

¹ INS_IEP: A MATLAB package for fitting peaks of Inelastic Neutron Scattering data in spin cluster systems

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⁷ Summary

⁸ Inelastic neutron scattering (INS) is a spectroscopic technique that can be used to measure
⁹ the magnetic excitations in materials with interacting electron spins. For samples composed of
¹⁰ finite-size clusters of magnetic moment-carrying atoms, such as single ions or molecular-based
¹¹ magnets, INS experiments yield inelastic excitation with energies that correspond to the
¹² eigenvalues of the spin Hamiltonian of the material being studied. Fitting a model spin system
¹³ to these experimental eigenvalues can be formulated as an inverse eigenvalue problem (IEP),
¹⁴ where the matrix formed is the spin Hamiltonian operator of the sample molecule. Solving
¹⁵ this IEP is less computationally expensive than fitting the full INS data, so can be used as an
¹⁶ initial proxy for the full fitting problem. **INS_IEP** is a MATLAB package that uses deflated
¹⁷ numerical optimisation methods to systematically find multiple solutions to this IEP. The
¹⁸ package requires and is fully compatible with **easyspin** ([Stoll & Schweiger, 2006](#)), a package
¹⁹ for solving fitting problems in electron paramagnetic resonance (EPR).

²⁰ Statement of need

²¹ Neutrons are an excellent bulk probe of material properties since they carry no charge and
²² therefore penetrate deeply into matter. Neutrons also carry a quantum spin of a half, making
²³ them a sensitive probe of magnetism ([Squires, 2012](#)). Reactors and spallation sources with
²⁴ dedicated high-flux neutron sources serve the international research community with neutron
²⁵ scattering experiment capabilities for material research. Inelastic neutron scattering (INS) is
²⁶ one such experimental technique that can be used to study magnetism. In an INS experiment,
²⁷ a sample under investigation is irradiated with a beam of neutrons and the scattered neutron
²⁸ energy and momentum transfer are detected. For samples composed of finite-size clusters of
²⁹ magnetic moment-carrying atoms, such as single ions or molecular-based magnets, the detected
³⁰ neutron energy transfer gives direct access to the quantum spin excitations ([M. Baker & Mutka,](#)
³¹ [2012](#); [Furrer & Waldmann, 2013](#); “[Spectroscopy Methods for Molecular Nanomagnets](#),” [2014](#)).

³² The energy of such excitations relates to the energy difference between eigenvalues of the
³³ Hamiltonian matrix that describes the quantum spin dynamics of the compound in question.
³⁴ Single-ion and molecular-based magnets are studied as prototype components (quantum bits,
³⁵ sensors) for quantum technologies. INS can therefore provide crucial information concerning
³⁶ the precise quantum properties of such systems. However, to relate the INS experimental
³⁷ results to the Hamiltonian that describes quantum spin dynamics requires parameterisation of
³⁸ matrix elements such that a set of eigenvalues and eigenstates matching the experiment are
³⁹ determined. This situation is known as the inverse eigenvalue problem.

⁴⁰ To date, this problem is addressed in an iterative process where parameters of the Hamiltonian

41 are varied manually, often one at a time, and the resultant eigenvalues compared to the
 42 experimental values - each such iteration requires an eigendecomposition of the Hamiltonian
 43 matrix. INS_IEP presents an elegant solution to solving this problem, using algorithms to
 44 calculate multiple parameter sets that minimise the difference in eigenvalues, reducing the
 45 number of Hamiltonian matrix diagonalisations, and providing a more robust method to reliably
 46 extract an accurate spin Hamiltonian model from INS experimental data.

47 Current State of the Art

48 There currently exists fitting software designed to solve this problem, such as:

- 49 ■ SPECTRE ([A. T. Boothroyd, 1990](#)) uses the NAG Fortran Library for minimisation and
 50 matrix diagonalisation.
 51 ■ FOCUS ([Fabi, 1995](#)) uses Monte-Carlo and general quasi-Newton methods to find the
 52 parameters.
 53 ■ PyCrystalField ([Scheie, 2021](#)) - fitting uses the SciPy minimize library.

54 All of these methods however only work for single ions (i.e. no exchange coupling) and only
 55 calculate one minimising system. INS_IEP however can fit INS data from samples with multiple
 56 spin centres, it can also find multiple minimising systems without multiple initial guesses to
 57 the parameter set.

