



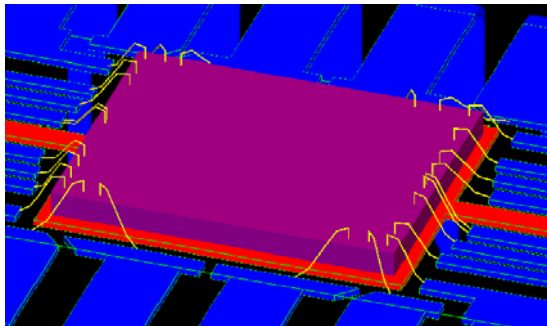
Passive reduced order multiport modelling: the Padé-Laguerre, Krylov-Arnoldi-SVD connection (and its application to FDTD)

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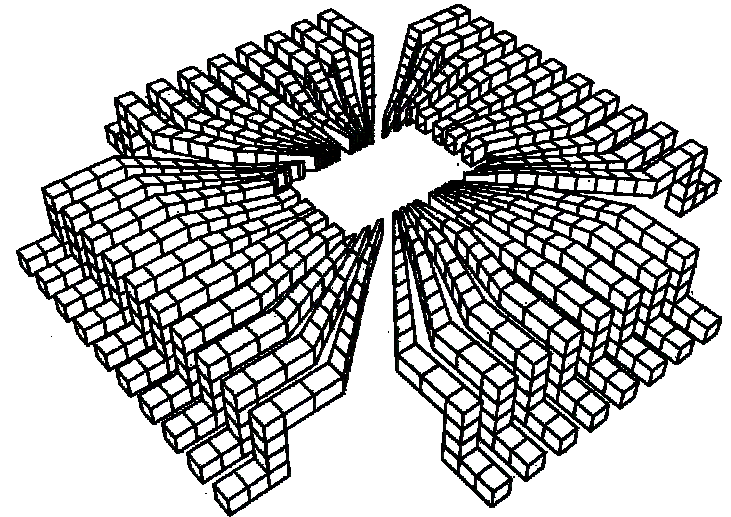


- Introduction
- Laguerre-SVD Reduced Order Modelling (ROM)
- Example: PEEC-circuit
- ROM and explicit reciprocity
- Example: two coupled lossy transmission lines
- Application to FDTD: automatic subcell generation
- Example: a grid of dielectric wires
- Conclusion

- The time-domain behaviour of a circuit or a physical device (e.g. a package) can be described by a state space description
- in electromagnetics this description typically emerges from numerical techniques such as FDTD, TLM or PEEC



package



perforated ground plane

- Typical representation as a set of first-order differential equations

$$\mathbf{C} \dot{\mathbf{x}} = -\mathbf{G} \mathbf{x} + \mathbf{B} \mathbf{u} \quad (\mathbf{C}, \mathbf{G}: N \times N)$$

$$\mathbf{y} = \mathbf{L}^T \mathbf{x} \quad (\mathbf{B}, \mathbf{L}: N \times P)$$

\mathbf{x} : a vector of N internal variables (e.g. currents and voltages*)

\mathbf{u} : a vector of P input (source) variables (e.g. port voltages)

\mathbf{y} : a vector of P output variables (e.g. port currents)

**Note: the variables could also be fields → FDTD!*

- The purpose of ROM techniques is to reformulate the previous state-space description as

$$\hat{\mathbf{C}} \dot{\mathbf{w}} = -\hat{\mathbf{G}} \mathbf{w} + \hat{\mathbf{B}} \mathbf{u} \quad (\hat{\mathbf{C}}, \hat{\mathbf{G}}: Q \times Q)$$

$$\mathbf{y} = \hat{\mathbf{L}}^T \mathbf{w} \quad (\hat{\mathbf{B}}, \hat{\mathbf{L}}: Q \times P)$$

\mathbf{w} : a **new** vector of Q ($Q < N$) internal variables

\mathbf{u} : the **same** vector of P input (source) variables (i.e. port voltages)

\mathbf{y} : the **same** vector of P output variables (i.e. port currents)

- 4 major techniques are found in literature
 - Asymptotic Waveform Evaluation (AWE)
 - Matrix Padé via Lanczos
 - Arnoldi-PRIMA
 - Congruence transformation
- these techniques mostly work for early time responses but are less performant for low frequencies!
- a new solution (based on an expansion in Laguerre polynomials) was developed by L. Knockaert, circumventing this low frequency problem.

Expansion in scaled Laguerre functions (1)



- In most techniques described in literature the transfer matrix $\mathbf{H}(s)$ with

$$\mathbf{H}(s) = \mathbf{L}^T(\mathbf{G} + s\mathbf{C})^{-1}\mathbf{B}$$

is used as the starting point to obtain Padé approximations by means of moment matching.

- Here we start from $\mathbf{h}(t) = \text{Lapl}^{-1}(\mathbf{H}(s))$ and expand $\mathbf{h}(t)$ in a set of **scaled Laguerre functions** La_n :

$$\mathbf{h}(t) = \sum_{n=0}^{\infty} \mathbf{F}_n \phi_n^{\alpha}(t)$$

with

$$\phi_n^{\alpha}(t) = (2\alpha)^{1/2} e^{-\alpha t} \text{La}_n(2\alpha t) \quad \text{with } \text{La}_n(t) = \frac{e^t}{n!} \frac{d^n}{dt^n} (e^{-t} t^n)$$

Expansion in scaled Laguerre functions (2)



- The Laplace transform of the scaled Laguerre functions is

$$\Phi_n^\alpha(s) = (2\alpha)^{1/2} \frac{1}{s + \alpha} \left(\frac{s - \alpha}{s + \alpha} \right)^n$$

- Consequently:

$$\mathbf{H}(s) = \mathbf{L}^T (\mathbf{G} + s\mathbf{C})^{-1} \mathbf{B}$$

$$= \sum_{n=0}^{\infty} \mathbf{F}_n \Phi_n^\alpha(s) = (2\alpha)^{1/2} \frac{1}{s + \alpha} \sum_{n=0}^{\infty} \mathbf{F}_n \left(\frac{s - \alpha}{s + \alpha} \right)^n$$

This implies that $\mathbf{H}(s)$ is written as the **product** of a simple **low-pass filter** and a weighted sum of **all-pass filters**!

- We now map the s-domain onto the u-domain with the bilinear transformation (right half-plane \rightarrow unit circle)

$$u = \frac{s - \alpha}{s + \alpha}$$

leading to

$$\mathbf{H}(u) = \mathbf{L}^T [(\alpha \mathbf{C} + \mathbf{G}) + u (\alpha \mathbf{C} - \mathbf{G})]^{-1} \mathbf{B}$$

Consequently:

- an **m-th order Padé approximation** in the u-domain of $\mathbf{H}(u)$
- is equivalent to**
- an **m-th order Laguerre approximation** of $\mathbf{H}(s)$ in the s-domain!

- The previous statement implies that $\mathbf{H}(s)$ can be optimally approximated in the Hardy space by the truncated Laguerre-expansion, i.e.

$$\mathbf{H}(s) \approx \sum_{n=0}^m \mathbf{F}_n \Phi_n^\alpha(s) = \mathbf{H}_m(s)$$

- The point-wise convergence $\mathbf{H}_m(i\omega) \rightarrow \mathbf{H}(i\omega)$ can be proven (see L. Knockaert and D. De Zutter, “Passive Reduced Multiport Modelling...”, Int. J. Electr. Commun. (AEÜ), pp. 254-260, no. 5, **53**, 1999. **[1]**)
- Mathematical arguments show that a good choice for α is

$$\alpha \approx 2\pi f_{\max}$$

with f_{\max} the bandwidth of the considered system.

