

EXERCISE 1 — STRUCTURE OPTIMIZATION

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- > At a (local) minimum of a crystal structure:
- The forces on all atoms are 0
- 2. The pressure is equal to the barostat (given) pressure
- ➤ QE has procedures for optimization (relaxation) of atomic positions and cell parameters until the above conditions are achieved

EXERCISE 1 — STRUCTURE OPTIMIZATION

- \triangleright calculation = 'relax' \longrightarrow optimize atomic positions
- \triangleright calculation = 'vc-relax' \rightarrow optimize atomic positions and cell parameters
- > An optimization cycle consists of two steps:
- 1. The forces and pressure are calculated from an SCF calculation
- 2. The atom positions (and cell parameters) are varied using a suitable algorithm (BFGS)
- > The above steps are repeated until the various convergence thresholds are satisfied:
 - Look at etot_conv_thr, forc_conv_thr and press_conv_thr in pw.x documentation

EXERCISE 2 — PHONONS

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- > This exercise consists of two examples:
- 1. Calculating the phonon frequencies at Γ point and determining LO-TO splitting
- 2. Plotting the phonon dispersion and DOS

- \triangleright The basic program for phonon calculations is $\underline{ph.x}$
- > An additional program dynmat.x is used for (e.g.) calculating LO-TO splitting
- \triangleright Programs q2r.x and matdyn.x are used for interatomic force constant calculations

EXERCISE 2 — PHONONS (Γ)

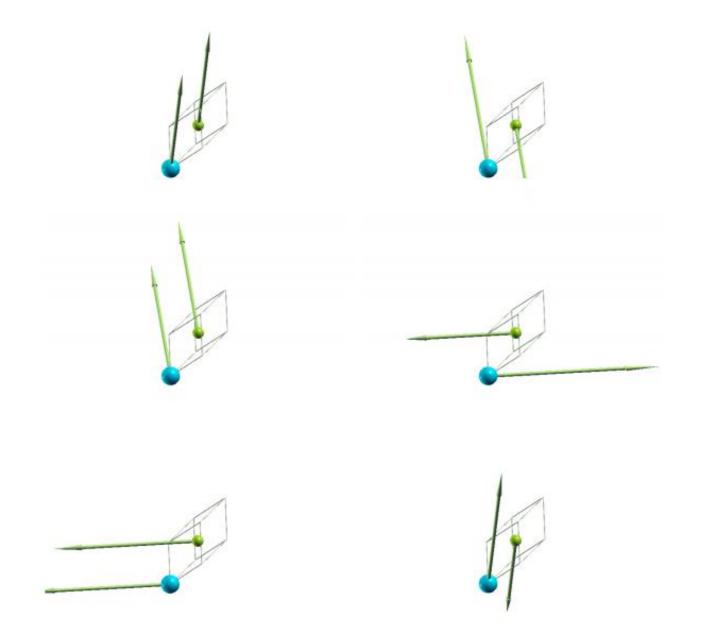
```
-0.196663 [THz] =
                                       -6.559986 [cm-1]
freq (
        2) =
                  -0.196663 [THz] =
                                       -6.559986 [cm-1]
freq (
freq (
                  -0.196663 [THz] =
                                       -6.559986 [cm-1]
        4) = 4.549038 [THz] = 151.739587 [cm-1]
freq (
freq (
        5) = 4.549038 [THz] =
                                      151.739587 [cm-1]
                   4.549038 [THz] =
                                      151.739587 [cm-1]
freq :
```

- \triangleright We have $3*N_{atoms}$ frequencies, one per displacement
- > Small (negative) values of the 3 acoustic modes <u>are numerical artefacts</u> (can be regarded as 0)
- \succ The optical modes are degenerate \longrightarrow should not happen for an ionic crystal such as NaCl

EXERCISE 2 — PHONONS (Γ)

> After adding the non-analytical term to the dynamical matrix and imposing the acoustic sum rule:

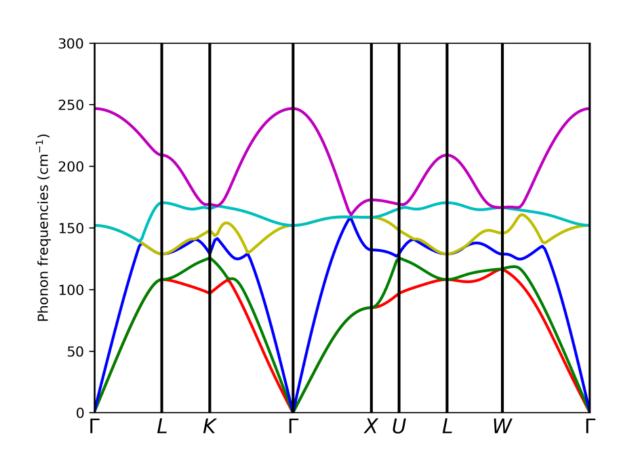
```
-0.0000000 [THz] =
                                         -0.000000 [cm-1]
freq
                                         0.000003 [cm-1]
                    0.0000000 [THz] =
freq (
                    0.000000 [THz] =
                                         0.000003 [cm-1]
freq (
freq (
                   4.555892 [THz] = 151.968191 [cm-1]
freq (
                   4.555892 [THz] = 151.968191 [cm-1]
         6) =
                    7.399029 [THz] =
                                        246.805025 [cm-1]
freq
```

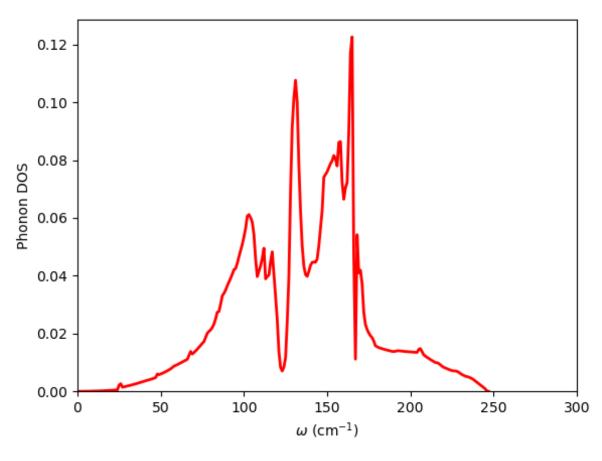


EXERCISE 2 — PHONONS (DISPERSION & DOS)

- Phonon frequencies for any wavenumber can be calculated in two ways:
- 1. Explicitly list a path in 1BZ at the end of ph.x input file and calculate as we did for Γ
- 2. Calculate frequencies for a finite wavenumber grid and extract interatomic force constants (IFC)
- Method 1 is rather slow
- > Once IFCs are known, calculating frequencies at any wavenumber is very fast

EXERCISE 2 — PHONONS (DISPERSION & DOS)





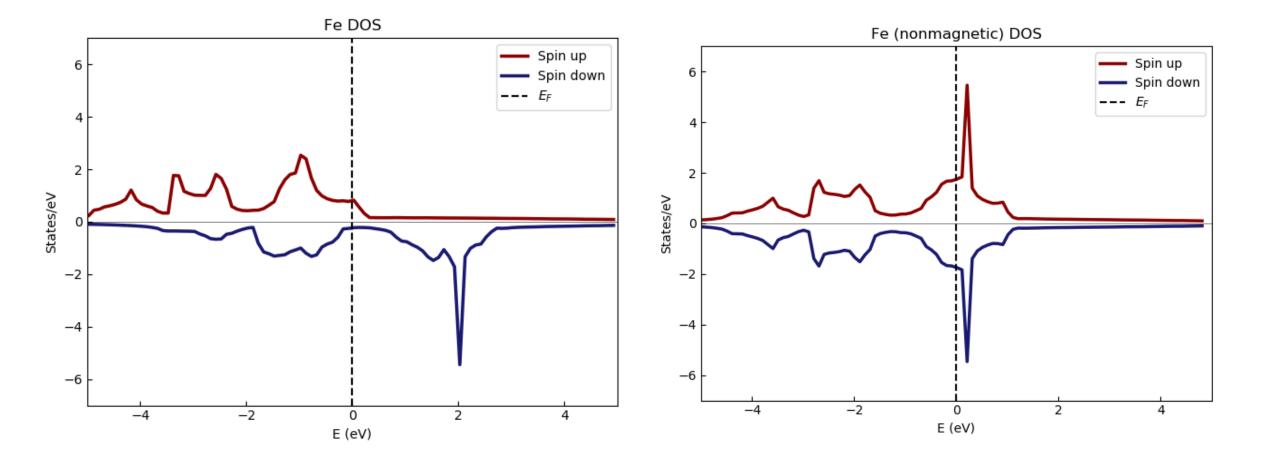
- > This exercise consists of two examples:
- Ferromagnetic system (Fe)
- 2. Antiferromagnetic system (NiO)