58 Key Concepts

59 The Spin Hamiltonian

60 The Spin Hamiltonian, H , is an approximation of the Hamiltonian that uses spin coordinates
 61 instead of orbital coordinates, and is widely used to model data arising from many spectroscopy
 62 techniques ([Launay & Verdaguier, 2014](#)). It can be modeled as a linear combination of
 63 interaction terms; in this package we will use the zero field interaction, H_{ZFI} , and the
 64 electron-electron interaction, H_{EEI} :

$$H = H_{ZFI} + H_{EEI}.$$

65 Both of these terms can themselves be modelled as the linear sum of other basis matrices.
 66 The zero field interaction can be written as:

$$H_{ZFI} = \sum_{-k \leq q \leq k} B_k^q O_k^q$$

67 where the O_k^q are Stevens Operators ([Rudowicz & Chung, 2004](#)), and B_k^q the associated
 68 parameter. When there are multiple spin centres it is necessary to take Kronecker products of
 69 the operator with identity matrices of the appropriate for each other spin centre.

70 When there are multiple spin centres it is also necessary to include an electron-electron
 71 interaction term, H_{EEI} . This term will be the sum of interaction terms between each pair of
 72 spin centres:

$$H_{EEI} = - \sum_{i \neq j} J_{ij} S_i \cdot S_j$$

73 where S_i is the vector of spin operators $S_i = [S_x, S_y, S_z]$ for the i -th spin centre, and J_{ij} is
 74 the parameter to be found that represents the strength of interaction between the two spin
 75 centres. Note that in the isotropic case J can be thought of as a scalar value, but in the
 76 anisotropic case will be a matrix where the off diagonals are skew symmetric (often the off

⁷⁷ diagonals are assumed to be zero). While the summation is in theory over all spin centre
⁷⁸ combinations, in practice many of these contributions will be negligible - often only the nearest
⁷⁹ neighbour interactions are significant.

⁸⁰ It is important to mention that these matrix operators can be very large. The size is defined
⁸¹ by the number of spin centres (n) and the spin (S_i) of each spin centre. The dimension of
⁸² matrices is given by:

$$\prod_i^n (2S_i + 1).$$

⁸³ The operators are however highly sparse, this means that it is possible to use eigensolvers that
⁸⁴ can take advantage of this sparsity.

⁸⁵ Inverse Eigenvalue Problem

⁸⁶ The INS experiments provide eigenvalues of the Spin Hamiltonian matrix of the sample, the
⁸⁷ task of calculating the matrix from the eigenvalues is an inverse eigenvalue problem:

⁸⁸ Let $A(x)$ be the affine family of matrices,

$$A(x) = A_0 + \sum_{i=1}^{\ell} x_i A_i,$$

⁸⁹ where $x \in \mathbb{R}^\ell$ and $A_0, \dots, A_\ell \in \mathbb{C}^{n \times n}$ are linearly independent Hermitian matrices, and denote
⁹⁰ the ordered eigenvalues of $A(x)$ as $\lambda_1(x) \leq \dots \leq \lambda_n(x)$. Then the least squares inverse
⁹¹ eigenvalue problem (LSIEP) is to find the parameters $x \in \mathbb{R}^\ell$ that minimises

$$F(x) = \frac{1}{2} \|r(x)\|_2^2 = \frac{1}{2} \sum_{i=1}^m (\lambda_i(x) - \lambda_i^*)^2$$

⁹² where $\lambda_1^* \leq \dots \leq \lambda_m^*$ are the experimental eigenvalues (Chu & Golub, 2005). In the case of INS
⁹³ fitting the A_i basis matrices will be a combination of Stevens operators and electron-electron
⁹⁴ exchange terms. The IEP described above is formulated as a least squares problem because
⁹⁵ the number of eigenvalues that can be probed by INS experiments is often a small subset of
⁹⁶ the full spectrum. Due to the low temperatures that these experiments are performed at (can
⁹⁷ be as low as 1K) it is generally the smallest eigenvalues that are involved. Note also that since
⁹⁸ it is the energy difference between the eigenvalues that is probed we actually have to modify
⁹⁹ the IEP - either by adding an additional parameter (an identity matrix) that shifts the values
¹⁰⁰ of the eigenvalues, or by changing the above formula for F to directly sum the difference in
¹⁰¹ eigenvalues thereby reducing the number of residual equations in $r(x)$ by one.