■ The Krylov-Arnoldi-SVD connection (1)



- Reduced order modelling via Krylov-Arnoldi

Starting point: original system (N x N)

$$\mathbf{C}\mathbf{x} = -\mathbf{G}\mathbf{x} + \mathbf{B}\mathbf{u}$$

$$\mathbf{y} = \mathbf{L}^T \mathbf{x}$$

- via the u-domain a modified \mathbf{A} and \mathbf{R} matrix can be defined:

$$\mathbf{A} = -(\alpha\mathbf{C} + \mathbf{G})^{-1}(\alpha\mathbf{C} - \mathbf{G}) \quad (\text{classical approach: } \mathbf{A} = -\mathbf{G}^{-1}\mathbf{C})$$

$$\mathbf{R} = (\alpha\mathbf{C} + \mathbf{G})^{-1} \mathbf{B} \quad (\text{classical approach: } \mathbf{R} = \mathbf{G}^{-1}\mathbf{B})$$

- **End point: reduced order system (Q x Q)**

$$\hat{\mathbf{C}}\mathbf{w} = -\hat{\mathbf{G}}\mathbf{w} + \hat{\mathbf{B}}\mathbf{u}$$

$$\mathbf{y} = \hat{\mathbf{L}}^T \mathbf{w}$$

with

- $\hat{\mathbf{C}} = \mathbf{X}^T \mathbf{C} \mathbf{X}$

- $\hat{\mathbf{G}} = \mathbf{X}^T \mathbf{G} \mathbf{X}$

- $\hat{\mathbf{B}} = \mathbf{X}^T \mathbf{B}$

- $\hat{\mathbf{L}} = \mathbf{X}^T \mathbf{L}$

■ The Krylov-Arnoldi-SVD connection (2)



- **X**: a $N \times Q$ matrix that orthogonalizes the columns of the $N \times Q$ Krylov matrix

$$\mathbf{K}_M = [\mathbf{R}, \mathbf{A}\mathbf{R}, \mathbf{A}^2\mathbf{R}, \dots, \mathbf{A}^{M-1}\mathbf{R}, (\mathbf{A}^M\mathbf{R})] \quad \text{with } Q = P \times M + r$$

- **X** can be generated using the block Arnoldi algorithm which is equivalent to so-called ‘thin’ QR factorization based on modified Gram-Schmidt orthogonalization.
- to avoid precision problems we use an SVD approach i.e. write \mathbf{K}_M as

$$\mathbf{K}_M = \mathbf{U} \Sigma \mathbf{V}^T \quad \text{implying that } \mathbf{X}^T \mathbf{X} = \mathbf{U}^T \mathbf{U}$$

and further implying that the role of **X can be replaced by U**.

(see [1])

■ The complete algorithm

- select the values for α ($\approx 2\pi f_{\max}$) and Q (reduction factor N/Q)
- solve $(\mathbf{G} + \alpha\mathbf{C})\mathbf{R}_0 = \mathbf{B}$
- for $k = 1, M - 1$ solve (with $M = Q/P$)
 $(\mathbf{G} + \alpha\mathbf{C})\mathbf{R}_k = (\mathbf{G} - \alpha\mathbf{C})\mathbf{R}_{k-1}$
- $\mathbf{U} \Sigma \mathbf{V}^T = \text{SVD}[\mathbf{R}_0, \mathbf{R}_1, \dots, \mathbf{R}_{M-1}]$
- finally,

$$\hat{\mathbf{C}} = \mathbf{U}^T \mathbf{C} \mathbf{U}$$

$$\hat{\mathbf{G}} = \mathbf{U}^T \mathbf{G} \mathbf{U}$$

$$\hat{\mathbf{B}} = \mathbf{U}^T \mathbf{B}$$

$$\hat{\mathbf{L}} = \mathbf{U}^T \mathbf{L}$$

■ Example: a PEEC-circuit



■ Typical example: a PEEC circuit

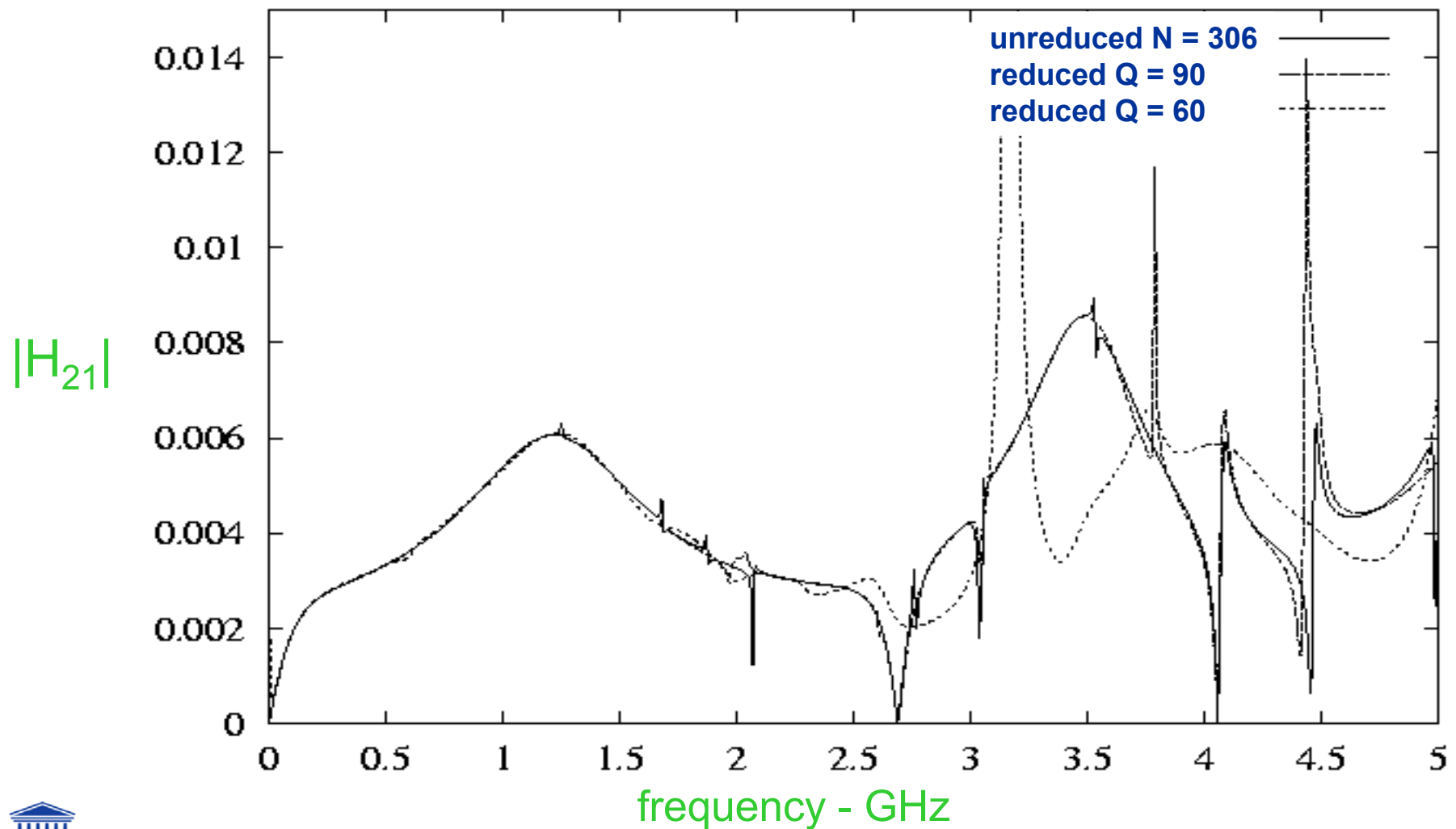
A.E. Ruehli, “Equivalent circuit models for three-dimensional multiconductor systems”, IEEE Trans. MTT, March 1974