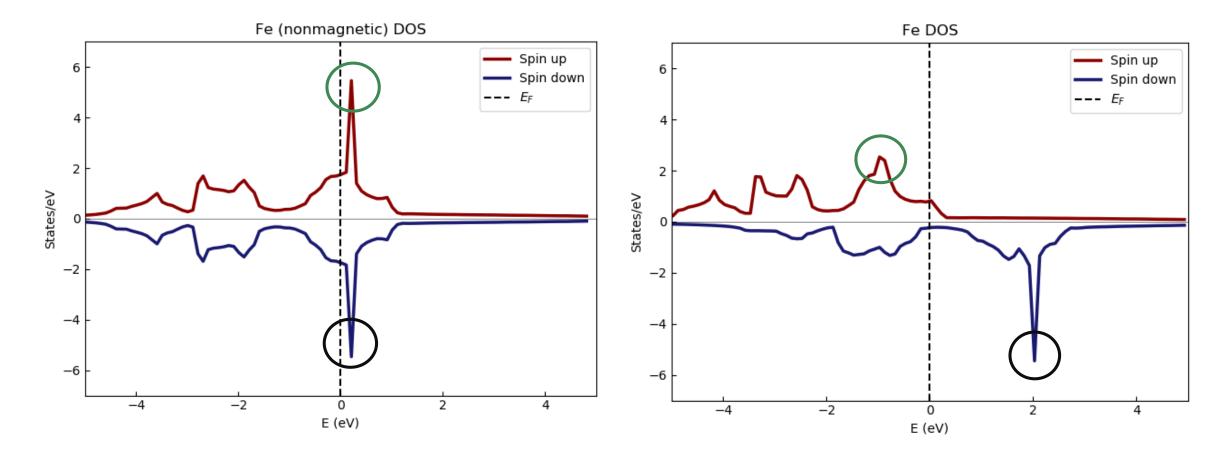
- > The total charge density becomes a sum of densities of subsystems with different spins
- Calculations on magnetic systems take at least twice as long compared to nonmagnetic

- Notice the appearance of two new variables in the SCF input file:
- $7. \quad nspin = tell QE to make a spin polarized calculation$
- 2. $starting_magnetization(i) = starting spin polarization on atomic type i$

- > SCF calculation will determine the final value of the magnetic moment
- > The magnetic moment can also be fixed with the tot_magnetization variable

- > k-points are written twice in the output because they are seperately calculated for the two subsystems
- > The total and absolute magnetization are now written in the output files
- \rightarrow If we put starting_magnetization(1) = 0.0 we always end up with a nonmagnetic state
- > The magnetic state is lower in energy by 0.04 Ry