¹⁰² As far as we are aware this is the first time that the fitting of INS data has been explicitly
¹⁰³ formulated as an IEP. An advantage of this formulation is that there are explicit formulas for
¹⁰⁴ the derivatives of $r(x)$. The first derivative (Jacobian) is:

$$J_r(x) = \begin{pmatrix} q_1(x)^T A_1 q_1(x) & \dots & q_1(x)^T A_\ell q_1(x) \\ \vdots & \ddots & \vdots \\ q_m(x)^T A_1 q_m(x) & \dots & q_m(x)^T A_\ell q_m(x) \end{pmatrix},$$

¹⁰⁵ and the second derivative (Hessian) is:

$$(H_r)_{ij} = 2 \sum_{k=1}^m (\lambda_k - \lambda_k^*) \sum_{\substack{t=1 \\ \lambda_t \neq \lambda_k}}^m \frac{(q_t^T A_i q_k)(q_t^T A_j q_k)}{\lambda_k - \lambda_t}.$$

¹⁰⁶ Another advantage is the number of constraints to fit is much smaller than fitting the spectrum
¹⁰⁷ itself, as it corresponds to fitting only the locations of the peaks of the spectrum.

108 **Methods**

109 All of the methods used are iterative schemes of the form $x^{k+1} = x^k + p^k$ where the step p^k
110 uniquely defines each algorithm:

- 111 ▪ Newton's method: $p^k = (J_r^T J_r + H_r r)^{-1} J_r^T r$ ([Nocedal & Wright, 2006](#))
112 ▪ Gauss-Newton method: $p^k = (J_r^T J_r)^{-1} J_r^T r$ ([Nocedal & Wright, 2006](#))
113 ▪ Lift and Projection Method: $p^k = B^{-1} J_r^T r$ ([Bloor Riley et al., 2025a](#))

114 Where the matrix B is the Gram matrix formed from the frobenius inner products of the basis
115 matrices: $B_{ij} = \langle A_i, A_j \rangle_F$. The Lift and Projection method is a Riemannian Gradient descent
116 method ([Bloor Riley et al., 2025a](#)), inspired by the Lift and Projection method ([Chu & Golub,](#)
117 [2005](#)), specifically designed for solving IEPs. In ([Bloor Riley et al., 2025a](#)) it is proven that the
118 method is a strictly descending algorithm, that is it reduces the value of the objective function
119 every step. Both the deflated Gauss-Newton method and the Riemannian Gradient descent
120 Lift and Projection method are new methods designed for this package ([Bloor Riley et al.,](#)
121 [2025a, 2025b](#)).

122 The m eigenvalues required to calculate J_r are calculated using MATLAB's `eigs` command,
123 which is an efficient method based on Krylov subspaces to calculate a small subset of the
124 eigenvalues of the matrix. This is invaluable in this setting because computing the full set of
125 eigenvalues by exact diagonalisation is in some cases infeasible ([Bloor Riley et al., 2025a](#)).

126 **Deflation**

127 The number of eigenvalues that can be probed via INS experiments varies depending on the
128 equipment and sample in question, meaning that the fitting problem is often under or even
129 over determined. The IEP is also highly nonlinear and due to the experimental nature of the
130 data may be ill-posed. One consequence of this is that the solution space may be very 'bumpy',
131 that is there may exist many local minimisers to the problem. For example in [Figure 1](#), there
132 are clearly 4 distinct solutions (for more details see Example 1 and the file `Example1_Mn12.m`
133 in the examples folder). We seek to solve the problem of multiple local minima by the use
134 of Deflation, a numerical technique used to find multiple solutions to systems of equations
135 ([Farrell et al., 2015](#)). Fortunately it is cheap to apply deflation for the above methods, it is
136 simply a change to the length of the step - notably this means that the direction of each step
137 does not change. It is proven in ([Bloor Riley et al., 2025b](#)) that the deflated methods will not
138 converge to deflated points. The usual requirements still apply to the convergence of the new
139 methods - that the initial guess is close enough to the new minimum, and that the Jacobian is
140 full rank in a neighbourhood around that minimum. The rate of convergence of the deflated
141 methods is also more complicated, although the number of iterations required to converge can
142 go up with the number of deflations this is not a strict correlation, as can be seen in [Figure 2](#).