and

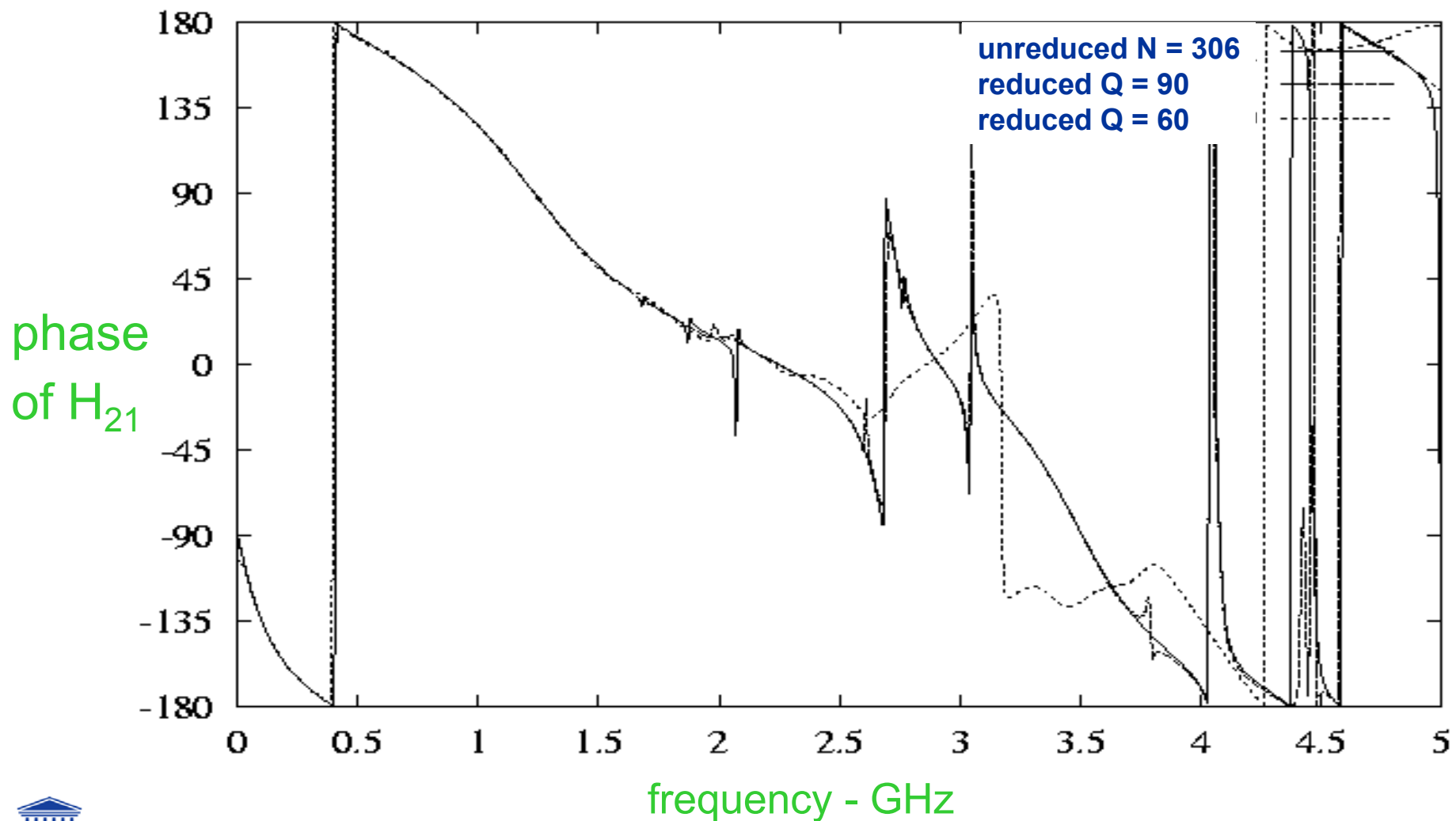
P. Feldmann and R.W. Freund, “Efficient linear circuit analysis by Padé approximation via the Lanczos process”, IEEE Trans. Computer-Aided Design, May 1995)

- 2100 capacitors
 - 127 inductors
 - 6990 mutual inductive couplings
 - **$N = 306$**
- reduced to **$Q < N = 60, 90$**

■ PEEC-circuit: numerical results (1)



■ PEEC-circuit: numerical results (2)



State-space description and reciprocity



- Suppose now that we start from the following state-space description

$$\begin{cases} \mathbf{C}\dot{\mathbf{x}} = -\mathbf{G}\mathbf{x} + \mathbf{L}u \\ \mathbf{y} = \mathbf{L}^T \mathbf{x} \end{cases}$$

with

$$\mathbf{C} = \mathbf{C}^T; \mathbf{G} = \mathbf{G}^T$$

Consequently, the transfer matrix $\mathbf{H}(s)$ (Laplace domain) is reciprocal, where

$$\mathbf{H}(s) = \mathbf{L}^T (\mathbf{G} + s\mathbf{C})^{-1} \mathbf{L}$$

and

$$\mathbf{H}^T(s) = \mathbf{H}(s)$$



Reversely:

- start from standard minimal state-space description:

$$\mathbf{H}(s) = \mathbf{L}^T (s\mathbf{I} - \mathbf{A})^{-1} \mathbf{R}$$

- suppose that $\mathbf{H}^T(s) = \mathbf{H}(s)$ (reciprocity)

- then Frobenius: $\mathbf{A} = -\mathbf{C}^{-1} \mathbf{G}$ with $\mathbf{C}^T = \mathbf{C}$ and $\mathbf{G}^T = \mathbf{G}$

- consequence: $\mathbf{H}(s) = \mathbf{L}^T (s\mathbf{C} + \mathbf{G})^{-1} \mathbf{C}\mathbf{R}$

- since $\mathbf{H}(s) = \mathbf{H}^T(s)$ we have $\mathbf{C}\mathbf{R} = \mathbf{L}$

- or $\mathbf{H}(s) = \mathbf{L}^T (s\mathbf{C} + \mathbf{G})^{-1} \mathbf{L}$

Reciprocity \Rightarrow explicitly reciprocal state space description

■ Congruent ROM and explicit reciprocity



Using congruent reduced order modeling:

$$\hat{\mathbf{G}} = \mathbf{U}^T \mathbf{G} \mathbf{U} \quad \hat{\mathbf{C}} = \mathbf{U}^T \mathbf{C} \mathbf{U} \quad \hat{\mathbf{L}} = \mathbf{U}^T \mathbf{L}$$

e.g. \mathbf{U} left factor of the SVD of the Krylov matrix (as in our method)

