

1.021, 3.021, 10.333, 22.00 Introduction to Modeling and Simulation
Spring 2018

Basic molecular dynamics

Lecture 2

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Massachusetts Institute of Technology

Content overview

I. Fundamentals of particle methods

1. Atoms, molecules, chemistry
2. Statistical mechanics
3. Molecular dynamics, Monte Carlo
4. Visualization and data analysis
5. Mechanical properties – application: how things fail (and how to prevent it)
6. Multi-scale modeling paradigm
7. Biological systems (simulation in biophysics) – how proteins work and how to model them

Lectures 1-12
February/March

II. Advanced topics in particle methods

1. Quantum Weirdness: The Theory of Quantum Mechanics
2. The Many-Body Problem: From Many-Body to Single-Particle
3. Quantum modeling of materials
4. From Atoms to Solids
5. Basic properties of materials
6. Advanced properties of materials
7. Materials Informatics

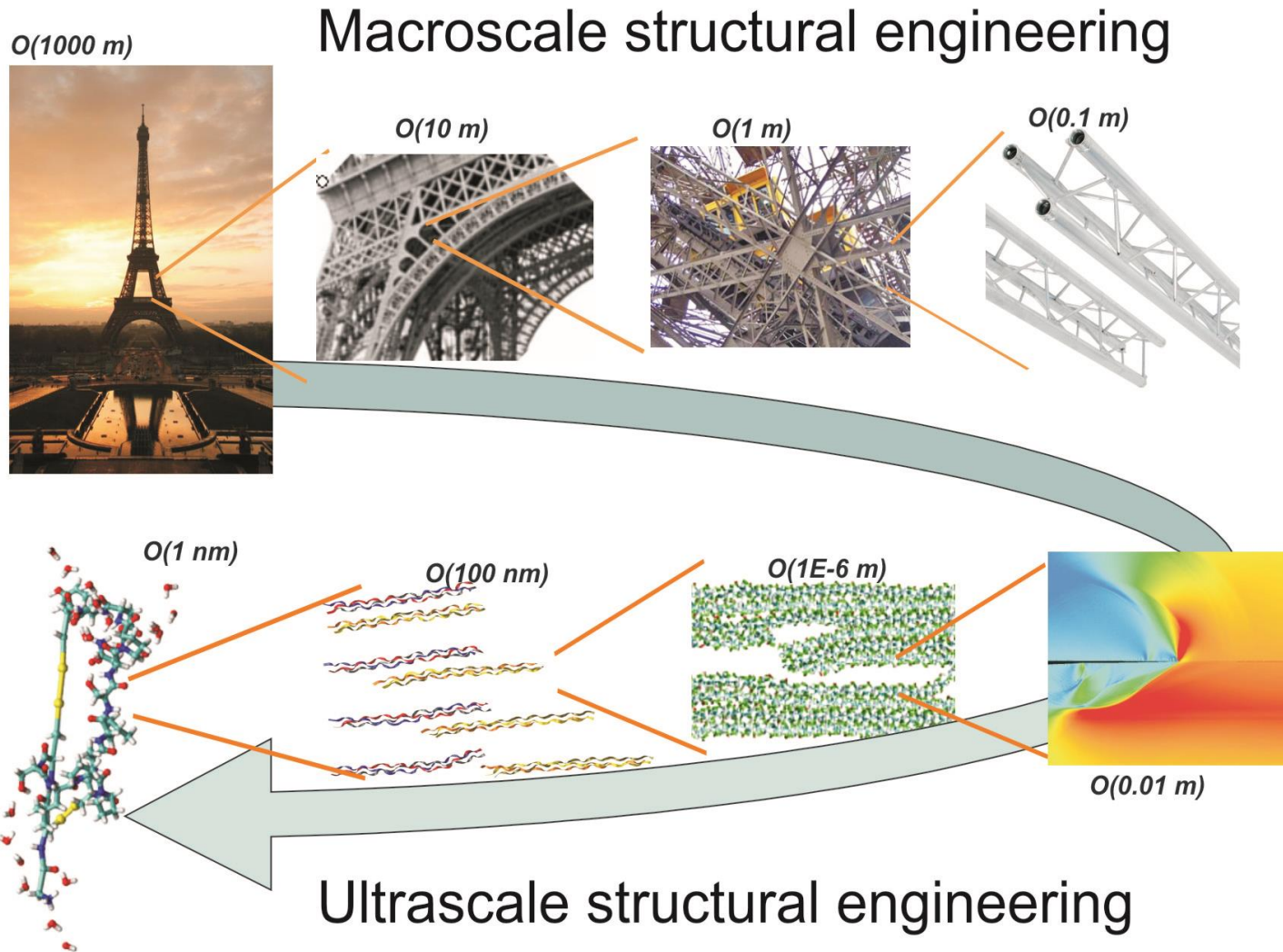
Lectures 13-24
March/April/May

Goals of part I (particle methods)

You will be able to ...

- **Carry out atomistic simulations** of various processes (diffusion, deformation/stretching, materials failure)
Nanowires, bulk metals, proteins, silicon crystals, etc.
- **Analyze** atomistic simulations (make sense of the numbers)
- **Visualize** atomistic/molecular data (bring data to life)
- Understand how to **link atomistic simulation results with continuum models** within a multi-scale scheme

Multi-scale view of materials



Lecture 2: Basic molecular dynamics

Outline:

1. Example application: Stiffness of materials
2. Introduction: Bottom-up vs. top-down
3. Case study: Diffusion
 - 3.1 Continuum model
 - 3.2 Atomistic model (to be continued)

Goals of today's lecture:

- Through case study of diffusion, illustrate the concepts of a continuum model and an atomistic model
- Develop appreciation for distinction of continuum and atomistic approach
- Develop equations/models for diffusion problem from both perspectives

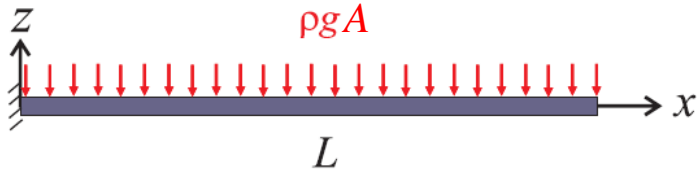
Next:

- Develop atomistic simulation approach (e.g. algorithm, pseudocode, etc.) and apply to describe diffusion (calculate diffusivity)
- Historical perspective on computer simulation with MD, examples from literature

1. Example application: Stiffness of materials (Young's modulus)

Objective: *Illustrate the significance of multiple scales for material behavior and introduce multi-scale modeling paradigm*

Beam deformation problem – continuum model



Question: Displacement field

Governing equation (PDE)

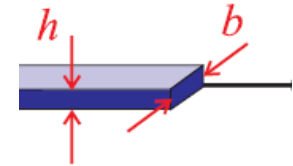
$$-EI_{zz} \frac{\partial^4 u_{z,0}}{\partial x^4} + q_z = 0$$



Integration & BCs

Geometry

$$I_{zz} = \frac{bh^3}{12}$$



BC - load:

$$\rho g A$$



$$u_z(x) = -\frac{\rho g A}{24EI_{zz}} x^4$$

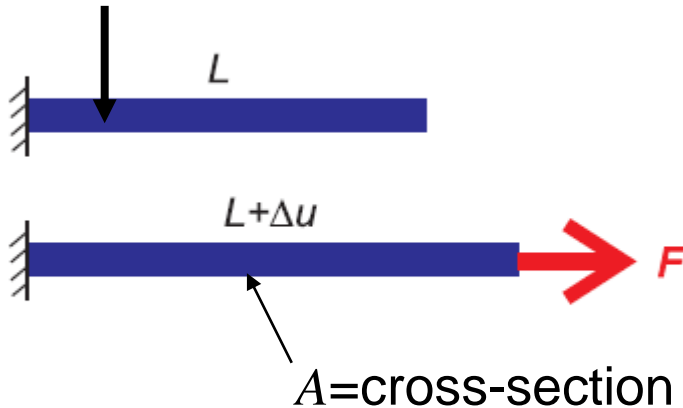
E = unknown parameter

E is parameter called “Young’s modulus” that relates how force and deformation are related (captures properties of material)

How to determine Young's modulus E ?

Measurement (laboratory):

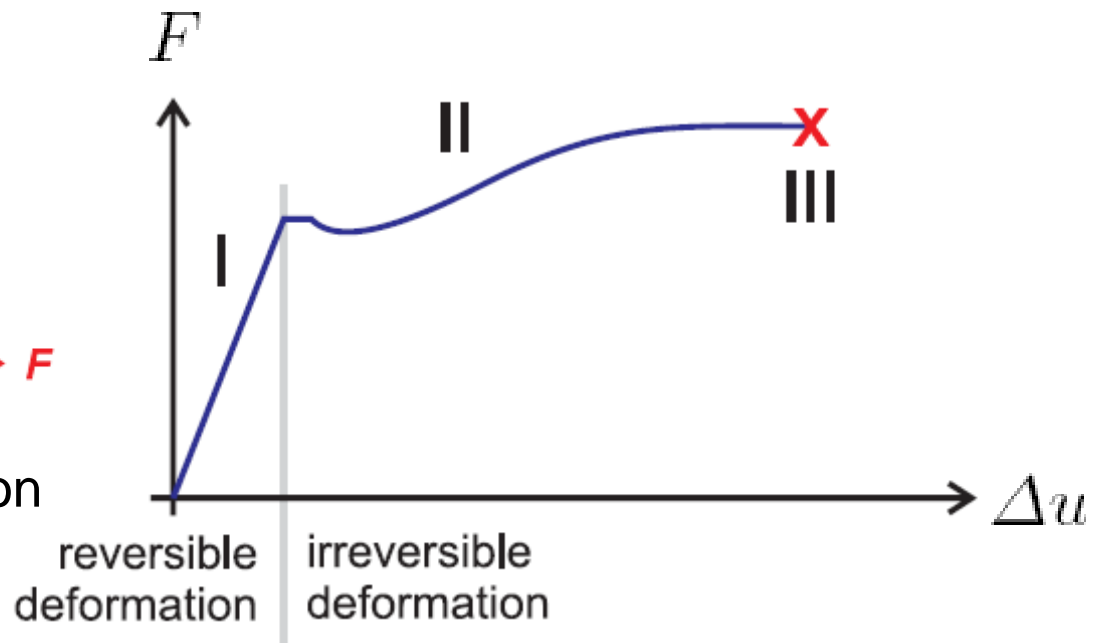
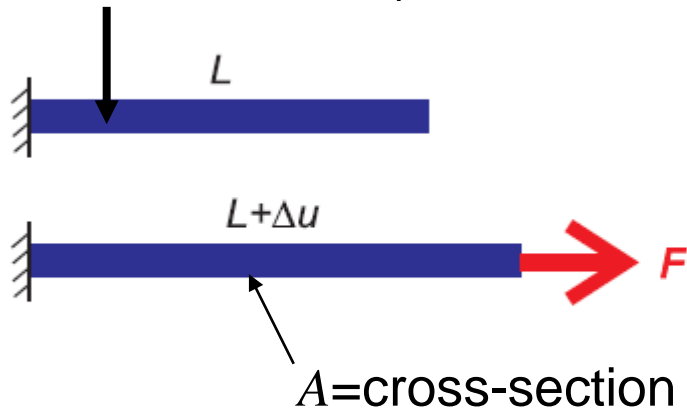
Rod/beam (e.g. plastic, metal, nanowire)



How to determine Young's modulus E ?

Measurement (laboratory):

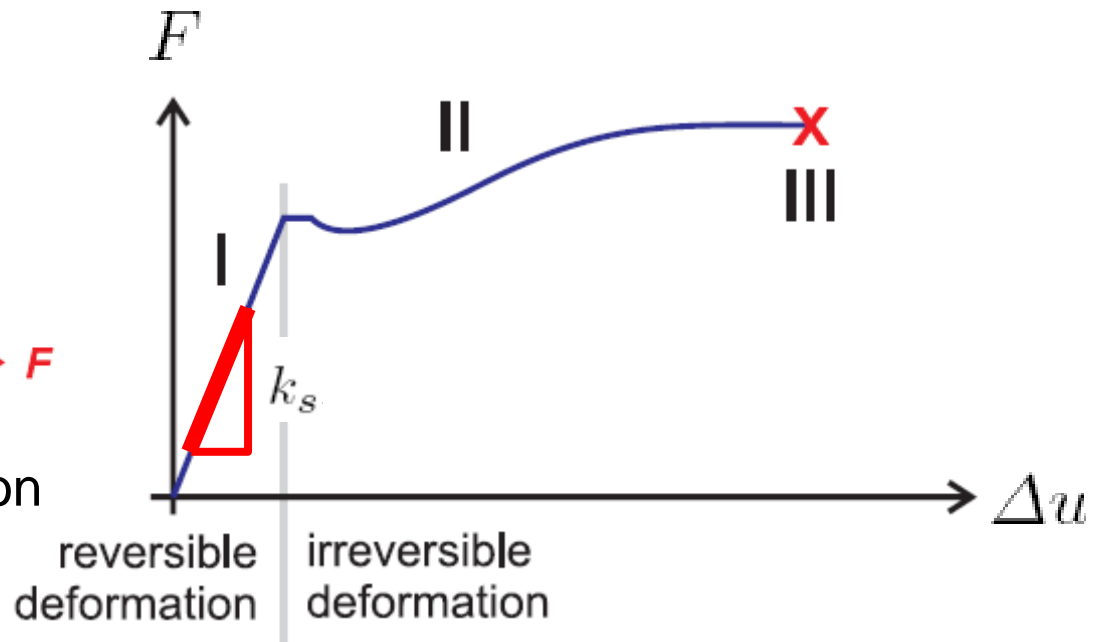
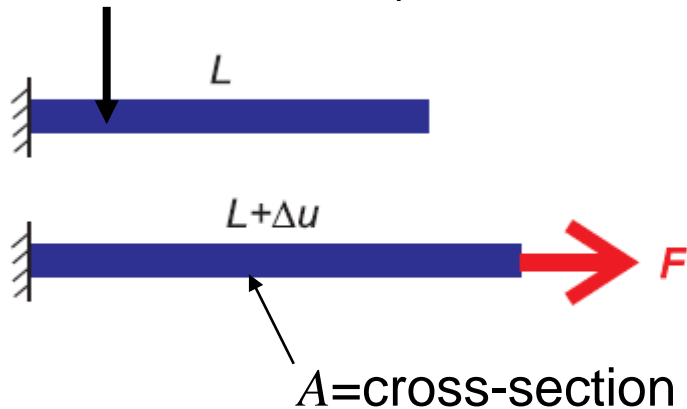
Rod/beam (e.g. plastic, metal, nanowire)



How to determine Young's modulus E ?

Measurement (laboratory):

Rod/beam (e.g. plastic, metal, nanowire)



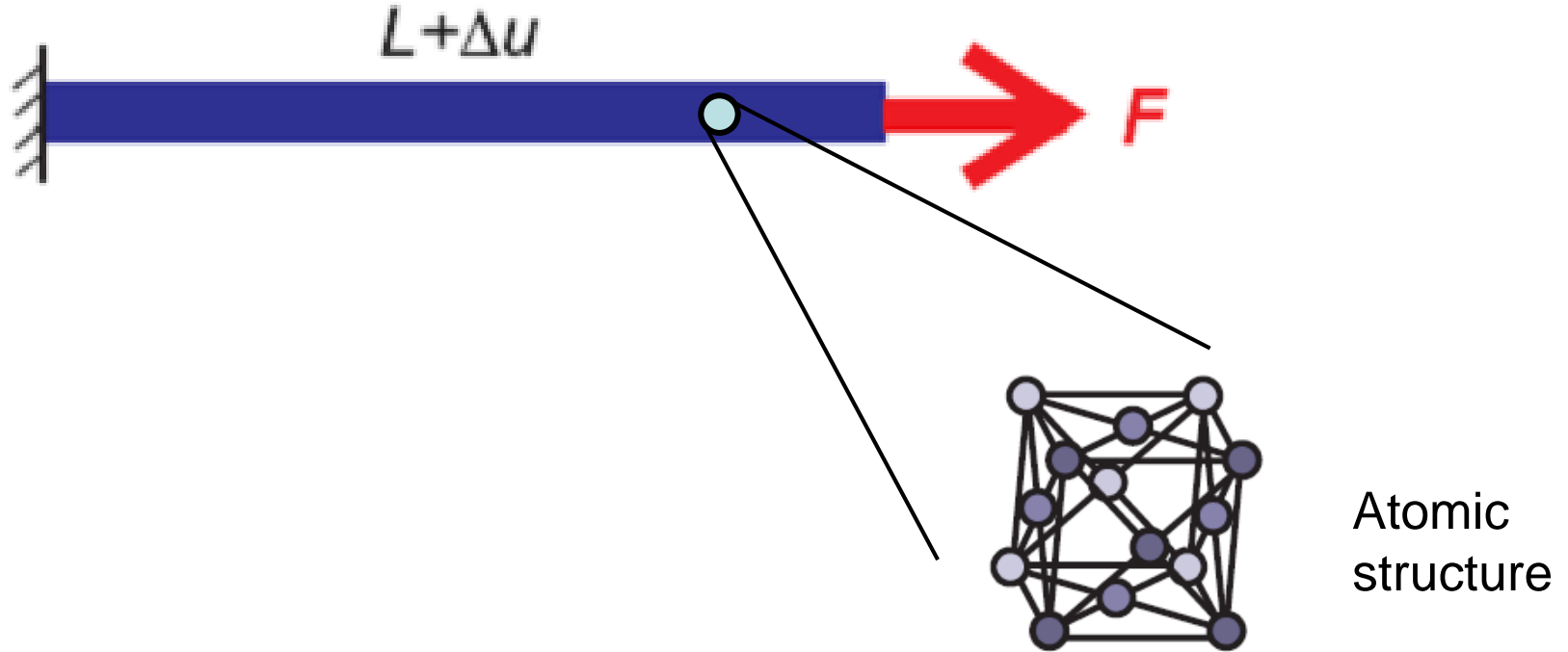
$$F = k_s \Delta u$$

$$E = \frac{L k_s}{A}$$

Young's modulus E (\sim stiffness= $\text{proportionality between force and displacement}$)

