

# Solids

Daniel Wysocki and Dylan McIntyre  
(aka  $D^2$ )

April 30, 2015

# Introduction

# Solid State

- some loosely bound, outermost valence electrons from each atom detached
- not subject to Coulomb field of a specific nucleus
  - subject to potential of entire crystal lattice

# Primitive Models

## ① electron gas theory of Sommerfeld

- ignores all forces except confining boundaries
- treats wandering electrons as free particles in a box ( $\infty$  cube well)

# Primitive Models

## ① electron gas theory of Sommerfeld

- ignores all forces except confining boundaries
- treats wandering electrons as free particles in a box ( $\infty$  cube well)

## ② Bloch's theory

- periodic potential representing electrical attraction of regularly spaced nuclei
- ignores electron-electron repulsion

# The Free Electron Gas

# Schrödinger Equation

- rectangular solid with dimensions  $\ell_x, \ell_y, \ell_z$

$$V(x, y, z) = \begin{cases} 0, & \text{if } 0 < x < \ell_x, 0 < y < \ell_y, \text{ and } 0 < z < \ell_z \\ \infty, & \text{otherwise} \end{cases}$$

$$-\frac{\hbar^2}{2m} \nabla^2 \psi = E \psi$$

$$\psi(x, y, z) = X(x) Y(y) Z(z)$$

# Schrödinger Equation (cont.)

$$k_x \equiv \frac{\sqrt{2mE_x}}{\hbar}, k_y \equiv \frac{\sqrt{2mE_y}}{\hbar}, k_z \equiv \frac{\sqrt{2mE_z}}{\hbar},$$

General solutions

$$X(x) = A_x \sin(k_x x) + B_x \cos(k_x x); \quad Y(y) = \dots; \quad Z(z) = \dots$$



# Schrödinger Equation (cont.)

Boundary conditions require

$$X(0) = Y(0) = Z(0) = 0$$

so

$$B_x = B_y = B_z = 0$$

and

$$X(\ell_x) = Y(\ell_y) = Z(\ell_z)$$

meaning

$$k_x \ell_x = n_x \pi; \quad k_y \ell_y = n_y \pi; \quad k_z \ell_z = n_z \pi$$

# Schrödinger Equation (cont.)

$$\psi_{n_x n_y n_z} = \sqrt{\frac{8}{\ell_x \ell_y \ell_z}} \sin\left(\frac{n_x \pi}{\ell_x} x\right) \sin\left(\frac{n_y \pi}{\ell_y} y\right) \sin\left(\frac{n_z \pi}{\ell_z} z\right)$$

$$E_{n_x n_y n_z} = \frac{\hbar^2 \pi^2}{2m} \left( \frac{n_x^2}{\ell_x^2} + \frac{n_y^2}{\ell_y^2} + \frac{n_z^2}{\ell_z^2} \right) = \frac{\hbar^2 k^2}{2m}$$

$$k = \|\mathbf{k}\|$$

$$\mathbf{k} = (k_x, k_y, k_z)$$

# $k$ -space

- 3D space with axes  $k_x$ ,  $k_y$ , and  $k_z$
- planes drawn at  $k_x = (\pi/\ell_x), (2\pi/\ell_x), \dots$ ,  $k_y = \dots$ ,  $k_z = \dots$
- each intersection represents a distinct (one-particle) stationary state
- each block occupies a volume in ( $k$ -space) of

$$k_x k_y k_z = \frac{\pi^3}{\ell_x \ell_y \ell_z} = \frac{\pi^3}{V_{\text{object}}}$$

- each block contains two electrons ( $\uparrow\downarrow$ )

$k$ -space

- all free electrons fill 1 octant of sphere in  $k$ -space, centered at origin
  - sphere has radius  $k_F$

$$V_{\text{octant}} = \frac{1}{8} \left( \frac{4}{3} \pi k_F^3 \right) = \frac{Nq}{2} \left( \frac{\pi^3}{V_{\text{object}}} \right)$$

$$k_F = (3\rho\pi^2)^{1/3}$$

$$\rho \equiv \frac{Nq}{V_{\text{object}}}$$

# Fermi surface

- boundary separating occupied and unoccupied states in  $k$ -space
- corresponding energy is Fermi energy,  $E_F$

$$E_F = \frac{\hbar^2}{2m} k_F^2 = \frac{\hbar^2}{2m} (3\rho\pi^2)^{2/3}$$

# Total Energy

- a shell of thickness  $dk$  has volume

$$\frac{1}{8}(4\pi k^2)dk$$

- number of electron states in shell

$$\frac{2[(1/2)\pi k^2 dk]}{\pi^3/V} = \frac{V}{\pi^2} k^2 dk$$

- each carries energy  $\hbar^2 k^2/2m$ , so energy of shell is

$$dE = \frac{\hbar^2 k^2}{2m} \frac{V}{\pi^2} k^2 dk$$

## Total Energy (cont.)

$$\begin{aligned}
 E_{\text{tot}} &= \int_0^{k_F} dE = \frac{\hbar^2 V}{2\pi^2 m} \int_0^{k_f} k^4 dk \\
 &= \frac{\hbar^2 k_F^5 V}{10\pi^2 m} = \frac{\hbar^2 (3\pi^2 Nq)^{5/3}}{10\pi^2 m} V^{-2/3}
 \end{aligned}$$

- total energy is inversely proportional to total volume

# Degeneracy pressure

- if the box expands by  $dV$ , the total energy decreases

$$dE_{\text{tot}} = -\frac{2}{3} \frac{\hbar^2 (3\pi^2 Nq)^{5/3}}{10\pi^2 m} V^{-5/3} dV = -\frac{2}{3} E_{\text{tot}} \frac{dV}{V}$$

