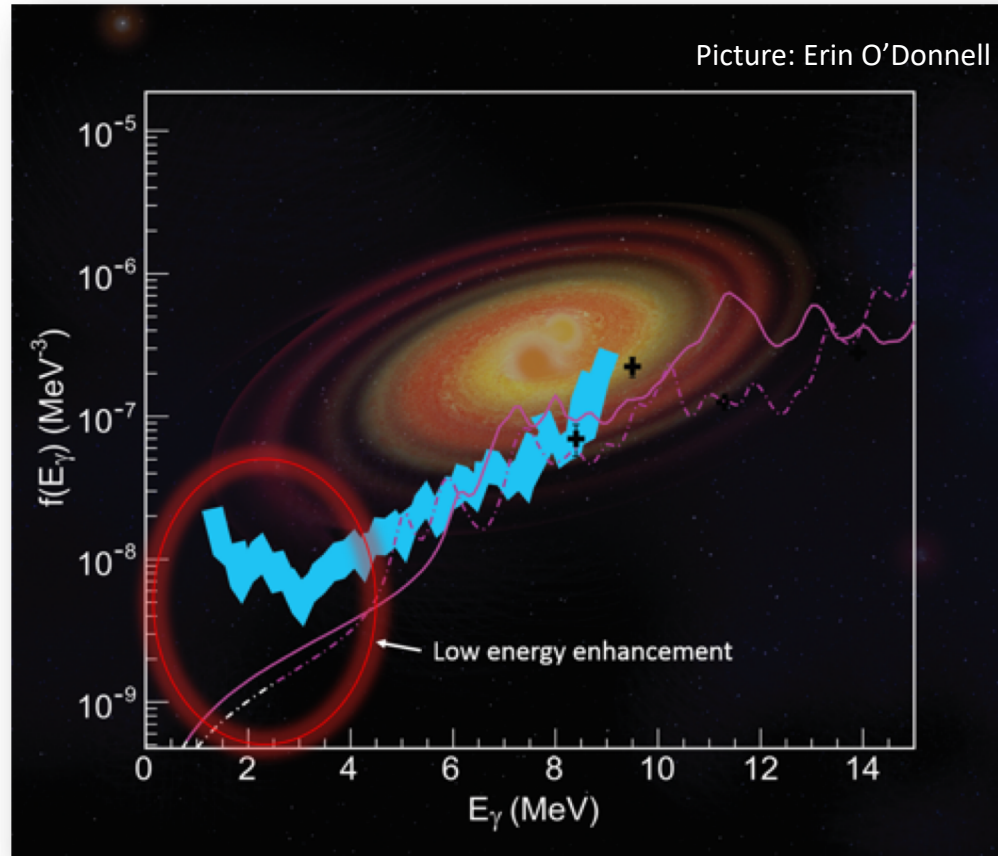


# Mini workshop on the beta-Oslo method: case study of $^{70}\text{Ni}$



Thursday, Aug 1, and Friday, Aug 2, 2019

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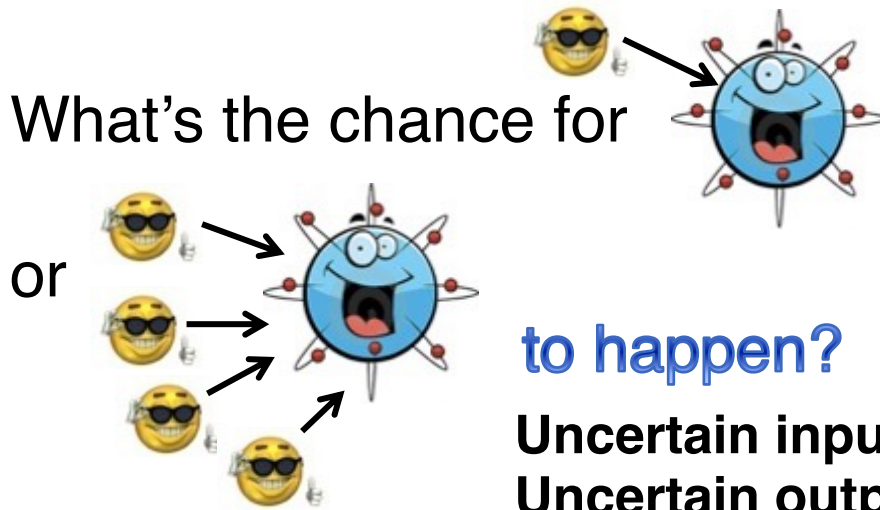


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7. Normalization of the  $\gamma$ -ray transmission coefficient and getting the dipole  $\gamma$ -ray strength function using *normalization.c*
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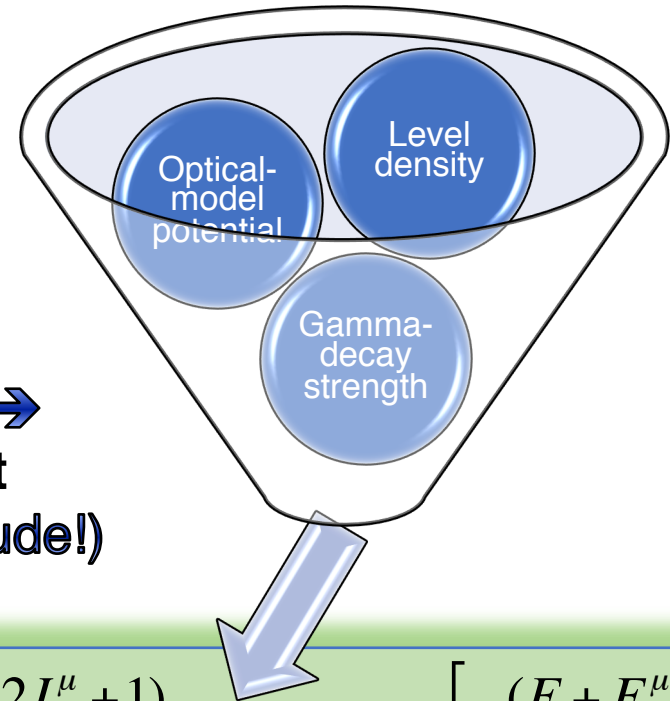
# 1. Introduction/motivation

# Nuclear input for radiative-capture rates



to happen?

**Uncertain input →  
Uncertain output  
(orders of magnitude!)**



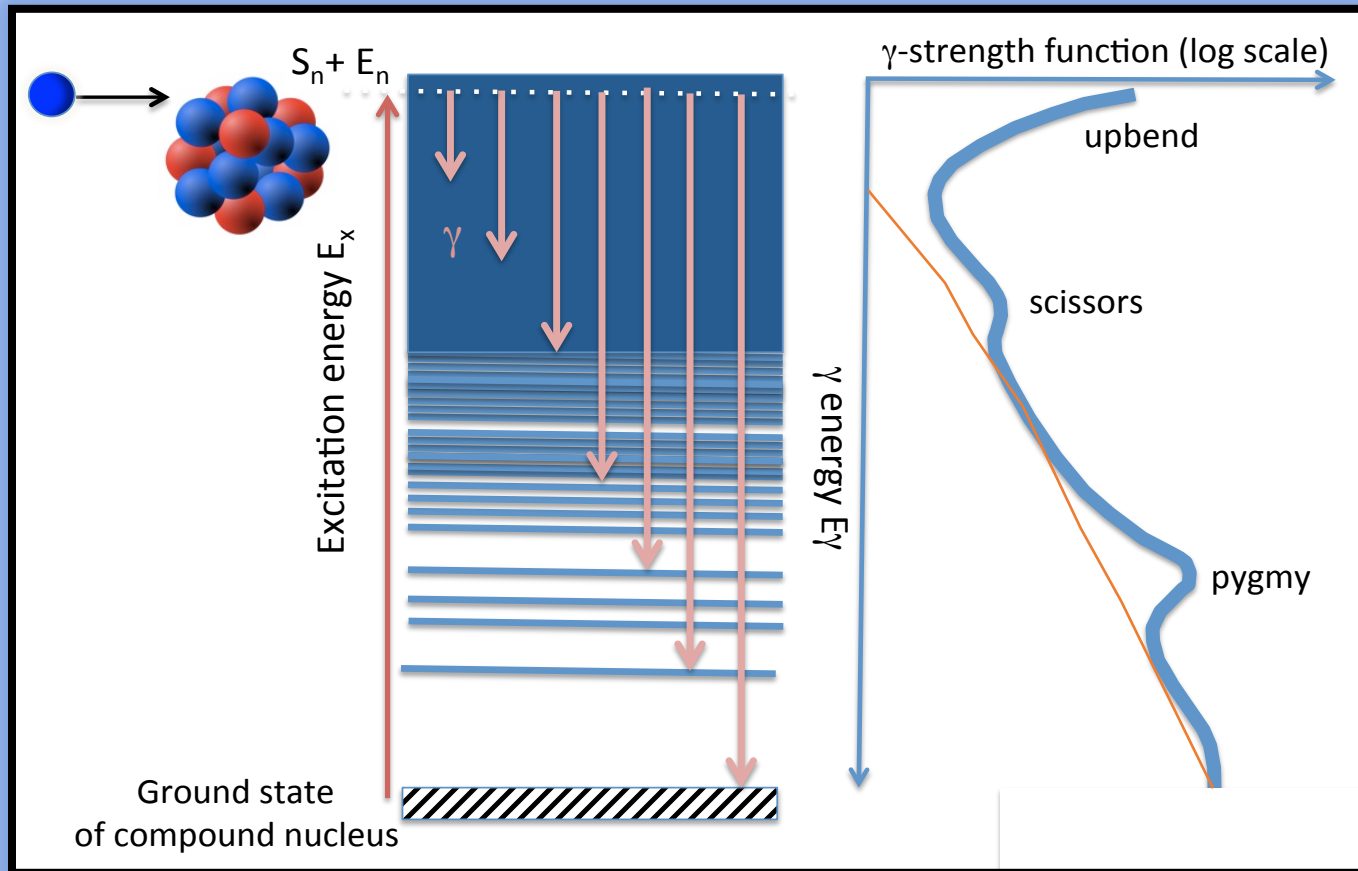
$$N_A \langle \sigma v \rangle (T) = \left( \frac{8}{\pi m} \right)^{1/2} \frac{N_A}{(kT)^{3/2} G(T)} \int_0^\infty \sum_\mu \frac{(2I^\mu + 1)}{(2I^0 + 1)} \sigma^\mu(E) E \exp \left[ -\frac{(E + E_x^\mu)}{kT} \right] dE$$

$$G(T) = \sum_\mu (2I^\mu + 1) / (2I^0 + 1) \exp(-E_x^\mu / kT)$$

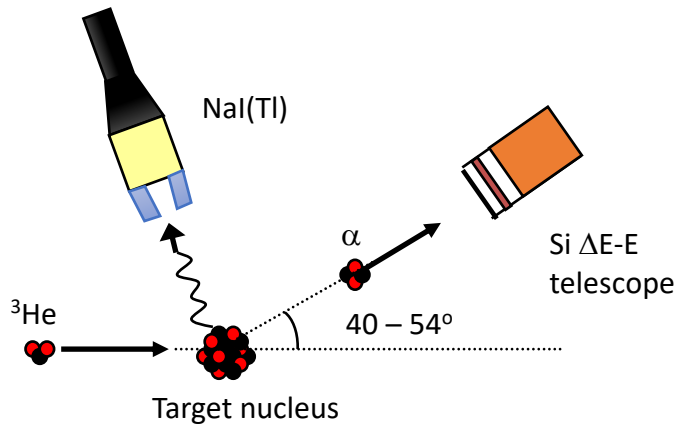
# Nuclear input for radiative-capture rates

**Level density:** Available quantum levels at a given excitation-energy bin per energy unit

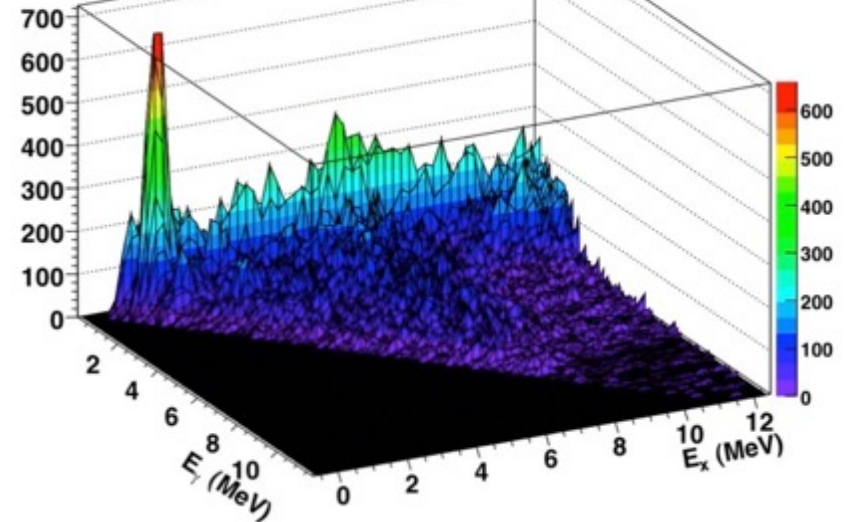
**Gamma strength:** The nucleus' ability to emit electromagnetic radiation when other known dependencies are removed (level density,  $E_g^3$  factor for dipole radiation) -> directly connected to reduced transition probabilities



# The Oslo method in a



$^{45}\text{Sc}(^3\text{He},\alpha\gamma)^{44}\text{Sc}$  data



0. Get a hold of an  $(E_\gamma, E_x)$  matrix ( $\sim 30\text{-}40\ 000$  coincidences)
1. Correct for the NaI response [Guttormsen et al., NIM A 374, 371 (1996)]
2. Extract *distribution of primary  $\gamma$ s* for each  $E_x$  [Guttormsen et al., NIM A 255, 518 (1987)]
3. Get level density and  $\gamma$ -strength from primary  $\gamma$ 's [Schiller et al., NIM A 447, 498 (2000)]
4. Normalize & evaluate systematic errors [Schiller et al., NIM A 447, 498 (2000), Larsen et al., PRC **83**, 034315 (2011)]

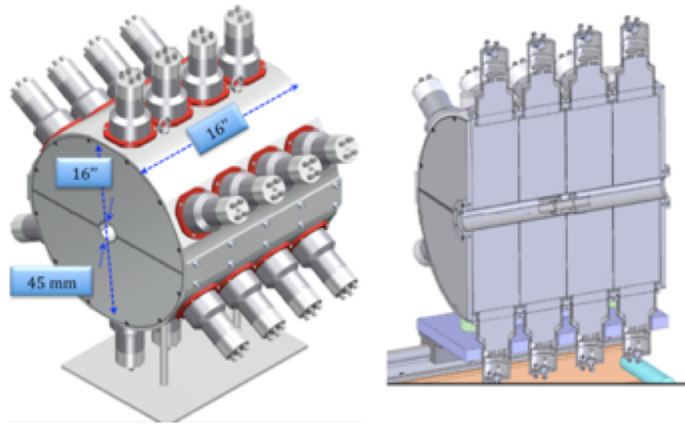
Data and references: [oel.uio.no/compilation/](http://oel.uio.no/compilation/)