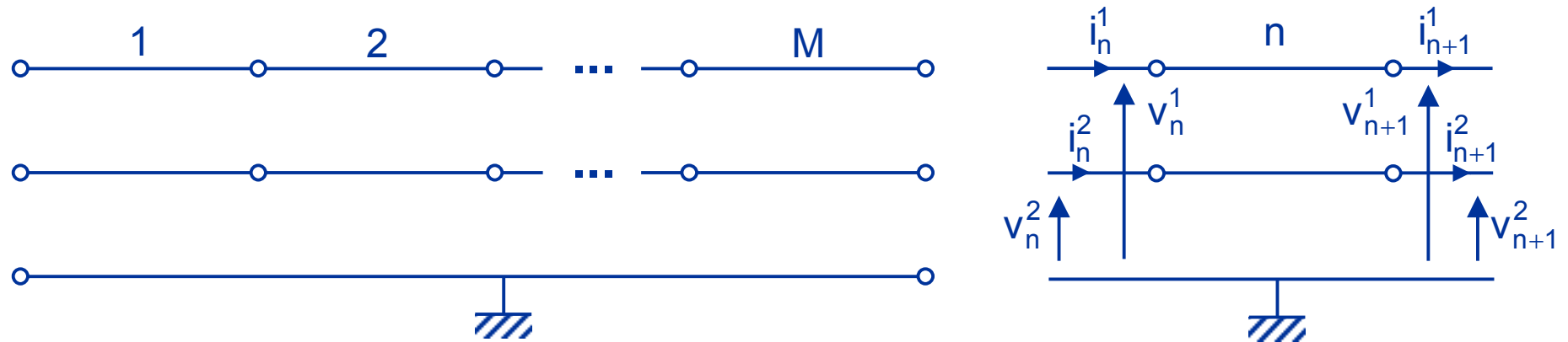
Then reduced state space description is still **explicitly** reciprocal:

$$\hat{\mathbf{H}}(s) = \hat{\mathbf{L}}^T (s\hat{\mathbf{C}} + \hat{\mathbf{G}})^{-1} \hat{\mathbf{L}} = \hat{\mathbf{H}}^T(s)$$

For the non-reciprocal form this is not the case!

$$\hat{\mathbf{H}}'(s) = \hat{\mathbf{L}}^T \mathbf{U} (s\mathbf{I} - \mathbf{U}^T \mathbf{A} \mathbf{U})^{-1} \mathbf{U}^T \mathbf{R} \neq \hat{\mathbf{H}}'^T(s)$$

Example: 2 coupled lines



2 coupled transmission lines represented by M discrete sections
total length: 5 cm

$$\mathbf{L} = \begin{bmatrix} 494.6 & 63.3 \\ 63.3 & 494.6 \end{bmatrix} \text{ nH/m} \quad \mathbf{R} = \begin{bmatrix} 0.1 & 0.02 \\ 0.02 & 0.1 \end{bmatrix} \Omega/\text{m}$$

$$\mathbf{C} = \begin{bmatrix} 62.8 & -4.9 \\ -4.9 & 62.8 \end{bmatrix} \text{ pF/m} \quad \mathbf{G} = \begin{bmatrix} 0.1 & -0.01 \\ -0.01 & 0.1 \end{bmatrix} \text{ S/m}$$

Normal format

Internal state space
variables: $\{\mathbf{v}_n, \mathbf{i}_n\}$

$$\mathbf{C}_n \frac{d\mathbf{v}_n}{dt} = -\mathbf{G}_n \mathbf{v}_n + \mathbf{i}_n - \mathbf{i}_{n+1}$$

$$\mathbf{L}_n \frac{d\mathbf{i}_n}{dt} = -\mathbf{R}_n \mathbf{i}_n + \mathbf{v}_{n-1} - \mathbf{v}_n$$

$$\underbrace{\begin{pmatrix} \mathbf{C}_1 & 0 & 0 & 0 & 0 \\ 0 & \mathbf{C}_2 & 0 & 0 & 0 \\ 0 & 0 & \mathbf{L}_1 & 0 & 0 \\ 0 & 0 & 0 & \mathbf{L}_2 & 0 \\ 0 & 0 & 0 & 0 & \mathbf{L}_3 \end{pmatrix}}_{\mathbf{C}}, \underbrace{\begin{pmatrix} \mathbf{G}_1 & 0 & -\mathbf{I} & \mathbf{I} & 0 \\ 0 & \mathbf{G}_2 & 0 & -\mathbf{I} & \mathbf{I} \\ \mathbf{I} & 0 & \mathbf{R}_1 & 0 & 0 \\ -\mathbf{I} & \mathbf{I} & 0 & \mathbf{R}_2 & 0 \\ 0 & -\mathbf{I} & 0 & 0 & \mathbf{R}_3 \end{pmatrix}}_{\mathbf{G}}, \underbrace{\begin{pmatrix} 0 & 0 \\ 0 & 0 \\ \mathbf{I} & 0 \\ 0 & 0 \\ 0 & -\mathbf{I} \end{pmatrix}}_{\mathbf{B}}$$

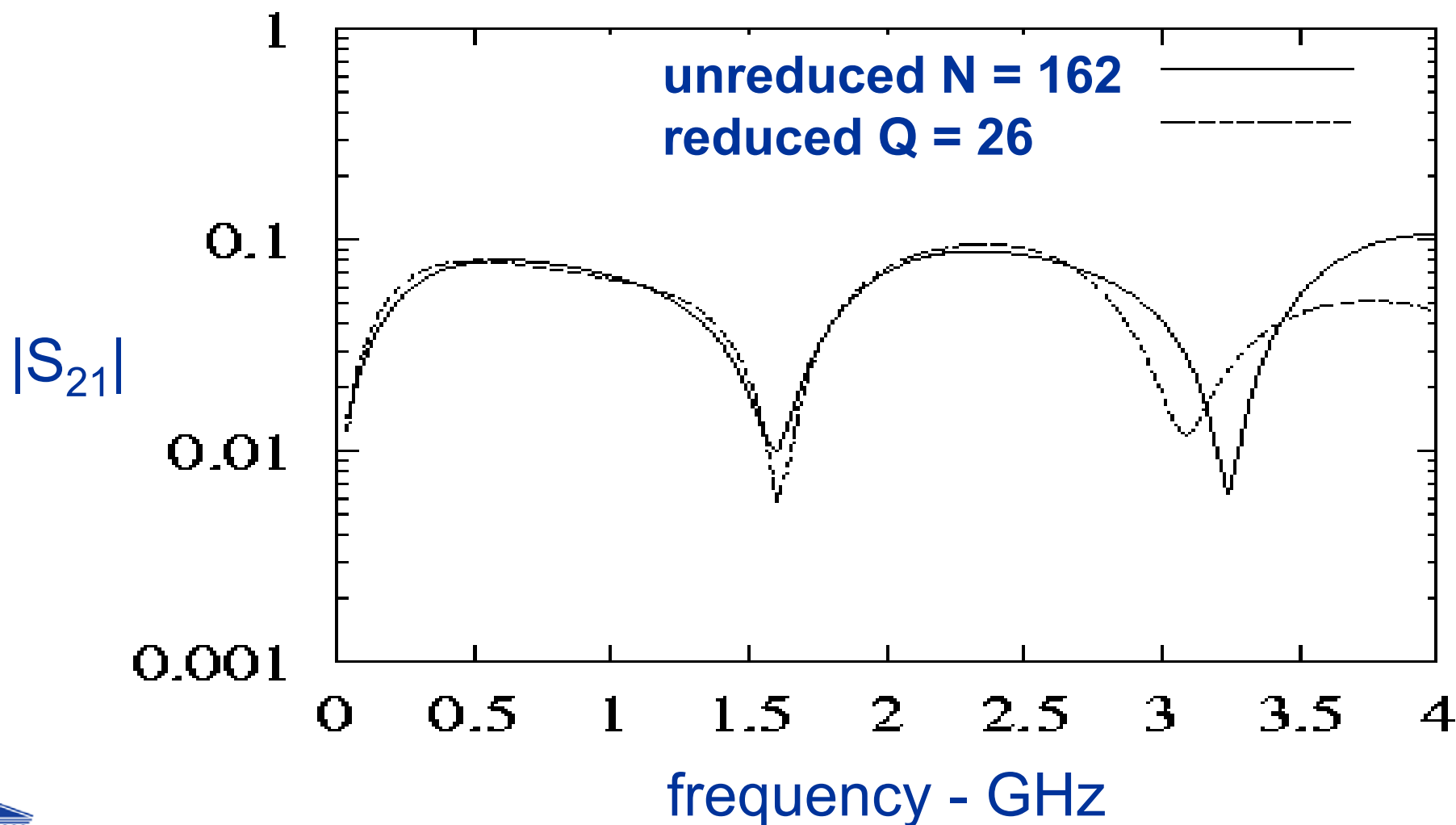
(for $M = 2$)