How to determine E ? - alternative approach to laboratory experiment

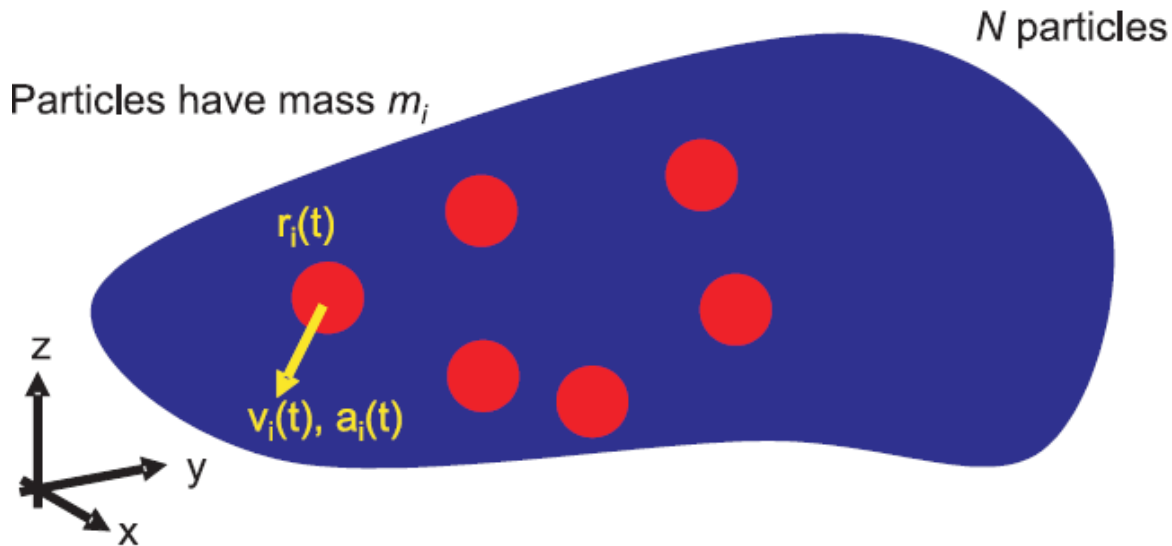
Atomistic simulation – *new engineering paradigm*



Concept: Consider the behavior of a collection of atoms inside the beam as deformation proceeds

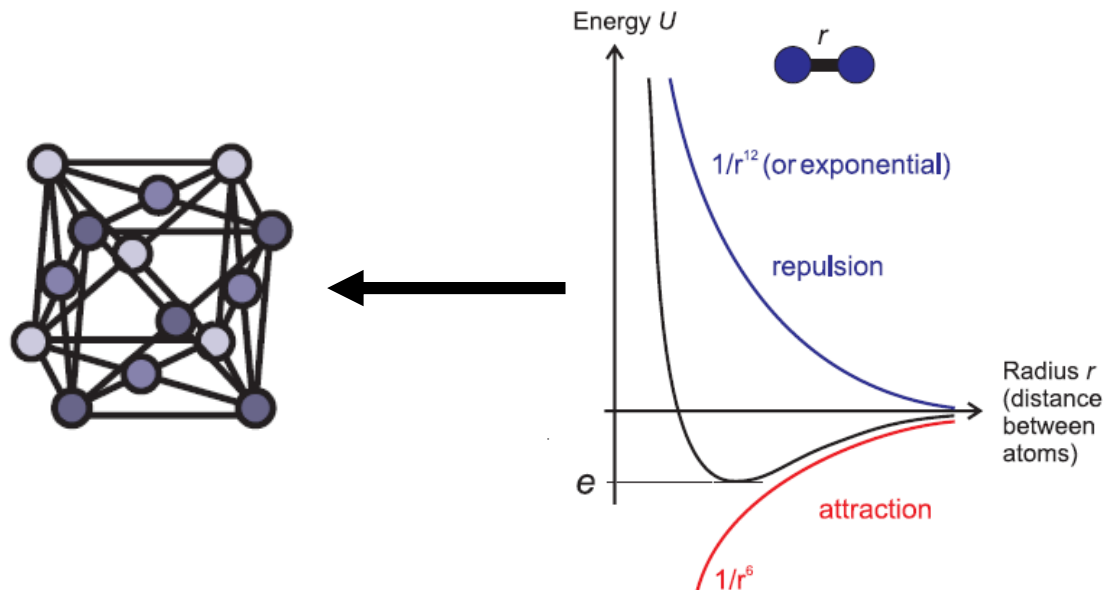
Molecular dynamics simulation

- Newton's laws: $F=ma$
- Chemistry: Atomic interactions – calculate interatomic forces from atomic interactions, that is, calculate F from energy landscape of atomic configuration (note that force and energy are related...)



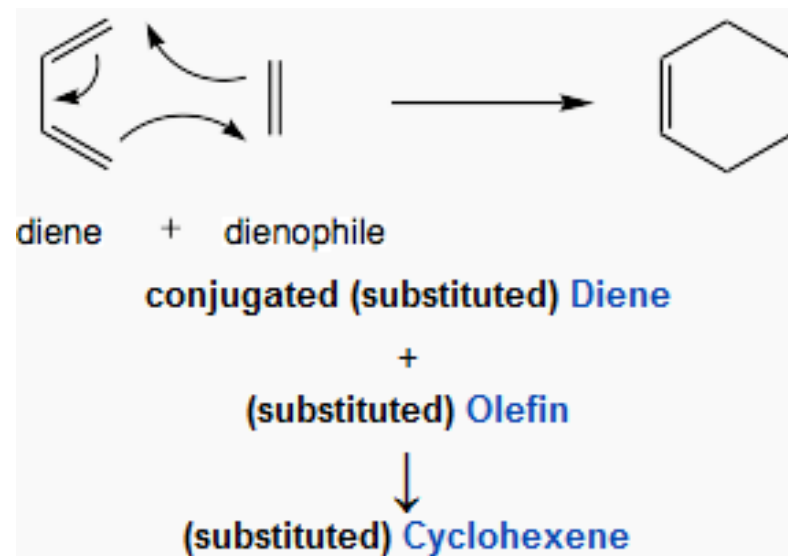
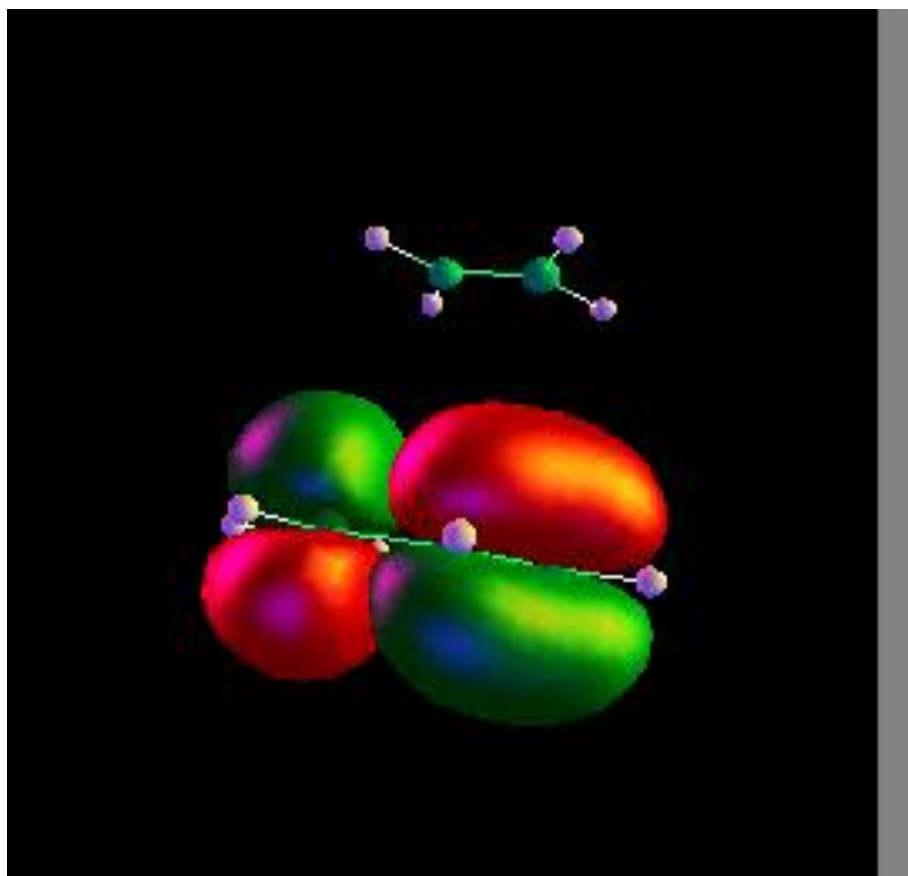
Linking atomistic and continuum perspective

- Atomistic viewpoint enables us to calculate how force and deformation is related, that is, we can predict E once we know the atomic structure and the type of chemical bonds
- Example: In metals we have metallic bonding and crystal structures – thus straightforward calculation of E
- *Atomistic models provide fundamental perspective, and thereby a means to determine (solely from the atomistic / chemical structure of the material) important parameters to be used in continuum models*



Quantum mechanics

- Deals with **fundamental view** of chemical bonding, based on electrons in atoms



“Schroedinger equation”

$$\frac{\partial^2 \psi}{\partial x^2} + \frac{8\pi^2 m}{h^2} (E - V) \psi = 0$$

Second derivative with respect to X

Shroedinger Wave Function

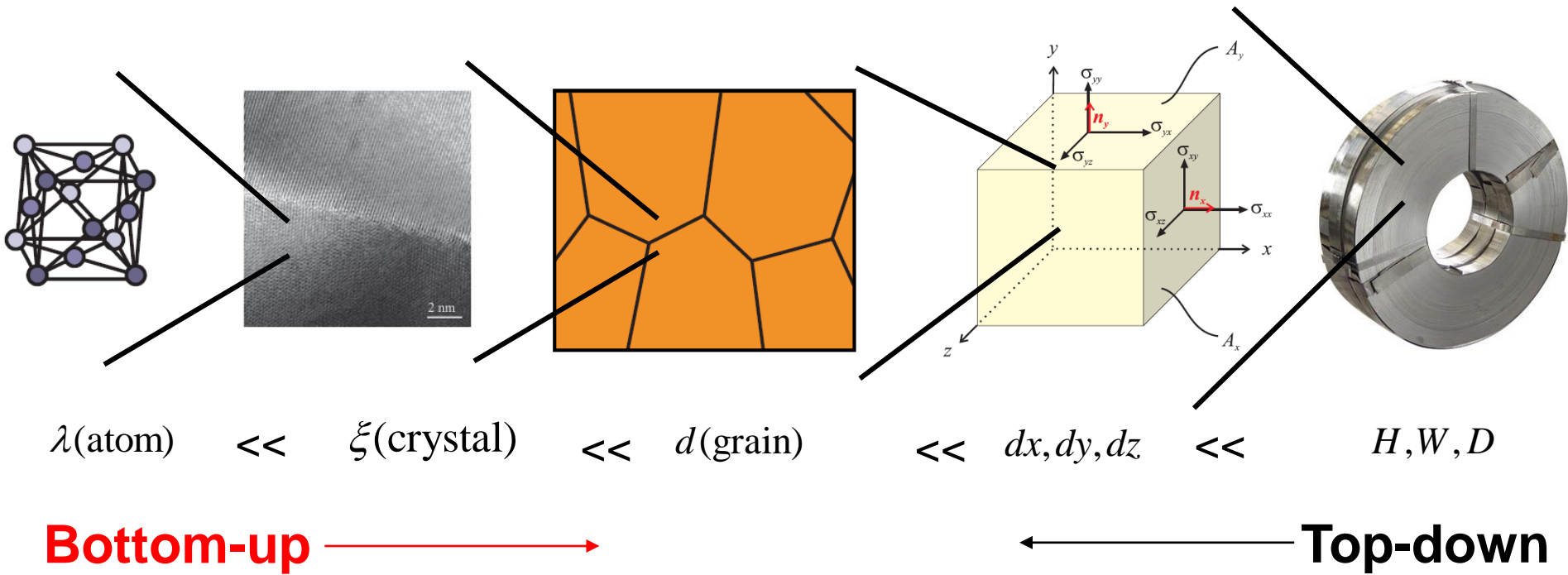
Position

Energy

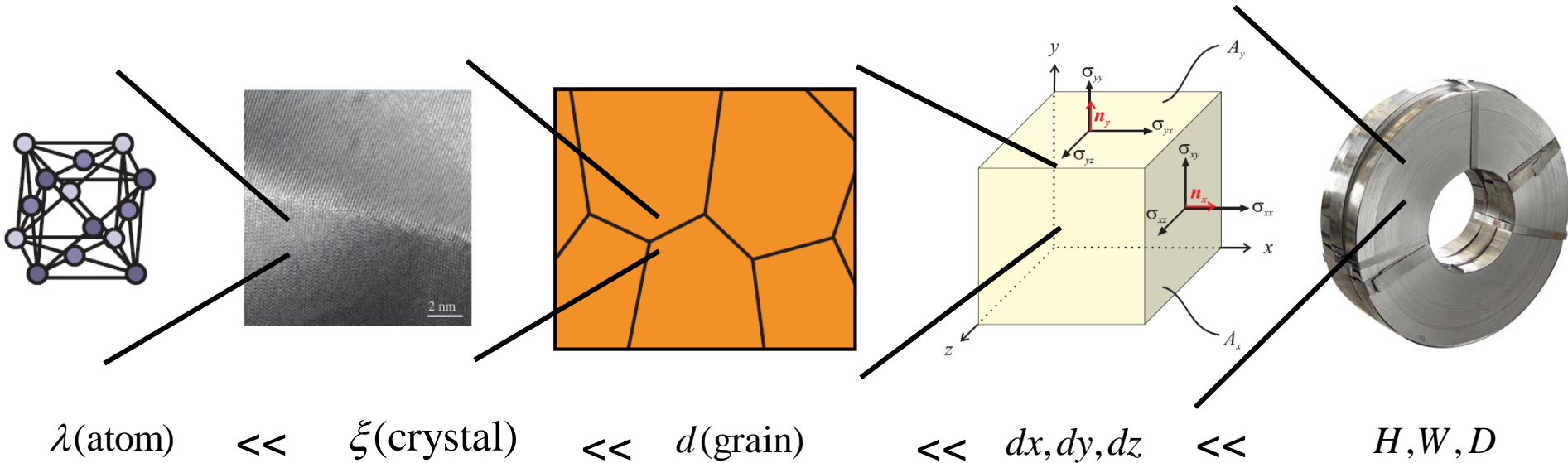
Potential Energy

2. Introduction: Bottom-up vs. top-down

Relevant scales in materials



Relevant scales in materials



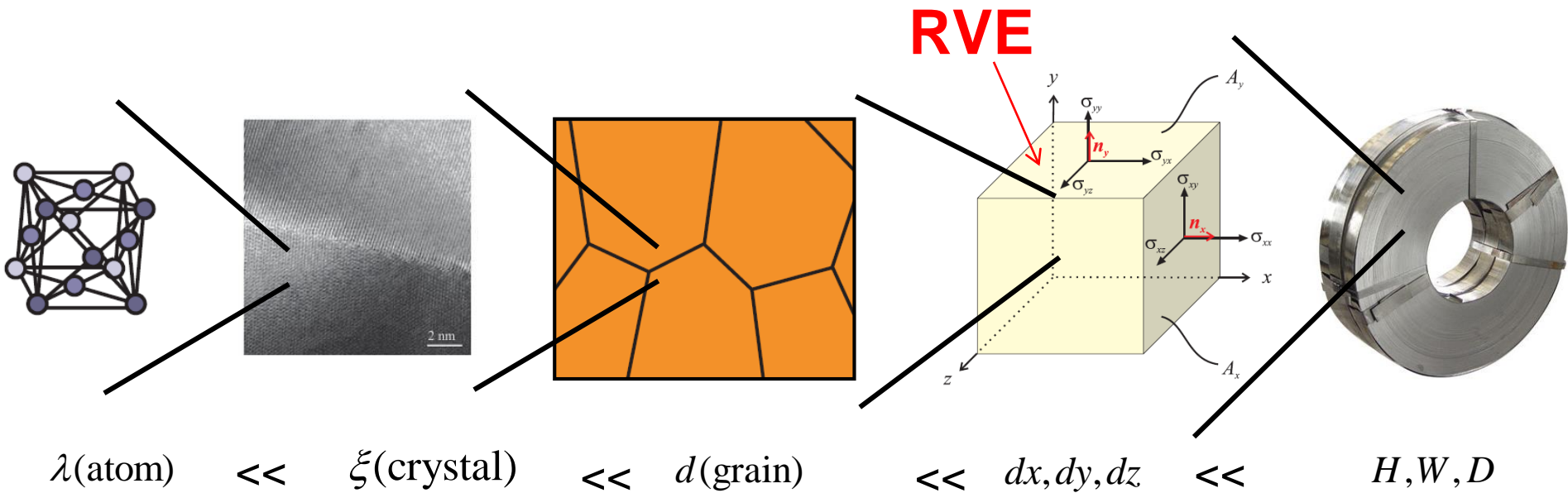
Bottom-up \longrightarrow

Atomistic viewpoint:

- Explicitly consider discrete atomistic structure
- Solve for atomic trajectories and infer from these about material properties & behavior
- Features internal length scales (atomic distance)

“Many-particle system with statistical properties”

Relevant scales in materials



← **Top-down**

Continuum viewpoint:

- Treat material as matter with no internal structure
- Develop mathematical model (governing equation) based on representative volume element (RVE, contains “enough” material such that internal structure can be neglected
- Features no characteristic length scales, provided RVE is large enough¹⁸

“PDE with parameters”

3. Case study: Diffusion

Continuum and atomistic modeling

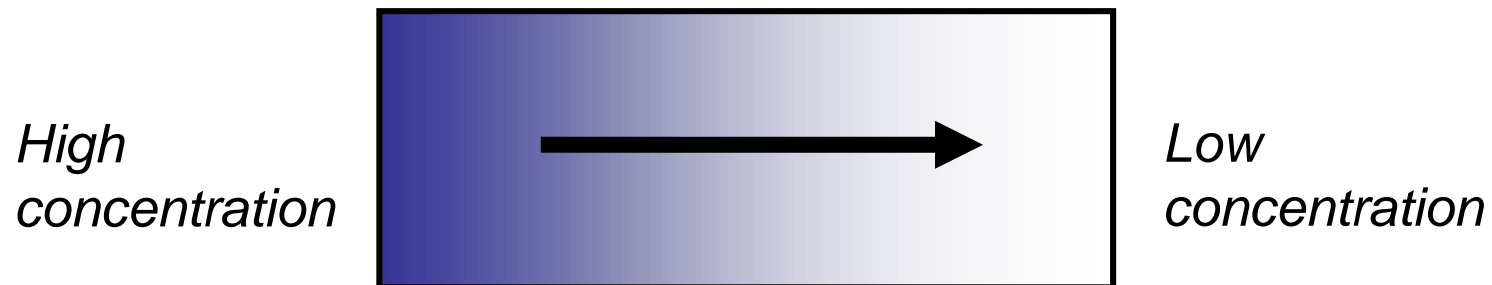
Goal of this section

- Diffusion as example
- **Continuum description** (top-down approach), partial differential equation
- **Atomistic description** (bottom-up approach), based on dynamics of molecules, obtained via numerical simulation of the molecular dynamics

Introduction: Diffusion

- Particles move from a domain with **high concentration** to an area of **low concentration**
- **Macroscopically**, diffusion measured by change in concentration
- **Microscopically**, diffusion is process of spontaneous net movement of particles

Result of random motion of particles (“Brownian motion”)



$$c = m/V = c(\vec{x}, t)$$

Ink droplet in water



hot

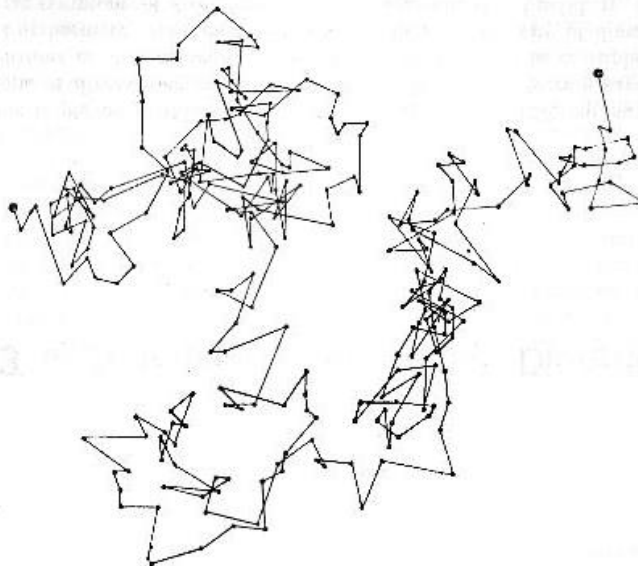
cold

(same time has passed)

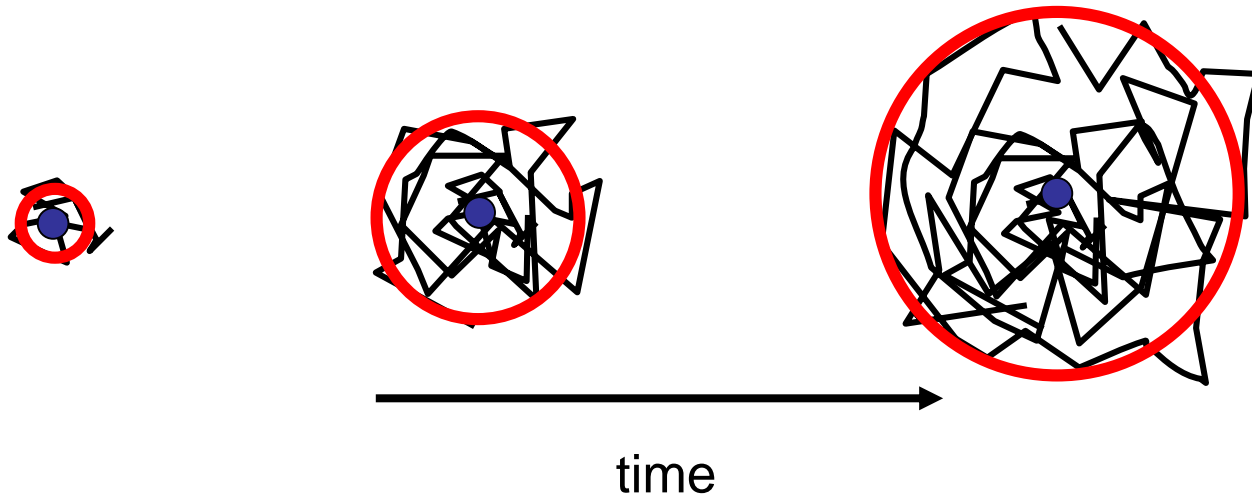
Microscopic observation of diffusion

Microscopic mechanism: “Random walk” – or Brownian motion

- Brownian motion was first observed (1827) by the British botanist **Robert Brown (1773-1858)** when studying pollen grains in water
- Initially thought to be sign of life, but later confirmed to be also present in inorganic particles
- The effect was finally explained in **1905 by Albert Einstein**, who realized it was caused by water molecules randomly smacking into the particles.



Brownian motion leads to net particle movement



Particle “slowly” moves away from its initial position

Robert Brown's microscope

1827



<http://www.brianjford.com/pbrownmica.jpg>

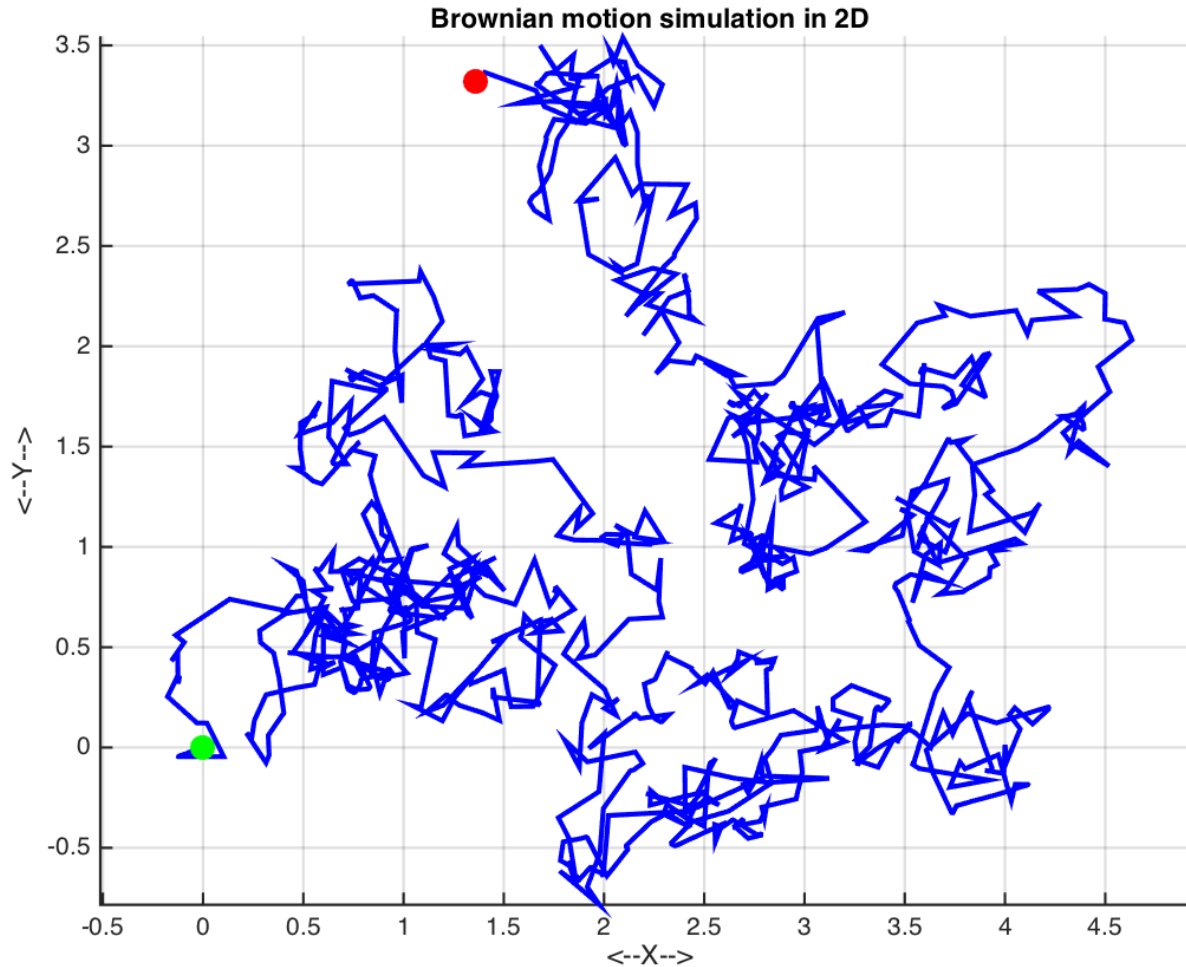
Robert Brown's Microscope

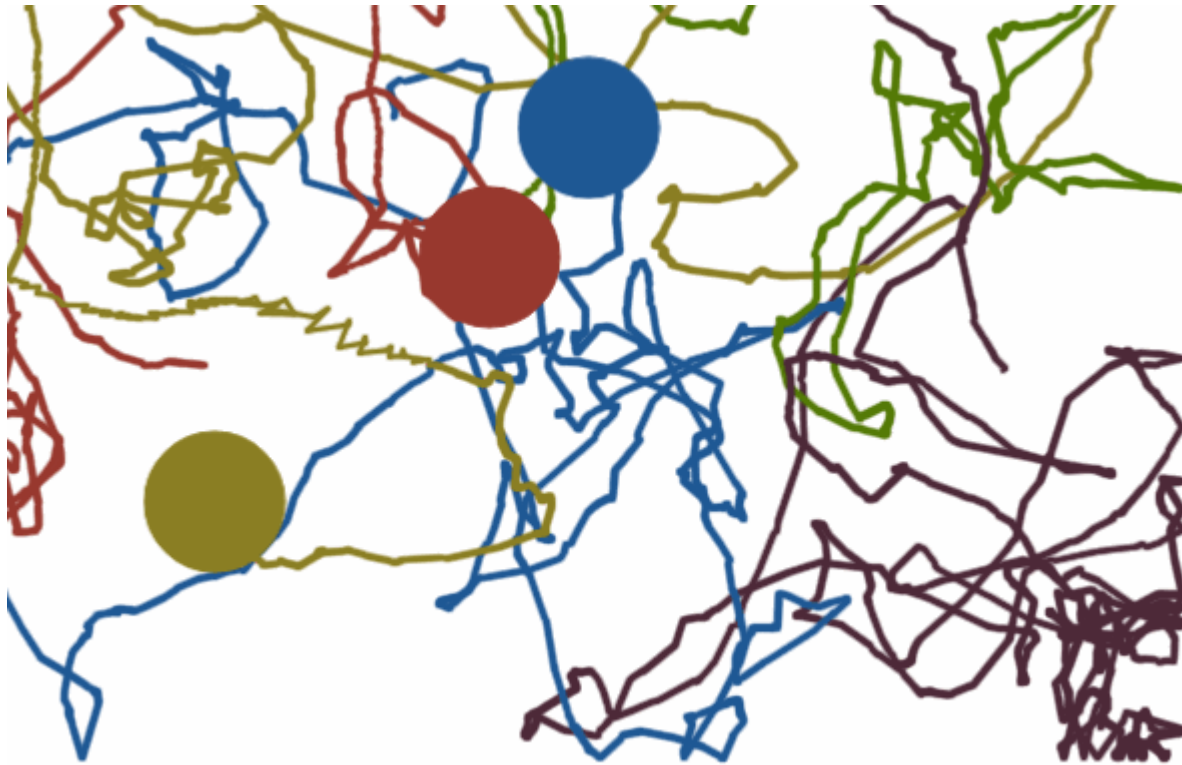
Instrument with which Robert Brown studied **Brownian motion** and which he used in his work on identifying the nucleus of the living cell

*Instrument is preserved at the **Linnean Society in London***

It is made of brass and is mounted onto the lid of the box in which it can be stored