Analysis codes and tools:

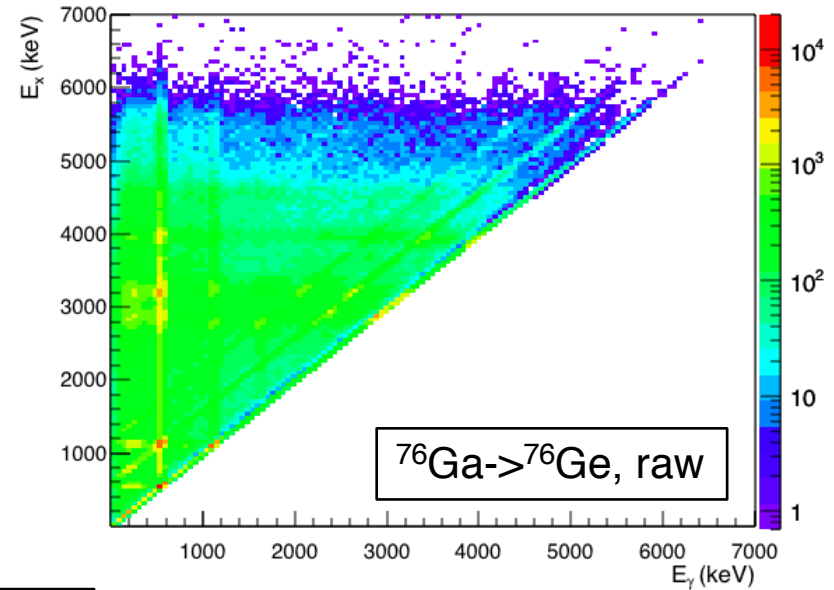
[github.com/oslocyclotronlab/oslo-method-software](https://github.com/oslocyclotronlab/oslo-method-software)

New Python Oslo-method package: **OMP**y by J.E. Midtbø, F.Zeiser, E. Lima

# The beta-Oslo method in a



[A. Simon, S.J. Quinn, A. Spyrou et al, NIM A 703, 16 (2013)]



## Recipe:

- 1) Implant a neutron-rich nucleus in a segmented total-absorption spectrometer (preferably with  $Q_{\text{beta}} \approx S_n$ )
- 2) Measure  $\beta$ -particle in coincidence with  $\gamma$  rays from the daughter nucleus
- 3) Apply Oslo method to the  $E_x$ - $E_\gamma$  matrix to extract level density &  $\gamma$ -ray strength

A. Spyrou, S.N. Liddick, A.C.Larsen, M. Guttormsen et al., Phys. Rev. Lett. **113**, 232502 (2014)



## 2. Case study: $^{70}\text{Ni}$



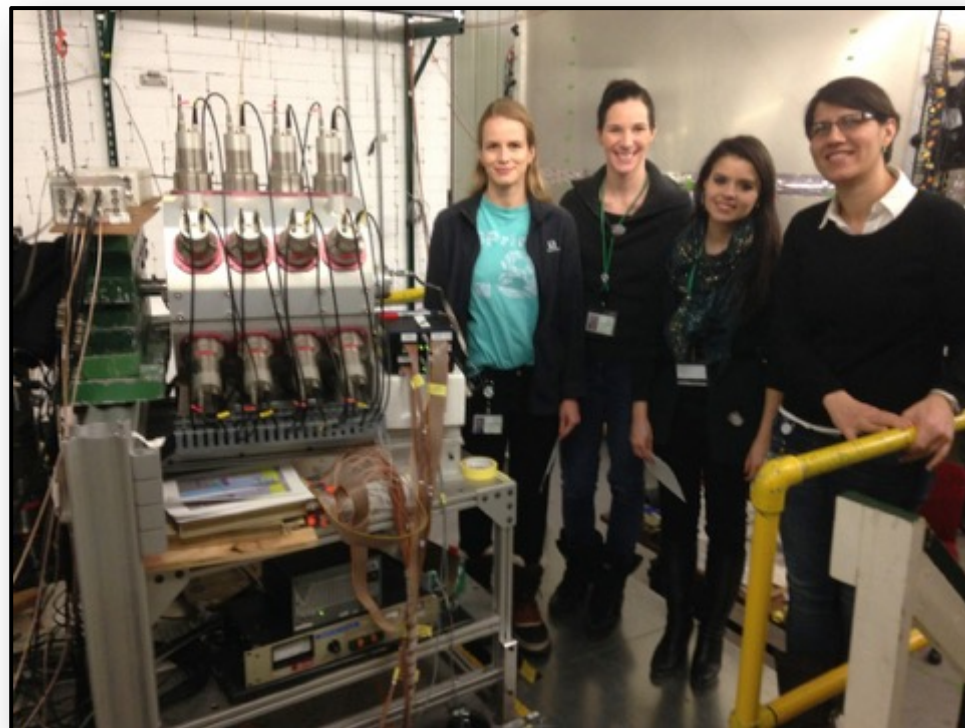
# Some experimental details, $^{70}\text{Ni}$

Discretionary beam time @ NSCL/MSU, February 2015;  $^{70}\text{Co} \rightarrow ^{70}\text{Ni}$   
 $^{86}\text{Kr}$  primary beam, 140 MeV/nucleon on thick Be target producing  $^{70}\text{Co} ++$   
A1900 mass separator optimized for  $^{70}\text{Co}$   
DSSD inside SuN for detecting  $^{70}\text{Co}$  &  $\beta^-$

$^{70}\text{Co}$   $T_{1/2}$ : 112 ms  
 $^{70}\text{Co}$   $|\pi = (6^-, 7^-)$  and isomer  $(2^+, 3^+)$   
Beta-decay Q-value: 12.3 MeV  
 $S_n, ^{70}\text{Ni}$ : 7.3 MeV

S.N. Liddick et al., Phys. Rev. Lett. **116**,  
242502 (2016)

A.C. Larsen et al., Phys. Rev. C **97**, 054329  
(2018)



# The $^{70}\text{Ni}$ analysis steps

1. Unfold the raw matrix on the excitation energy axis
2. Unfold the  $E_x$ -unfolded matrix on the  $\gamma$ -energy axis
3. Extract primary  $\gamma$  rays for each excitation-energy bin
4. Extract level density and  $\gamma$ -ray transmission coefficient (which is directly proportional to the  $\gamma$ -strength function) from the primary  $\gamma$ -ray matrix
5. Normalize the level density to known discrete levels and Hartree-Fock-Bogoliubov calculations of Stephane Goriely et al. at high  $E_x$  (consider the populated spin range...)
6. Normalize the  $\gamma$ -ray transmission coefficient and get the  $\gamma$ -ray strength function
7. Use the level-density and  $\gamma$ -strength data to guide models to be used as input in the nuclear reaction code TALYS
8. Calculate the  $^{69}\text{Ni}(n,\gamma)^{70}\text{Ni}$  cross section and reaction rate with the data-constrained input for the level density and  $\gamma$  strength

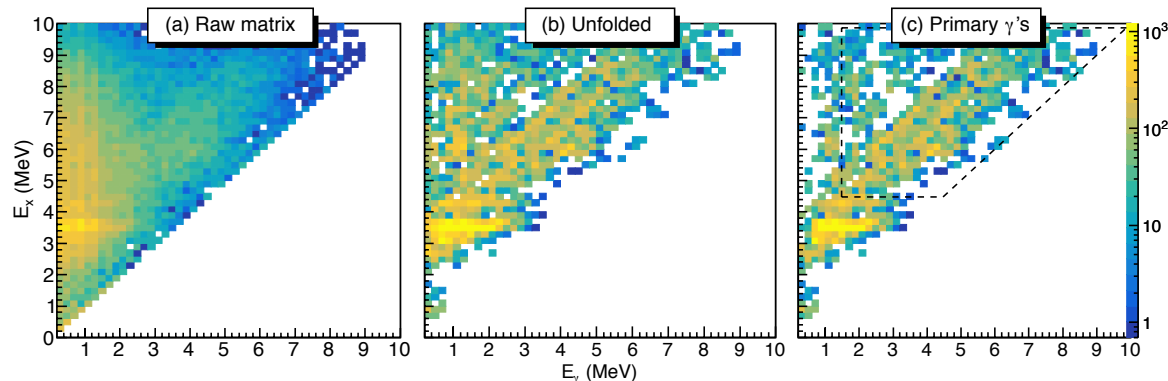
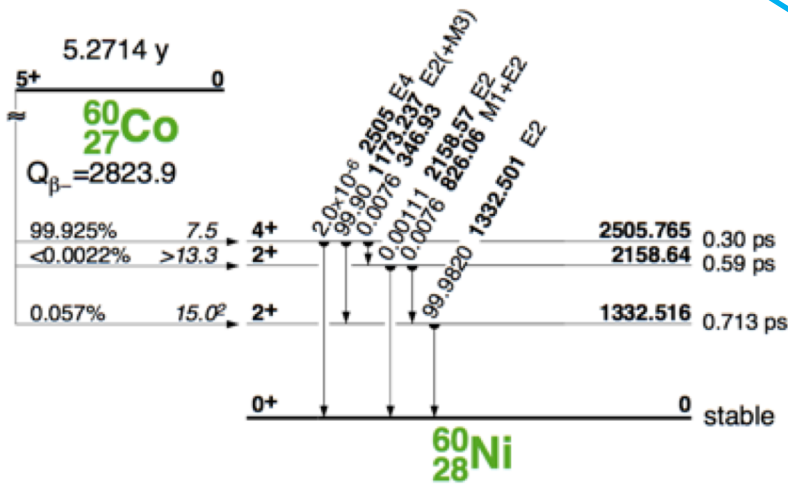


Figure from A.C. Larsen et al., Phys. Rev. C **97**, 054329 (2018)

### 3. $E_x$ and $E_\gamma$ unfolding using MaMa

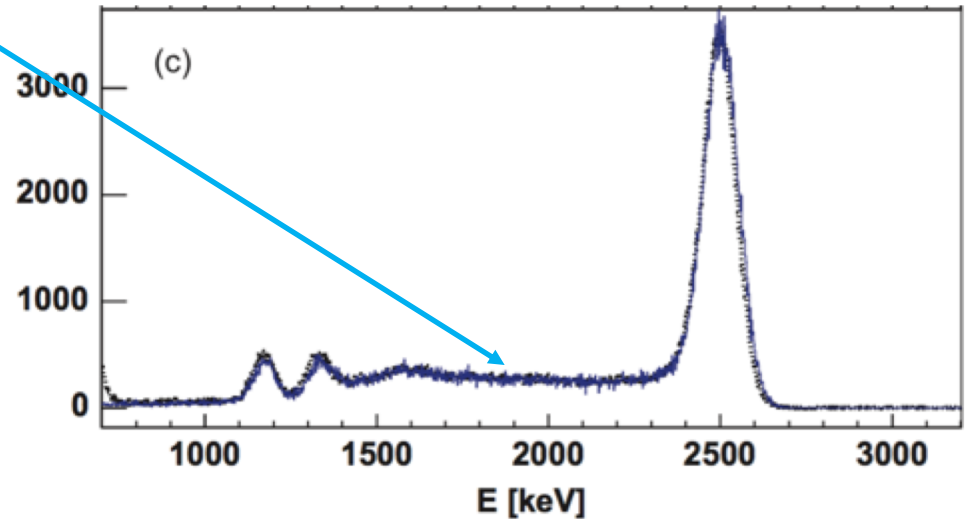
# Incomplete summing: what is it?

Note the “tail” towards low  $E_x$ !



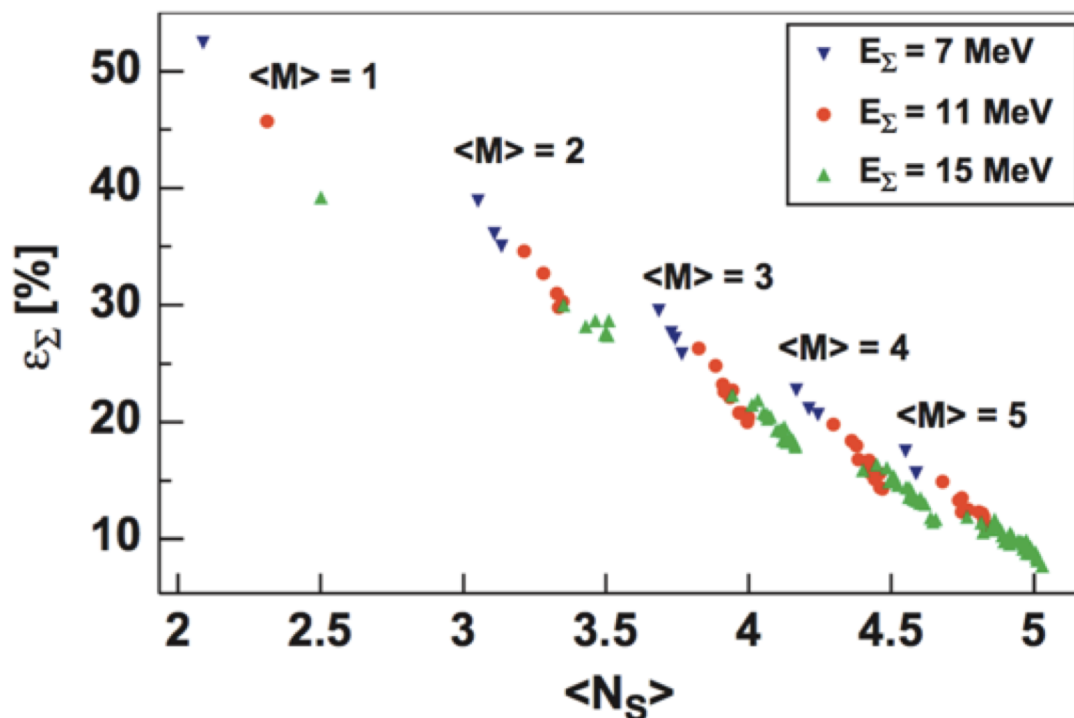
[From Table of isotopes, R.B. Firestone]

Sum of all segments,  $^{60}\text{Co}$  source (1173keV+1332keV)



[A. Simon, S.J. Quinn, A. Spyrou et al, NIM A 703, 16 (2013)]

# Summing efficiency and multiplicity



Summing efficiency of the SuN detector as a function of the average number of hits  $\langle N_S \rangle$  for various sum-peak energies