■ Explicitly reciprocal format

Internal state space variables: $\{-\mathbf{v}_n, \mathbf{i}_n\}$

$$\underbrace{\begin{pmatrix} -\mathbf{C}_1 & 0 & 0 & 0 & 0 \\ 0 & -\mathbf{C}_2 & 0 & 0 & 0 \\ 0 & 0 & \mathbf{L}_1 & 0 & 0 \\ 0 & 0 & 0 & \mathbf{L}_2 & 0 \\ 0 & 0 & 0 & 0 & \mathbf{L}_3 \end{pmatrix}}_{\mathbf{C}}, \underbrace{\begin{pmatrix} -\mathbf{G}_1 & 0 & -\mathbf{I} & \mathbf{I} & 0 \\ 0 & -\mathbf{G}_2 & 0 & -\mathbf{I} & \mathbf{I} \\ -\mathbf{I} & 0 & \mathbf{R}_1 & 0 & 0 \\ \mathbf{I} & -\mathbf{I} & 0 & \mathbf{R}_2 & 0 \\ 0 & \mathbf{I} & 0 & 0 & \mathbf{R}_3 \end{pmatrix}}_{\mathbf{G}}, \underbrace{\begin{pmatrix} 0 & 0 \\ 0 & 0 \\ \mathbf{I} & 0 \\ 0 & 0 \\ 0 & -\mathbf{I} \end{pmatrix}}_{\mathbf{B}}$$

■ Coupled lossy lines: numerical result

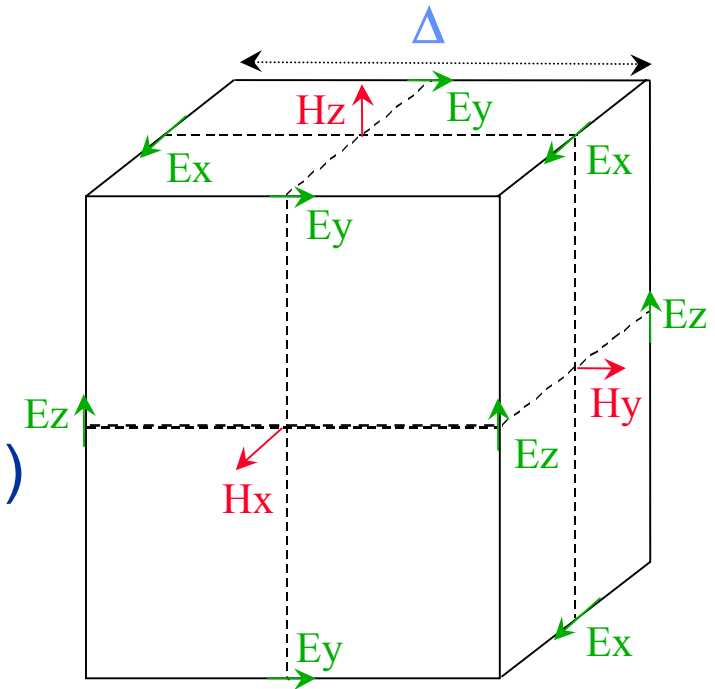


- we have discussed a reduced order modeling technique which remains accurate for low frequencies i.e. late time responses
- we have drawn attention to the fact that reciprocity can be explicitly represented in the state-space representation
- we have shown how this explicit reciprocity can be maintained when applying a ROM technique

2-D FDTD Automatic subcell generation

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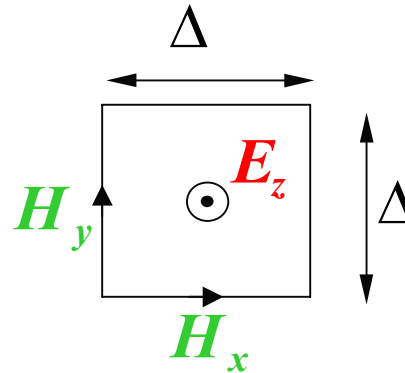
- In normal FDTD: cells of constant size
- Space step Δ (size of cell) is dependent on
 - wavelength (unavoidable)
 - smallest geometrical feature
- Techniques to choose Δ independent of small geometrical features
 - subcell models (thin wire, thin slot,...)
 - non-uniform orthogonal grid
 - subgridding
 - ...



Yee cell

2D FDTD - the TM-case

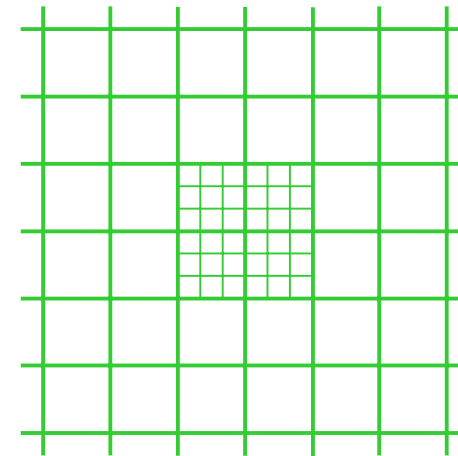
- 2-D FDTD, TM-case, no variation in z-direction
- field variables E_z , H_x , H_y
- Typical cell:
- equations:



$$\partial E_z / \partial x = \mu \partial H_y / \partial t$$

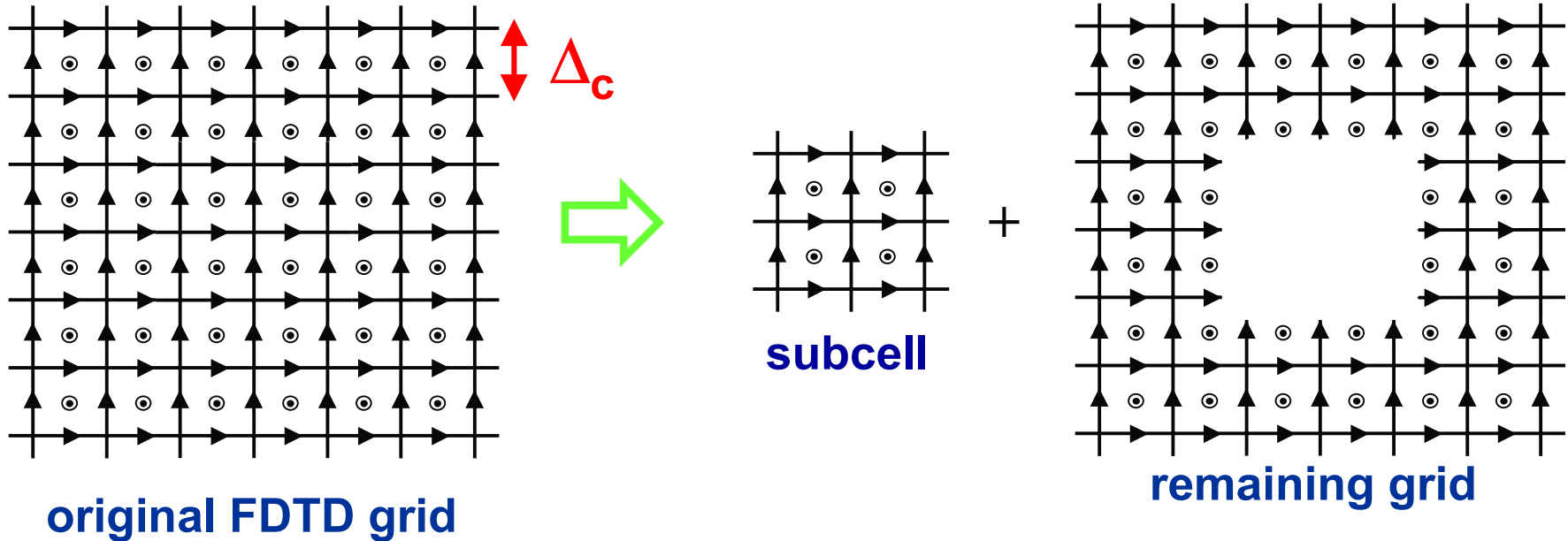
$$\partial E_z / \partial y = -\mu \partial H_x / \partial t$$

$$\partial H_y / \partial x - \partial H_x / \partial y = -\sigma E_z + \epsilon \partial E_z / \partial t$$

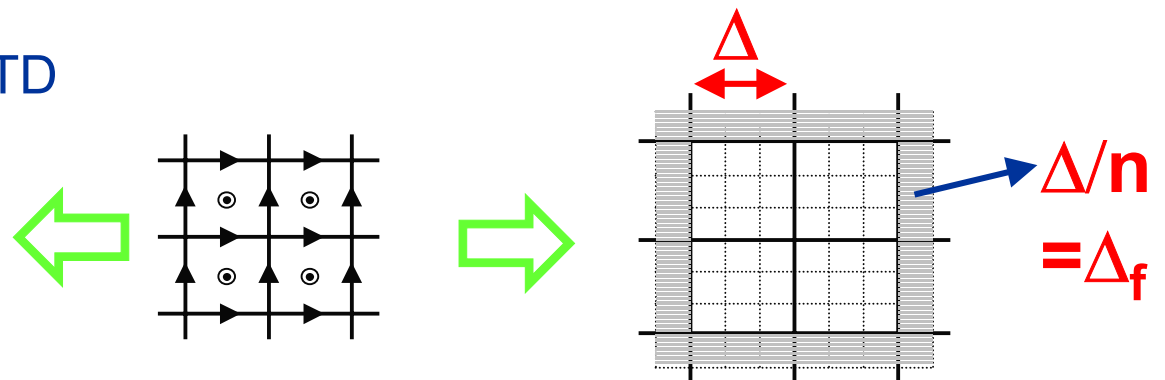


- subindex c refers to coarse grid
subindex f refers to fine grid

■ The subcell idea (1)



- state-space model through FDTD on fine grid
- only the space derivatives are discretised in the FDTD-way

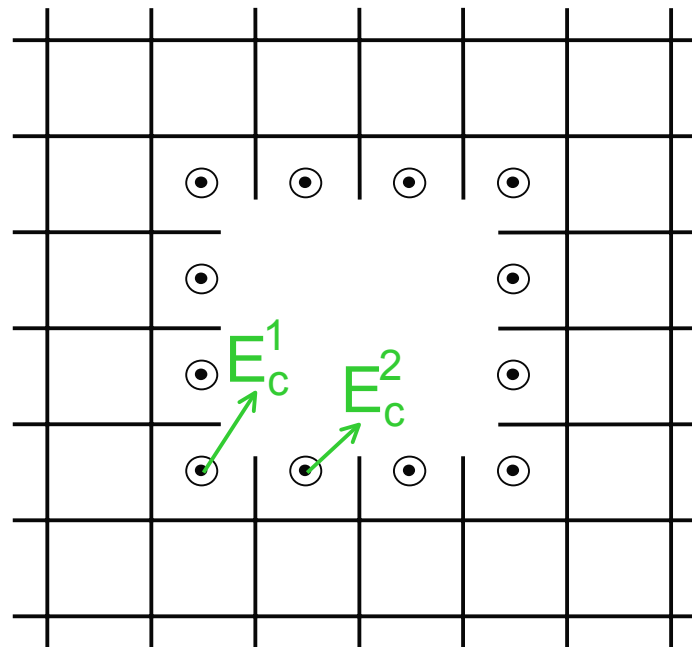
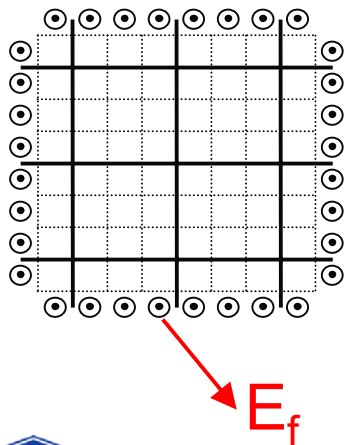
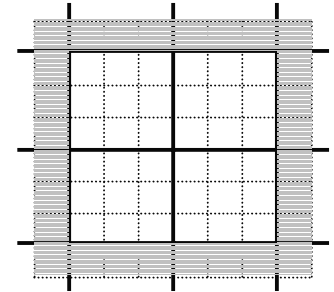


■ The subcell idea (2)

- State-space description of subcell $\begin{cases} \mathbf{C} \dot{\mathbf{x}} = -\mathbf{G} \mathbf{x} + \mathbf{B} \mathbf{u} \\ \mathbf{y} = \mathbf{L}^T \mathbf{x} \end{cases}$

\mathbf{x} : internal E_z , H_x , H_y fields

\mathbf{u} : inputs \rightarrow fine grid E_z -values derived from E_z fields of coarse grid through a suitable interpolation