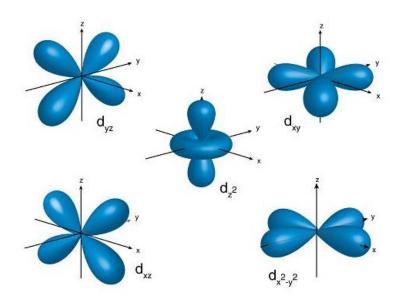


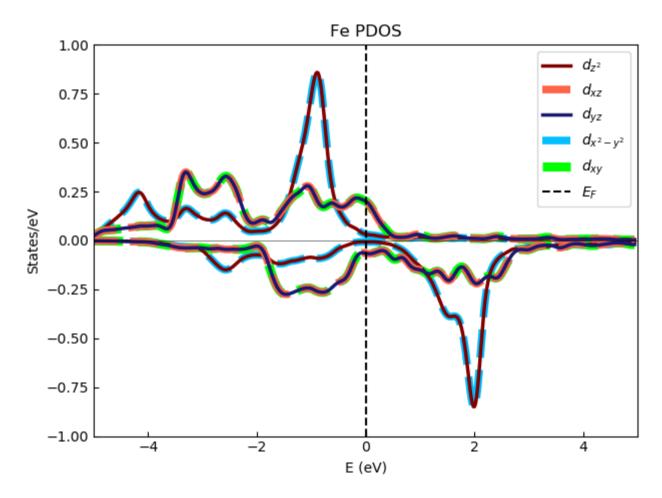
- > Hund rule: some d-orbitals of an isolated Fe atom are occupied only with electrons of the same spin orientation (broken spin degeneracy of the d-orbitals)
- > Let's see how the crystal field (other Fe atoms) affects the d-orbitals hybridization

- > Kohn-Sham wave functions can be projected onto localized atomic orbitals using projwfc.x
- > The orbitals are read from the pseudopotential files

```
Atomic states used for projection
    (read from pseudopotential files):
    state # 1: atom 1 (Fe ), wfc 1 (l=0 m= 1)
    state # 2: atom 1 (Fe ), wfc 2 (l=1 m= 1)
    state # 3: atom 1 (Fe ), wfc 2 (l=1 m= 2)
    state # 4: atom 1 (Fe ), wfc 2 (l=1 m= 3)
    state # 5: atom 1 (Fe ), wfc 3 (l=2 m= 1)
    state # 6: atom 1 (Fe ), wfc 3 (l=2 m= 2)
    state # 7: atom 1 (Fe ), wfc 3 (l=2 m= 3)
    state # 8: atom 1 (Fe ), wfc 3 (l=2 m= 4)
    state # 9: atom 1 (Fe ), wfc 3 (l=2 m= 5)
   state # 10: atom 1 (Fe ), wfc 4 (l=0 m= 1)
   state # 11: atom 1 (Fe ), wfc 5 (l=1 m= 1)
   state # 12: atom 1 (Fe ), wfc 5 (l=1 m= 2)
   state # 13: atom 1 (Fe ), wfc 5 (l=1 m= 3)
==== e( 1) = -70.05016 eV ====
   psi = 0.954*[# 1]+0.046*[# 10]
   |psi|^2 = 1.000
==== e( 2) = -35.84389 eV ====
   psi = 0.009*[# 2]+0.009*[# 3]+0.009*[# 4]
   |psi|^2 = 0.026
==== e( 3) = -35.84389 eV ====
   psi = 0.641*[# 2]+0.641*[# 3]+0.641*[# 4]+0.011*[# 11]+0.011*[# 12]
        +0.011*[# 13]
   |psi|^2 = 1.954
==== e( 4) = -35.84389 eV ====
   psi = 0.334*[# 2]+0.334*[# 3]+0.334*[# 4]+0.006*[# 11]+0.006*[# 12]
        +0.006*[# 13]
   |psi|^2 = 1.019
==== e( 5) = 9.60552 eV ====
   psi = 0.954*[# 10]+0.046*[# 1]
   |psi|^2 = 1.000
==== e( 6) = 15.75113 eV ====
   psi = 0.331*[# 6]+0.331*[# 7]+0.331*[# 9]
   |psi|^2 = 0.994
==== e( 7) = 15.75113 eV ====
   psi = 0.331*[# 6]+0.331*[# 7]+0.331*[# 9]
   |psi|^2 = 0.994
```

- \triangleright We are interested in d-orbitals (l=2)
- In PDOS/Fe.pdos_atm#1 (Fe)_wfc#3(d) the projected density of states (PDOS(E)) is written for spin-up and spin-down components in the order:
- 1. d_{z^2}
- d_{xz}
- d_{yz}
- 4. $d_{x^2-y^2}$
- d_{xy}