# $E_x$ unfolding, $^{70}\text{Ni}$

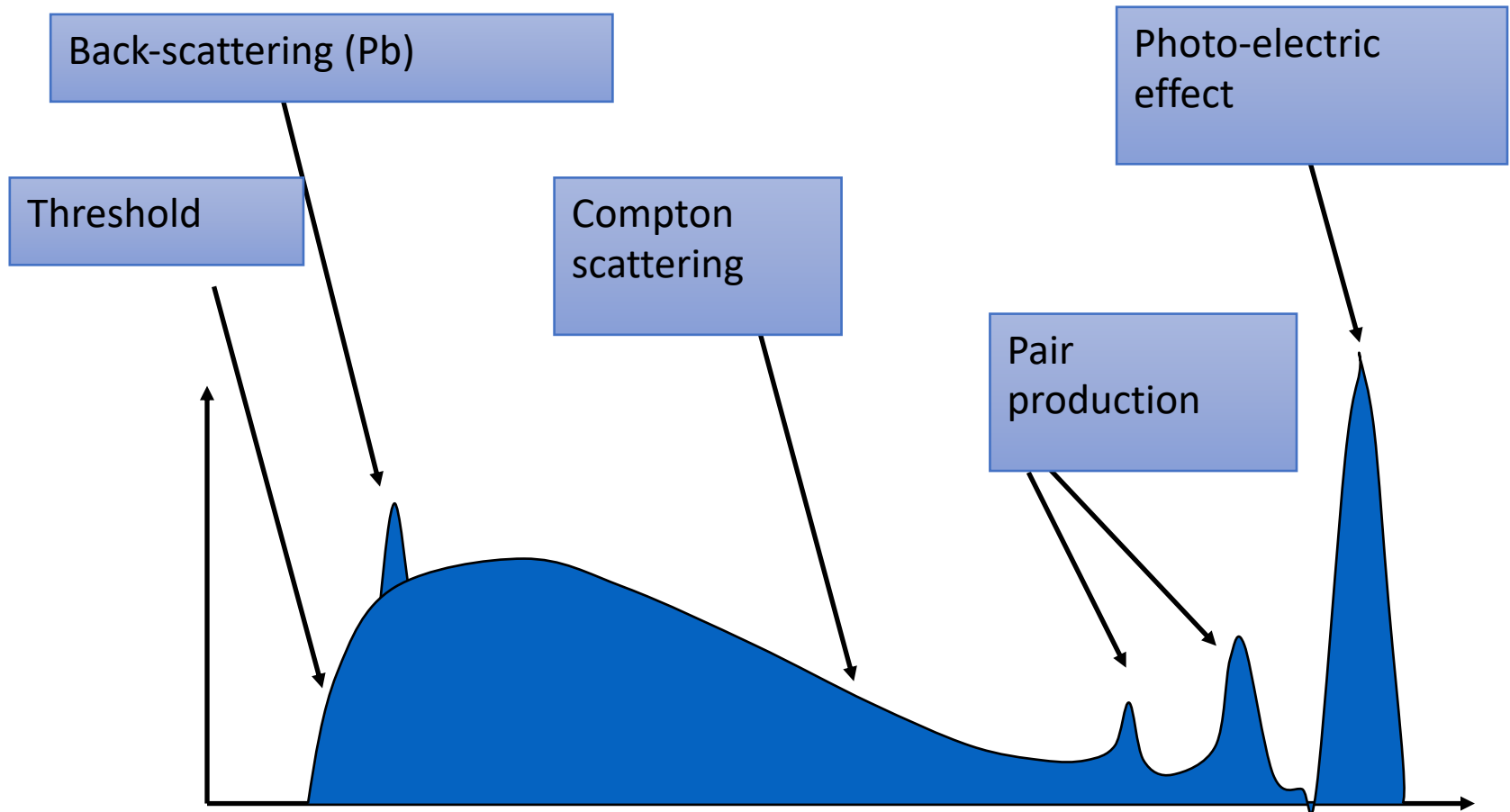
In the  $\beta$ -Oslo method, the sum of all segments gives the initial  $E_x$ , but:

- > if incomplete summing, the obtained  $E_x$  is not correct
- > the higher the  $\gamma$  multiplicity, the lower the SuN efficiency
- > the high  $Q_\beta$  value gives a background from electrons interacting with SuN

⇒ We need to unfold the  $E_x$  axis

The next slides are from Magne's presentation at a nuclear-physics group meeting in Oslo, 2017

# One crystal NaI response function



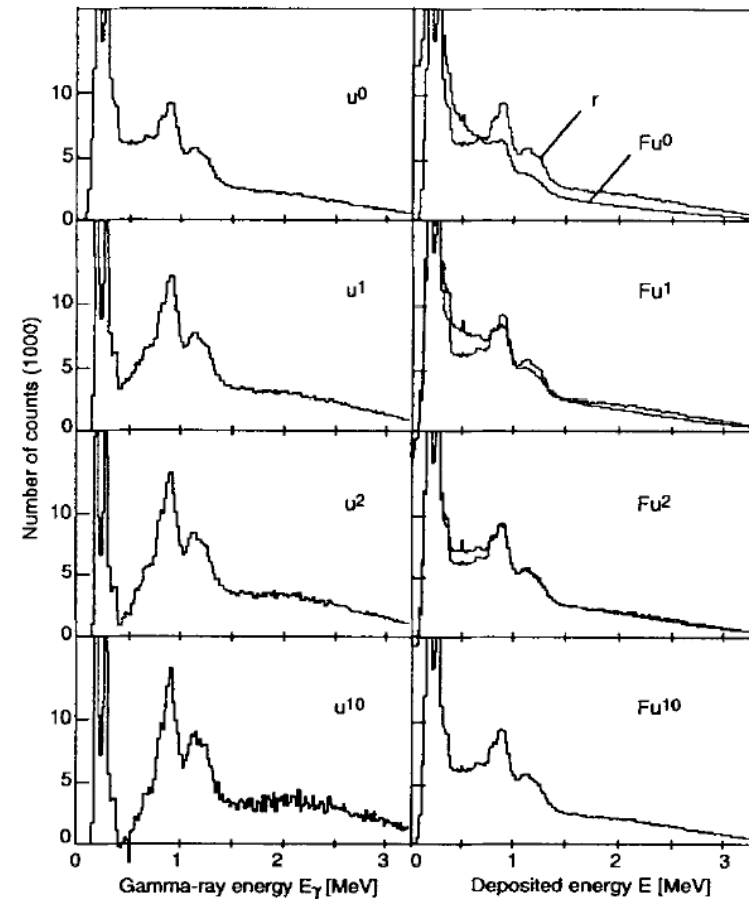
# The unfolding algorithm

We know how to fold:

$$\mathbf{f} = \mathbf{R}\mathbf{u}$$

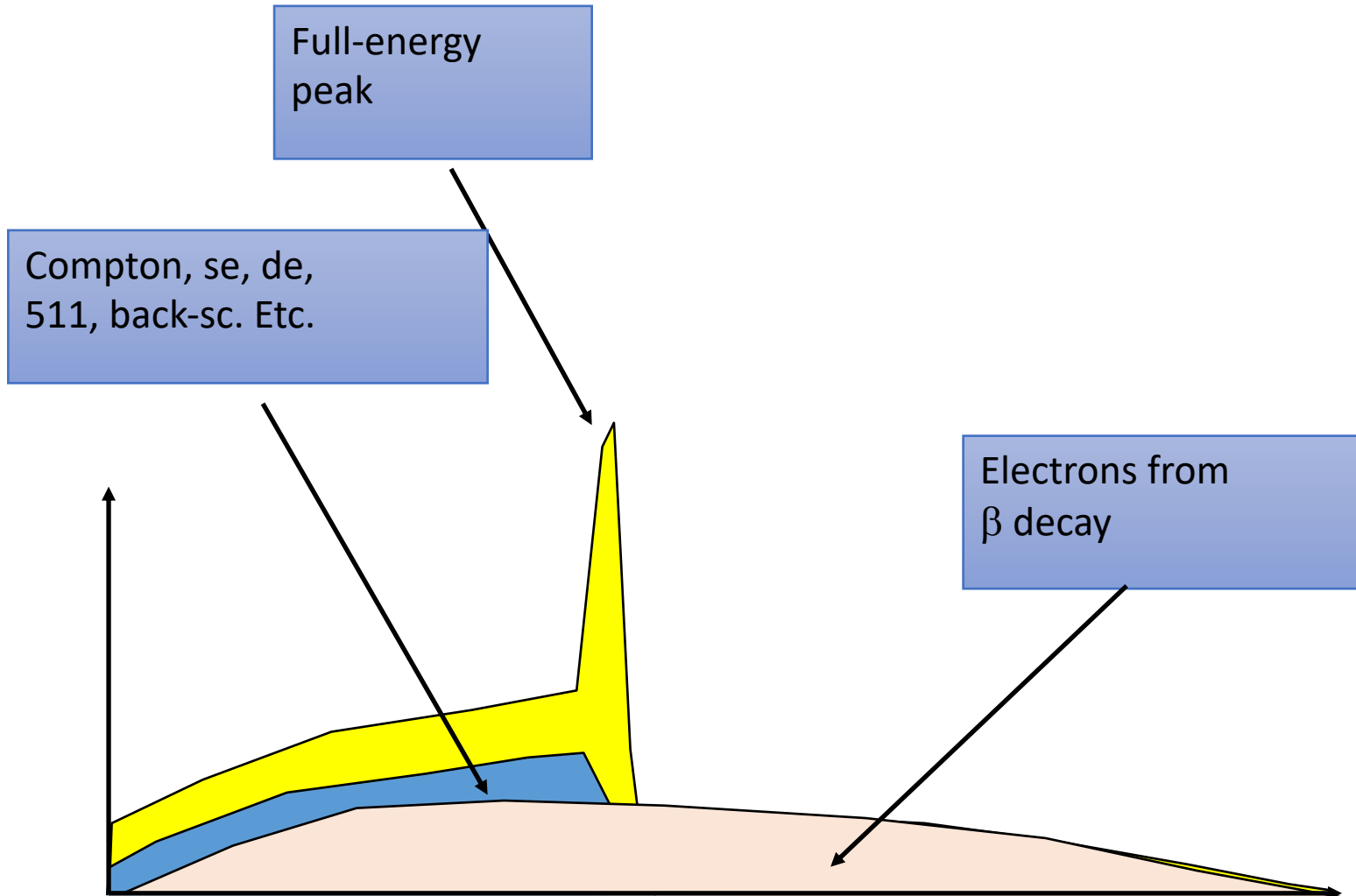
- (i) First trial function:  $\mathbf{u}^0 = \mathbf{r}$
  - (ii) First folded spectrum:  $\mathbf{f}^0 = \mathbf{R}\mathbf{u}^0$
  - (iii) Correct for how much we fail:  $\mathbf{u}^1 = \mathbf{u}^0 + (\mathbf{r} - \mathbf{f}^0)$
  - (iv) Second folded spectrum:  $\mathbf{f}^1 = \mathbf{R}\mathbf{u}^1$
  - (v) The third trial function:  $\mathbf{u}^2 = \mathbf{u}^1 + (\mathbf{r} - \mathbf{f}^1)$
- and so on until  $\mathbf{f}^i \approx \mathbf{r}$ .

Guttormsen et al., NIM A 374, 371 (1996)



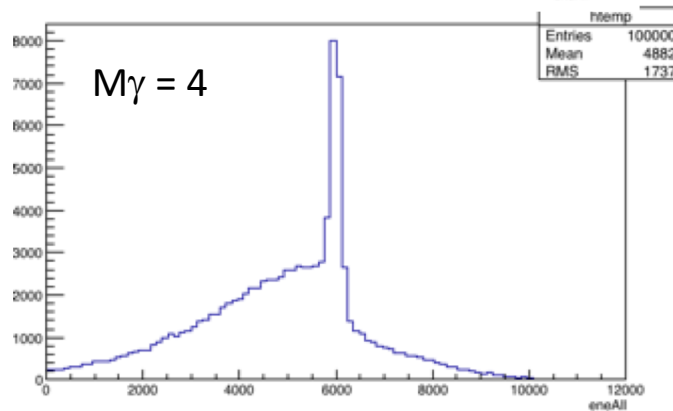
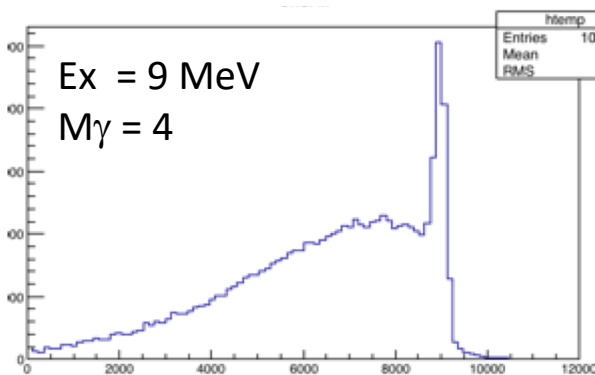
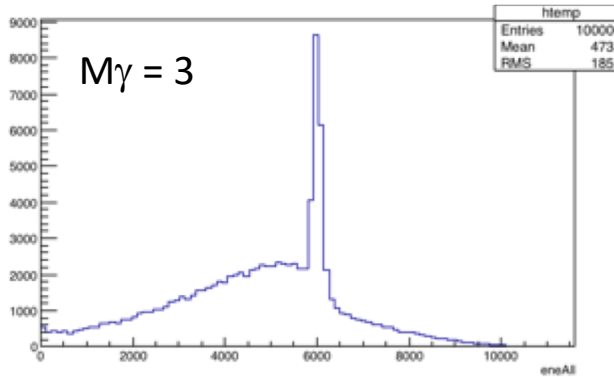
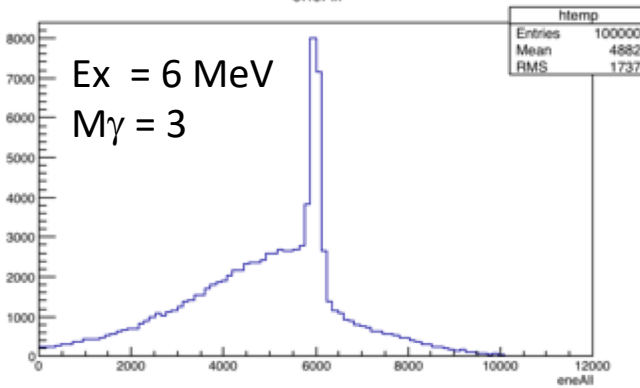
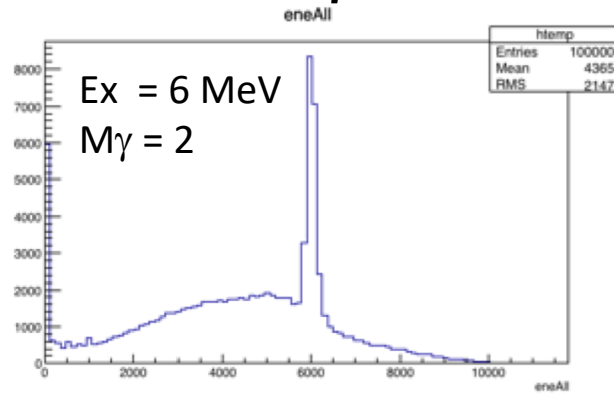
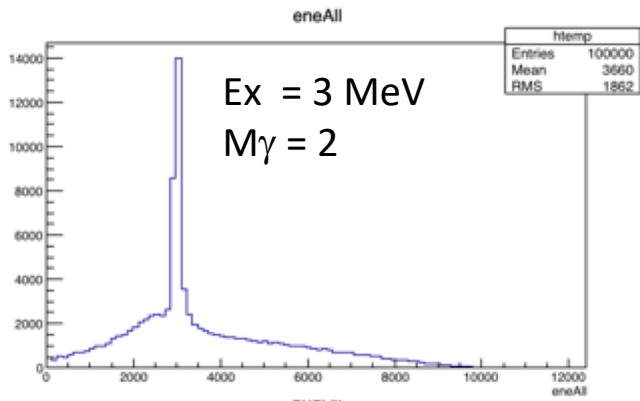


$$E_x = \sum E_\gamma \text{ NaI response}$$



# Dependence on $E_x$ and $M_\gamma$

( GEANT4 simulations  
by Artemis Spyrou )