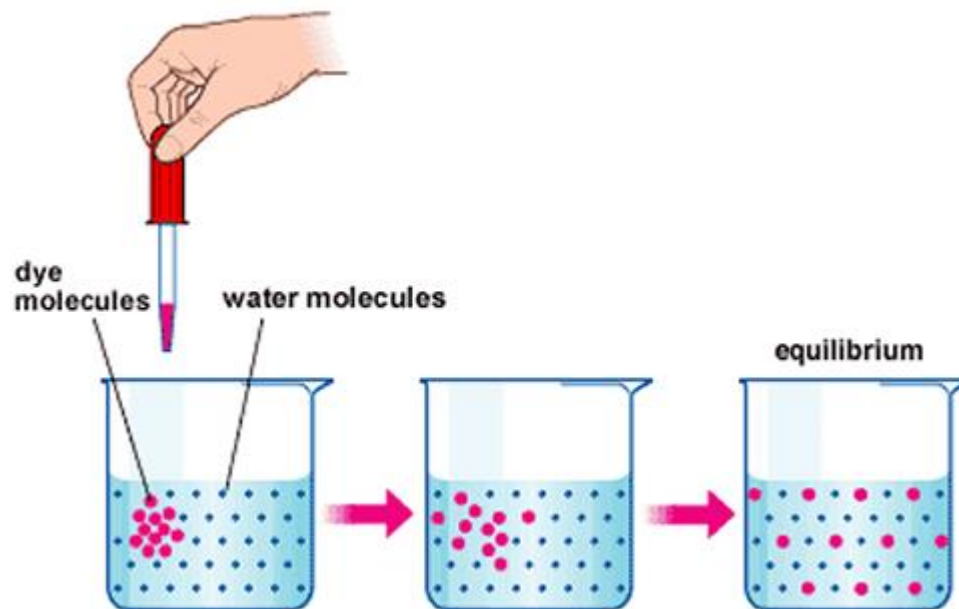
Simulation of Brownian motion



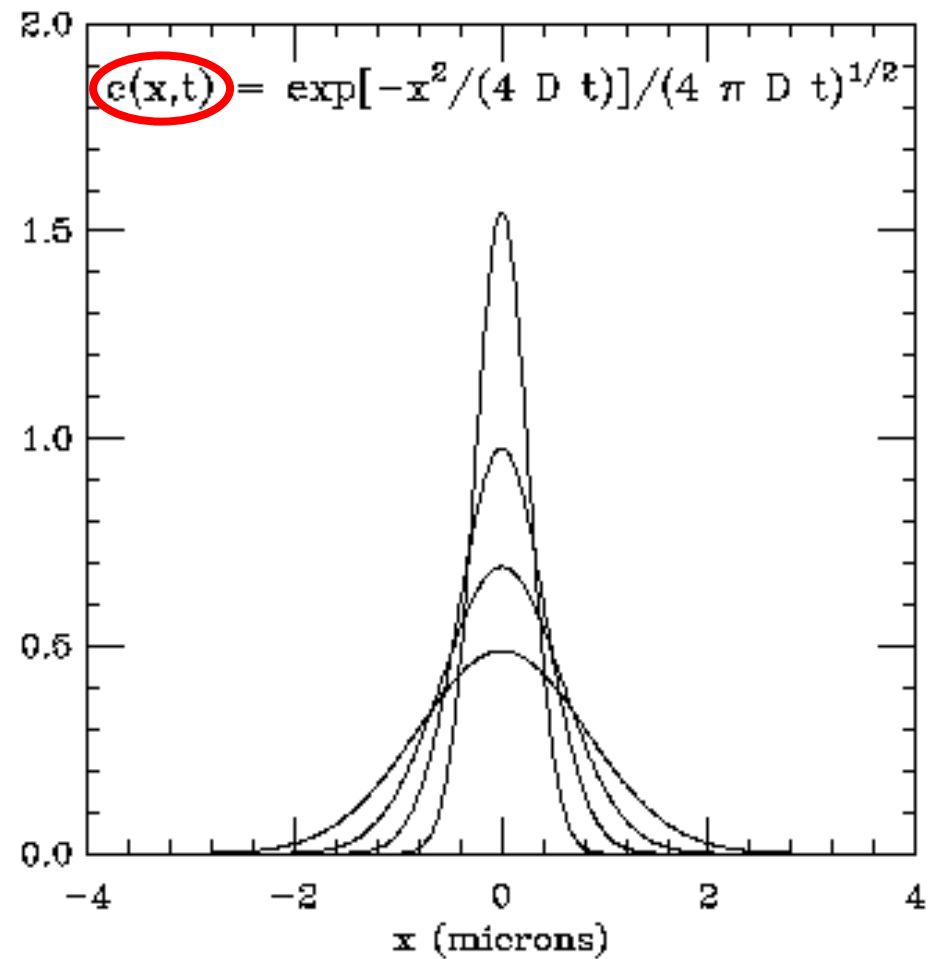
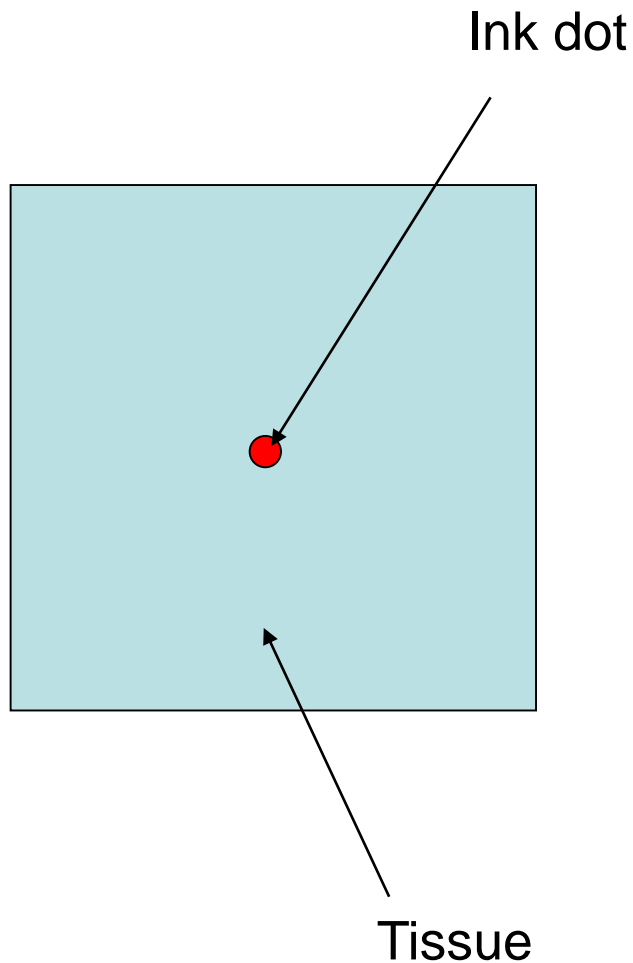




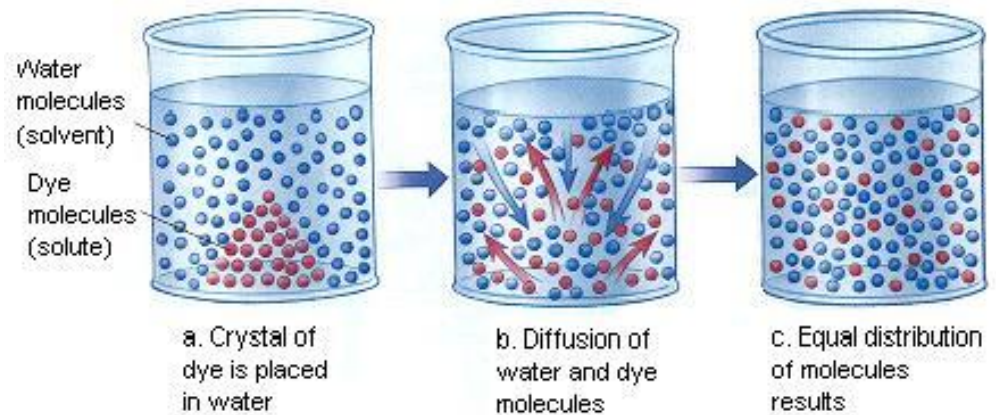
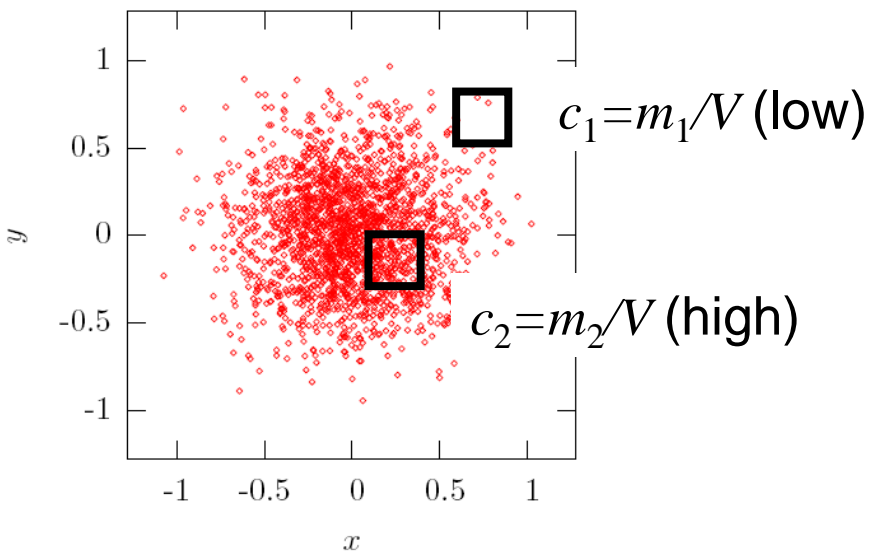
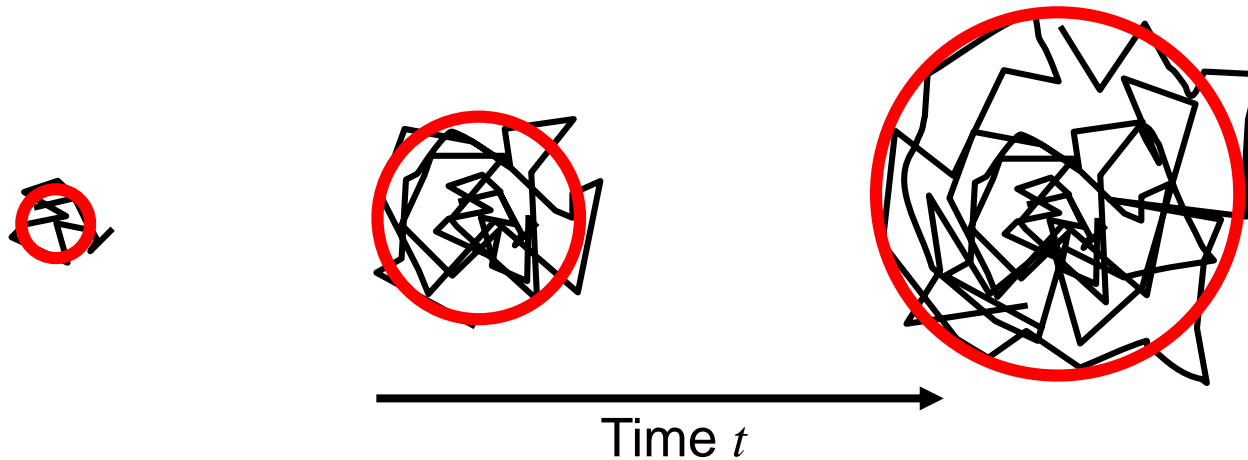
Macroscopic observation of diffusion



Macroscopic observation: concentration change



Brownian motion leads to net particle movement

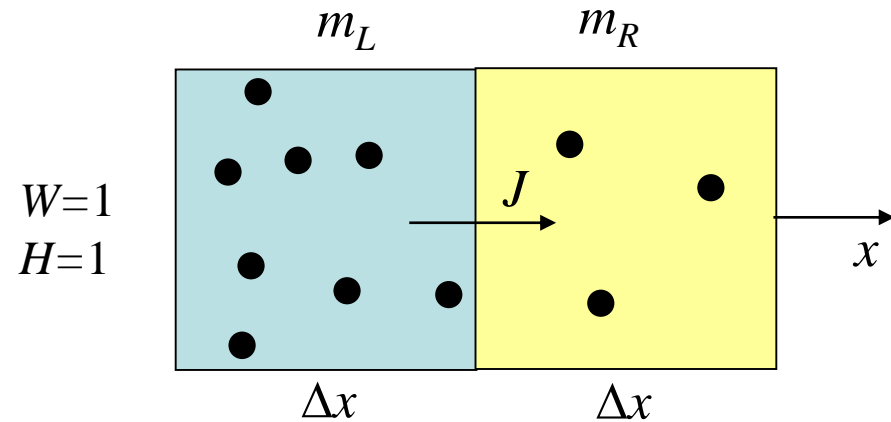


3.1 Continuum model

How to build a continuum model to describe the physical phenomena of diffusion?

Approach 1: Continuum model

- Develop differential equation based on differential element

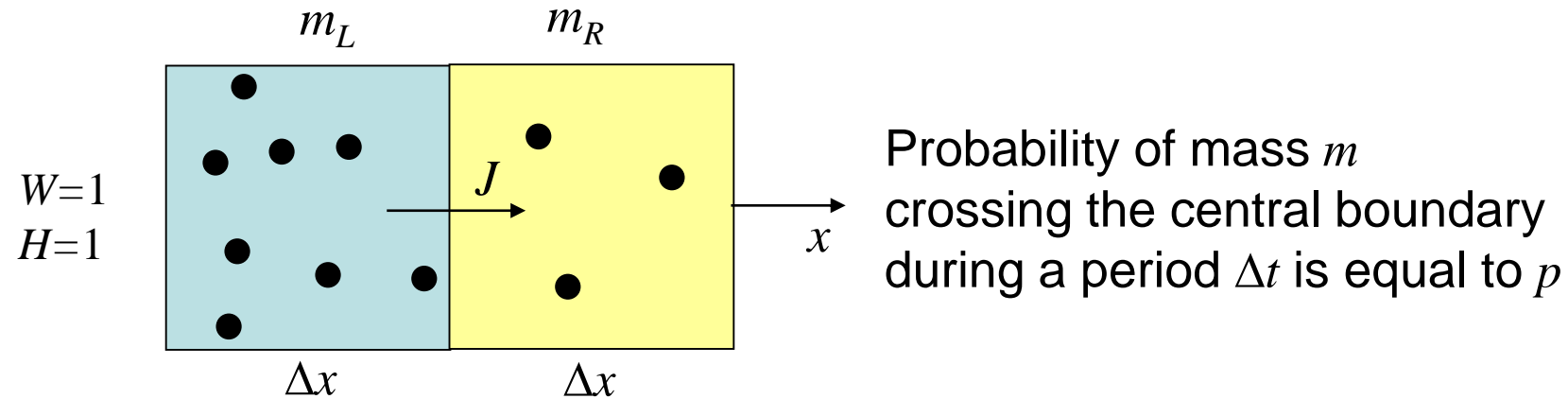


J : Mass flux (mass per unit time per unit area)

Concept: Balance mass [here], force etc. in a differential volume element; much greater in dimension than inhomogeneities (“sufficiently large RVE”)

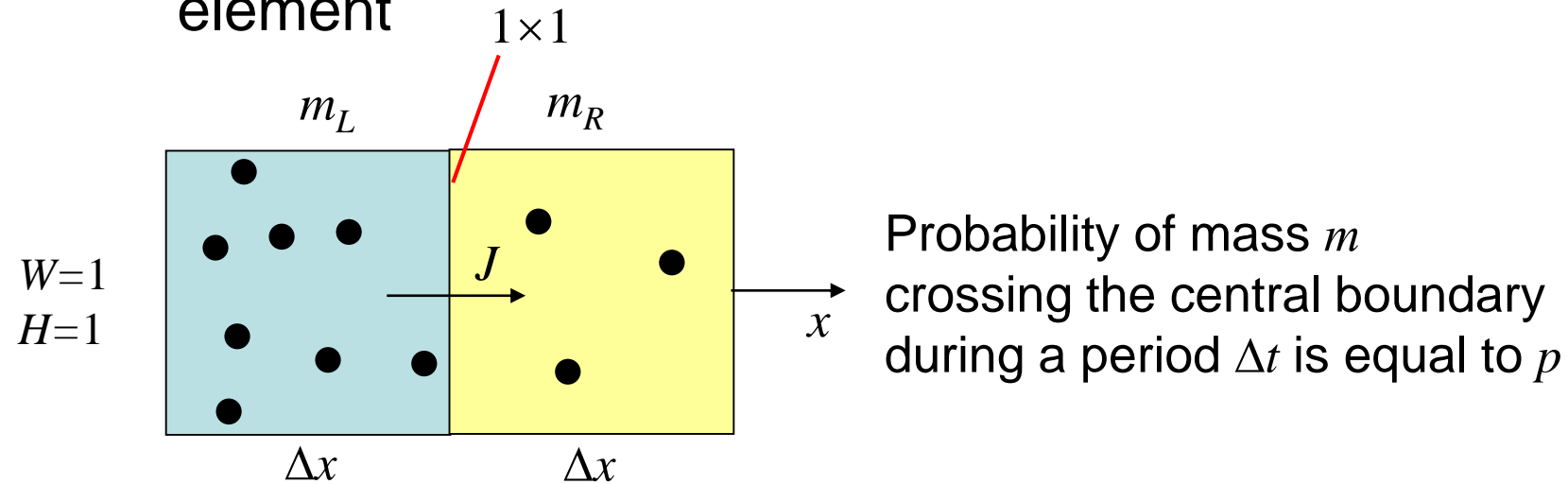
Approach 1: Continuum model

- Develop differential equation based on differential element



Approach 1: Continuum model

- Develop differential equation based on differential element



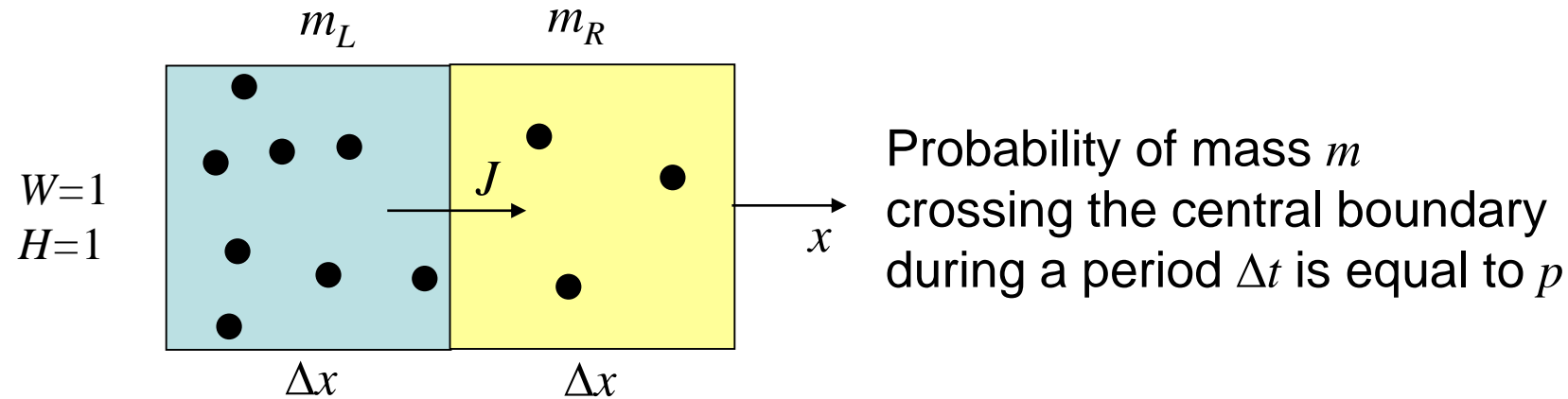
$$J_L = \frac{1}{1 \times 1} \frac{p}{\Delta t} m_L \quad \text{Mass flux from left to right}$$

$$J_R = \frac{1}{1 \times 1} \frac{p}{\Delta t} m_R \quad \text{Mass flux from right to left}$$

$[J]$ = mass per unit time per unit area

Approach 1: Continuum model

- Develop differential equation based on differential element



$$J_L = \frac{1}{1 \times 1} \frac{p}{\Delta t} m_L \quad \text{Mass flux from left to right}$$

$$J_R = \frac{1}{1 \times 1} \frac{p}{\Delta t} m_R \quad \text{Mass flux from right to left}$$