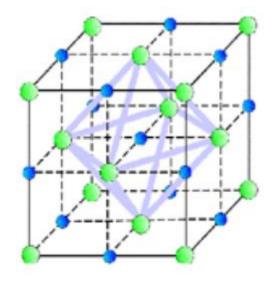


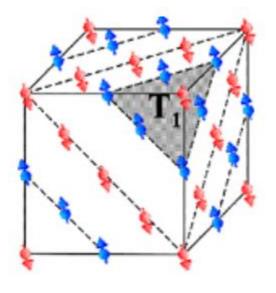
 $ightarrow d_{z^2}$ and $d_{x^2-y^2}$ are weakly hybridized (sharp peaks) so the Hund rules break spin degeneracy

The other orbitals are strongly hybridized (wide peaks) because they point towards the other Fe atoms at the vertices of the BCC cube

- \triangleright **Note:** The calculations in this exercise (ex2.afm) take quite some time
- > The results are therefore already prepared in the reference directory
- Following the instructions in *README.md*, you can copy those results and take a look at the input/output files

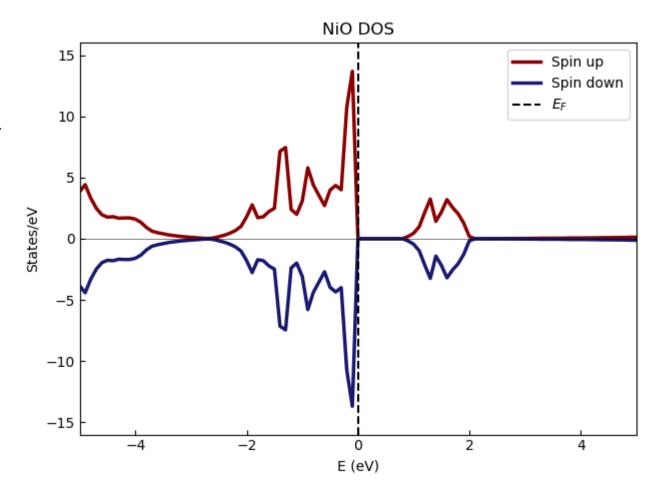
- The Ni ions of the NiO crystal have alternating magnetic moment
- We have to treat them as different atom types with the same pseudopotential but different starting_magnetization





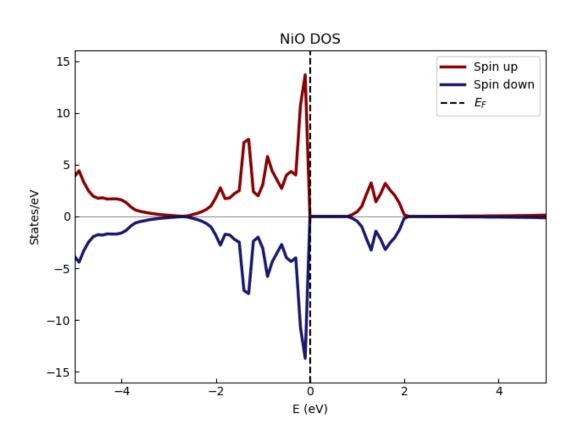
```
Magnetic moment per site (integrated on atomic sphere of radius R) atom 1 (R=0.206) charge= 5.6606 magn= -0.0000 atom 2 (R=0.206) charge= 5.6606 magn= -0.0000 atom 3 (R=0.206) charge= 15.6783 magn= 1.3342 atom 4 (R=0.206) charge= 15.6783 magn= -1.3342
```

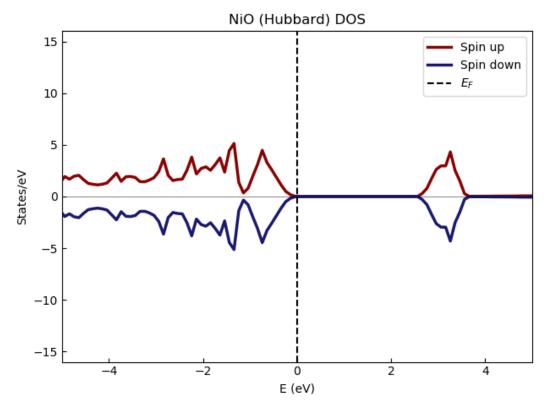
- > The gap is about ~0.84 eV
- > The experimental gap is 3.6-4 eV
- > A significant discrepancy!