# Multiplicity interpolation

Weights:

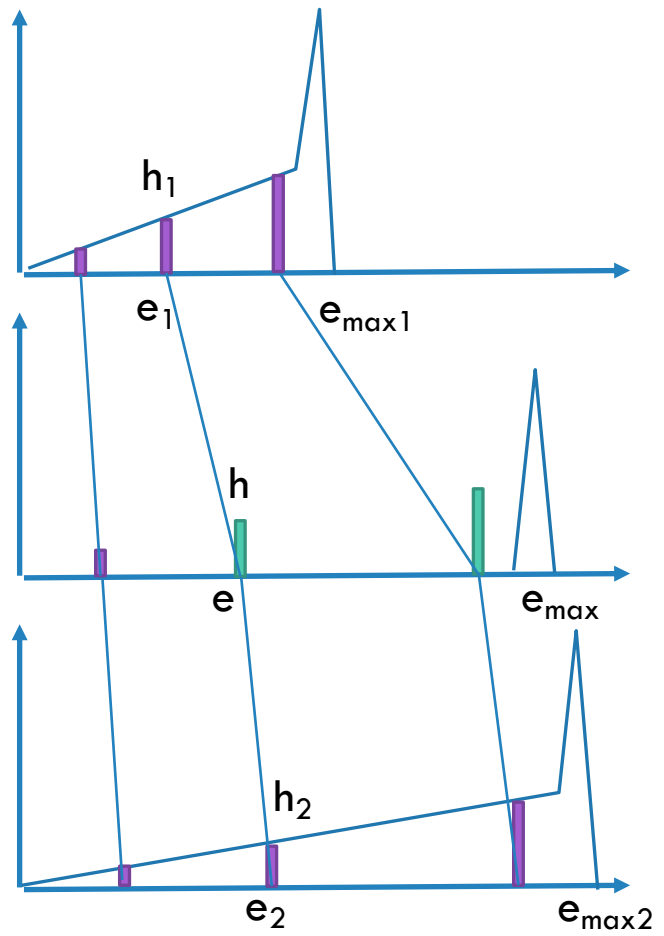
$$wm_1 = (m_2 - m)/(m_2 - m_1)$$

$$wm_2 = 1 - wm_1$$

For all h's at the same Ex:

$$h(e) = wm_1 h_{m_1}(e) + wm_2 h_{m_2}(e)$$

# $E_x$ interpolation below the full-energy peak



Channel energy:

$$e_1 = e (e_{max1} / e_{max})$$

$$e_2 = e (e_{max2} / e_{max})$$

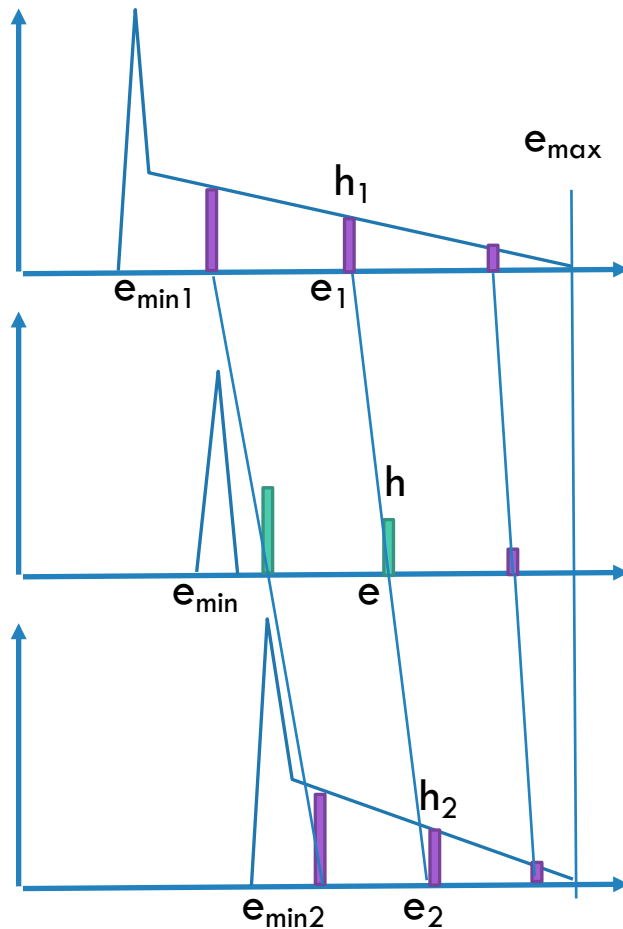
Weights:

$$w_{e1} = (e_{max2} - e) / (e_{max2} - e_{max1})$$

$$w_{e2} = 1 - w_{e1}$$

$$h(e) = w_{e1} h_1(e_1) + w_{e2} h_2(e_2)$$

# $E_x$ interpolation above the full-energy peak



$$e_{\max} = 12 \text{ MeV}$$

Channel energy:

$$e_1 = e_{\min 1} + (e - e_{\min})(e_{\max} - e_{\min 1}) / (e_{\max} - e_{\min})$$

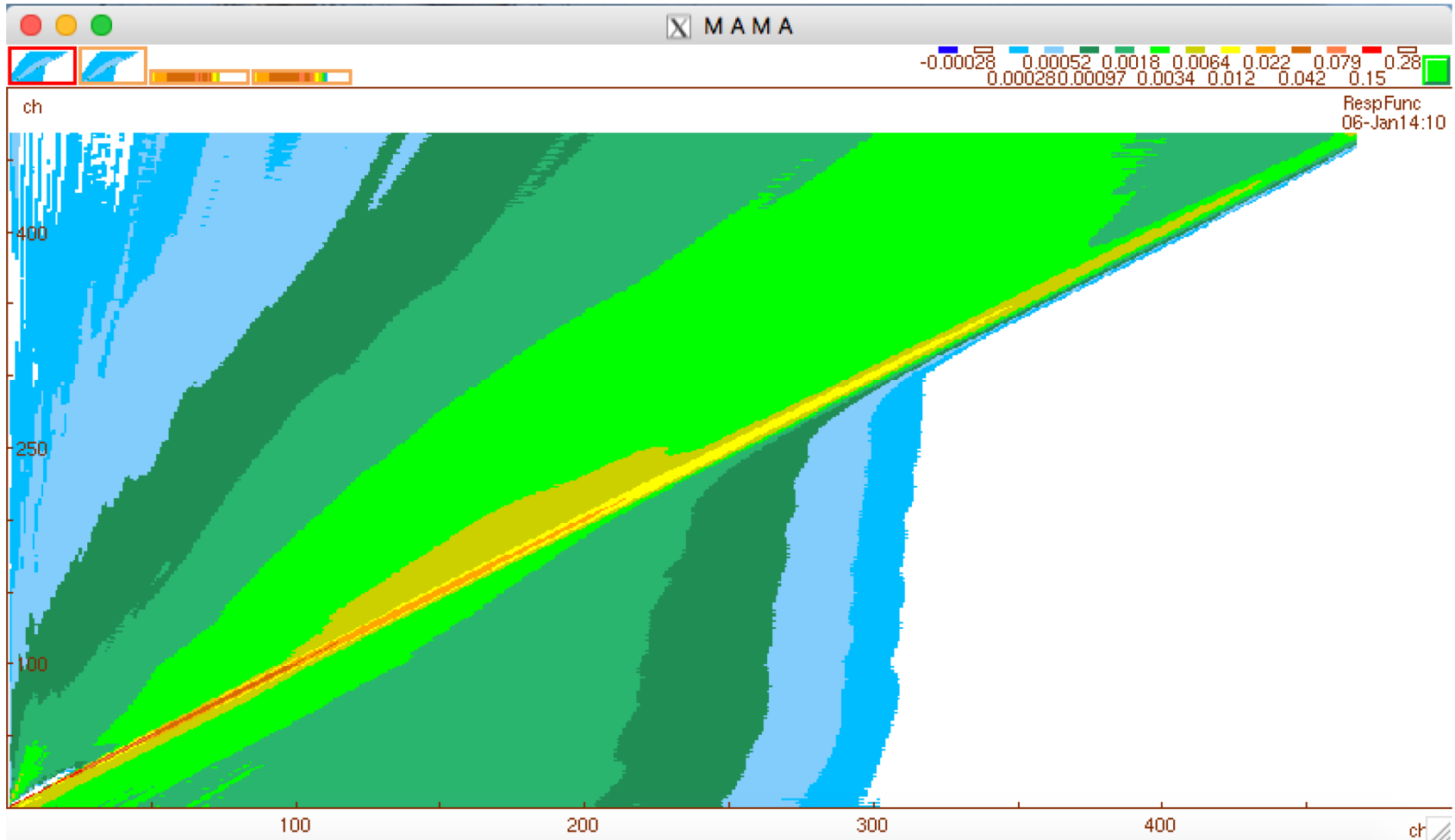
$$e_2 = e_{\min 2} + (e - e_{\min})(e_{\max} - e_{\min 2}) / (e_{\max} - e_{\min})$$

$$h(e) = w e_1 h_1(e_1) + w e_2 h_2(e_2)$$

# The $E_x$ response matrix

The y-axis is the true  $E_x$  populated. The x-axis is the  $E_x$  values detected by SuN.

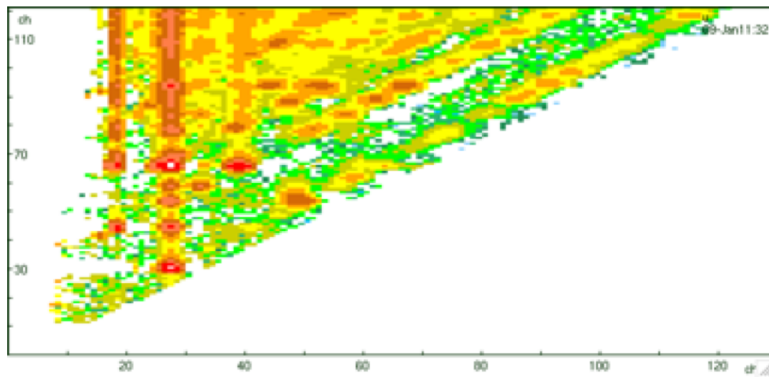
Note that we observe values below (incomplete summing) and above (electron energies) the true  $E_x$  value.



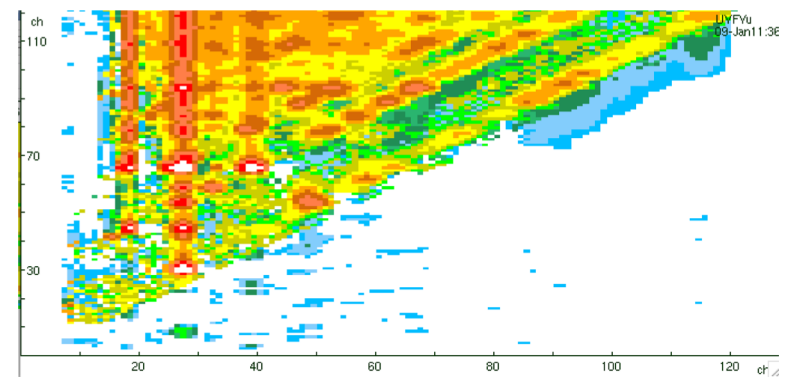
# Folding and unfolding the $E_x$ axis

Starts with an unfolded  $^{92}\text{Zr}$  matrix  $U_0$  from (p,p) reaction with Oslo method. Then fold along y-axis  $F_y(U_0)$  and finally unfold back again, so that  $U_0 = U_y(F_y(U_0))$ . Note all the “curtains” hanging down below the  $E_x = E_g$  diagonal.

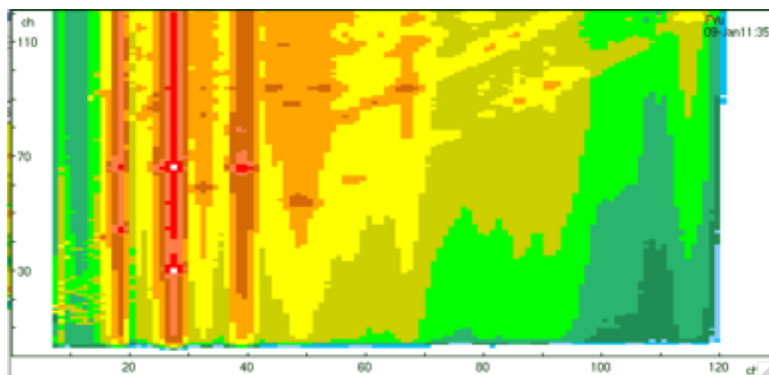
$U_0$



$U_y(F_y(U_0))$



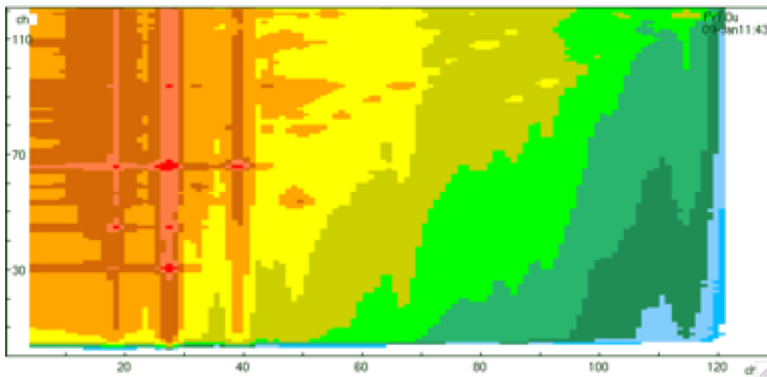
$F_y(U_0)$



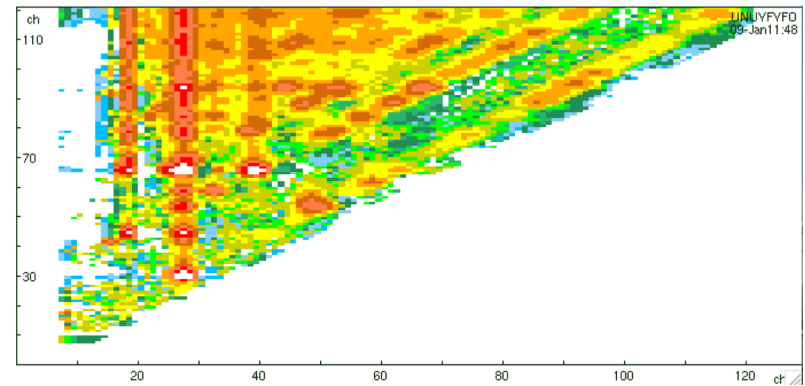
# Folding and unfolding $E_x$ and $E_y$ axis

Starts with  $U_0$  folded along  $E_g$  and  $E_x$  axis:  $F_y(F_x(U_0))$ . Then unfolding along  $E_x$ -axis:  $U_y(F_y(F_x(U_0)))$ . Then unfold  $E_g$ -axis in order to obtain the original the  $U_0$  matrix: so that  $U_0 = U_x(U_y(F_y(F_x(U_0))))$ .