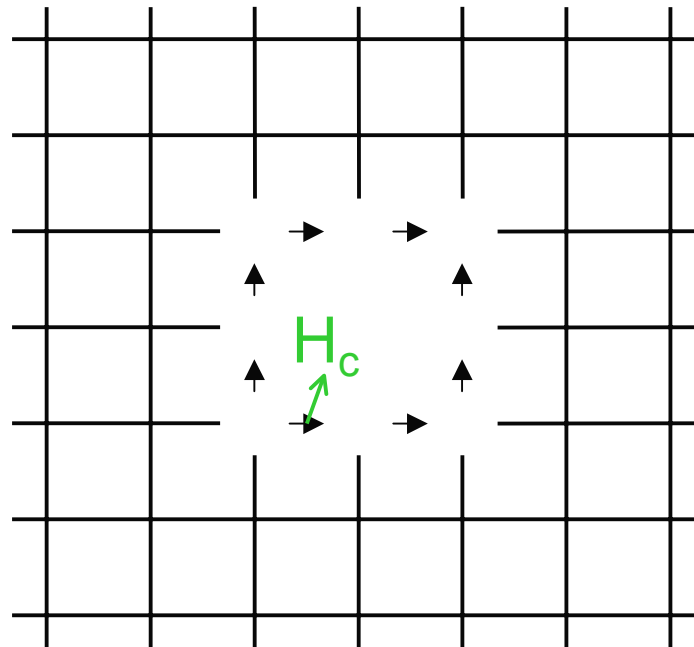
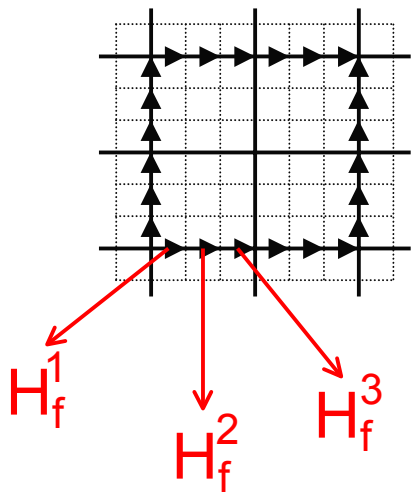
$$r = \Delta_c / \Delta_f = 3$$

$$E_f = \frac{2}{3} E_c^1 + \frac{1}{3} E_c^2$$

■ The subcell idea (3)

- State-space description of subcell
$$\begin{cases} \mathbf{C} \dot{\mathbf{x}} = -\mathbf{G} \mathbf{x} + \mathbf{B} \mathbf{u} \\ \mathbf{y} = \mathbf{L}^T \mathbf{x} \end{cases}$$

\mathbf{y} : outputs \rightarrow fine grid H-values leading to
H-fields of coarse grid through suitable averaging



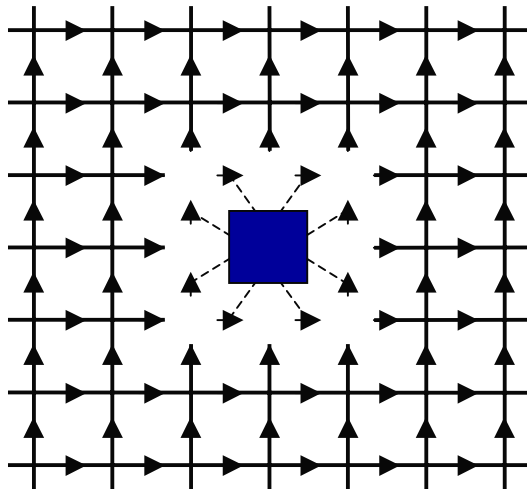
$$r = \Delta_c / \Delta_f = 3$$

$$H_c = \frac{1}{3}(H_f^1 + H_f^2 + H_f^3)$$

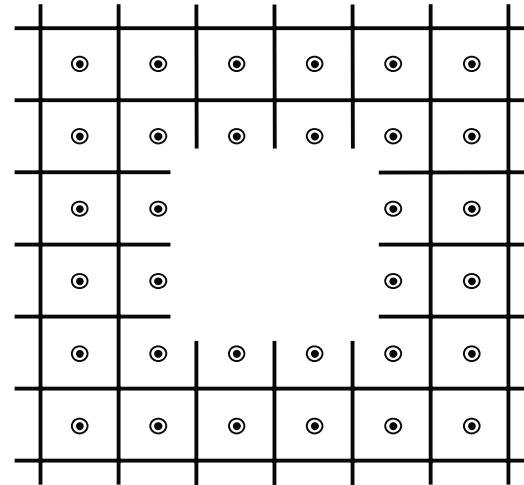
■ The subcell idea (4) - algorithm

- generate MIMO state-space description of subcell $\begin{cases} \mathbf{C} \dot{\mathbf{x}} = -\mathbf{G} \mathbf{x} + \mathbf{B} \mathbf{u} \\ \mathbf{y} = \mathbf{L}^T \mathbf{x} \end{cases}$
- apply ROM-algorithm $\begin{cases} \hat{\mathbf{C}} \dot{\mathbf{w}} = -\hat{\mathbf{G}} \mathbf{w} + \hat{\mathbf{B}} \mathbf{u} \\ \mathbf{y} = \hat{\mathbf{L}}^T \mathbf{w} \end{cases}$
- discretise time in the reduced state-space description
- solve the overall problem using FDTD with updating as below

