

Dirac Notation

Dylan J. McIntyre and Nicholas C. Jira

March 5, 2015

Overview

- 1 Vectors
- 2 Matrix Algebra
- 3 Dirac Notation
- 4 Bibliography

Vectors

- A vector A can be described by specifying its components with respect to a set of axes

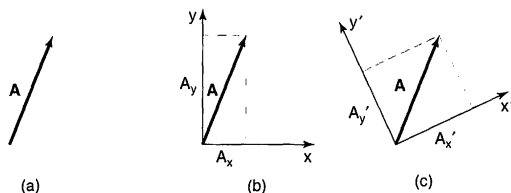


Figure 1: (a) vector, (b) x,y axis, (c) x',y' axis [Griffiths, 2005]

- $A : A_x = \hat{i} \cdot A, A_y = \hat{j} \cdot A$ or $A : A'_x = \hat{i}' \cdot A, A'_y = \hat{j}' \cdot A$
- Same vector A , different choice of basis
- The vector lives in a space independent of coordinates

The Wave Function

- The state of a system is also represented by a vector, $|S\rangle$
- This vector lives in Hilbert space
- Represented with respect to different bases

$$\Psi(x, t) = \langle x|S(t)\rangle \quad (1)$$

- $|x\rangle$ is the position eigenfunction with eigenvalues x

$$\Phi(p, t) = \langle p|S(t)\rangle \quad (2)$$

- $|p\rangle$ is the momentum eigenfunction with eigenvalues p
- Same vector S expanded over different bases of eigenfunctions

The Wave Function

- S may also be expanded in the energy eigenfunction basis (assuming discrete spectrum for simplicity)

$$c_n(t) = \langle n | S(t) \rangle \quad (3)$$

- In this case, $|n\rangle$ is the n th eigenfunction of \hat{H}
- In all cases, it is the same state S
- Ψ , Φ and c_n contain the same exact information and are simply different ways of describing the same vector:

$$\begin{aligned} \Psi(x, t) &= \int \Psi(y, t) \delta(x - y) dy = \int \Phi(p, t) \frac{1}{\sqrt{2\pi\hbar}} e^{ipx/\hbar} dp \\ &= \sum c_n e^{-iE_n t/\hbar} \psi_n(x) \end{aligned} \quad (4)$$

Linear Transformations

- Operators transform one vector into another:

$$|\beta\rangle = \hat{Q}|\alpha\rangle \quad (5)$$

- A vector is represented by it's components with respect to a basis $|e_n\rangle$ by:

$$|\alpha\rangle = \sum_n a_n |e_n\rangle \text{ with } a_n = \langle e_n | \alpha \rangle \quad (6)$$

- Operators are expressed by their matrix elements:

$$\langle e_m | \hat{Q} | e_n \rangle \equiv Q_{mn} \quad (7)$$

- With this notation, 5 becomes:

$$\sum_n b_n |e_n\rangle = \sum_n a_n \hat{Q} |e_n\rangle \quad (8)$$

Linear Transformations

- 8 can be rewritten by using the inner product with $|e_m\rangle$

$$\sum_n b_n \langle e_m | e_n \rangle = \sum_n a_n \langle e_m | \hat{Q} | e_n \rangle \quad (9)$$

- Which gives

$$b_m = \sum_n Q_{mn} a_n \quad (10)$$

- The elements of the matrix describe how the components transform
- We might encounter systems with a finite number of linearly independent states
- Then, we have an N -dimensional vector space with operators in the form of $(N \times N)$ matrices

Dirac - bra and ket

- Dirac split the inner product bracket notation into two pieces: the bra, $\langle\alpha|$ and ket, $|\beta\rangle$
- The ket is a vector, but the bra is a linear function of vectors
- When the bra acts on the ket, it results in a complex number (the inner product)
- Similar to an operator, except an operator acting on a vector produces another vector, whereas the bra acting on a vector produces a number
- We can think of the bra as an instruction telling us to integrate:

$$\langle f| = \int f^*[\dots]dx$$

- Where the $[\dots]$ represents the function encountered in the ket

Dirac - bra and ket

- A vector in a finite-dimensional vector space can be expressed in columns as

$$|\alpha\rangle = \begin{pmatrix} a_1 \\ a_2 \\ \vdots \\ a_n \end{pmatrix} \quad (11)$$

- The bra corresponding to this vector is:

$$\langle\alpha| = (a_1^* a_2^* \cdots a_n^*) \quad (12)$$

- The collection of all bras describes a vector space known as *dual space*

Projection Operator

- If $|\alpha\rangle$ is normalized, we can define the *projection operator* to be

$$\hat{P} \equiv |\alpha\rangle\langle\alpha| \quad (13)$$

- which picks out the portion of any vector along $|\alpha\rangle$:

$$\hat{P}|\beta\rangle = \langle\alpha|\beta\rangle|\alpha\rangle \quad (14)$$

- Lets consider a discrete orthonormal basis $|e_n\rangle$ such that

$$\langle e_m|e_n\rangle = \delta_{mn} \quad (15)$$

- Then the identity operator is

$$\sum_n |e_n\rangle\langle e_n| = 1 \quad (16)$$

Dirac Orthonormalized Continuous Basis

- If this operator acts on a vector $|\alpha\rangle$ we obtain

$$\sum_n |e_n\rangle \langle e_n | \alpha \rangle = |\alpha\rangle \quad (17)$$

- The same can be said for a Dirac orthonormalized continuous basis $|e_z\rangle$:

$$\langle e_z | e_{z'} \rangle = \delta(z - z') \quad (18)$$

- Then

$$\int |e_z\rangle \langle e_z| dz = 1 \quad (19)$$

References



David J. Griffiths (2005)

Introduction to Quantum Mechanics, 2nd Edition

Pearson Prentice Hall