Effective mass flux

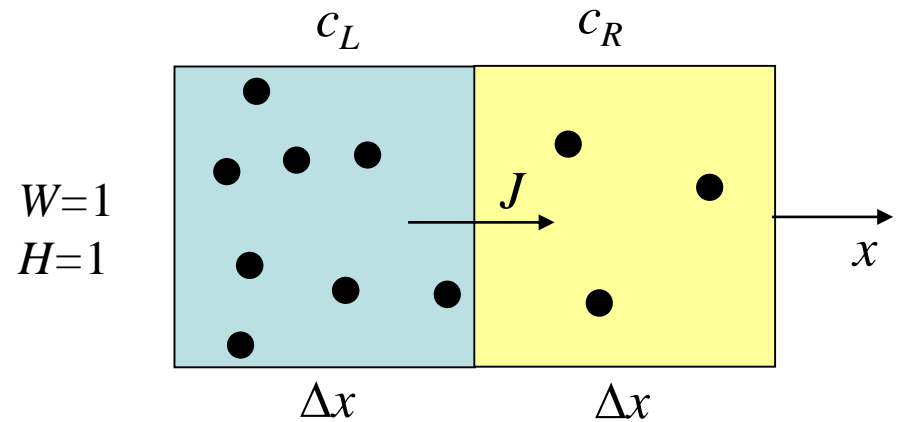
$$J = \frac{1}{1 \times 1} \frac{p}{\Delta t} (m_L - m_R)$$

More mass, more flux (m_L is \sim to number of particles)

Continuum model of diffusion

Express in terms of mass concentrations

$$c = \frac{m}{V} \quad m = cV \quad J = \frac{p}{\Delta t} (m_L - m_R)$$

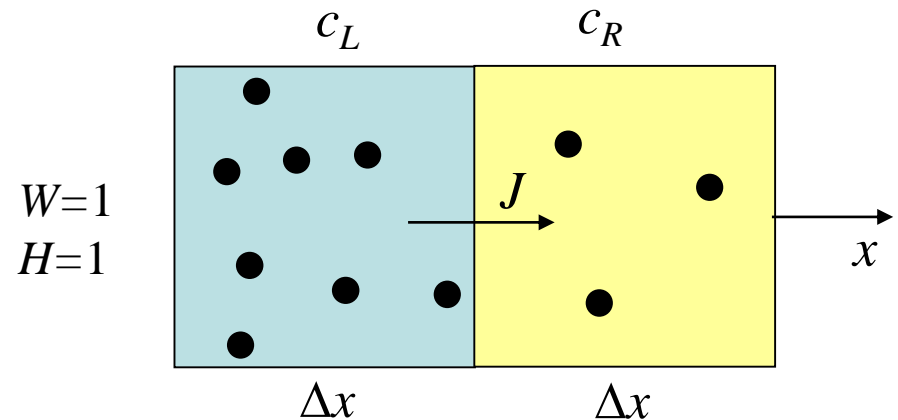


Continuum model of diffusion

Express in terms of mass concentrations

$$c = \frac{m}{V} \quad m = cV \quad J = \frac{p}{\Delta t} (m_L - m_R)$$

$$J = \frac{1}{1 \times 1} \frac{p}{\Delta t} \underbrace{(c_L - c_R)}_{= -\Delta c} \cdot \underbrace{\Delta x \times 1 \times 1}_{= V}$$



Continuum model of diffusion

Express in terms of mass concentrations

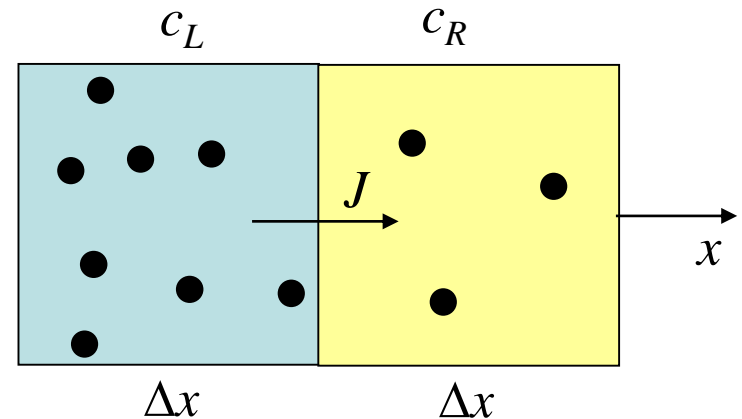
$$c = \frac{m}{V}$$

$$J = \frac{1}{\cancel{1 \times 1}} \frac{p}{\Delta t} \underbrace{(c_L - c_R)}_{= -\Delta c} \cdot \underbrace{\Delta x \times \cancel{1 \times 1}}_{= V}$$

$$J = -\frac{p}{\Delta t} \Delta c \Delta x = -\frac{p}{\Delta t} \Delta x^2 \frac{\Delta c}{\Delta x}$$

Concentration
gradient

expand Δx



Continuum model of diffusion

Express in terms of mass concentrations

$$c = \frac{m}{V}$$

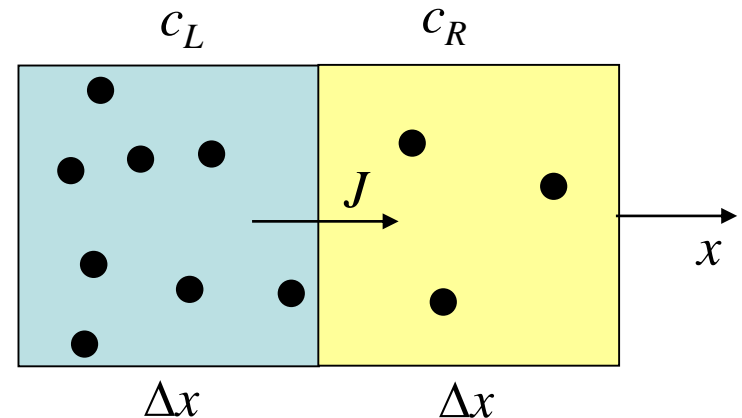
$$J = \frac{1}{\cancel{1 \times 1}} \frac{p}{\Delta t} \underbrace{(c_L - c_R)}_{= -\Delta c} \cdot \underbrace{\Delta x \times \cancel{1 \times 1}}_{= V}$$

$$J = -\frac{p}{\Delta t} \Delta c \Delta x = -\frac{p}{\Delta t} \Delta x^2 \frac{\Delta c}{\Delta x}$$

Concentration
gradient

$$W=1$$

$$H=1$$



$$J = -\left[\frac{p}{\Delta t} \Delta x^2 \right] \frac{dc}{dx} = -D \frac{dc}{dx}$$

$$D = p \frac{\Delta x^2}{\Delta t}$$

Parameter that measures how “fast” mass moves (in square of distance per unit time)

Diffusion constant & 1st Fick law

Reiterate: Diffusion constant D describes the how much mass moves per unit time

Movement of mass characterized by square of displacement from initial position

Flux

$$J = -\frac{p}{\Delta t} \Delta x^2 \frac{dc}{dx} = -D \frac{dc}{dx}$$

1st Fick law
(Adolph Fick, 1829-1901)

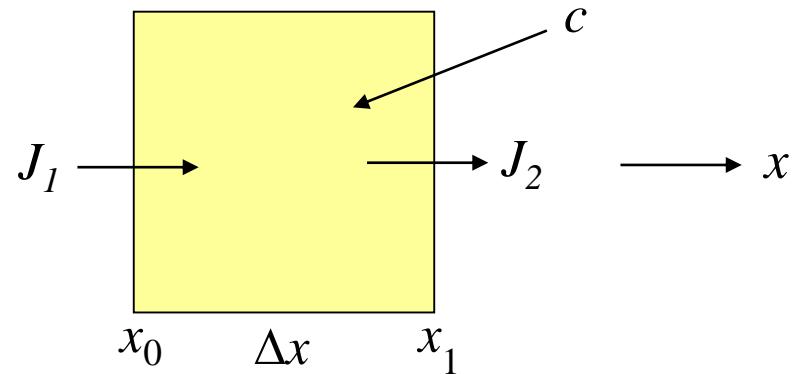
$$D = p \frac{\Delta x^2}{\Delta t}$$

$$D \sim p$$

Diffusion constant relates to the ability of mass to move a distance Δx^2 over a time Δt (strongly temperature dependent, e.g. Arrhenius)

2nd Fick law (time dependence)

$$J = -D \frac{dc}{dx} \quad \text{1st Fick law}$$

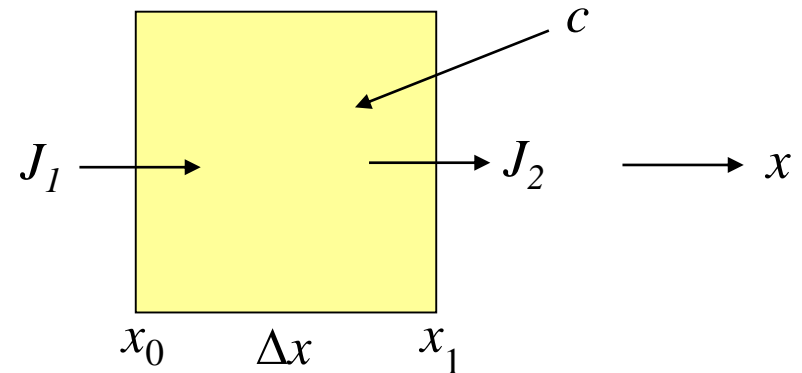


$$\frac{\Delta c}{\Delta t} = \frac{(J_1 - J_2) \times \cancel{1 \times 1}}{\Delta x \times \cancel{1 \times 1}}$$

J : Mass flux (mass per unit time per unit area)

2nd Fick law (time dependence)

$$J = -D \frac{dc}{dx} \quad \text{1st Fick law}$$

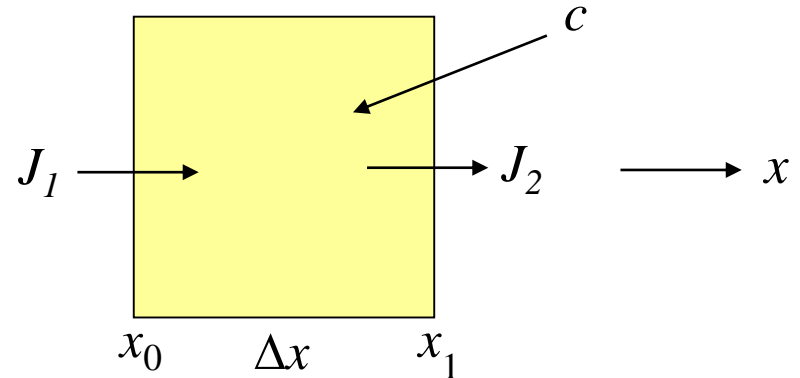


$$\frac{\Delta c}{\Delta t} = \frac{J_1 - J_2}{\Delta x} = \frac{1}{\Delta x} \left(-D \frac{dc}{dx} \Big|_{x=x_0} - \left[-D \frac{dc}{dx} \Big|_{x=x_1} \right] \right)$$

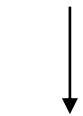
$$J_1 = J(x = x_0) = -D \frac{dc}{dx} \Big|_{x=x_0}$$

2nd Fick law (time dependence)

$$J = -D \frac{dc}{dx} \quad \text{1st Fick law}$$



$$\frac{\Delta c}{\Delta t} = \frac{1}{\Delta x} \left(\overbrace{-D \frac{dc}{dx} \Big|_{x=x_0} + D \frac{dc}{dx} \Big|_{x=x_1}}^{\Delta J} \right)$$



$$\frac{\partial c}{\partial t} = -\frac{d}{dx}(J) = -\frac{d}{dx} \left(-D \frac{dc}{dx} \right)$$

Change of concentration in time equals change of flux with x (mass balance)

$$\frac{\partial c}{\partial t} = D \frac{d^2 c}{dx^2}$$

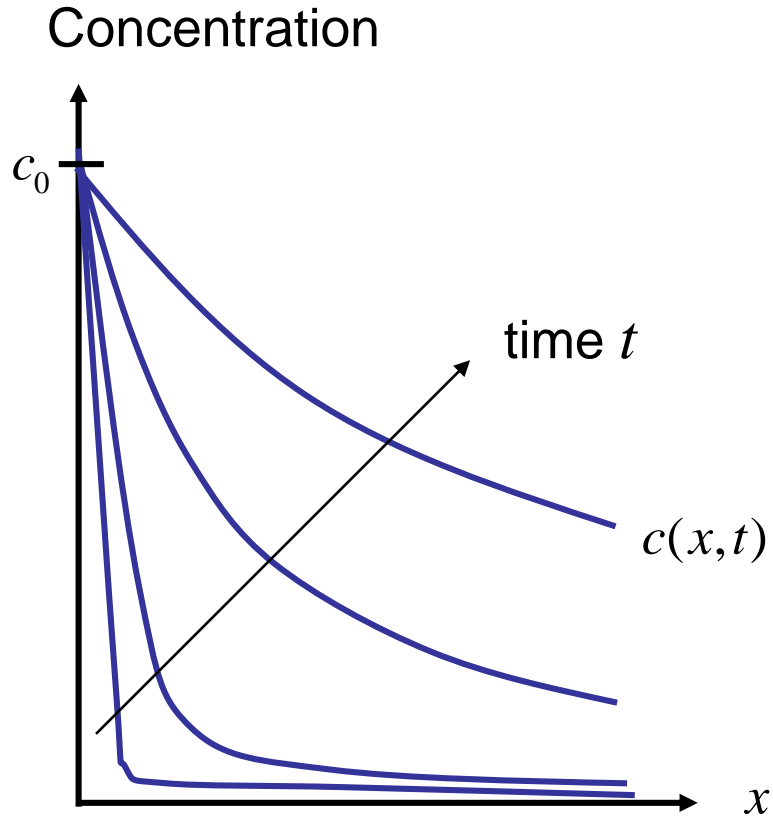
2nd Fick law

PDE

Solve by applying ICs and BCs...

Example solution – 2nd Fick's law

$$\frac{\partial c}{\partial t} = D \frac{d^2 c}{dx^2}$$



Boundary condition (BC): $c(x=0) = c_0$

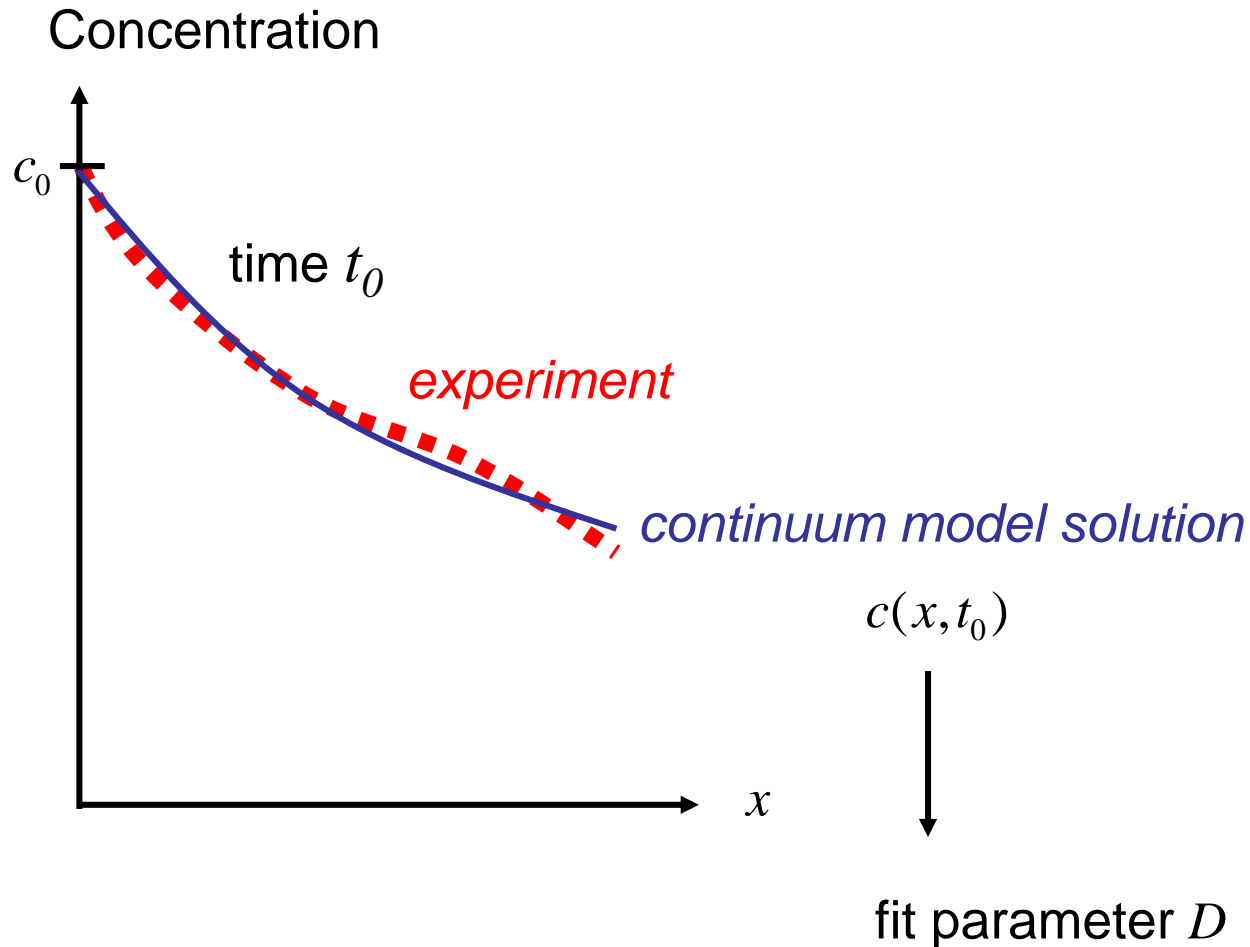
Initial condition (IC): $c(x > 0, t=0) = 0$

Need diffusion coefficient to solve for distribution!