- shows up as work done on outside ( $dW = PdV$ ) by quantum pressure  $P$

$$P = \frac{2}{3} \frac{E_{\text{tot}}}{V} = \frac{2}{3} \frac{\hbar^2 k_F^5}{10\pi^2 m} = \frac{(3\pi^2)^{2/3} \hbar^2}{5m} \rho^{5/3}$$

- stabilizing internal pressure, independent of  $e$ - $e$  & thermal repulsion



# Band Structure

# Introduction

- we improve on the free electron model by including forces exerted by nuclei
  - regularly spaced
  - positively charged
  - stationary

# Dirac Comb

- model the potential by a 1D Dirac comb
  - evenly spaced delta function spikes
- periodic in  $x$  at intervals of  $a$

$$V(x + a) = V(x)$$

# Bloch's Theorem

the solutions to the Schrödinger equation

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V(x)\psi = E\psi$$

can be taken to satisfy the condition

$$\psi(x+a) = e^{iKa}\psi(x)$$

for some constant  $K$  (independent of  $x$  but not necessarily  $E$ )

# Displacement Operator

$$Df(x) = f(x + a)$$

- for a periodic potential,  $D$  commutes with the Hamiltonian

$$[D, H] = 0$$

- free to choose eigenfunctions of  $H$  that are eigenfunctions of  $D$

$$D\psi = \lambda\psi \implies \psi(x + a) = \lambda\psi(x)$$

# Wave Function

- considering the loop model,  $0 < x < a$  yields the wave function:

$$\psi(x) = A \sin(kx) + B \cos(kx), \quad (0 < x < a)$$

- Bloch's theorem allows us to write for final cell:

$$\psi(x) = e^{-iKa} [A \sin k(x + a) + B \cos k(x + a)], \quad (-a < x < 0)$$

# Wave Function (cont.)

- from boundary conditions:

$$B = e^{-iKa} [A \sin(ka) + B \cos(ka)],$$

whose derivative is discontinuous proportional to magnitude of  $\delta$  fn

- solving for  $A \sin(ka)$  gives

$$A \sin(ka) = [e^{iKa} - \cos(ka)] B$$

# Wave Function (cont.)

- derivative of the wave function is given by

$$\Delta\left(\frac{d\psi}{dx}\right) = -\frac{2m\alpha}{\hbar^2}\psi(0)$$

- plugging into  $B$  given by boundary conditions:

$$kA - e^{-iKa} k[A \cos(ka) - B \sin(ka)] = \frac{2m\alpha}{\hbar^2} B$$



# Wave Function (cont.)

- plugging into  $A \sin(ka)$  gives

$$\left[ e^{iKa} - \cos(ka) \right] \left[ 1 - e^{-iKa} \cos(ka) \right] + e^{-iKa} \sin^2(ka) = \frac{2m\alpha}{\hbar k} \sin(ka)$$

- simplifies to

$$\cos(Ka) = \cos(ka) + \frac{m\alpha}{\hbar^2 k} \sin(ka)$$

# Discrete Energy States

$$\cos(Ka) = \cos(ka) + \frac{m\alpha}{\hbar^2 k} \sin(ka)$$

- allows us to find the possible values of  $k$  and allowed energies
- we define this dimensionless notation:

$$z \equiv ka, \quad \beta \equiv \frac{m\alpha a}{\hbar^2}$$

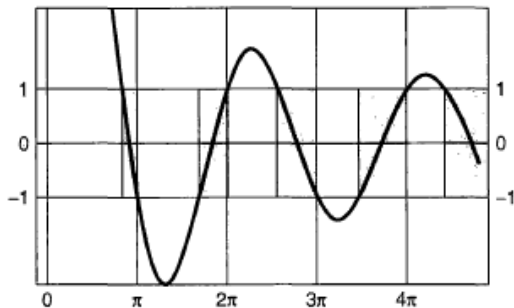
- and rewrite the top as

$$f(z) \equiv \cos(z) + \beta \frac{\sin(z)}{z}$$

- $\beta$  is the “strength” of the delta function

# Forbidden Energies

- plot  $f(z)$  using  $\beta = 10$

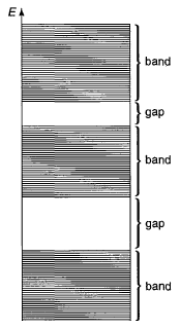


# Forbidden Energy Considerations

- $f(z)$  exists outside the range  $(-1, +1)$
- unsolvable in these regions
  - $\cos(Ka)$  is confined between  $(-1, +1)$
- gives rise to forbidden energy regions
  - energy gaps
- gaps are separated by allowed regions
  - bands
- practically any energy allowed within given band ( $Ka = 2\pi n/N$ )

# Energy Bands

- draw  $N$  horizontal lines on previous graph at values of  $\cos(2\pi n/N)$
- intersection points represent energy levels within each band



# Multielectron Systems

- we have only considered one electron in the potential
- in reality there are  $Nq$  electrons
- Pauli exclusion principle dictates at most 2 electrons may occupy one state
- with  $N$  states, we have the following possibilities
  - $q = 1 \implies$  first band half filled
  - $q = 2 \implies$  first band completely filled
  - $q = 3 \implies$  second band half filled
  - etc.

# System Classifications

- entirely filled band requires relatively large energy to excite electron
  - electrical **insulators**
- partly filled band requires very small energy to excite electron
  - electrical **conductors**

# Doped Insulators

- we dope an insulator with some atoms of a different  $q$
- two cases arise:
  - 1  $q_2 > q_1 \implies$  obtain extra electrons in next higher band
  - 2  $q_2 < q_1 \implies$  create holes in previously filled band
- in both cases, small electric currents may flow
  - **semiconductors**



# Thank You