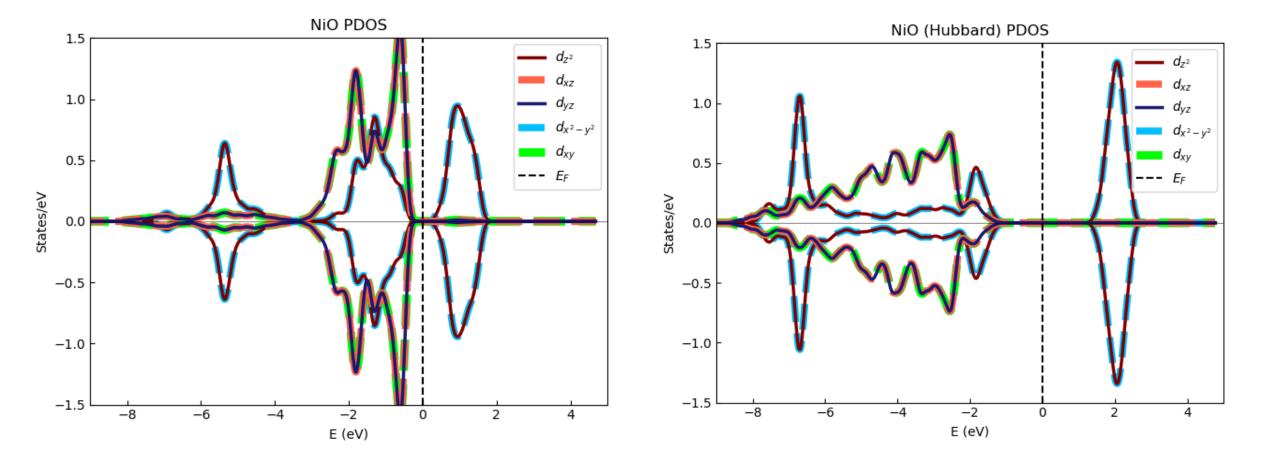


- > The problem arises because PBE (GGA/LDA in general) underestimates band gaps
- Ways to deal with insulators/semiconductors:
- Band insulators (gap due to crystal potential) <u>GW approximation (YAMBO)</u>
- 2. Strongly correlated systems (Mott/charge-transfer insulators) **DFT+U**
- > A term representing the repulsion between electrons of different spins is added to the functional
- DFT+U favours integer orbital occupancy (localization)

- > DFT+U is included in the SCF calculation by the Ida_plus_u and Hubbard_U(i) variables
- Hubbard correction is applied to the states defined in Modules/set_hubbard_n.f90 and Modules/set_hubbard_1.f90
- The values of Hubbard_U(i) need to be separately computed:
- 1. Physical Review B 71.3 (2005): 035105.
- 2. Physical Review B 84.11 (2011): 115108

