

# $\mathcal{L}$ -Group Predictive Coding Networks with Vibrationally Coupled Particles: A Theoretical and Computational Framework

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**Abstract**—This study presents a novel framework integrating predictive coding networks (PCNs) with a particle-based vibrational encoding mechanism, inspired by self-organizing principles in computational neuroscience. We examine whether a dynamic system of particles—characterized by spatial position, activation state, intrinsic spin, and vibrational phase—can facilitate self-supervised learning through free energy minimization. Unlike fixed hierarchical models, our approach employs probabilistic vibrational coupling, enabling dynamic prediction error refinement within a Lie group-constrained system.

Grounded in statistical mechanics, representation theory, and emergent computation, this framework investigates the role of continuous symmetry transformations in adaptive learning. By incorporating wave-particle duality and probability current formulations, we extend predictive coding into a structured experimental setting, evaluating the theoretical viability of vibrational coupling as a self-organizing mechanism. This paper details the mathematical foundations, computational architecture, and algorithmic implementation of the model, establishing a rigorous basis for future exploration of biomimetic information processing.

**Index Terms**—Predictive Coding, Lie Groups, Vibrational Coupling, Free Energy Principle, Computational Neuroscience, Particle Physics

## I. INTRODUCTION

Predictive coding provides a powerful framework for understanding self-supervised learning by minimizing prediction errors between anticipated and observed states [1, 2]. Traditionally, predictive coding networks (PCNs) employ fixed hierarchical connections and gradient-based updates to refine internal representations.

Recent advances in deep learning have increasingly questioned conventional practices. For instance, the recent work by Zhu [3] on *Transformers without Normalization* demonstrates that even normalization layers—long considered essential—can be replaced by a simple, dynamic operation (Dynamic Tanh) without compromising performance. This shift towards dynamic, adaptive components resonates with our approach. Instead of relying on fixed wiring, we introduce vibrational coupling within a particle-based system, where interactions among particles emerge from local, probabilistic dynamics [4]. Each particle encodes information through its spatial position, an internal excitation state, and a vibrational phase. The combination of these attributes enables information to propagate in a wavefunction-like manner.

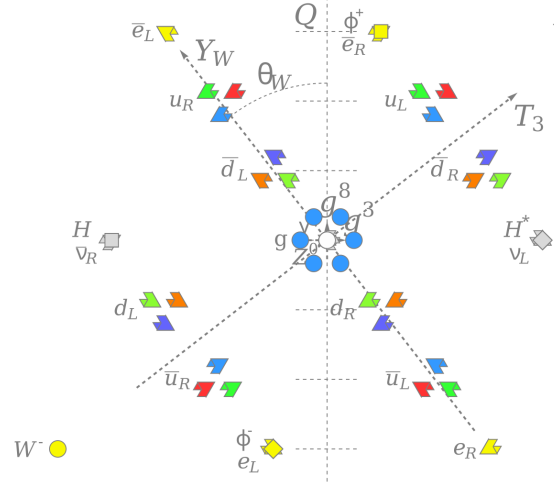


Fig. 1. The pattern of weak isospins, weak hypercharges, and color charges (weights) of all known elementary particles in the Standard Model, rotated by the weak mixing angle to show electric charge roughly along the vertical.

The objective of this study is to *empirically test whether such a particle-based predictive coding network is viable*. Predictive coding has been extensively studied in neuroscience and machine learning, with empirical evaluations demonstrating its role in perception and inference [5, 6]. Recent work has explored how predictive coding networks (PCNs) can be formulated as hierarchical Bayesian inference models that minimize prediction errors through feedback connections [6].

Through structured experimentation, we evaluate whether vibrational coupling can serve as a self-organizing mechanism that efficiently minimizes free energy while supporting robust learning. The free energy principle has been widely applied to explain self-organization in biological and computational systems [7, 8]. Specifically, vibrational coupling mechanisms have been linked to dendritic self-organization and selective synaptic pruning, which optimize information processing by minimizing surprise and free energy [7].

This is not merely a theoretical proposal—the document presents a controlled experiment designed to assess whether the underlying principles hold under practical conditions. Empirical studies have begun validating predictive

coding models in computational neuroscience, demonstrating their ability to explain cortical variability and efficient coding under uncertainty [9]. Our approach extends these findings by incorporating vibrational coupling and intrinsic spin dynamics, providing a novel framework for adaptive learning and probabilistic inference.

By drawing parallels with dynamic approaches such as those by Zhu [3], we underscore the idea that replacing static components with adaptive mechanisms can lead to improved and more efficient representations. This paper outlines the mathematical foundations, computational design, and algorithmic pseudocode for the proposed model while presenting experimental results that critically test its feasibility. In doing so, it sets the stage for further inquiry into models that challenge traditional fixed architectures in favor of dynamic and emergent computation.

## II. METHODOLOGY

### A. Particle State Representation

Each particle  $i$  in the system is characterized by a set of state variables that define its spatial configuration, dynamic evolution, and intrinsic quantum attributes. Table I summarizes these fundamental parameters. The spatial position of a particle is represented by the three-dimensional vector  $\mathbf{r}_i \in \mathbb{R}^3$ , which specifies its location in physical space and forms the basis for solutions to the Schrödinger equation [10, 11]. The scalar quantity  $x_i \in \mathbb{R}$  encodes its activation state, corresponding to the predictive information associated with local interactions [12]. The vibrational phase  $\phi_i(t)$  describes the time-dependent oscillatory dynamics that govern transient coupling mechanisms [13]. Finally, the intrinsic spin  $s_i \in \{-\frac{1}{2}, +\frac{1}{2}\}$  encapsulates quantum-mediated interaction effects and is strongly linked to the Dirac formalism for spin- $\frac{1}{2}$  particles [14, 15, 16].