How to obtain diffusion coefficient?

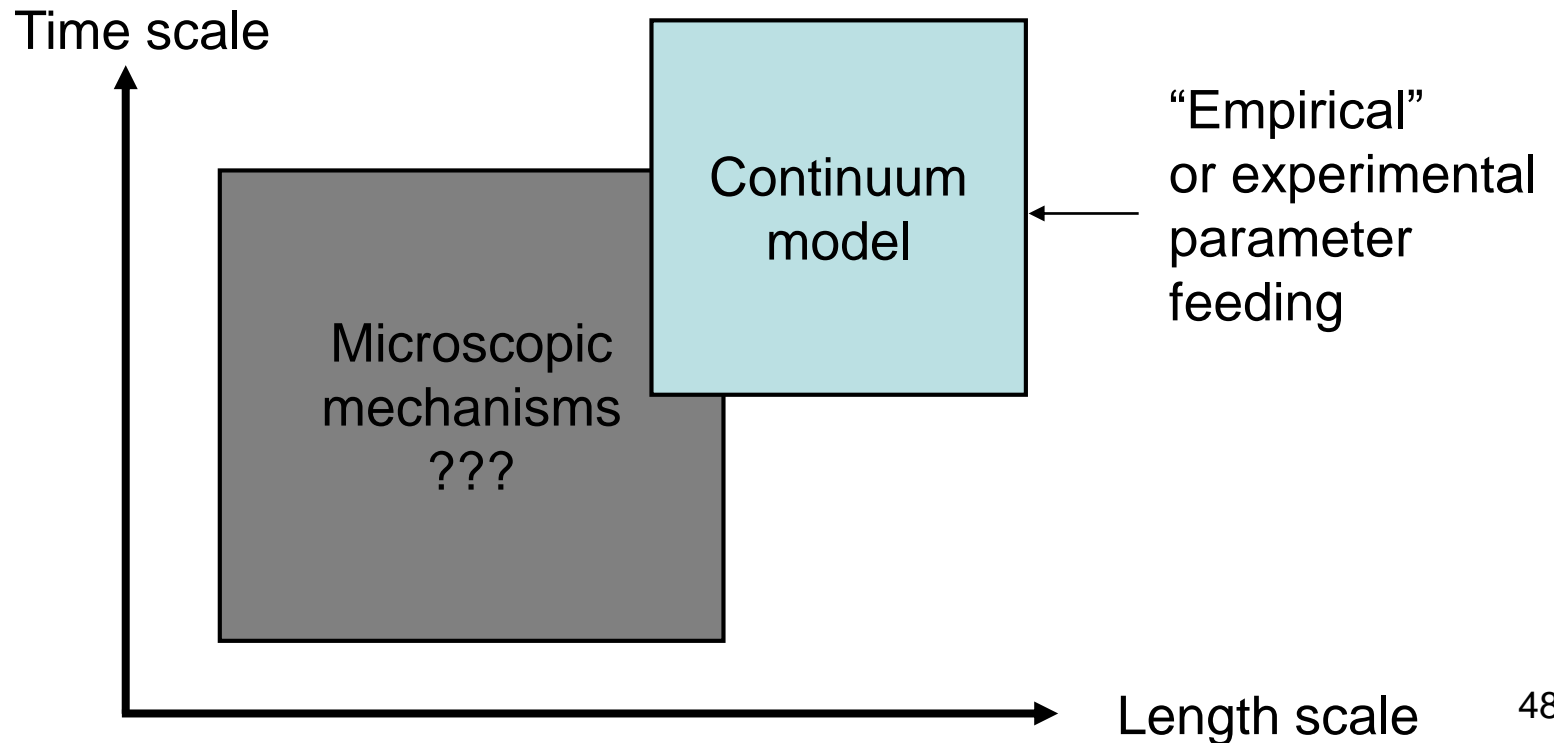
- Laboratory experiment
- Study distribution of concentrations (previous slide)
- Then “fit” the appropriate diffusion coefficient so that the solution matches
- Approach can then be used to solve for more complex geometries, situations etc. for which no lab experiment exists
- “Top down approach”

Matching with experiment (parameter identification)



Summary

- Continuum model requires parameter that describes microscopic processes inside the material
- Typically need experimental measurements to calibrate



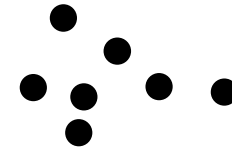
3.2 Atomistic model

How to build an atomistic bottom-up model to describe the physical phenomena of diffusion?

Approach 2: Atomistic model

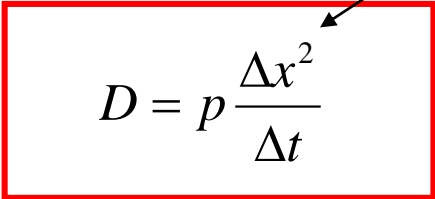
- Atomistic model provides an alternative approach to describe diffusion
- Enables us to directly calculate the diffusion constant from the trajectory of atoms (“microscopic definition”)

- Approach: Consider set of atoms/molecules



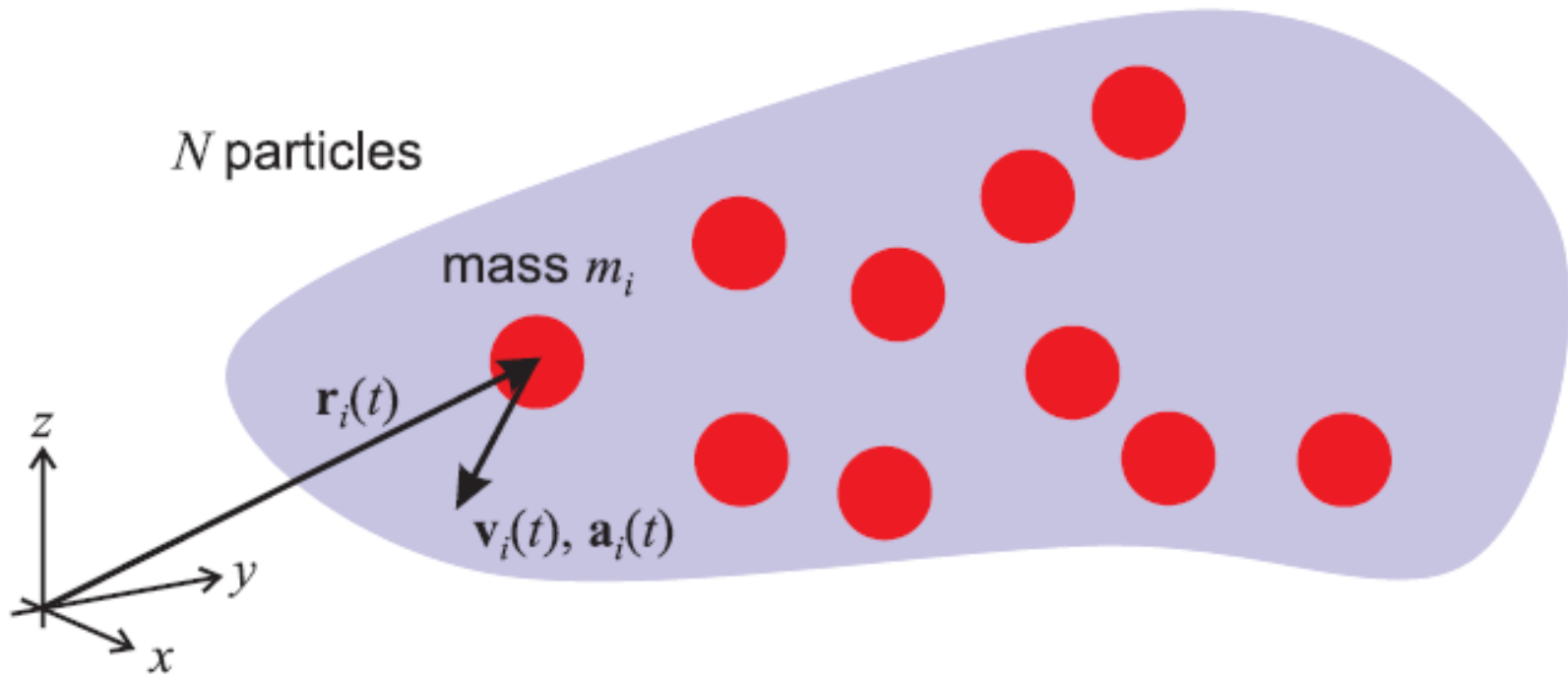
- Follow their trajectory and calculate how fast atoms leave their initial position

Follow this quantity over time


$$D = p \frac{\Delta x^2}{\Delta t}$$

Recall: Diffusion constant relates to the “ability” of particle to move a distance Δx^2 over a time Δt

Molecular dynamics – simulate trajectory of atoms



Goal: Need an algorithm to predict positions, velocities, accelerations as function of time

Solving the equations: What we want

To solve those equations: Discretize in time (n steps), Δt time step:

$$r_i(t_0) \rightarrow r_i(t_0 + \Delta t) \rightarrow r_i(t_0 + 2\Delta t) \rightarrow r_i(t_0 + 3\Delta t) \rightarrow \dots \rightarrow r_i(t_0 + n\Delta t)$$

Solving the equations

To solve those equations: Discretize in time (n steps), Δt time step:

$$r_i(t_0) \rightarrow r_i(t_0 + \Delta t) \rightarrow r_i(t_0 + 2\Delta t) \rightarrow r_i(t_0 + 3\Delta t) \rightarrow \dots \rightarrow r_i(t_0 + n\Delta t)$$

Recall: Taylor expansion of function f around point a

$$f(x) = f(a) + f'(a)(x-a) + \frac{f''(a)}{2!}(x-a)^2 + \frac{f^{(3)}(a)}{3!}(x-a)^3 + \dots + \frac{f^{(n)}(a)}{n!}(x-a)^n + \dots$$

Solving the equations

To solve those equations: Discretize in time (n steps), Δt time step:

$$r_i(t_0) \rightarrow r_i(t_0 + \Delta t) \rightarrow r_i(t_0 + 2\Delta t) \rightarrow r_i(t_0 + 3\Delta t) \rightarrow \dots \rightarrow r_i(t_0 + n\Delta t)$$

Recall: Taylor series expansion of function f around point a

$$f(x) = f(a) + f'(a)(x-a) + \frac{f''(a)}{2!}(x-a)^2 + \frac{f^{(3)}(a)}{3!}(x-a)^3 + \dots + \frac{f^{(n)}(a)}{n!}(x-a)^n + \dots$$

Taylor series expansion $r_i(t)$ around

$$a = t_0 \quad x = t_0 + \Delta t$$

$$x - a = t_0 + \Delta t - t_0 = \Delta t$$

Solving the equations

To solve those equations: Discretize in time (n steps), Δt time step:

$$r_i(t_0) \rightarrow r_i(t_0 + \Delta t) \rightarrow r_i(t_0 + 2\Delta t) \rightarrow r_i(t_0 + 3\Delta t) \rightarrow \dots \rightarrow r_i(t_0 + n\Delta t)$$

Recall: Taylor series expansion of function f around point a

$$f(x) = f(a) + f'(a)(x-a) + \frac{f''(a)}{2!}(x-a)^2 + \frac{f^{(3)}(a)}{3!}(x-a)^3 + \dots + \frac{f^{(n)}(a)}{n!}(x-a)^n + \dots$$

Taylor series expansion $r_i(t)$ around

$$a = t_0 \quad x = t_0 + \Delta t$$

$$x - a = t_0 + \Delta t - t_0 = \Delta t$$

$$r_i(t_0 + \Delta t) = r_i(t_0) + v_i(t_0)\Delta t + \frac{1}{2}a_i(t_0)\Delta t^2 + \dots$$

Taylor expansion of $r_i(t)$

$$r_i(t_0 + \Delta t) = r_i(t_0) + v_i(t_0)\Delta t + \frac{1}{2}a_i(t_0)\Delta t^2 + \dots$$

$$r_i(t_0 - \Delta t) = r_i(t_0) - v_i(t_0)\Delta t + \frac{1}{2}a_i(t_0)\Delta t^2 + \dots$$

$$a = t_0 \quad x = t_0 + \Delta t \quad x - a = t_0 + \Delta t - t_0 = \Delta t$$

$$a = t_0 \quad x = t_0 - \Delta t \quad x - a = t_0 - \Delta t - t_0 = -\Delta t$$

Taylor expansion of $r_i(t)$

$$r_i(t_0 + \Delta t) = r_i(t_0) + v_i(t_0)\Delta t + \frac{1}{2}a_i(t_0)\Delta t^2 + \dots$$
$$+ \left[r_i(t_0 - \Delta t) = r_i(t_0) - v_i(t_0)\Delta t + \frac{1}{2}a_i(t_0)\Delta t^2 + \dots \right]$$



$$r_i(t_0 - \Delta t) + r_i(t_0 + \Delta t) = 2r_i(t_0) - \cancel{v_i(t_0)\Delta t} + \cancel{v_i(t_0)\Delta t} + a_i(t_0)\Delta t^2 + \dots$$



$$r_i(t_0 + \Delta t) = \underbrace{2r_i(t_0)}_{\text{Positions at } t_0} - \underbrace{r_i(t_0 - \Delta t)}_{\text{Positions at } t_0 - \Delta t} + \underbrace{a_i(t_0)\Delta t^2}_{\text{Accelerations at } t_0} + \dots$$

Positions
at t_0

Positions
at $t_0 - \Delta t$

Accelerations
at t_0

Physics of particle interactions

Laws of Motion of Isaac Newton (1642 – 1727):

1. Every body continues in its state of rest, or of uniform motion in a right line, unless it is compelled to change that state by forces impressed upon it.
2. The change of motion is proportional to the motive force impresses, and is made in the direction of the right line in which that force is impressed.
3. To every action there is always opposed an equal reaction: or, the mutual action of two bodies upon each other are always equal, and directed to contrary parts.

$$f = m \frac{d^2 x}{dt^2} = ma$$

2nd law

Verlet central difference method

$$r_i(t_0 + \Delta t) = \underbrace{2r_i(t_0)}_{\substack{\text{Positions} \\ \text{at } t_0}} - \underbrace{r_i(t_0 - \Delta t)}_{\substack{\text{Positions} \\ \text{at } t_0 - \Delta t}} + \underbrace{a_i(t_0)\Delta t^2}_{\substack{\text{Accelerations} \\ \text{at } t_0}} + \dots$$

How to obtain
accelerations?

$$f_i = ma_i$$
$$a_i = f_i / m$$

Need forces on atoms!

Verlet central difference method

$$r_i(t_0 + \Delta t) = 2\underbrace{r_i(t_0)}_{\text{Positions at } t_0} - \underbrace{r_i(t_0 - \Delta t)}_{\text{Positions at } t_0 - \Delta t} + \underbrace{f_i(t_0)/m\Delta t^2}_{\text{Forces at } t_0} + \dots$$

Positions
at t_0

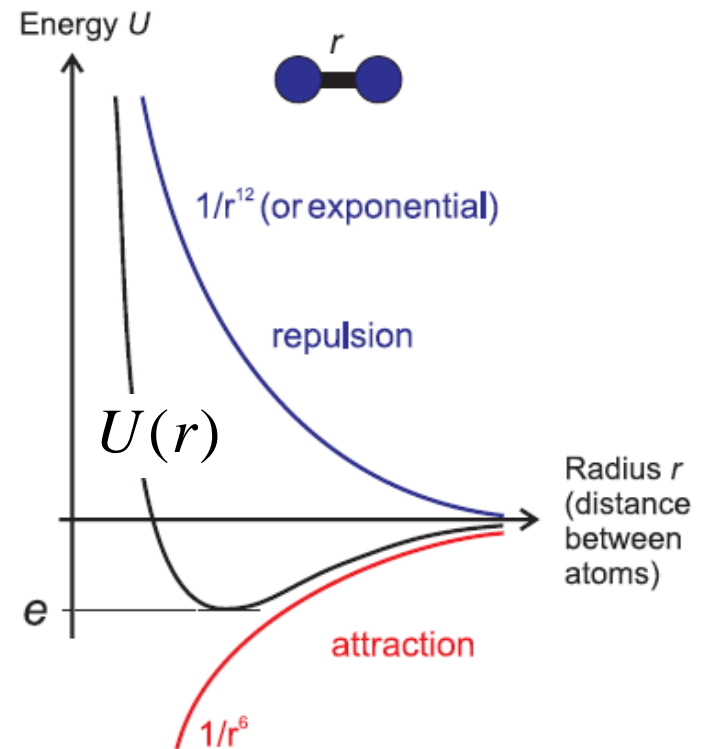
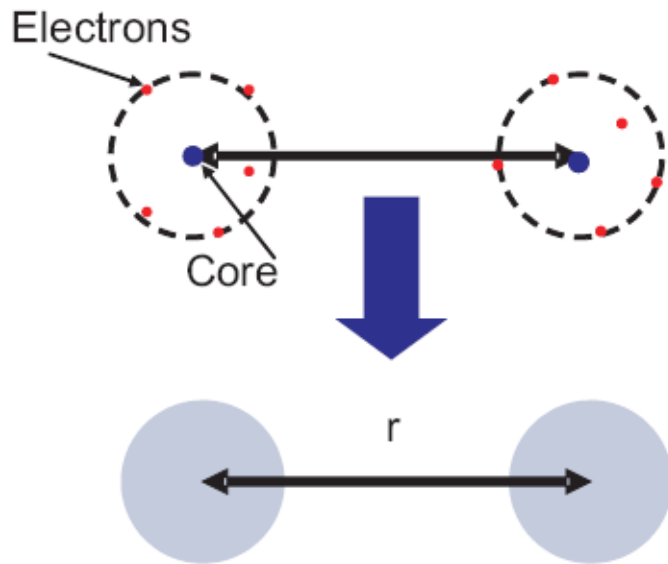
Positions
at $t_0 - \Delta t$

Forces
at t_0

Lecture 3

Preview: Forces on atoms

- Consider energy landscape due to chemical bonds



Attraction: Formation of chemical bond by sharing of electrons

Repulsion: Pauli exclusion (too many electrons in small volume)

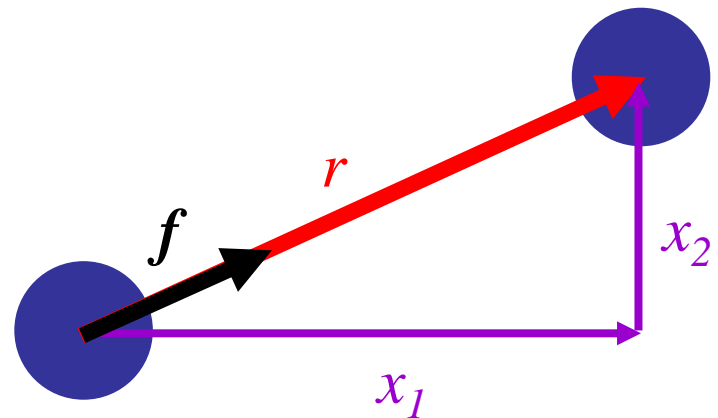
How are forces calculated?