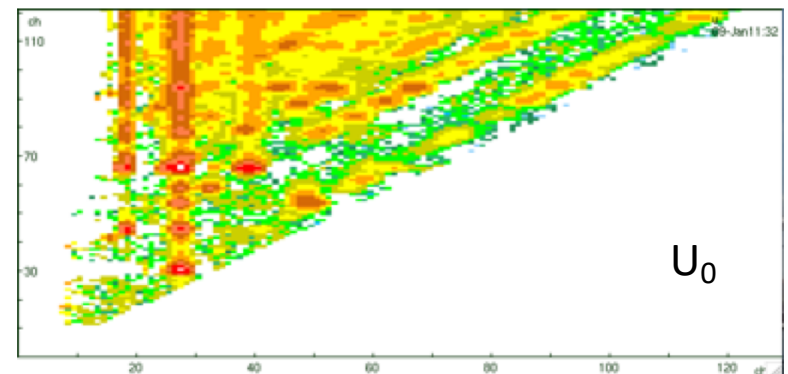
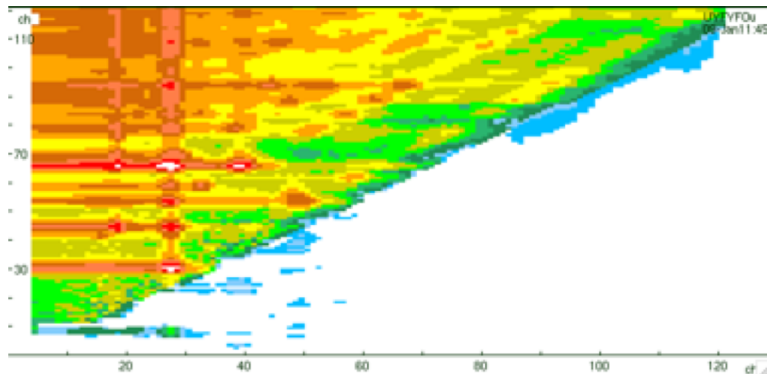
$$F_y(F_x(U_0))$$



$$U_x(U_y(F_y(F_x(U_0))))$$

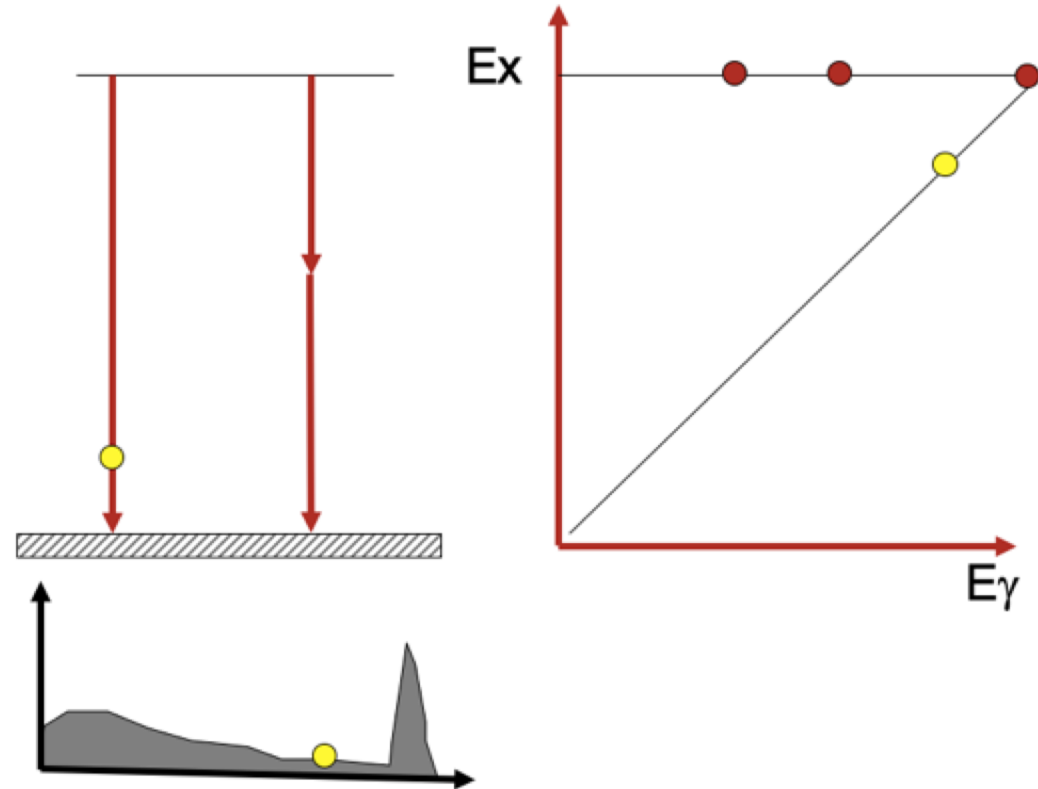
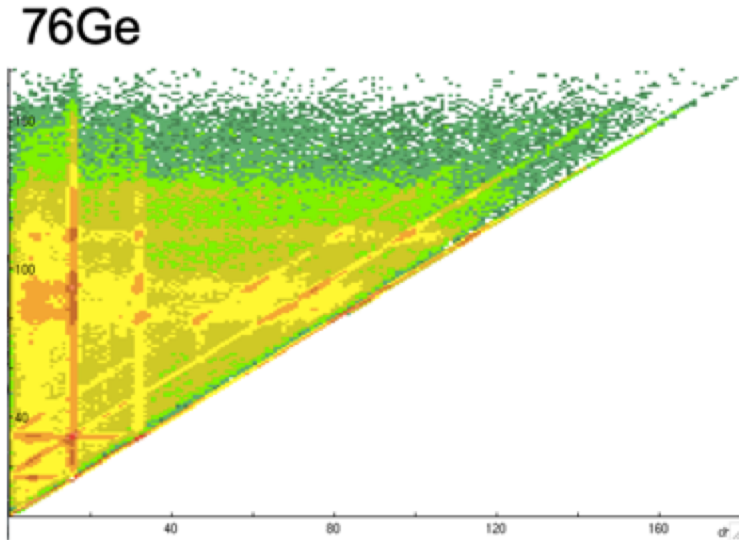


$$U_y(F_y(F_x(U_0)))$$





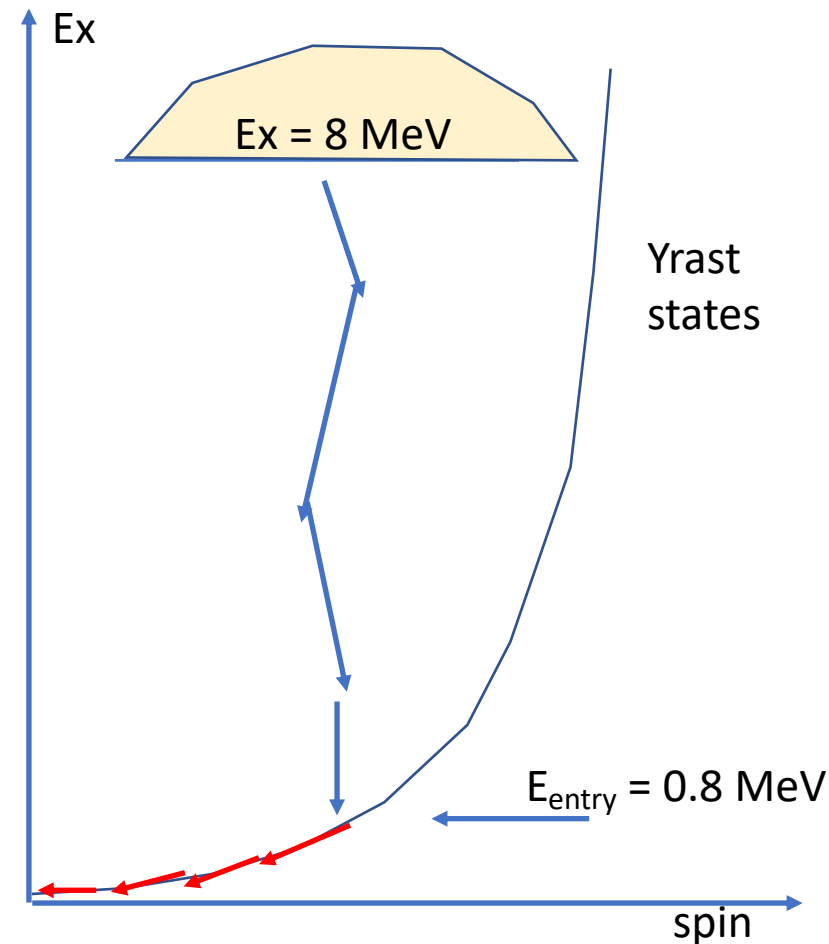
# Correlations between $E_x$ and $E_\gamma$ – not taken into account so far



## 4. Extraction of 1<sup>st</sup> generation $\gamma$ rays using MaMa

# The 1<sup>st</sup>-generation method

- We want to isolate the distribution of the primary  $\gamma$  rays from all possible decay cascades at a given excitation-energy bin (i.e. branching ratios) [M. Guttormsen, T. Ramsøy, and J. Rekstad, NIM A255, 518 (1987)]



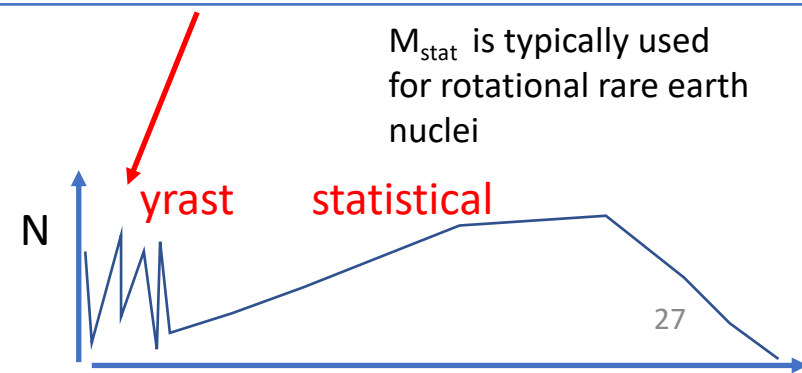
Total gamma-multiplicity:

$$M_{\text{tot}} = M_{\text{stat}} + M_{\text{yrast}} = 4 + 4 = 8$$

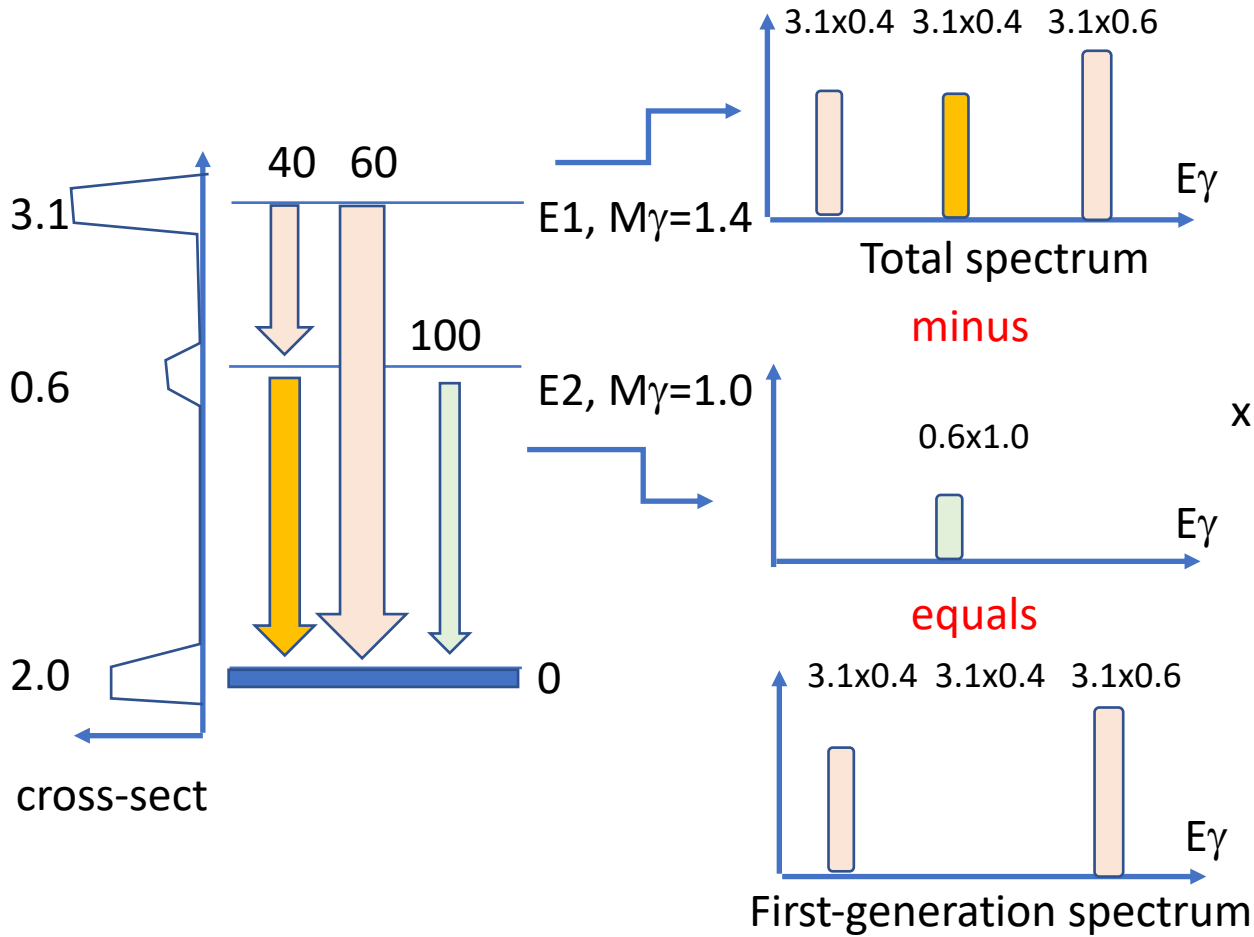
$$M_{\text{tot}} = Ex / \langle E_{\gamma} \rangle = 8 / 1 = 8$$

$$M_{\text{stat}} = (Ex - E_{\text{entry}}) / \langle E_{\gamma} \rangle_{\text{stat}} = (8 - 0.8) / 1.8 = 4$$

We only use  $M_{\text{stat}}$  with an artificial ground state of  $E_{\text{entry}}$  (instead of  $M_{\text{tot}}$ ) if we cannot determine accurately the intensities of the yrast transitions.