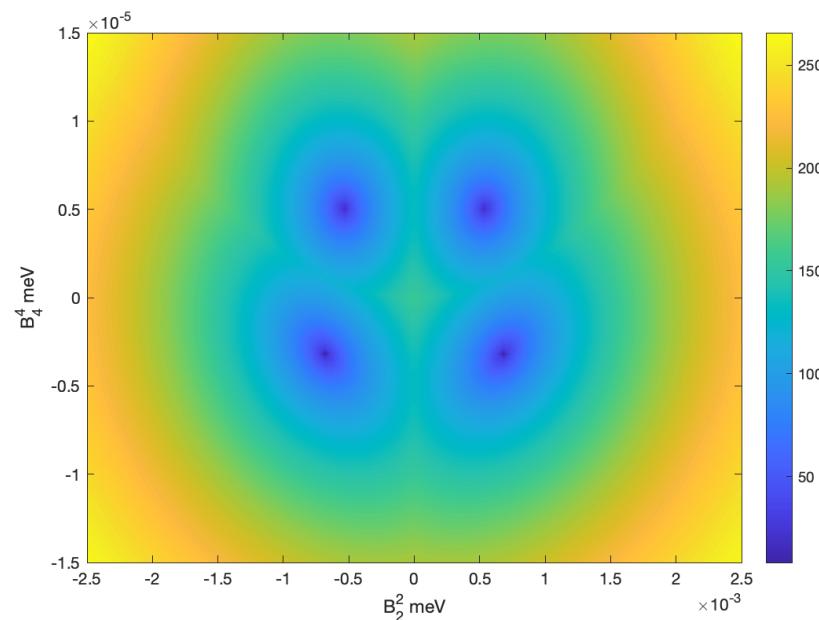


Figure 1: Contour plot of how F varies with the two parameters B_2^2 and B_4^4 for the molecule Mn_12 as described in Example 1. There are four locally minimising parameter pairs corresponding to the four blue regions

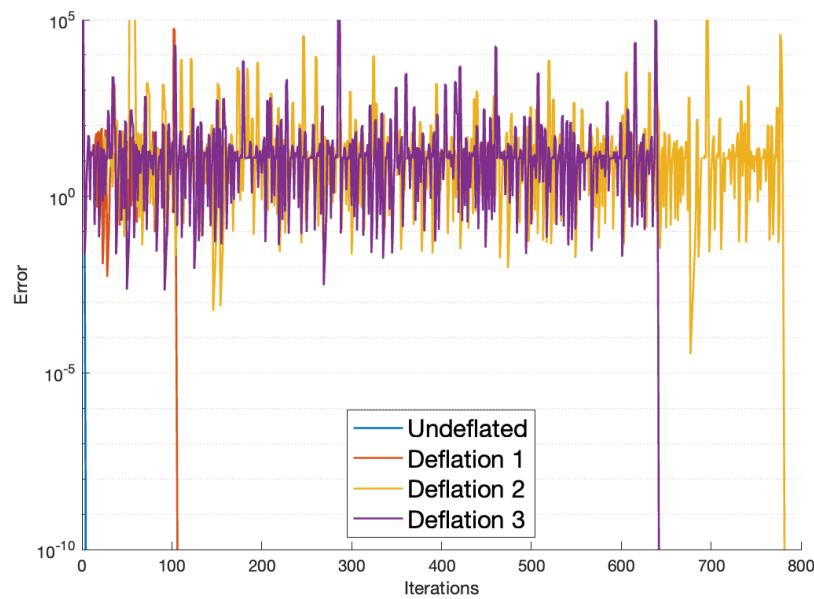


Figure 2: Comparison of the convergence behaviour for computing each solution in Example 1. The gradient of the lines as the method approaches the solution shows how the methods are quadratically convergent local to a minimum.

143 Examples

144 Example 1 - Mn12

145 The first example we will look at is Manganese-12-acetate. This is a well known example in the
 146 INS and magnetism community, as one of the first molecules that behaves like a nano-sized
 147 magnet with a molecular magnetic coercivity as well as its role in the research of quantum
 148 tunnelling of magnetisation (Friedman et al., 1996; Sessoli et al., 1993).