Force magnitude: Derivative of potential energy with respect to atomic distance

$$f = -\frac{dU(r)}{dr}$$

To obtain force vector f_i , take projections into the three axial directions

$$f_i = f \frac{x_i}{r}$$

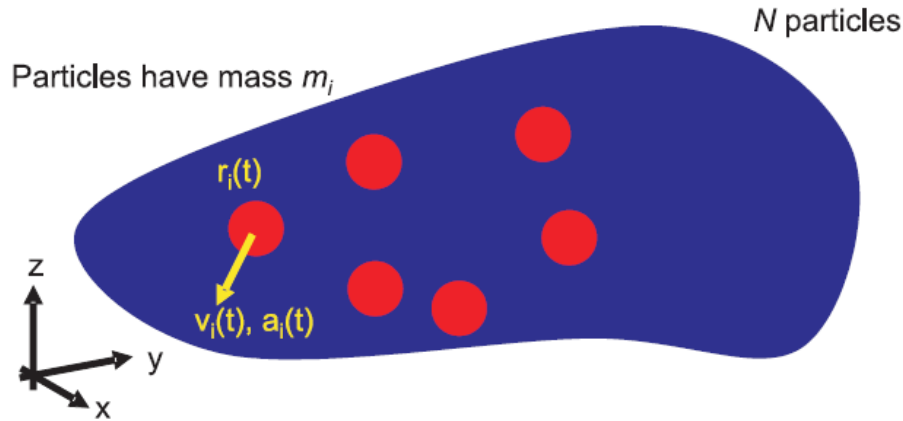


Often: Assume pair-wise interaction between atoms

Note on force calculation

- Forces can be obtained from a variety of models for interatomic energy, e.g.
 - Pair potentials (e.g. LJ, Morse, Buckingham)
 - Multi-body potentials (e.g. EAM, CHARMM, UFF, DREIDING)
 - Reactive potentials (e.g. ReaxFF)
 - Quantum mechanics (e.g. DFT) – **part II**
 - Tight-binding
 - ...
- **...will be discussed in next lectures**

Molecular dynamics



Follow trajectories of atoms

“Verlet central difference method”

$$r_i(t_0 + \Delta t) = \underbrace{2r_i(t_0)}_{\text{Positions at } t_0} - \underbrace{r_i(t_0 - \Delta t)}_{\text{Positions at } t_0 - \Delta t} + \underbrace{a_i(t_0)\Delta t^2}_{\text{Accelerations at } t_0} + \dots$$

Positions
at t_0

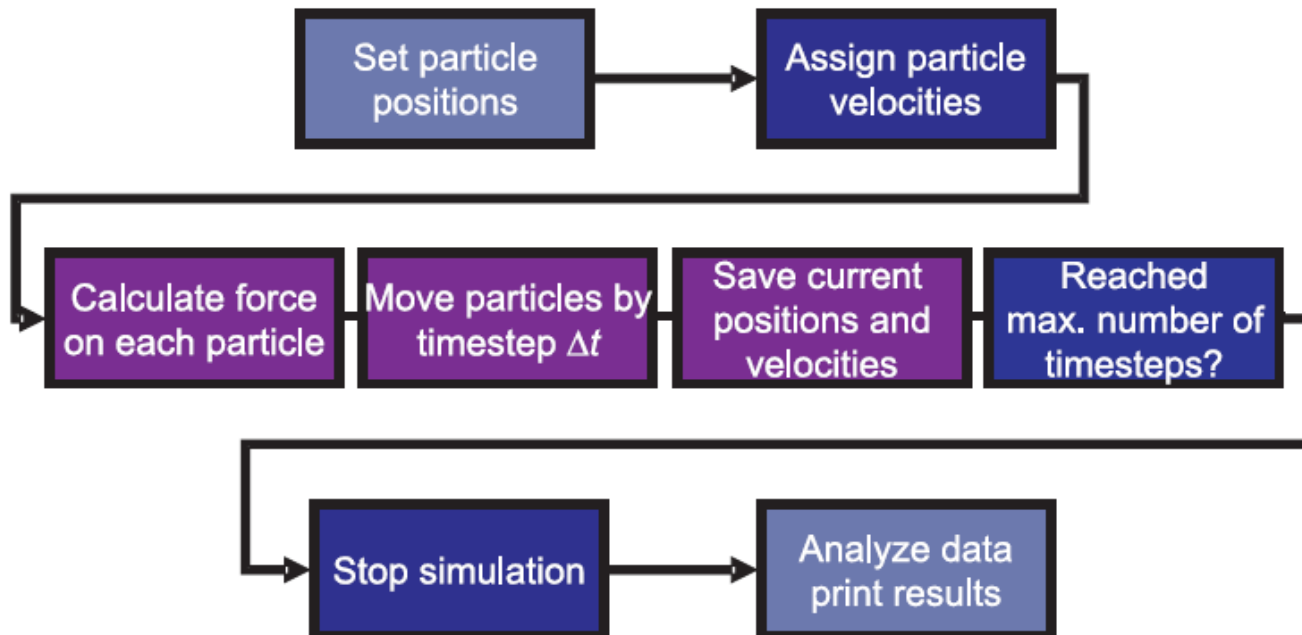
Positions
at $t_0 - \Delta t$

Accelerations
at t_0

$$a_i = f_i / m$$

Summary: Atomistic simulation – numerical approach “molecular dynamics – MD”

- Atomistic model; requires atomistic microstructure and atomic position at beginning
- Step through time by integration scheme
- Repeated force calculation of atomic forces
- Explicit notion of chemical bonds – captured in interatomic potential



Pseudocode

Set particle positions (e.g. crystal lattice)

Assign initial velocities

For (all time steps):

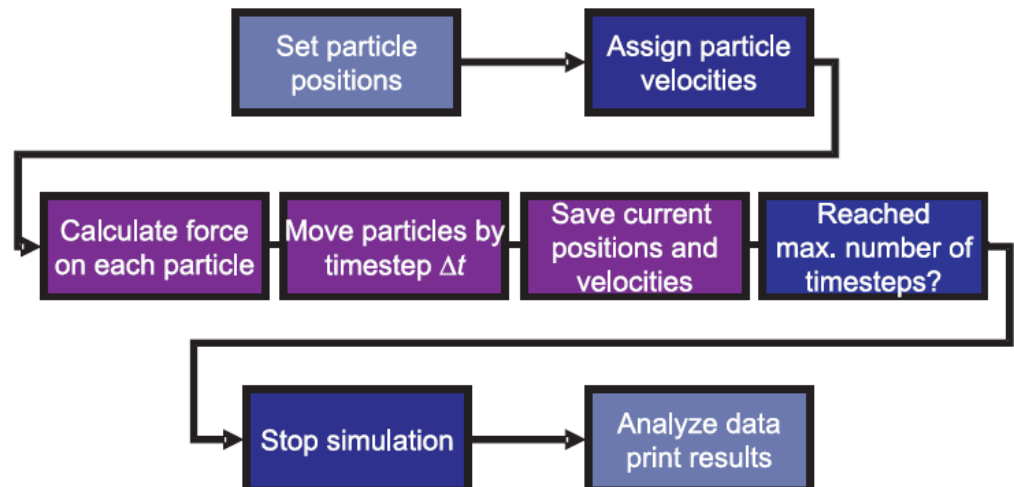
 Calculate force on each particle (subroutine)

 Move particle by time step Δt

 Save particle position, velocity, acceleration

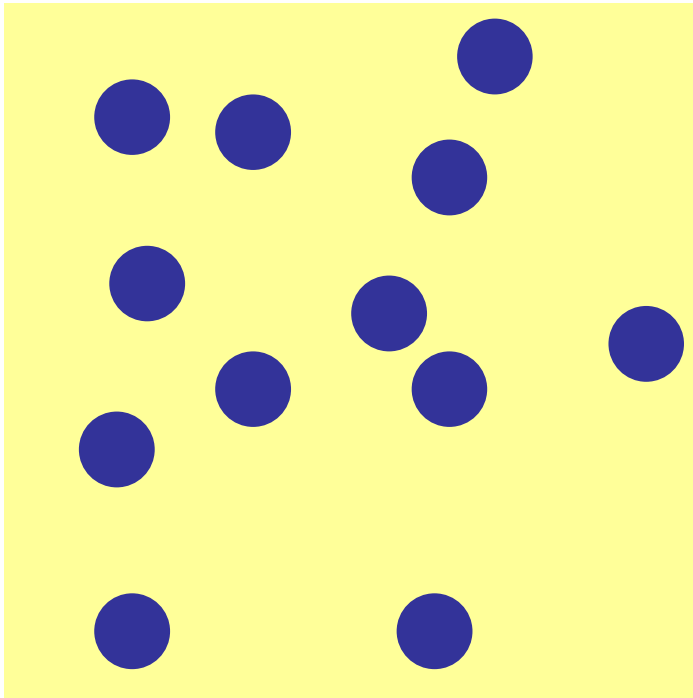
Save results

Stop simulation

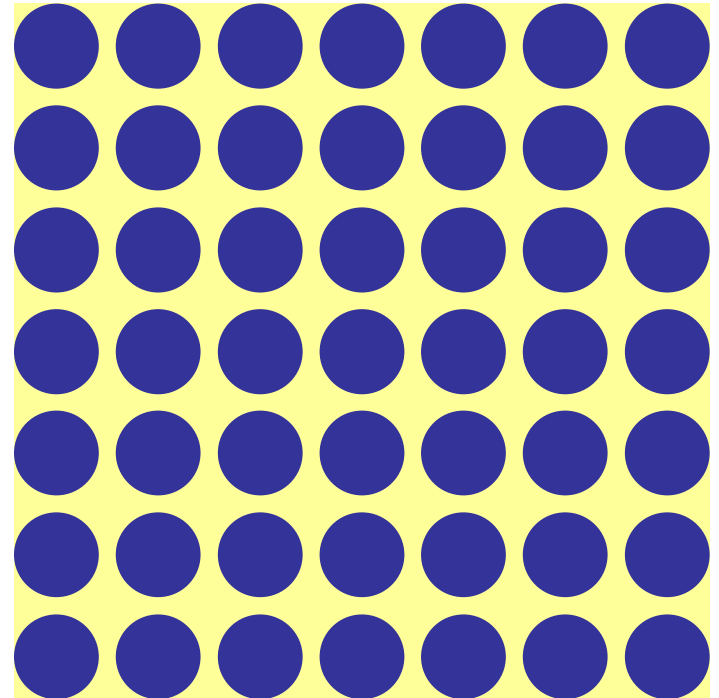


Atomic positions (initial conditions)

- Typically, have cubical cell in which particles are placed in a regular or irregular manner



“gas (liquid)”



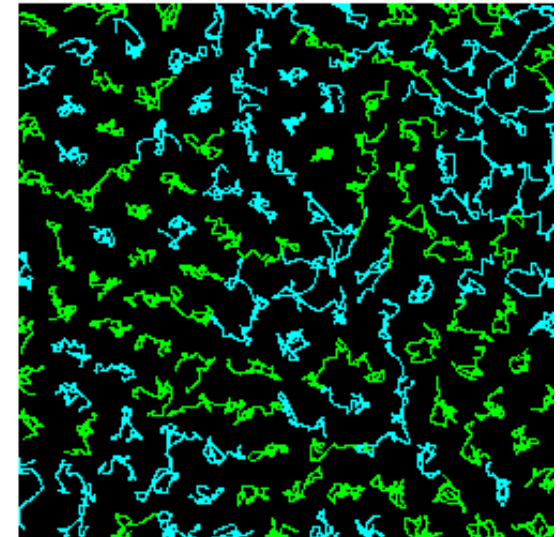
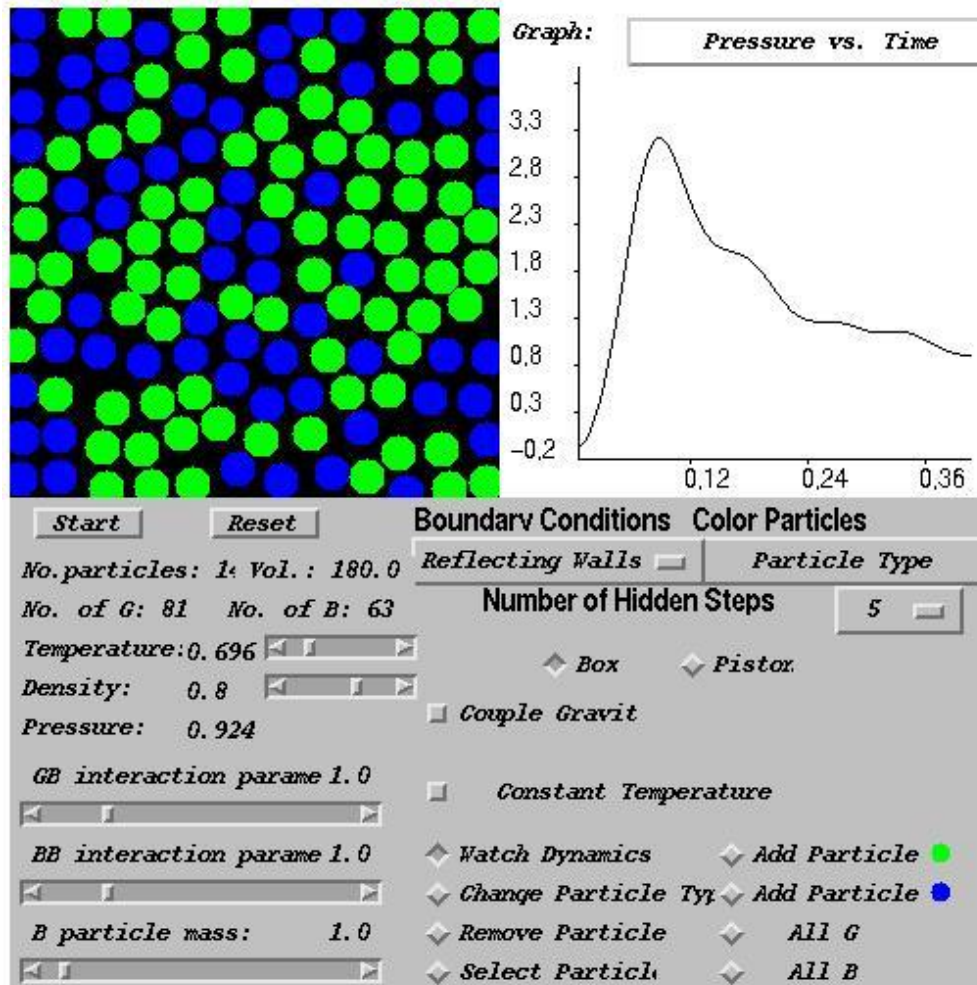
“solid - crystal”

Atomistic description

Back to the application of diffusion problem...