# Simple example, 1<sup>st</sup> gen. method



Cross sections:  
 Factor =  $3.1 \times 0.4 / 0.6 \times 1$   
 = 2.07

Multiplicity:  
 Factor =  $(3.1 \times 1.4 - 3.1 \times 1) / 0.6 \times 1$   
 = 2.07

# What about spin population?

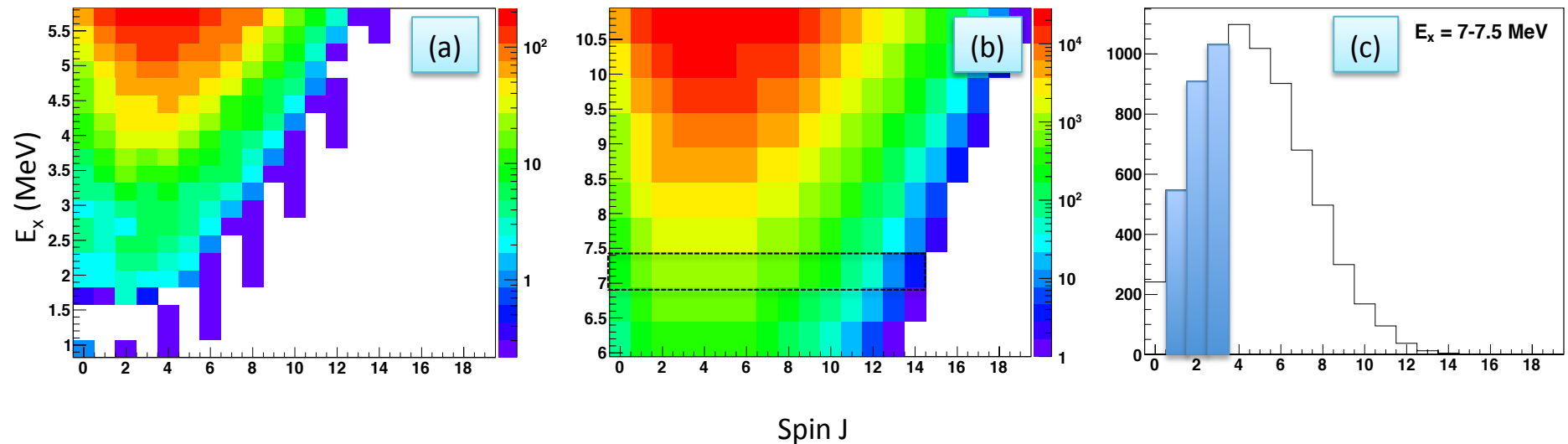
Assumption behind the first-generation method:

[M. Guttormsen, T. Ramsøy, and J. Rekstad, NIM A255, 518 (1987)]

**The present method is based on the assumption that states populated after the first  $\gamma$ -transition have the same decay properties as states populated directly in the particle reaction at that excitation energy. Using the**

# A possible culprit

- Populated spins in the standard Oslo method:  $J \approx 0-10$  both directly and from decay from above (and still we have some trouble!!)
- Beta-decay populates a few spins, much more selective (mainly Gamow-Teller, i.e. same parity as the mother nucleus and spins  $J_{\text{initial}} = J_{\text{mother}} - 1, J_{\text{mother}}, J_{\text{mother}} + 1$ )



Spin distribution for  $^{76}\text{Ge}$  from Goriely et al. [PRC 78, 064307 (2008)]  
(a) for  $E_x \approx 1-5.8$  MeV, (b) for  $E_x \approx 6-10.9$  MeV, and  
(c) for a projection for  $E_x \approx 7-7.5$  MeV. **Blue** histograms:  $J_{\text{initial}}$  of  $^{76}\text{Ge}$ .

# Information saved to figegaout.dat

```

figegaout.dat
Parameters used:
First generation spectra extracted for
excitation energies between 11000.- 0. keV
corresponding to y-channels 220 - 0
Number of spectra= 221
Ax0 = 0.0 Ax1 = 50.0 Ay0 = 0.0 Ay1 = 50.0
AxW0= 0.0 AxW1= 0.0 AyW0= 0.0 AyW1= 0.0
Weighting: 0 Level density parameter a= 7.6 Exponent n=4.2
Normalization= 2 Stat/Tot= 2 Areacorr.= 1
Experimental lower gamma threshold: 300.0 keV (ch= 6)
Upper threshold for statistical gammas: 300.0 keV (ch= 6)
Average energy entry point in g.s. band: 300.(stat) 2000.(tot)
Sliding threshold given by Ex*R, with R = 0.30
Multiplicities: MA = Af/Afg and ME = (Ex-ExEntry)/<Eg>

Iteration number: 1
Y-ch Ex Af Ag Afg Mult Sing Alpha MATot METot MASTa MESTa
18 220 11000. 0. 0. 0. 0.00 0 1.15 0.00 0.00 0.00 0.00
19 219 10950. 48. 26. 22. 2.22 1000 0.89 2.22 2.22 2.22 2.64
20 218 10900. 0. 0. 0. 0.00 1000 0.89 0.00 0.00 0.00 0.00
21 217 10850. 45. 30. 15. 1.74 1000 0.85 2.98 1.74 2.98 2.07
22 216 10800. 51. 29. 22. 2.34 1000 0.99 2.34 2.34 2.34 2.79
23 215 10750. 0. 0. 0. 0.00 1000 0.99 0.00 0.00 0.00 0.00
24 214 10700. 64. 27. 38. 3.71 1000 1.15 1.71 3.71 1.71 4.44
25 213 10650. 50. 27. 22. 2.06 1000 0.85 2.23 2.06 2.23 2.47
26 212 10600. 109. 60. 49. 2.06 1000 0.85 2.22 2.06 2.22 2.47
27 211 10550. 37. 21. 16. 2.30 1000 0.98 2.30 2.30 2.30 2.76
28 210 10500. 140. 76. 65. 2.08 1000 0.85 2.17 2.08 2.17 2.50
29 209 10450. 112. 62. 50. 2.26 1000 0.95 2.26 2.26 2.26 2.71
30 208 10400. 218. 124. 94. 1.96 1000 0.85 2.31 1.96 2.31 2.36
31 207 10350. 3. 1. 2. 8.04 1000 1.15 1.23 8.04 1.23 9.68
32 206 10300. 41. 26. 15. 1.73 1000 0.85 2.76 1.73 2.76 2.08
33 205 10250. 39. 32. 7. 1.33 1000 0.85 5.82 1.33 5.82 1.60
34 204 10200. 142. 84. 58. 2.45 1000 1.13 2.45 2.45 2.45 2.96
35 203 10150. 118. 71. 47. 1.81 1000 0.85 2.52 1.81 2.52 2.18
36 202 10100. 130. 69. 62. 2.12 1000 0.88 2.12 2.12 2.12 2.56
    
```

Average gamma multiplicity – this gets better with larger Ex bins (200 keV/ch instead of 50 keV/ch)

Correction factor to the weighting function. Can only deviate from 1 by 15%. If it is at its limits (0.85 or 1.15), it indicates that the primary spectrum at this Ex is not fully reliable

5. Getting the level density and  $\gamma$ -ray transmission coefficient using *rhosigchi.f*



# The principle

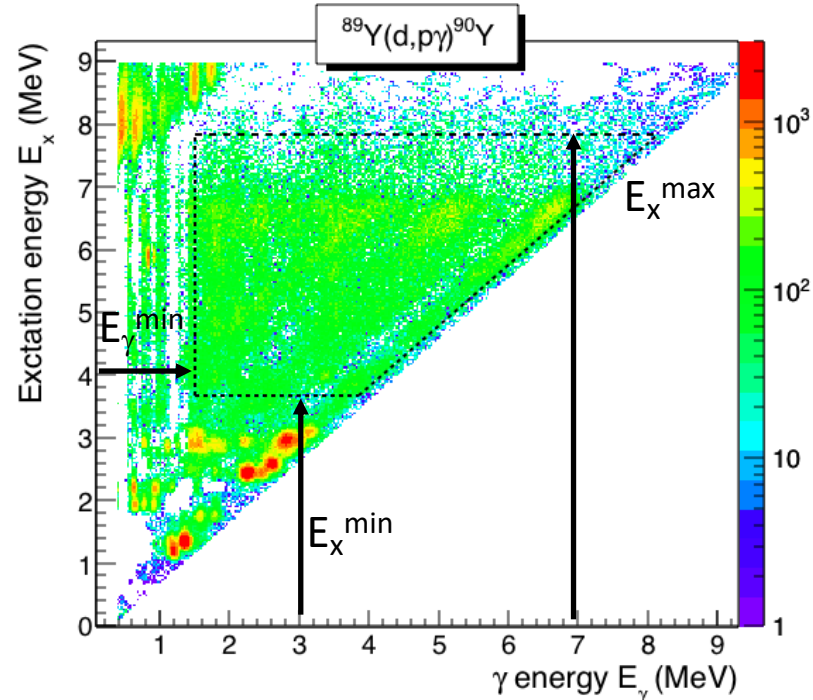
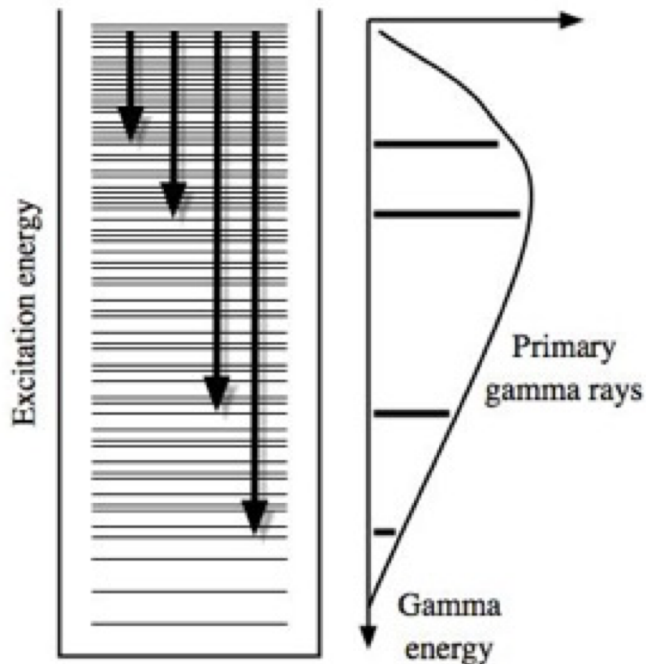
**Ansatz:** primary  $\gamma$  matrix can be factorized into two independent functions (vectors)

[Schiller et al., NIM A 447, 498 (2000)]

Assumes: (i) Compound-state *decay*  
(ii) The Brink hypothesis

$$P(E_i, E_\gamma) \propto \rho(E_i - E_\gamma) \tau(E_\gamma)$$

$$f(E_\gamma) = \tau(E_\gamma) / 2\pi E_\gamma^3$$



# Formalism, *rhosigchi.f*

- Normalize  $P(E_i, E_\gamma)$  so that  $\sum_{E_\gamma=E_\gamma^{\min}}^{E_i} P(E_i, E_\gamma) = 1$ .
- Theoretical estimate of experimental primary  $\gamma$  matrix:

$$P_{th} = \frac{\tau(E_\gamma)\rho(E_i - E_\gamma)}{\sum_{E_\gamma=E_\gamma^{\min}}^{E_i} \tau(E_\gamma)\rho(E_i - E_\gamma)}$$

- First trial function:

$$\rho^{(0)} = 1,$$

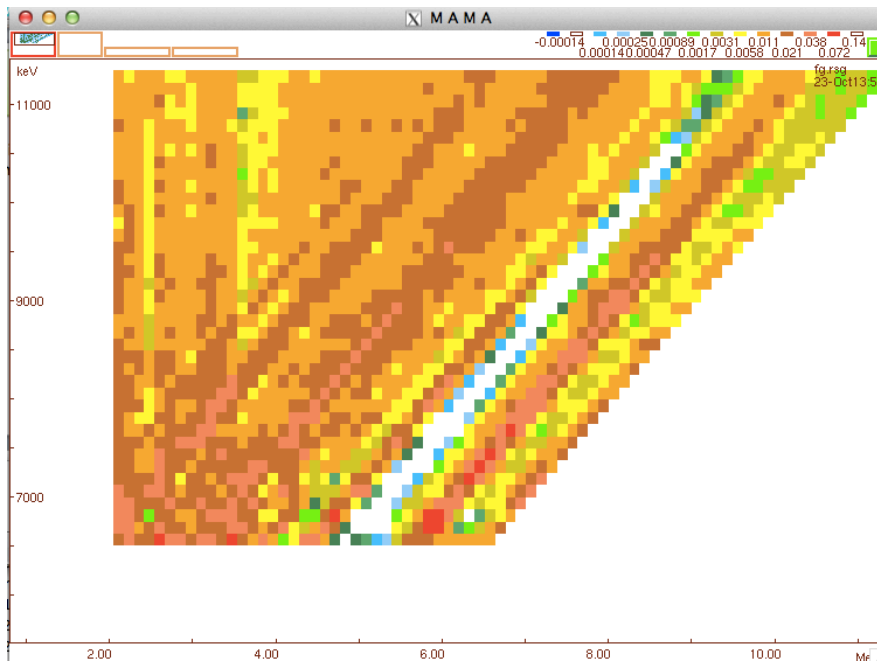
$$P(E_i, E_\gamma) = \frac{\tau^{(0)}(E_\gamma)}{\sum_{E_\gamma=E_\gamma^{\min}}^{E_i} \tau^{(0)}(E_\gamma)}$$

Note: there is no *a priori* assumption that the level density has a Fermi gas or constant-temperature shape!

