ToyFlow MC tool for FIT

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Introductions

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Motivation

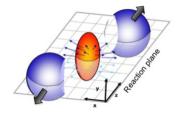
- Because of the initial pressure gradients and fluctuations, the particle azimuthal distribution is not isotropic.
- This is why it is customary to expand it in a Fourier series:

$$E\frac{\mathrm{d}^3N}{\mathrm{d}^3p}\propto 1+\sum_{n=1}^{\infty}2v_n\cos\left(n\left(\phi-\Psi_{\mathrm{RP}}\right)\right)$$

where

$$v_n = \langle \cos \left[n \left(\phi_i - \Psi_{\mathrm{RP}} \right) \right] \rangle$$

and angled brackets mean an average over all particles in all events.

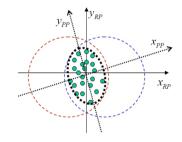


- Problem: $\Psi_{\rm RP}$ is unknown when using real data.
- Define the event flow vector $\mathbf{Q_n} \in \mathbb{C}$:

$$Q_{n,\mathsf{x}} = \sum_{i \in \mathrm{particles}} w_i \cos\left(n\phi_i\right) = |\mathbf{Q_n}| \cos\left(n\Psi_n\right)$$

$$Q_{n,y} = \sum_{i \in \text{particles}} w_i \sin(n\phi_i) = |\mathbf{Q_n}| \sin(n\Psi_n),$$

where w_i are the weights and optimally $w_i \approx v_n(p_T, y) \propto p_T$. Figure: From with small p_T . Many times it is set $w_i = 1$ for all i.



arXiv:0809.2949v2

Now we have

$$\Psi_n = rac{1}{n} \arctan \left(rac{Q_{n,y}}{Q_{n,x}}
ight)$$

and now the observed v_n can be calculated using

$$v_n^{\text{obs}} = \langle \cos \left[n \left(\phi_i - \Psi_n \right) \right] \rangle$$

 Autocorrelations need to be removed when using this equation.

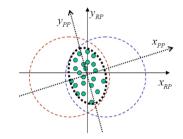


Figure: From arXiv:0809.2949v2

• It is easily shown that there is a relation

$$v_n^{\text{obs}} = v_n \cos \left[n \left(\Psi_n - \Psi_{\text{RP}} \right) \right].$$

To help, a resolution parameter is defined as

$$R_{n,true} = \langle \cos \left[n \left(\Psi_n - \Psi_{\text{RP}} \right) \right] \rangle$$

so that

$$v_n = \frac{v_n^{\text{obs}}}{R_{n,true}}.$$

The above definition of R_n can only be used with simulations as the $\Psi_{\rm RP}$ is unknown.

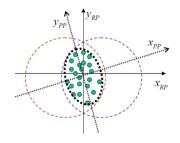


Figure: From arXiv:0809.2949v2

 \bullet R_n can also be written as

$$R_{k}\left(\chi\right) = \frac{\sqrt{\pi}}{2}\chi e^{-\frac{\chi^{2}}{2}}\left[I_{\frac{k-1}{2}}\left(\frac{\chi^{2}}{2}\right) + I_{\frac{k+1}{2}}\left(\frac{\chi^{2}}{2}\right)\right],$$

where I is the modified Bessel function and

$$\chi = v_n \sqrt{M}$$

where M is the multiplicity of the event. When comparing the same harmonic number correlation and event plane, k=1 always.

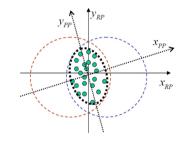


Figure: From arXiv:0809.2949v2

 By splitting the event into two equal size independent sub-events

$$R_{n,sub} = \sqrt{\left\langle \cos\left[n\left(\Psi_{n}^{\mathrm{A}} - \Psi_{n}^{\mathrm{B}}\right)\right]
ight
angle},$$

and then χ_{sub} can be solved iteratively from the previous page equation.

• Because $\chi \propto \sqrt{M}$, the full event $\chi = \sqrt{2}\,\chi_{\rm sub}$, and so finally full event resolution

$$R_n = R_{k=1} \left(\chi \right)$$

and

$$v_n = \frac{v_n^{\text{obs}}}{R_n}$$

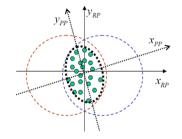


Figure: From arXiv:0809.2949v2

• Summary table:

Table: Is possible to use.

	Simulation	Data
$\Psi_{ m RP}$	×	
Ψ_n	×	X
$R_{n,true}$	×	
R_n	×	X

• When using a simulation the methods can be validated by comparing to the true values of reaction plane and $R_{n,true}$.