149 The Spin Hamiltonian of this system, using the giant spin approximation, can be represented
 150 as a 21×21 matrix modelled using 4 Stevens operators (Bircher et al., 2004):

$$H = B_2^0 O_2^0 + B_4^0 O_4^0 + B_2^2 O_2^2 + B_4^4 O_4^4 \in \mathbb{R}^{21 \times 21}$$

151 We utilise the same spin system syntax as easyspin, so to set up the problem we first set up
 152 the model, along with initial guesses for the parameters:

```
%One spin centre (because giant spin approximation)
Sys0.S=10;
%Four Stevens operators
Sys0.B2 = [-100,0,-1000,0,0];
Sys0.B4 = [-1,0,0,0,-1,0,0,0,0];
```

153 Then we input the experimental eigenvalues - these are typically shifted such that the smallest
 154 eigenvalue is zero - and define which parameters to fit. Note that all values given must be in
 155 gigahertz, so it may be useful to use conversions.

```
rcm = 29979.2458; meV = rcm*8.065; %Conversions values
%Input calculated eigenvalues:
Exp.ev = [0,0,1.24,1.24,2.3,2.3,3.18,3.18,3.91,3.91,4.5,4.5,
           4.97,4.97,5.32,5.32,5.54,5.59,5.69,5.75,5.78].*meV;
%Note that these eigenvalues are simulated from the parameters given in [@bircher_trans]

%Vary all non zero parameters (no Fixed parameters):
Vary = Sys0;
```

156 Then all that is required is to call `INS_IEP` with these three inputs:

```
SysOut = INS_IEP(Sys0,Vary,Exp);
```

157 If we wish to find all four solutions as shown in Figure 1 then we use the additional option:

```
Opt.NDeflations = 4;
SysOut = INS_IEP(Sys0,Vary,Exp,Opt);
```

158 In this case `SysOut` will be an array of four spin structures each containing a distinct
 159 locally optimal solution. It is possible to access information about the convergence of each
 160 deflation by using `SysOut.Output`. For example by utilising the iterates recorded, stored in
 161 `SysOut.Output.Iterates` it is possible to plot a graph of convergence, as can be seen in
 162 Figure 2. The Output structure also contains the value of F at the final point, as well as the
 163 number of iterations it took to get there.

164 A full list of options is provided in the help of `INS_IEP`.

165 Example 2 - Chromium(iii) Horseshoes

166 The second example concerns antiferro-magnetically coupled chain of six chromium(III) ions (M.
 167 L. Baker et al., 2011). Because there are multiple spin centres an electron-electron interaction
 168 term is required. The spin hamiltonian is a 4096×4096 matrix composed of two Stevens

169 operators and one interaction term, since it is known a priori that each spin centre will have
 170 the same value parameters we pin the parameters here, by setting the initial guess as the same
 171 value:

```
Sys0.S = [1.5 1.5 1.5 1.5 1.5 1.5];
Sys0.B2 = [1 0 -1 0 0;
            1 0 -1 0 0;
            1 0 -1 0 0;
            1 0 -1 0 0;
            1 0 -1 0 0;
            1 0 -1 0 0];
Sys0.J = [100,0,0,0,0,100,0,0,0,100,0,0,100,0,100];
Vary = Sys0;
Exp.ev = [0,0.355,0.457,0.497,1.576,1.577,1.592,1.629,1.632,
           2.97,2.98,3.002,3.004,3.01,3.038,3.821,3.824,3.827,
           3.837,3.856,3.879,3.888,3.895,3.903].*meV;
```

172 Note that only 24 eigenvalues were found experimentally, so this will form a partial LSIEP. To
 173 find the solution system is as simple as:

```
Opt.GradientTolerance = 1e-3; %Use additional stopping criterion.
SysOut= INS_IEP(Sys1,Vary1,Exp,Opt);
```

174 It is possible to find multiple minimising systems even if they do not make any sense physically,
 175 however due to the scaling of the problem a change in the default deflation parameters is
 176 necessary:

```
Opt = struct('NDeflations',5,'Sigma',1e-7,'StepTolerance',1e-3 );
SysOut= INS_IEP(Sys0,Vary,Exp,Opt);
```

177 The output contains five different spin systems that all have the same eigenvalues as input.
 178 Like in example one some of the minima found differ only by a sign change of the B_2^2 parameter,
 179 but all minima have the same exchange term.

180 Additional examples can be found in the Examples folder.

181 Note also that using mint ([M. L. Baker, 2022](#)), which is fully compatible with INS_IEP it is
 182 possible to simulate the INS spectrum of any calculated system - which can then be compared
 183 to the experimental spectrum. A description of how to do this can also be found in the
 184 examples folder.

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