# Formalism, *rhosigchi.f*

- Higher-order estimates through a least  $\chi^2$ - minimization:

$$\chi^2 = \frac{1}{N_{free}} \sum_{E_i=E_{min}}^{E_{max}} \sum_{E_\gamma=E_\gamma^{min}}^{E_i} \left[ \frac{P_{th}(E_i, E_\gamma) - P(E_i, E_\gamma)}{\Delta P(E_i, E_\gamma)} \right]^2$$



$^{56}\text{Fe}$ , fg.rsg

[see also Larsen et al., JPhysG (2017)]

Each vector element in  $\rho$  and  $\tau$  is treated as a free parameter

$^{56}\text{Fe}(p,p')$  example:

Data points (“pixels”): 2052

Free parameters: 184

$$N_{free} \ll N_{data}$$

Typically  $\approx 10$ - $20$  iterations, but converges after  $\sim 4$ - $5$  iterations

# Comparison, input/output

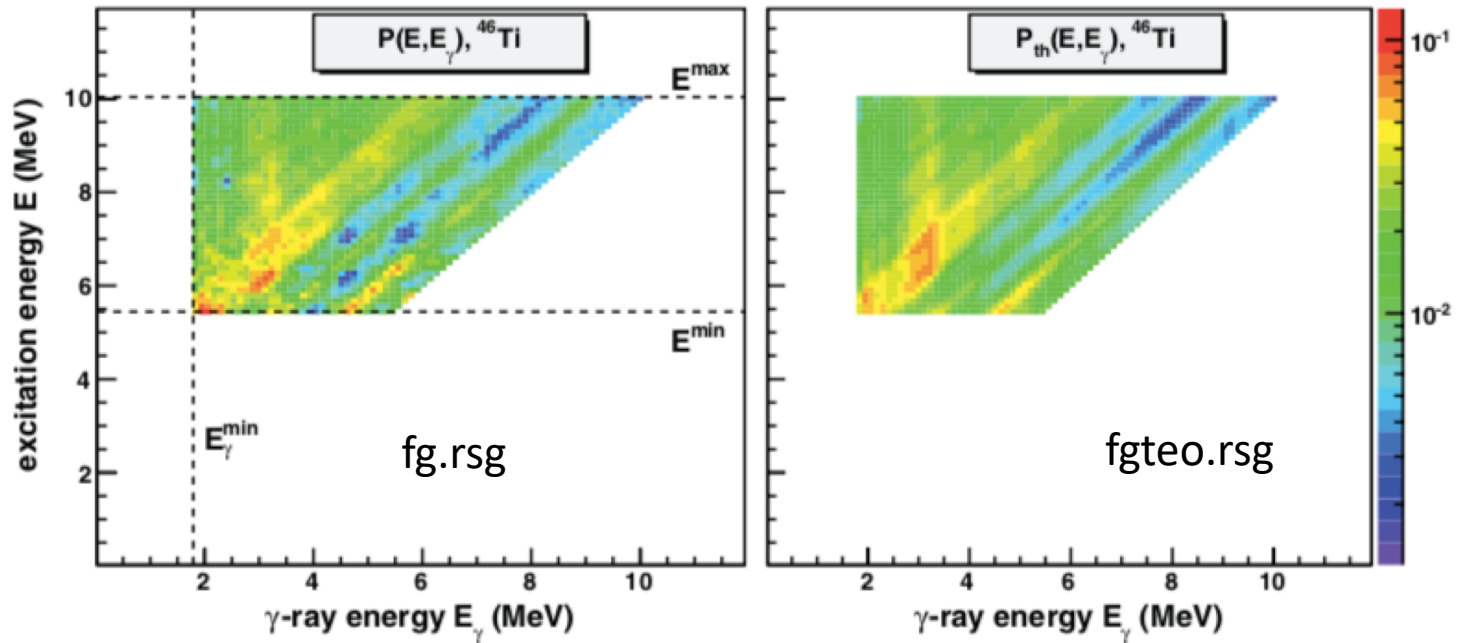
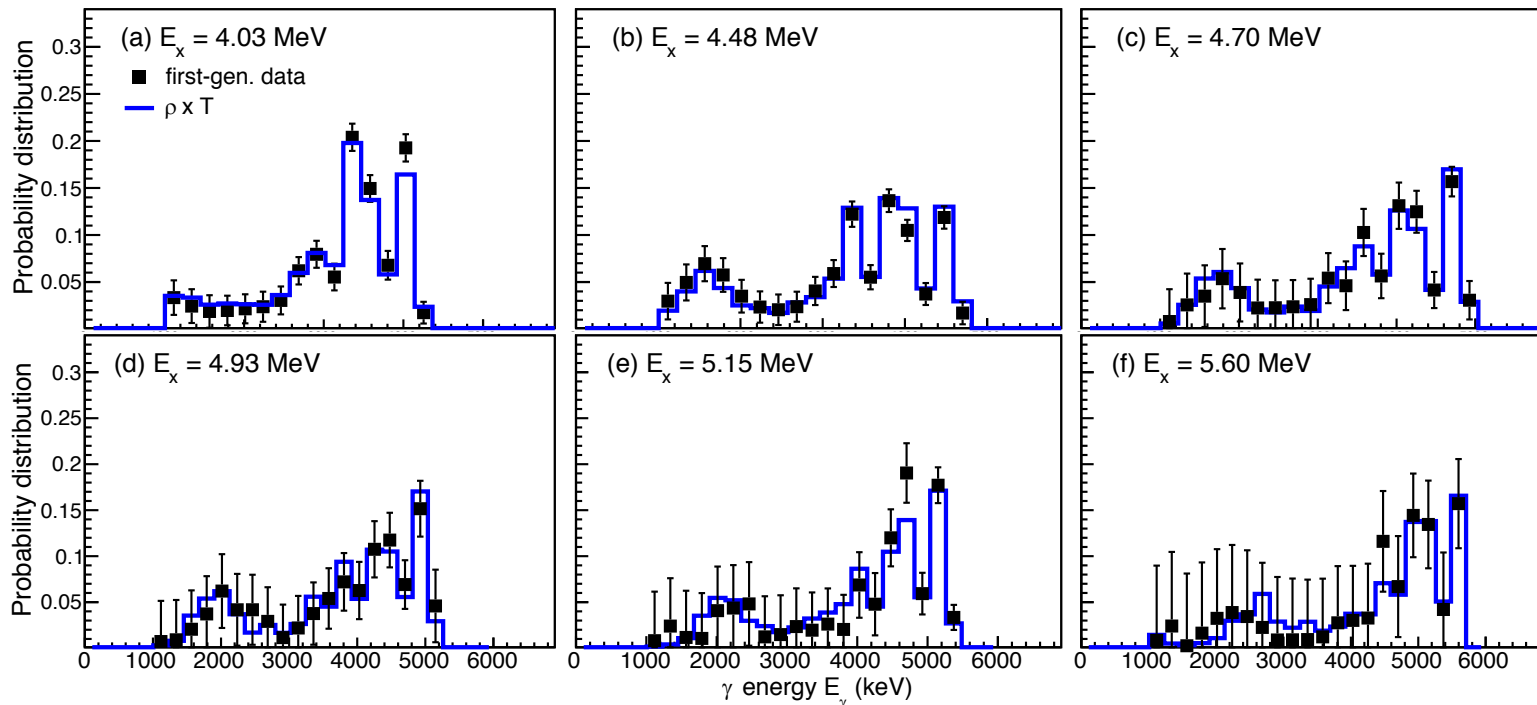


FIG. 5. (Color online) Experimental first-generation matrix  $P(E, E_\gamma)$  (left) and the calculated  $P_{\text{th}}(E, E_\gamma)$  (right) of  ${}^{46}\text{Ti}$  from the iteration procedure of Schiller *et al.* [25]. The dashed lines show the limits set in the experimental first-generation matrix for the fitting procedure. The data are taken from the experiment presented in Ref. [15].

# Does it work? $^{76}\text{Ge}$ example

- Extracted level density and  $\gamma$ -ray trans.coeff. from the whole region within  $E_x^{\min}$ ,  $E_x^{\max}$ ,  $E_\gamma^{\min}$  tested against primary  $\gamma$  spectra from individual bins (see Root script `does_it_work_51Ti_NSCL.cpp`)



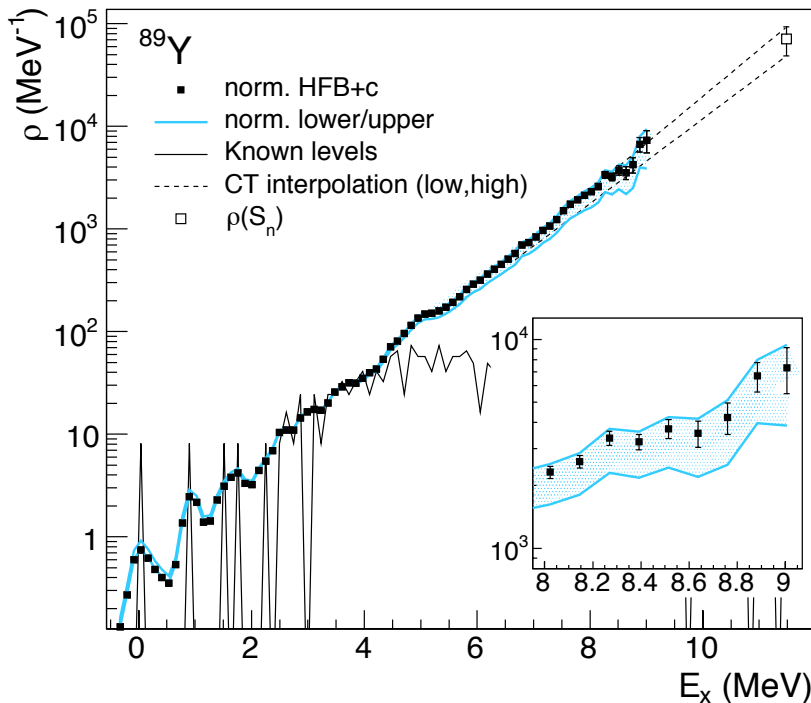
6. Normalization of the level density (and the slope of the  $\gamma$ -ray transmission coefficient) using *counting.c*

# For nuclei at/near stability

Low  $E_x$ : known, discrete levels

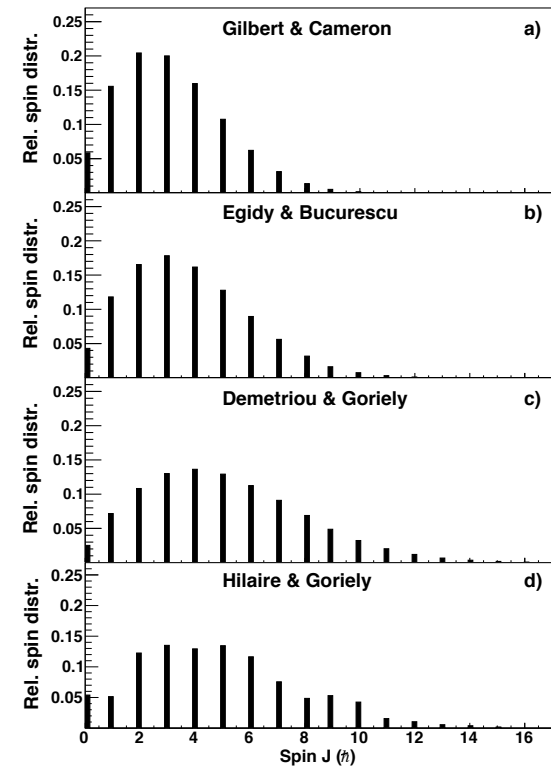
High  $E_x$ : calculate total level density from the neutron-resonance spacing  $D_0$

From Ericson (1960): 
$$g(E_x, J) \cong \frac{2J+1}{2\sigma^2} \exp[-(J+1/2)^2 / 2\sigma^2]$$



[From A.C. Larsen et al., PRC **93**, 045810 (2016)]

**CHALLENGE:**  
usually no  
experimental  
data on the  
spin  
distribution at  
high  $E_x$ !

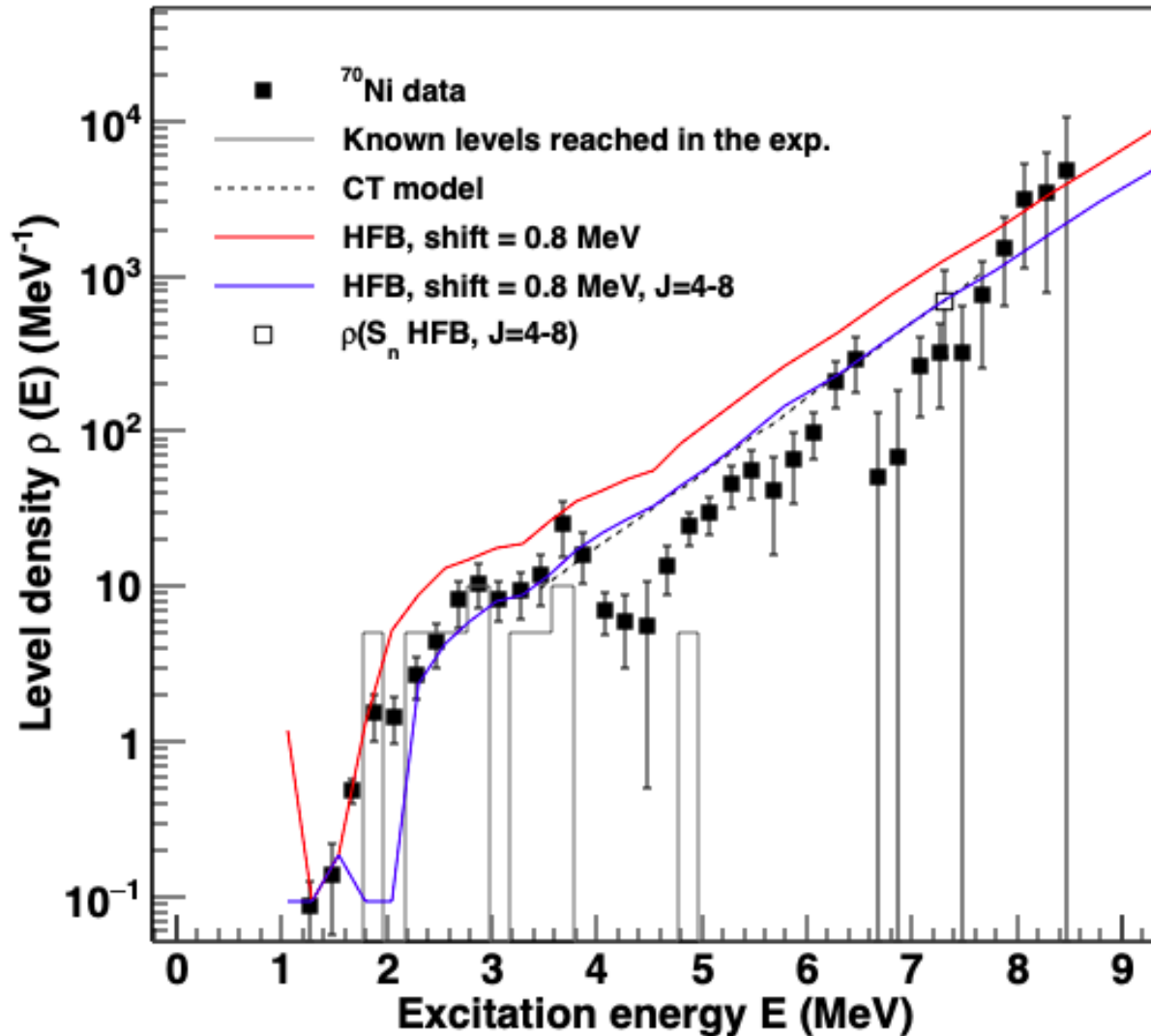


# Level-density normalization, $^{70}\text{Ni}$

- Discrete levels for both parities within the spin range  $J=4-8$  ( $J_{\text{initial}} = 5^-, 6^-, 7^-$ , assuming one dipole transition to reach both parities and spins 4-8)
- Using Goriely's HFB tables (calculations from 2008), adding together both parities for spins  $J=4-8$ , and shift them to match the discrete levels
- This gives the level density for this *restricted* spin range => gives the correct slope for the  $\gamma$ -transmission coefficient
- Note that the *total* level density should be used in the TALYS calculations (TALYS figures out the spins reached in n-capture, and needs the total level density for all spins as input)