$t = n\Delta t$

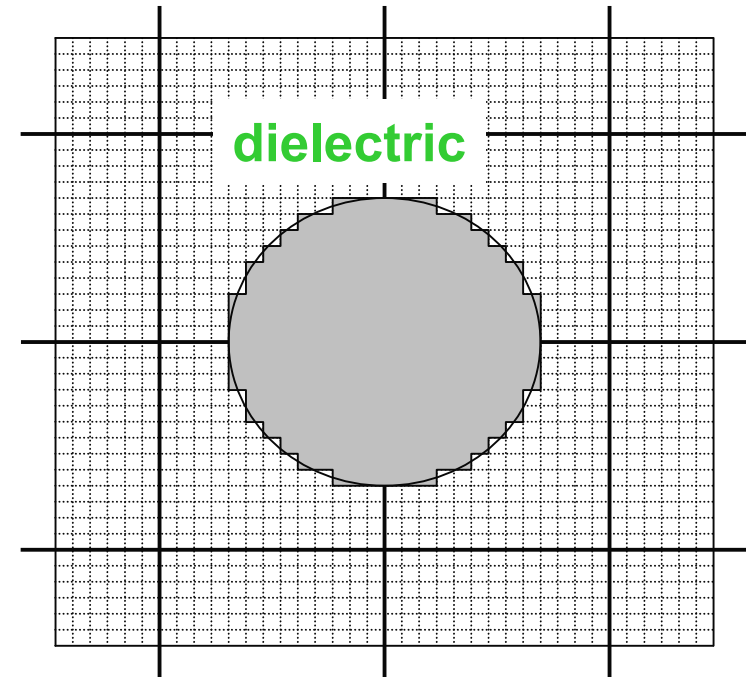


$t = (n+0.5)\Delta t$



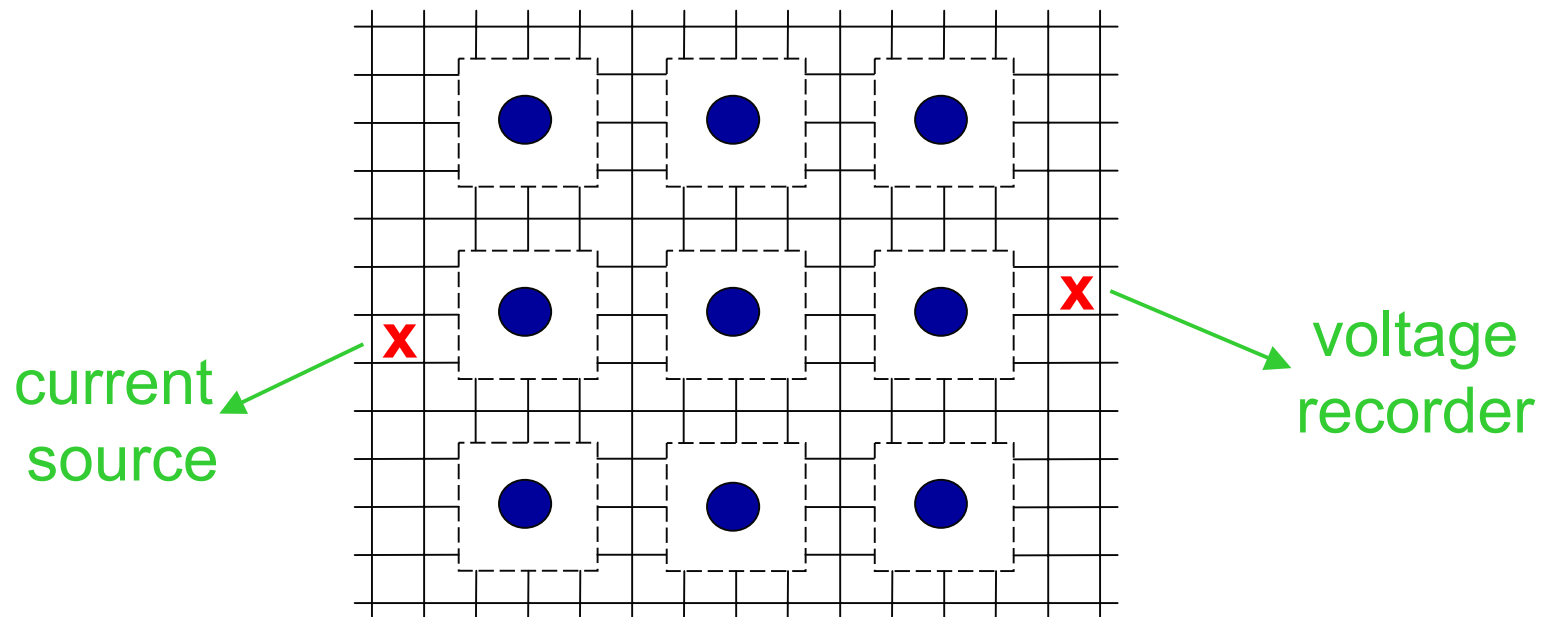
Example: subcell model for a dielectric wire (1)

- dielectric wire of radius = 0.9mm and $\varepsilon_r = 10$
- $\Delta_c = 1.3\text{mm}$ $\Delta_f = 0.1\text{mm}$ $r = 13$
- original number of state-space variables, $N = 4408$
- number of variables after ROM: either 12, 24 or 48
- subcell consists of 2×2 coarse cells

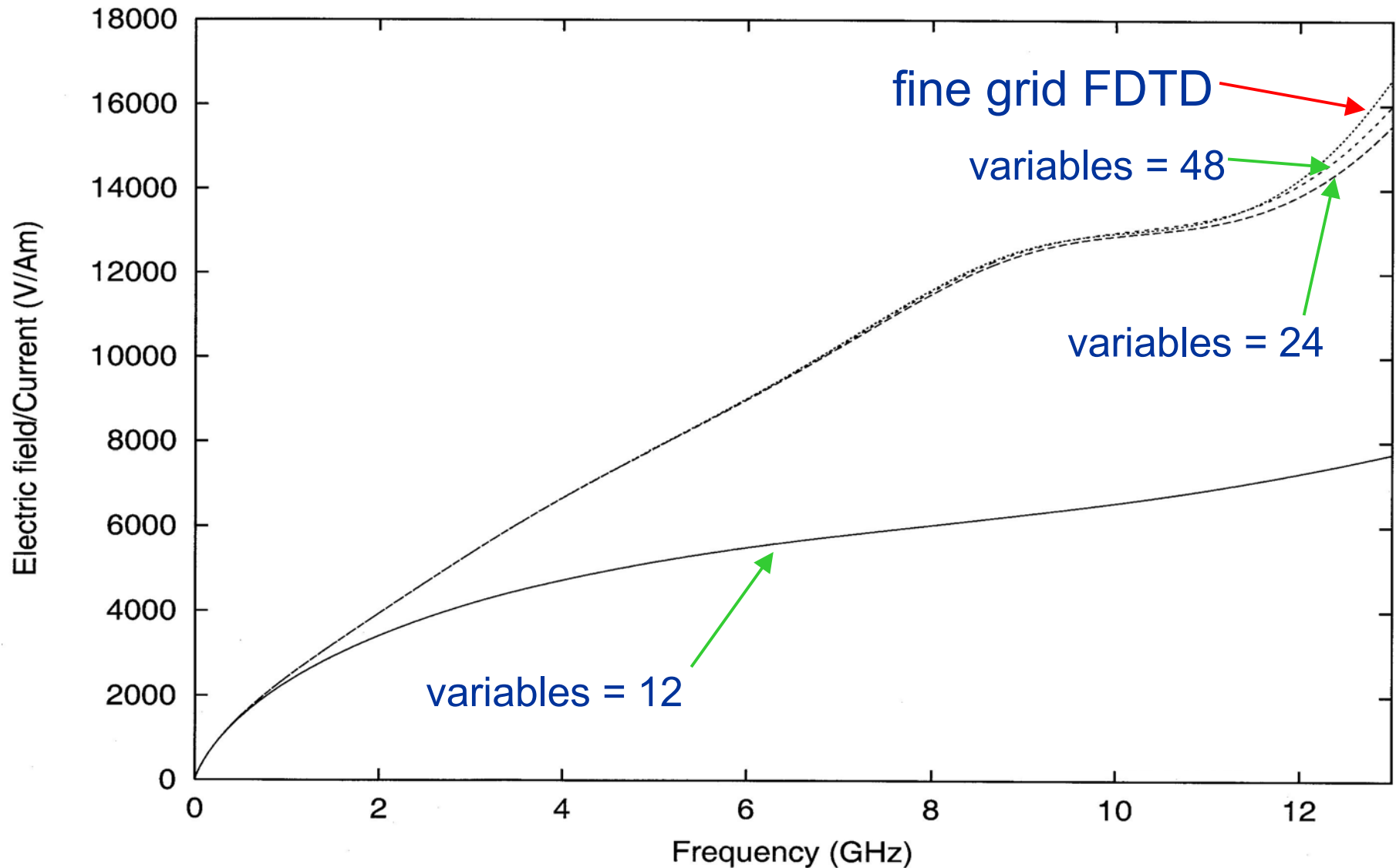


■ Example: dielectric wire (2)

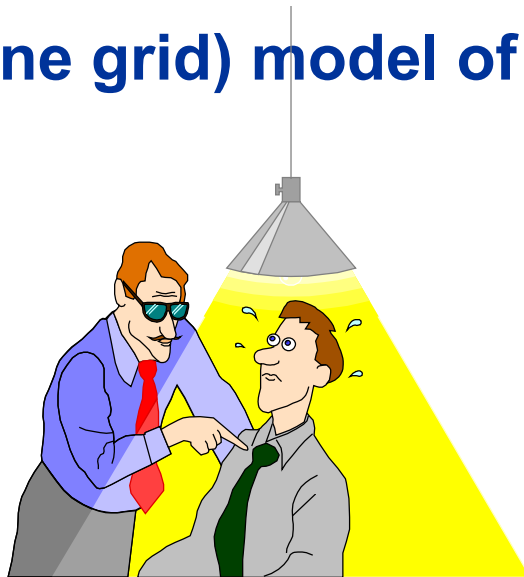
- reuse the subcell model e.g. in a periodic arrangement
- line source as excitation
- electric field recorded at other side of periodic structure
- comparison: normal FDTD with fine grid



Example: dielectric wire (3)



- A new reduced order modelling (ROM) technique based on representing time-domain signals in terms of a sum of Laguerre functions was proposed.
- It was shown how ROM can be used in EM field simulations using (2D) FDTD to combine an overall coarse grid with a subcell (fine grid) model of arbitrary objects.



Questions?