceSL(curl(u) * ds) * curl(v) * ds
```

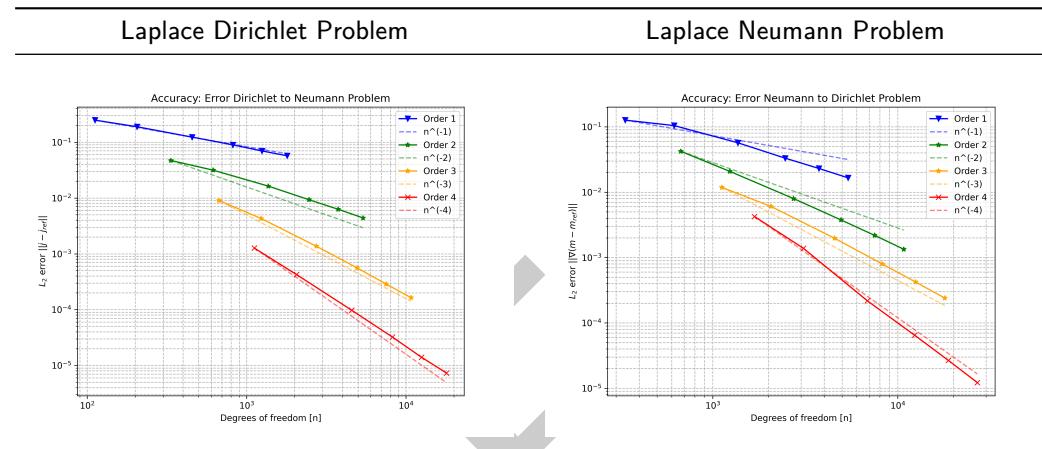
<sup>55</sup> The linear system with Galerkin matrix  $D$  is solved using a standard iterative solver.

<sup>56</sup> Single and double layer operators for the Helmholtz and Maxwell equations are provided  
<sup>57</sup> analogously.

## 58 Accuracy

59 Accurate integration is central to the BEM, particularly because singular kernels must be  
60 treated with care. NGSBEM employs the specialized quadrature rules of Sauter and Schwab  
61 ([Sauter & Schwab, 2011](#)), which support curved elements and enable high-order convergence.

62 The numerical experiments below compare the computed Cauchy data with the known exact  
63 solution for the Laplace equation. The results demonstrate optimal convergence rates.



64 Beyond scalar-valued problems, NGSBEM supports vector finite elements and BEM operators  
65 for electromagnetics, with numerical evidence of optimal convergence documented in ([Wegeler,](#)  
66 [2025](#)).

67 Another challenge is the accurate evaluation of potentials, especially in the near field. NGSBEM  
68 implements techniques developed by Gumerov and others ([Nail A. Gumerov & Duraiswami,](#)  
69 [2021](#); [Kaneko et al., 2023](#)). Depending on the geometric configuration, triangles are either  
70 subdivided or a line-based approximation is applied, ensuring stable and accurate potential  
71 evaluation.

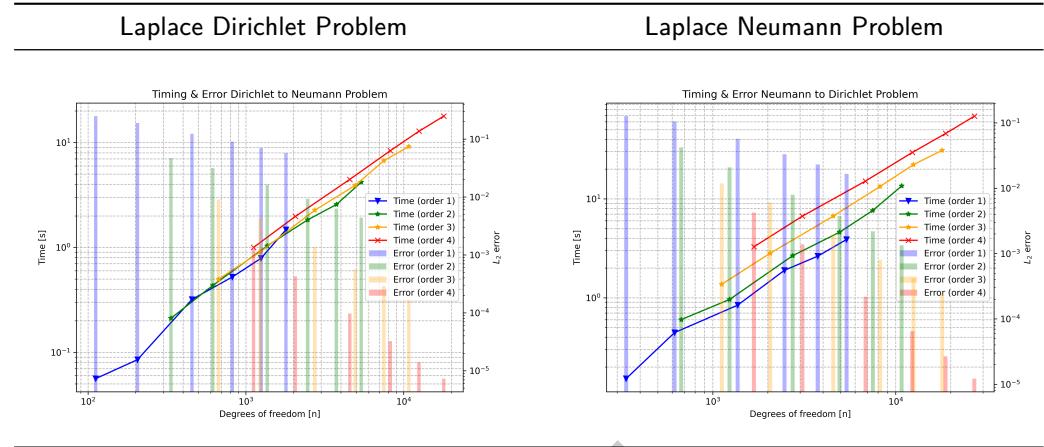
## 72 Efficiency

73 Galerkin matrices are generated on the fly using the multi-level fast multipole method (MLFMM)  
74 of Greengard and Rokhlin ([Greengard & Rokhlin, 1987](#)), with efficient computation of translation  
75 and rotation coefficients following the work of Gumerov and Duraiswami ([Dorigo, 2025](#); [Nail](#)  
76 [A. Gumerov & Duraiswami, 2001](#); [Nail A. Gumerov & Duraiswami, 2015](#)).

77 Performance is further enhanced by:

- 78 ▪ the use of AVX vectorization,
- 79 ▪ full parallelization across all major computational stages.

80 To assess efficiency, we compare accuracy and runtime across different polynomial orders.  
81 Higher-order elements provide significantly better accuracy at comparable runtime, despite the  
82 increased number of unknowns.



- 83 For example, in Laplace Dirichlet Problem with efficiency plot shown on the left:
- 84     ■ order 1 at 1250 DOFs (0.789 s runtime, error 0.07), vs.,  
 85     ■ order 4 at 1120 DOFs (1.000 s runtime, error 0.001)
- 86 This corresponds to an error reduction of approximately 98.6% at comparable runtime.
- 87 Similar improvements appear across other configurations, demonstrating the practical efficiency  
 88 of higher-order BEM.
- 89 All numerical experiments were run on a MacBook Pro equipped with an Apple M4 Pro (14  
 90 cores: 10 performance + 4 efficiency) and 48 GB of RAM. Timings were measured on the  
 91 host machine under macOS.

## 92 Example

93 Solving a PDE with the BEM proceeds in two steps:

- 94     1. solve an integral equation for the unknown boundary trace(s),  
 95     2. evaluate the representation formula to obtain the solution in the domain.

96 Consider a plate capacitor:

---

$$\Omega^c$$

$$-\Delta\phi = 0, \quad \text{in } \Omega^c,$$

$$\gamma_0\phi = m, \quad \text{on } \Gamma,$$

$$\lim_{\|x\| \rightarrow \infty} \phi(x) = \mathcal{O}\left(\frac{1}{\|x\|}\right), \quad \|x\| \rightarrow \infty.$$



---

97 Using [Equation 1](#), the unknown Neumann data  $j$  is computed via a BEM system:

```
j = GridFunction(fesL2)

# generate potential operators V, M and K
with TaskManager():
    SLPotential = LaplaceSL(u*ds(bonus_intorder=3))
    DLPotential = LaplaceDL(uH1*ds(bonus_intorder=3))
    V = SLPotential*v*ds(bonus_intorder=3)
```

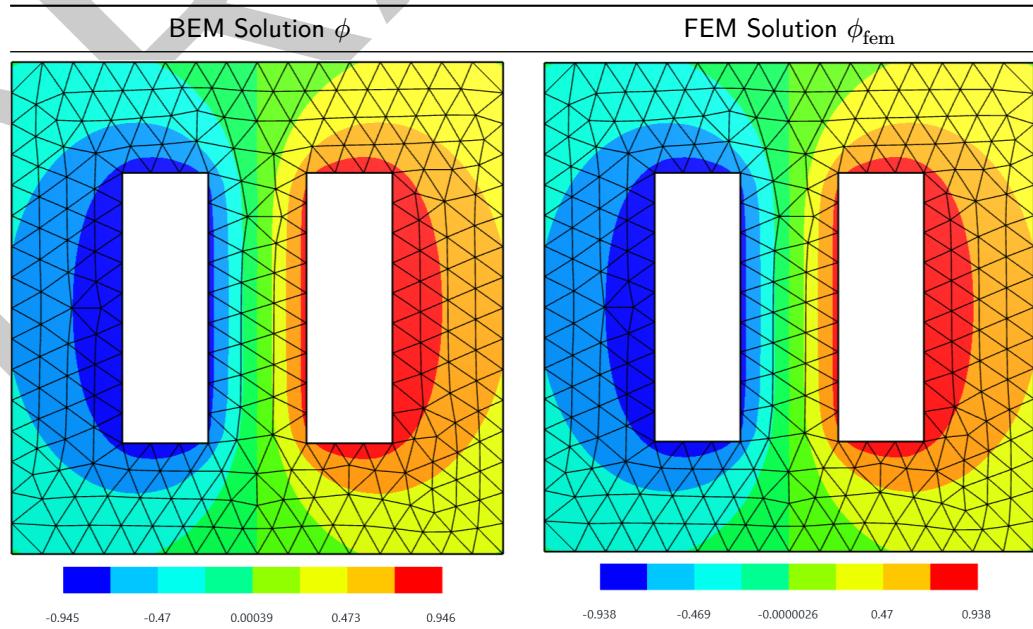
```
K = DLPotential*v*ds(bonus_intorder=3)
M = BilinearForm(uH1*v*ds(bonus_intorder=3)).Assemble()
```

```
# solve the linear system  $V \ j = (-1/2 M + K) m$ 
with TaskManager():
    pre = BilinearForm(u.Trace() * v.Trace() * ds, diagonal=True).Assemble().mat.Inverse()
    rhs = ((-0.5 * M.mat + K.mat) * m).Evaluate()
    CG(
        mat = V.mat, pre=pre, rhs = rhs, sol=j.vec, tol=1e-8, maxsteps=200,
        initialize=False, printrates=False
    )
```

98 Afterwards, the potential  $\phi$  can be evaluated anywhere. For example, on a screen cutting  
 99 through the plates:

```
screen = WorkPlane(
    Axes((0,0,0), X, Z)).RectangleC(4, 4).Face()
    - Box((-1.1,-1.1,0.4), (1.1,1.1,1.1))
    - Box((-1.1,-1.1,-1.1), (1.1,1.1,-0.4))
)
mesh_screen = Mesh(OCCGeometry(screen).GenerateMesh(maxh=0.25)).Curve(1)
fes_screen = H1(mesh_screen, order=3)
phi_on_screen = GridFunction(fes_screen)
print ("ndofscreen=", fes_screen.ndof)
with TaskManager():
    phi_on_screen.Set(
        -SLPotential(j)+DLPotential(utop)+DLPotential(ubot),
        definedon=mesh_screen.Boundaries(".*"), dual=False
    )
Draw (gf_screen)
```

100 The BEM solution (left) and a FEM reference solution (right) are shown here:



101 Because the problem is exterior, the FEM requires either infinite elements or an artificial  
 102 boundary with prescribed data. The BEM, in contrast, handles the unbounded domain

103 naturally and without mesh refinement at edges or corners. The maximum discrepancy between  
104 FEM and BEM is around 2.5%, occurring near corner singularities.  
105 Further examples, including FEM–BEM coupling, mixed formulations, and acoustic or electro-  
106 magnetic scattering, can be found in the online documentation ([Weggler, 2025](#)).  
107 Details on the theoretical background for high order BEM for electrostatics and electromagnetics,  
108 can be found in ([Weggler, 2011, 2012](#)).

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