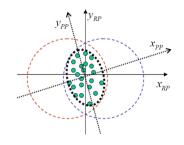


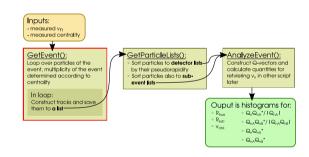
Figure: From arXiv:0809.2949v2

Toy Monte Carlo

- Toy MC to study the resolution parameter and flow coefficients
- \bullet Uses measured data for v_n and multiplicity from arXiv:1804.02944 and arXiv:1509.07299
- Available at https://github.com/hrytkone/ToyFlow

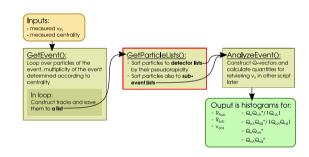
How the event is formed

- Event multiplicity taken from data according to centrality
- Then for each particle in the event
 - \bullet ϕ is taken from the Fourier series
 - ullet η is taken from multiplicity distribution



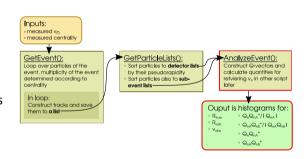
How the event is formed

- Particles in full event are divided into detector specific lists according to pseudorapidity
 - Detectors are TPC, FT0-A, FT0-C and FV0
- Particles are also divided into two sub events

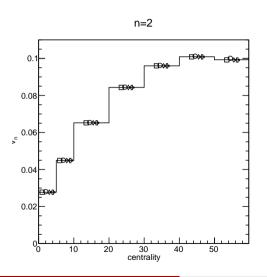


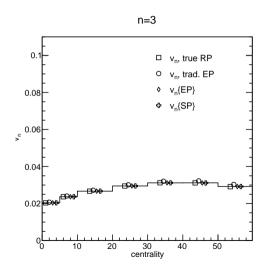
How the event is formed

- Q-vectors are then calculated
- Relevant information for retrieving the flow coefficients is saved to the histograms

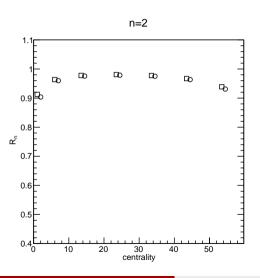


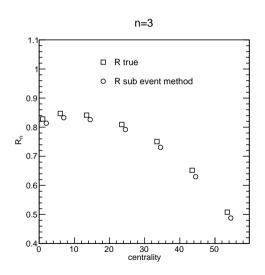
Validating





Validating (if needed)





FIT V0+

- Ideal vs. gran
- Compare to Maciej

TODO

- TODO:
 - Subevent handling
 - In FIT?
 - Realistic geometry.
- In the future:
 - Realistic simulation of FIT.

The end

Thank you for listening

Questions?

Event Plane (Scalar Product) method

• Another way to calculate v_n :

$$\left\langle \mathbf{Q_{n}} \frac{\mathbf{Q_{n,A}}^{*}}{|\mathbf{Q_{n,A}}|} \right\rangle = \left\langle \mathbf{Q_{n}} e^{-in\Psi_{n}} \right\rangle \left\langle \frac{\mathbf{Q_{n,A}}}{|\mathbf{Q_{n,A}}|} e^{-in\Psi_{n}} \right\rangle^{*} = \nu_{n} R_{n,A}$$

where the last equality is valid when Q_n and $Q_{n,A}$ are uncorrelated, except for the common Ψ_n .

• For two subevents with equal multiplicities

$$\left\langle \frac{\mathbf{Q}_{\mathsf{n},\mathsf{A}}}{|\mathbf{Q}_{\mathsf{n},\mathsf{A}}|} \frac{\mathbf{Q}_{\mathsf{n},\mathsf{B}}^*}{|\mathbf{Q}_{\mathsf{n},\mathsf{B}}|} \right\rangle = \left\langle \frac{\mathbf{Q}_{\mathsf{n},\mathsf{A}}}{|\mathbf{Q}_{\mathsf{n},\mathsf{A}}|} e^{-in\Psi_n} \right\rangle \left\langle \frac{\mathbf{Q}_{\mathsf{n},\mathsf{B}}}{|\mathbf{Q}_{\mathsf{n},\mathsf{B}}|} e^{-in\Psi_n} \right\rangle^* = \left| \left\langle \frac{\mathbf{Q}_{\mathsf{n},\mathsf{A}}}{|\mathbf{Q}_{\mathsf{n},\mathsf{A}}|} e^{-in\Psi_n} \right\rangle \right|^2 = R_{n,\mathsf{A}}^2$$

Event Plane (Scalar Product) method

• Thus the v_n can be calculated using event-plane method by

$$v_n\left\{\mathrm{EP}\right\} = \left\langle \mathbf{Q_n} \frac{\mathbf{Q_{n,A}}^*}{|\mathbf{Q_{n,A}}|} \right\rangle / \sqrt{\left\langle \frac{\mathbf{Q_{n,A}}}{|\mathbf{Q_{n,A}}|} \frac{\mathbf{Q_{n,B}}^*}{|\mathbf{Q_{n,B}}|} \right\rangle} \ .$$

- This means that to calculate v_n , one does not necessarily need the event plane information explicitly. Only the flow vectors are necessary.
- By comparing the v_n from different methods, the validation is better.