- Atomistic description provides alternative way to predict D
- Simple solve equation of motion
- Follow the trajectory of an atom
- Relate the average distance as function of time from initial point to diffusivity
- **Goal: Calculate how particles move “randomly”, away from initial position**

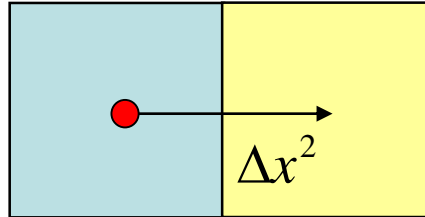
JAVA applet



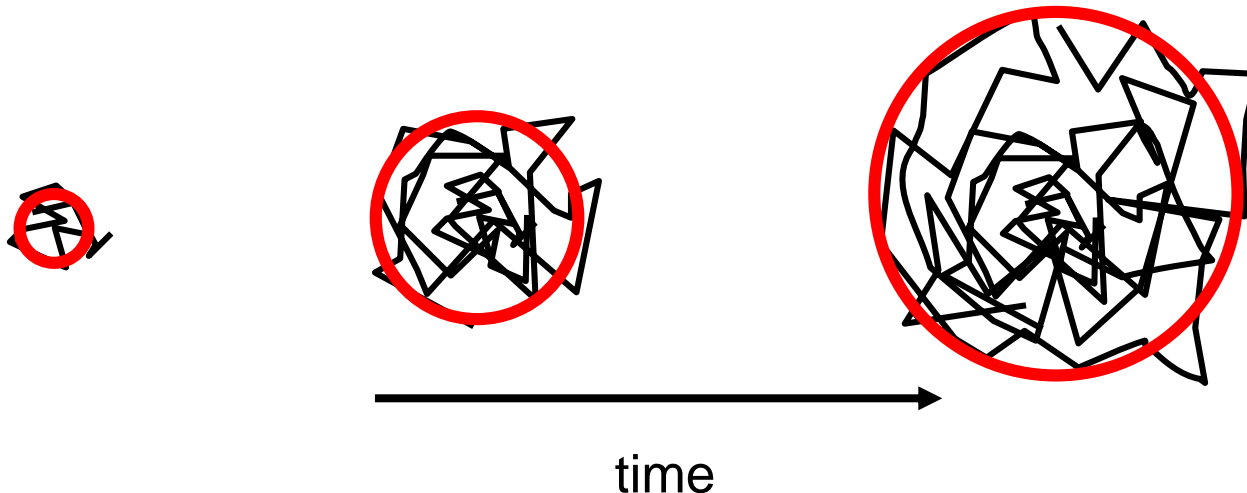
Link atomistic trajectory with diffusion constant (1D)

Diffusion constant relates to the “ability” of a particle to move a distance Δx^2 (from left to right) over a time Δt

$$D = p \frac{\Delta x^2}{\Delta t}$$



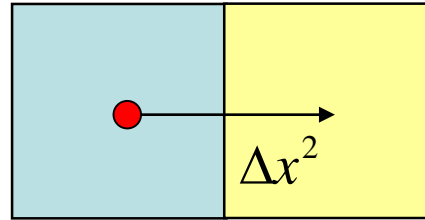
Idea – Use MD simulation to measure square of displacement from initial position of particles, $\Delta r^2(t)$:



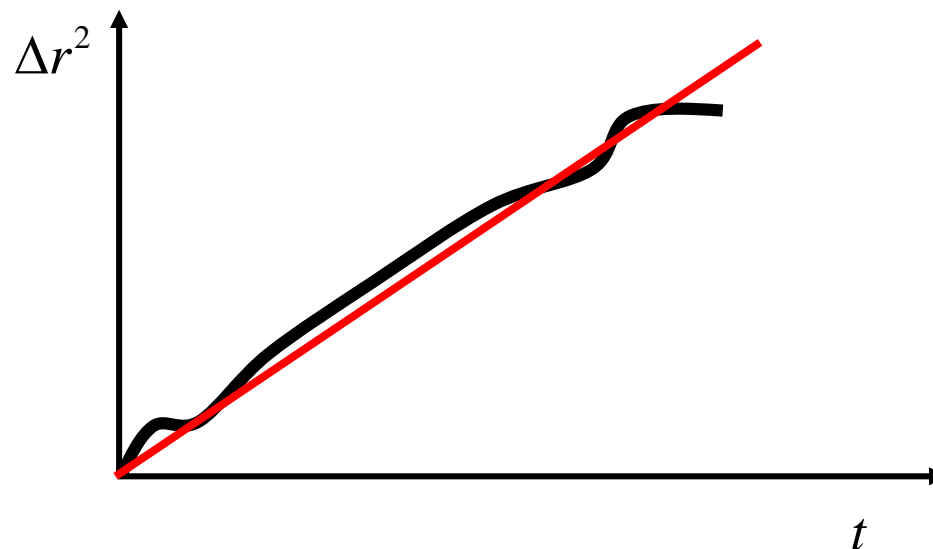
Link atomistic trajectory with diffusion constant (1D)

Diffusion constant relates to the “ability” of a particle to move a distance Δx^2 (from left to right) over a time Δt

$$D = p \frac{\Delta x^2}{\Delta t}$$



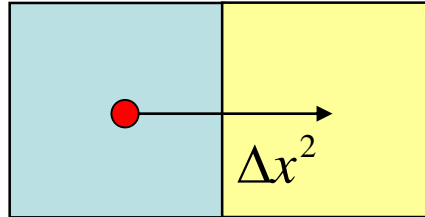
MD simulation: Measure square of displacement from initial position of particles, $\Delta r^2(t)$:



Link atomistic trajectory with diffusion constant (1D)

Diffusion constant relates to the “ability” of a particle to move a distance Δx (from left to right) over a time Δt

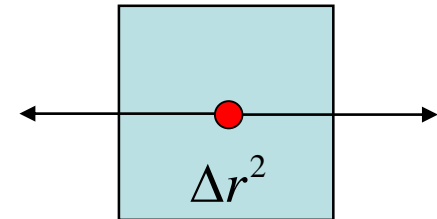
$$D = p \frac{\Delta x^2}{\Delta t}$$



MD simulation: Measure square of displacement from initial position of particles, $\Delta r^2(t)$ and not $\Delta x^2(t)$

Replace

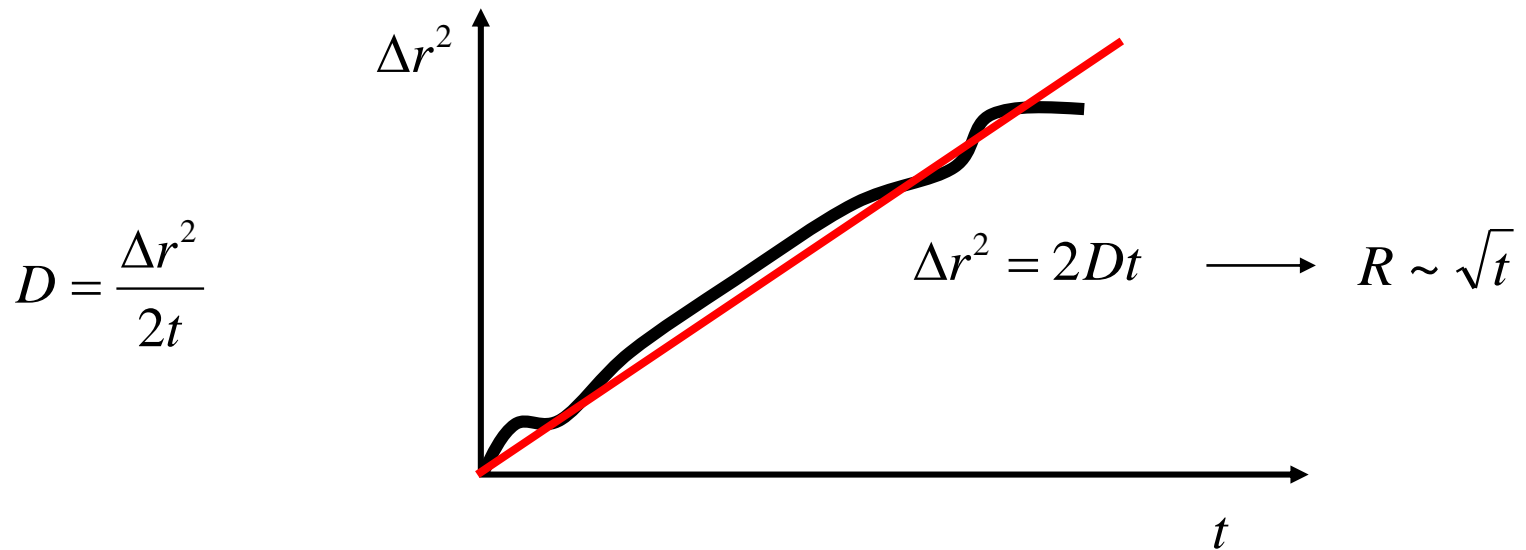
$$D = p \frac{\Delta x^2}{\Delta t} \longrightarrow D = \frac{1}{2} \frac{\Delta r^2}{\Delta t}$$



Factor 1/2 = no directionality in (equal probability to move forth or back)

Link atomistic trajectory with diffusion constant (1D)

MD simulation: Measure square of displacement from initial position of particles, $\Delta r^2(t)$:



Link atomistic trajectory with diffusion constant (2D/3D)

$$D = p \frac{\Delta x^2}{\Delta t}$$



Higher dimensions

$$D = \frac{1}{2} \frac{1}{d} \frac{\Delta r^2}{\Delta t}$$

Factor 1/2 = no directionality in (forth/back)

Factor $d = 1, 2, \text{ or } 3$ due to 1D, 2D, 3D (dimensionality)

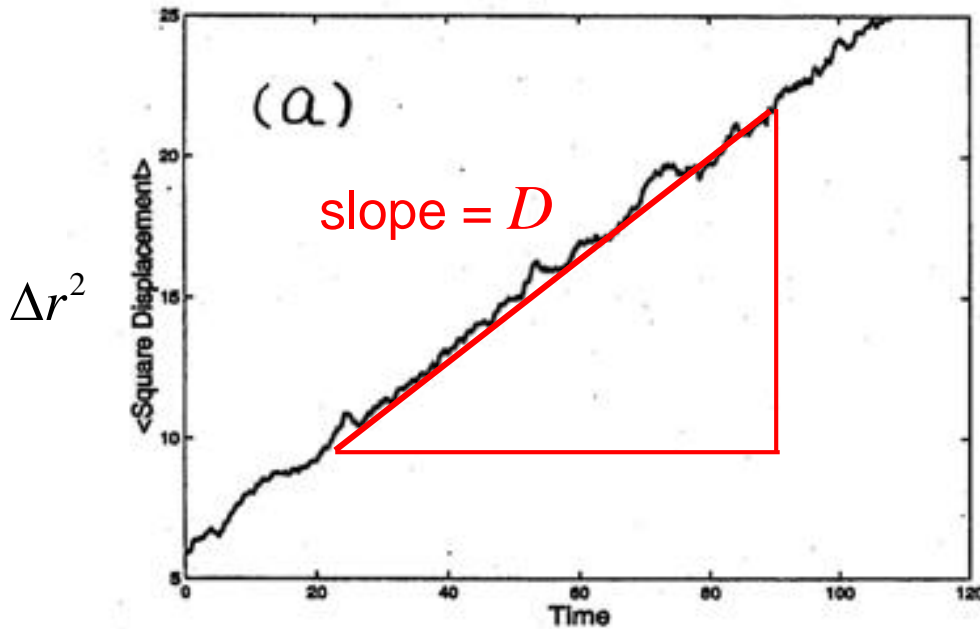
Since:

$$2dD\Delta t \sim \Delta r^2$$

$$2dD\Delta t + C = \Delta r^2$$

$C = \text{constant}$ (does not affect D)

Example: MD simulation



source: S. Yip, lecture notes

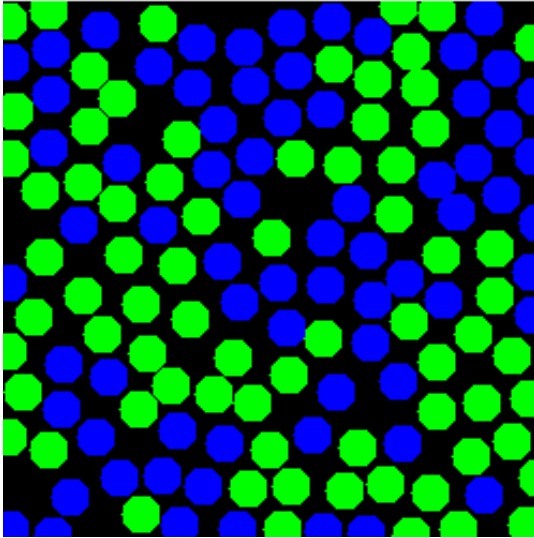
$$D = \frac{1}{2d} \lim_{t \rightarrow \infty} \frac{d}{dt} (\Delta r^2)$$

↑
1D=1, 2D=2, 3D=3

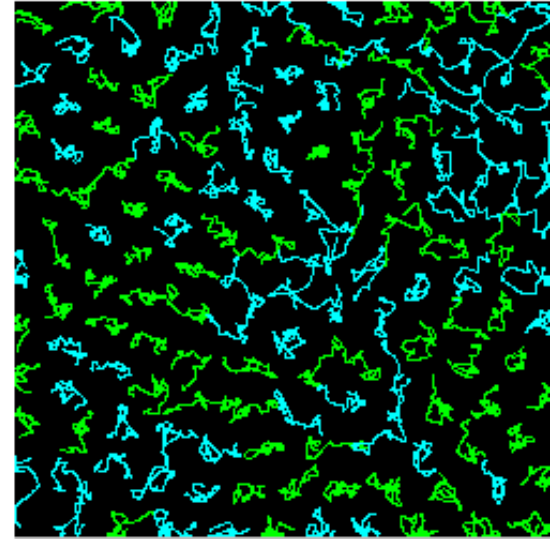
$$D = \frac{1}{2d} \lim_{t \rightarrow \infty} \frac{d}{dt} \langle \Delta r^2 \rangle$$

$\langle .. \rangle$ = average over all particles

Example molecular dynamics



Particles

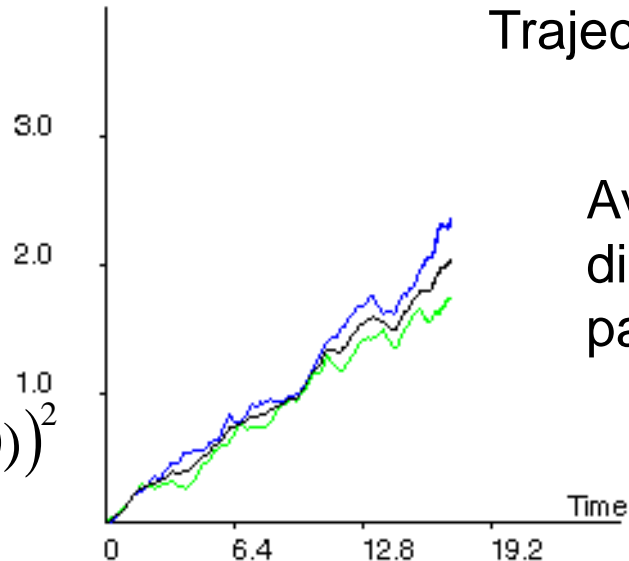


Trajectories

**Mean Square
Displacement function**

$$\langle \Delta r^2(t) \rangle = \frac{1}{N} \sum_i (\mathbf{r}_i(t) - \mathbf{r}_i(t=0))^2$$

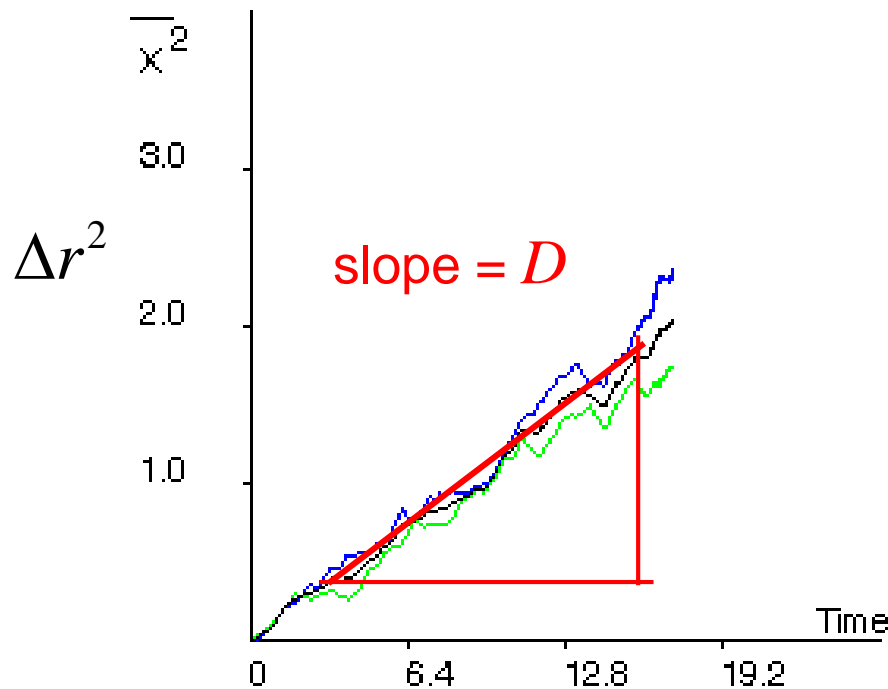
Δr^2



Average square of
displacement of all
particles

Example calculation of diffusion coefficient

$$\langle \Delta r^2(t) \rangle = \frac{1}{N} \sum_i \underbrace{(r_i(t))}_{\text{Position of atom } i \text{ at time } t} - \underbrace{r_i(t=0)}_{\text{Position of atom } i \text{ at time } t=0}$$



$$D = \frac{1}{2d} \lim_{t \rightarrow \infty} \frac{d}{dt} \langle \Delta r^2 \rangle$$

↑
1D=1, 2D=2, 3D=3

Summary: “Multiscale simulation paradigm”

- Molecular dynamics provides a powerful approach to relate the diffusion constant that appears in continuum models to atomistic trajectories
- Outlines multi-scale approach: Feed parameters from atomistic simulations to continuum models