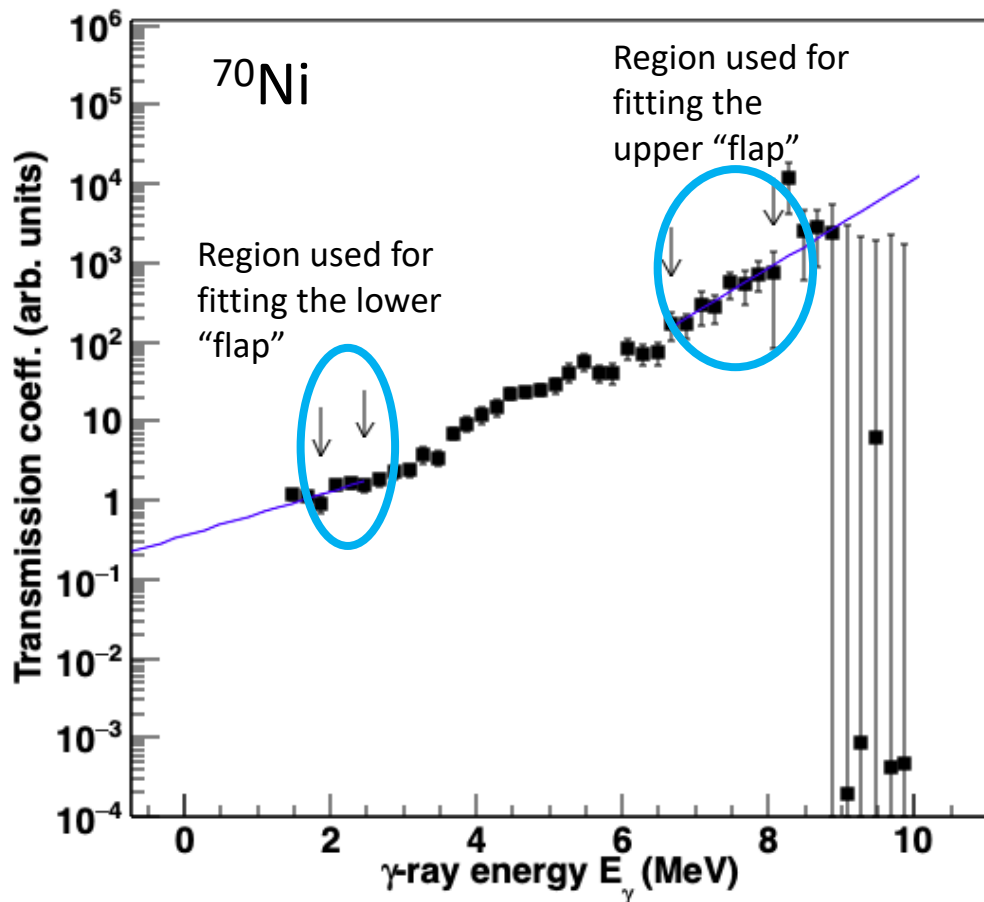


# Level-density normalization, $^{70}\text{Ni}$



*counting.c* also normalizes the slope of the  $\gamma$ -ray transmission coefficient

- ... and makes “flaps” [credit: Stephanie Lyons 😊] on “both sides” of the experimental data



**Note:**

This figure was made with an older version of *counting.c*. The current one on Github gives a steeper decline on the lower “flap”, to get an approx. exponential low-energy component in the  $\gamma$  strength function

7. Normalization of the  $\gamma$ -ray transmission coefficient and getting the dipole  $\gamma$ -ray strength function using *normalization.c*

# For nuclei near/at stability

We only got the slope from the level density normalization, and need to determine the scaling factor B by normalizing to the average, total radiative width measured for s-wave neutron resonances:

$$\begin{aligned} & \langle \Gamma_\gamma(S_n, I_t \pm 1/2, \pi_t) \rangle \\ &= \frac{B}{4\pi\rho(S_n, I_t \pm 1/2, \pi_t)} \int_{E_\gamma=0}^{S_n} dE_\gamma \mathcal{T}(E_\gamma) \\ & \quad \times \rho(S_n - E_\gamma) \sum_{J=-1}^1 g(S_n - E_\gamma, I_t \pm 1/2 + J), \end{aligned}$$

We assume (and have measured!) dominance of dipole radiation, so that

$$B \mathcal{T}(E_\gamma) = B \sum_{XL} \mathcal{T}_{XL}(E_\gamma) \approx B[\mathcal{T}_{E1}(E_\gamma) + \mathcal{T}_{M1}(E_\gamma)].$$

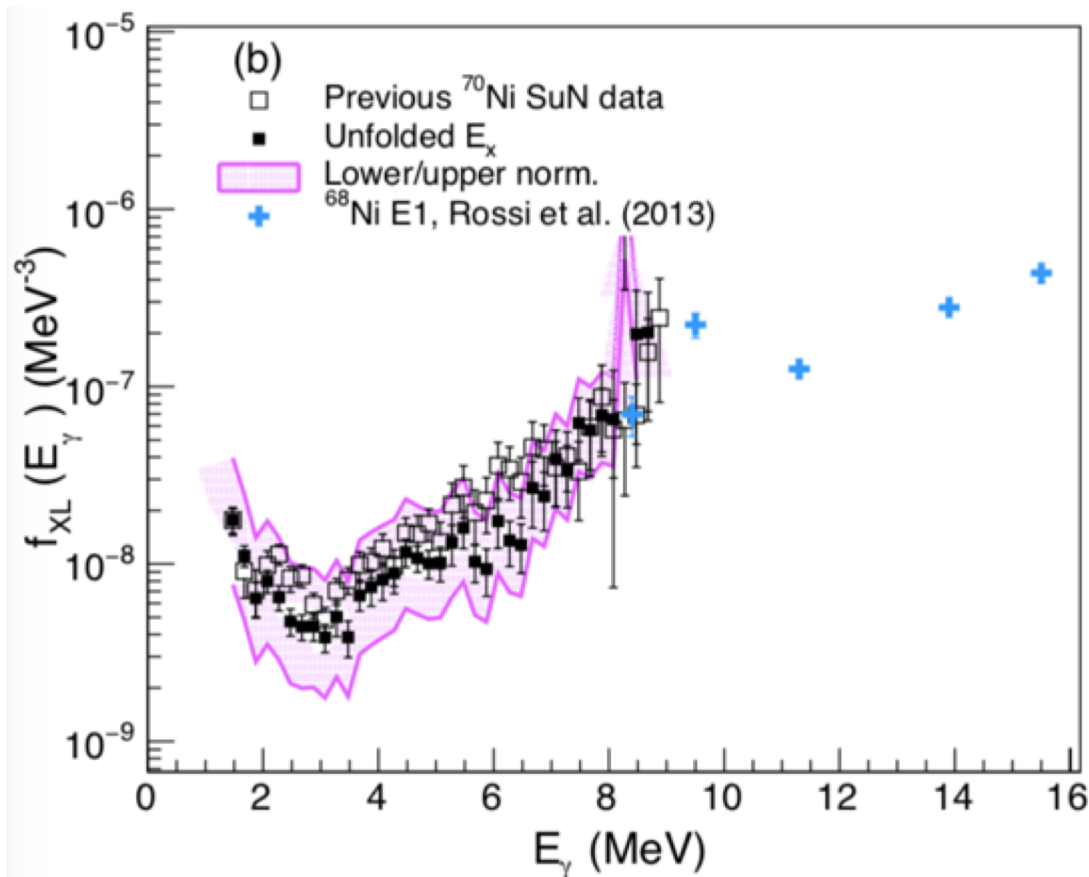
and 
$$f(E_\gamma) = \frac{1}{2\pi E_\gamma^3} B \mathcal{T}(E_\gamma)$$

# Normalizing the $\gamma$ strength of $^{70}\text{Ni}$

No neutron-resonance data available (obviously)

⇒ Used Coulomb-dissociation data of  $^{68}\text{Ni}$  measured at GSI (Rossi et al., PRL (2013))

⇒ Scaled the radiative width (used as a free parameter in normalization.c) until best match with the GSI data at  $\approx 8.5$  MeV

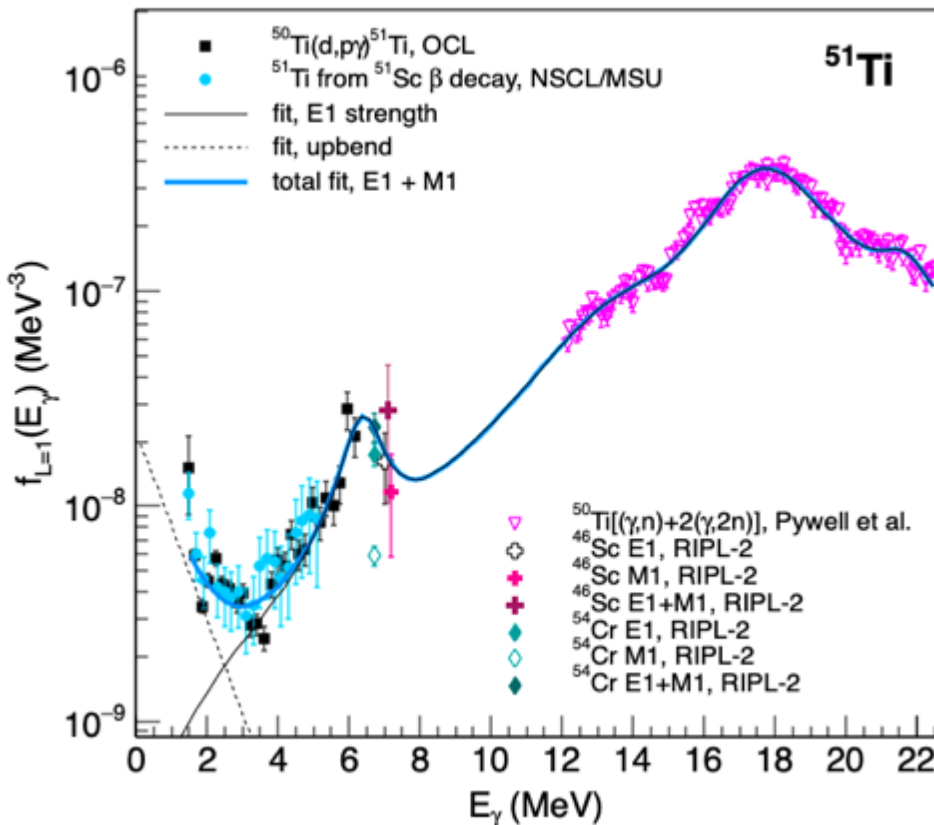


8. Fit of the  $\gamma$ -ray strength function to generate TALYS E1 and M1 tables  
(Root script)

# Using $^{51}\text{Ti}$ as an example

Root script fitexample\_gSF\_51Ti.cpp (input files also provided) makes a fit of

- the  $^{50}\text{Ti}(d,p\gamma)^{51}\text{Ti}$  data from Oslo
- the  $^{51}\text{Sc} \rightarrow ^{51}\text{Ti}$  beta-decay data from NSCL
- $(\gamma,n) + 2(\gamma,2n)$  photonuclear data from Pywell et al
- E1 and M1 strengths from RIPL-2 for  $^{46}\text{Sc}$  and  $^{54}\text{Cr}$



The fit includes:

- Generalized Lorentzian (two components), E1
- One Standard Lorentzian around  $E_\gamma = 14$  MeV, E1
- One Standard Lorentzian around  $E_\gamma = 6$  MeV, E1
- Upbend as an exponential function (assumed to be M1)

# E1 and M1 strength functions printed to files (TALYS format)

```
529 // Print E1 and M1 strengths to file, to use in the TALYS calculations
530 // REMEMBER that the TALYS functions are given in mb/MeV (Goriely's tables)
531 // so we must convert with the factor const double factor = 8.6737E-08; // const. factor in mb(-1) MeV(-2)
532 FILE *E1file, *M1file;
533
534 E1file = fopen("E1_gsf_51Ti_TALYSformat.txt","w");
535 fprintf(E1file," Z= 22 A= 51\n");
536 fprintf(E1file," U[MeV] fE1[mb/MeV]\n");
537 double dummy = 0.1;
538 //cout << " E1 strength:" << endl;
539 for(i=0;i<300;i++){
540     fprintf(E1file,"%9.3f%12.3E\n",dummy,(plot_strength_E1->Eval(dummy))/factor);
541     //cout << dummy << " " << (plot_strength_E1->Eval(dummy)) << endl;
542     dummy += 0.1;
543 }
544 fclose(E1file);
545
546 M1file = fopen("M1_gsf_51Ti_TALYSformat.txt","w");
547 fprintf(M1file," Z= 22 A= 51\n");
548 fprintf(M1file," U[MeV] fM1[mb/MeV]\n");
549 dummy = 0.1;
550 //cout << " M1 strength:" << endl;
551 for(i=0;i<300;i++){
552     fprintf(M1file,"%9.3f%12.3E\n",dummy,(plot_strength_upbend->Eval(dummy))/factor);
553     //cout << dummy << " " << (plot_strength_upbend->Eval(dummy)) << endl;
554     dummy += 0.1;
555 }
556 fclose(M1file);
557
558 cout << " Modeled strength functions for 51Ti written in TALYS format. " << endl;
```



# 9. TALYS input file for calculating $(n,\gamma)$ cross sections and rates

# Using again $^{51}\text{Ti}$ as an example

input\_ng\_51Ti\_example.txt

```
#
# TALYS input file, 50Ti(n,g)51Ti
# Date: Tue 13 March, 2018
# Recommended normalisation
# Additional comments added on Aug 6, 2019
# Cecilie

projectile n
ejectiles g
element ti
mass 50
energy energies.txt
massmodel 2

transeps 1.00E-15
xseps 1.00E-25
popeps 1.00E-25
preequilibrium y
fileresidual y

# CT model with parameters T = 1.227 MeV, E0 = 0.007 MeV
# I have copied the table generated by counting.c found in the
# output file talys_nld_cnt.txt into the Goriely tables contained
# in /talys/structure/density/ground/goriely/Ti.tab for Z= 22 A= 51
ldmodel 4
ptable 22 51 0.0
ctable 22 51 0.0

# Use first 14 discrete levels
Nlevels 22 51 14

# Gamma strength: tables for the E1 and M1 components
# generated with the Root script fitexample_gSF_51Ti.cpp
E1file 22 51 E1_gsf_51Ti_TALYSformat.txt
M1file 22 51 M1_gsf_51Ti_TALYSformat.txt
gnorm 1.
```

# 10. Some Root <-> MaMa scripts

To ease the conversion between Root and MaMa formats:

- `th22mama.C`: taking a Root TH2 and converting to a MaMa matrix
- `th22mama_hist.C`: taking a Root TH1 and converting to a MaMa spectrum
- `mamatoroot_70Niexample.cpp`: takes a MaMa matrix and converts to Root, writes to a Root file

# Remaining challenges 🤔

1. We need a better way to unfold on the Ex axis to include correlations between Ex and Ey
  2. The spin distribution of the initial and final levels is different -> problems with the 1<sup>st</sup> generation method?
  3. How can we reliably normalize the level density and  $\gamma$  strength far away from stability? [Some ideas are on the way...]
  4. ...
  5. ...
- ... and probably many more 😊



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