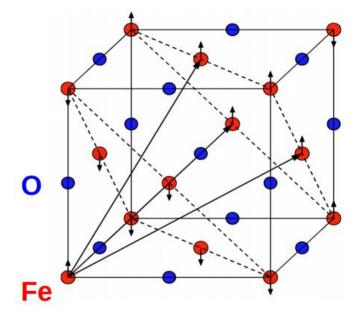


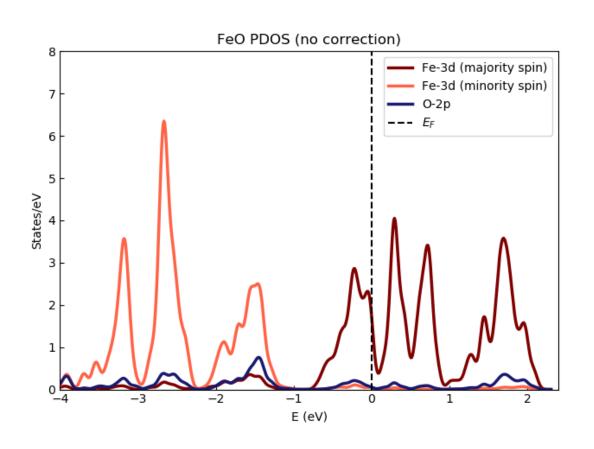


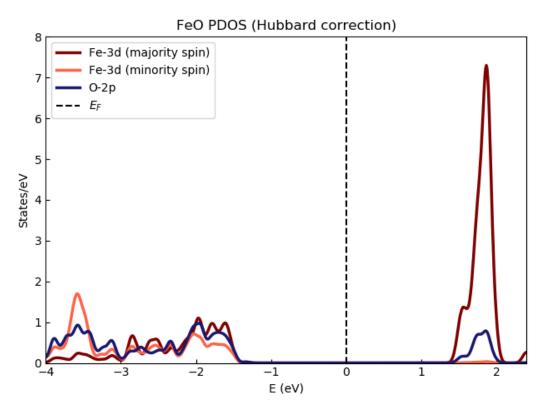


> The Hubbard correction changes the order of the d-orbitals and increases the gap

- > An even more drastic example of DFT gap underestimation is FeO
- > The system is in reality insulating, but DFT predicts it to be a metal



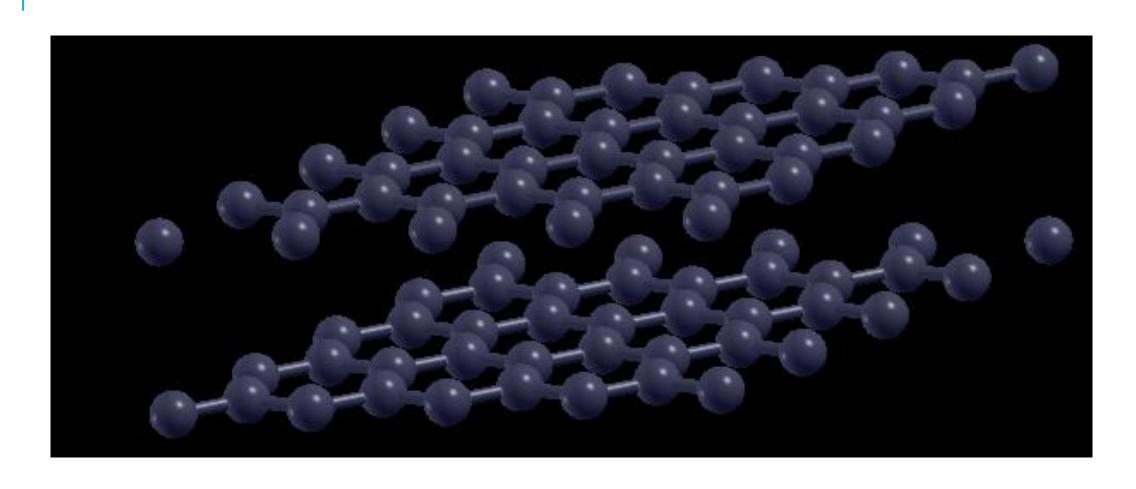




- > In older QE versions, a metallic state would be obtained even with DFT+U
- > The DFT+U SCF would converge to a local minimum of state occupations
- In these cases, it's necessary to force starting occupations with the starting_ns_eigenvalue variable

- > The input_dft variable controls the XC functional used in (N)SCF calculations
- > The functional specified in pseudpotential files is used by default
- Available functionals are listed in Modules/funct.f90
- In this exercise we will use the nonlocal 'vdw-df-cx' functional

- We use the pseudopotential for the closest GGA (PBE)
- There are no PPs for non-local functionals



> The optimized interlayer distances are:

- 1. PBE = 3.55 Å
- 2. vdw-df-cx = 3.22 Å
- 3. LDA = 3.19 Å
- > The experimental interlayer distance is 3.336 Å