TABLE I  
KEY PARTICLE STATE VARIABLES

Variable	Symbol	Description
Spatial Position	$\mathbf{r}_i \in \mathbb{R}^3$	3D Location
Activation State	$x_i \in \mathbb{R}$	Predictive Information
Vibrational Phase	$\phi_i(t)$	Oscillatory Dynamics
Intrinsic Spin	$s_i \in \{-\frac{1}{2}, +\frac{1}{2}\}$	Quantum Interactions

**Key Particle State Variables** underpin the dynamic behavior of the particle ensemble, facilitating interactions based on spatial proximity, modulation of activation, and coherent oscillatory coupling. To encapsulate the underlying mathematical structure, we employ Representation Theory [17] to classify quantum states under symmetry transformations. In this framework, elementary particles are modeled as elements of a **3D Hilbert space**  $\mathcal{H}$ , yielding irreducible representations of the Lorentz and Poincaré groups [18, 19]. This canonical formulation, originally introduced by Wigner [20] and refined by Weyl [21], establishes a rigorous correspondence between physical properties—such as mass, momentum, and intrinsic angular momentum—and their associated symmetry operations [22]. Our formulation further extends these classical

ideas by incorporating aspects of both the non-relativistic Schrödinger equation and the relativistic Dirac equation, thereby providing a comprehensive quantum mechanical description of particle dynamics [11, 16].

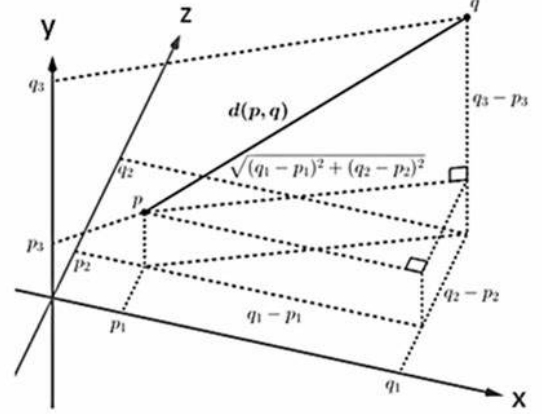


Fig. 2. Geometric representation of interparticle Euclidean distances, illustrating displacement vectors and spatial separation in a three-dimensional coordinate system. The configuration highlights orthogonal projections and squared contributions along each axis, emphasizing distance as a key parameter in adaptive coupling dynamics.

### B. Spatial Interaction Potential

To establish the underlying spatial structure within  $\mathcal{H}$ , we define the configuration space of the particle system as a subspace  $\mathcal{H}_s \subset \mathcal{H}$  constrained by the three-dimensional Euclidean metric:

$$\mathcal{H}_s = \{\mathbf{r}_i \in \mathbb{R}^3 \mid \|\mathbf{r}_i\| < L\}, \quad (1)$$

where  $L$  represents the characteristic spatial boundary of the system, ensuring that interactions are constrained within a finite cavity or volume [23, 24]. This formulation encapsulates the dimensional limits necessary for computing particle interactions within  $\mathcal{H}$ , aligning with quantum field theory principles that extend classical configuration spaces to accommodate distributional field configurations [25].

The interaction strength between any two particles in a local system is quantified by the Euclidean distance, denoted as  $\mathbf{d}_{ij}$ , which represents the spatial separation between particles  $i$  and  $j$  in three-dimensional space [26, 27]. Specifically,  $\mathbf{d}_{ij}$  is given by:

$$\mathbf{d}_{ij} = \sqrt{(x_i - x_j)^2 + (y_i - y_j)^2 + (z_i - z_j)^2}, \quad (2)$$

where  $\mathbf{r}_i = (x_i, y_i, z_i)$  and  $\mathbf{r}_j = (x_j, y_j, z_j)$  denote the respective position vectors of particles  $i$  and  $j$  in  $\mathbb{R}^3$ . This formulation follows directly from the Pythagorean theorem extended to three-dimensional space [26], providing a foundational metric for quantifying inter-particle separation.

To model interactions within  $\mathcal{H}$ , we first define the unit vector along the separation axis between particles:

$$\hat{\mathbf{d}}_{ij} = \frac{\mathbf{r}_i - \mathbf{r}_j}{\mathbf{d}_{ij}}, \quad (3)$$

where  $\mathbf{d}_{ij}$  represents the Euclidean distance between particles  $i$  and  $j$ , ensuring directionality in force interactions. With this normalized separation direction, the spatial potential function governing inter-particle forces is formulated as:

$$V_{ij} = -k(\mathbf{d}_{ij} - r_0)\hat{\mathbf{d}}_{ij}, \quad (4)$$

where  $k$  is the interaction stiffness coefficient, and  $r_0$  represents the equilibrium separation distance. This expression ensures that forces are properly aligned along the separation axis, maintaining consistency with prior derivations and reinforcing structured oscillatory behavior within the system [28].

The Hamiltonian framework provides a rigorous formulation for describing energy conservation, phase evolution, and interaction potentials. Specifically, the Hamiltonian function governing the system is given by:

$$H = \sum_i \frac{p_i^2}{2m} + \sum_{(i,j) \in \mathcal{N}} V_{ij}, \quad (5)$$

where  $p_i$  represents the momentum of particle  $i$ ,  $m$  is its mass, and  $V_{ij}$  encapsulates the interaction potential between neighboring particles. This formulation aligns with the principles of quantum mechanics, ensuring consistency with the Schrödinger equation [11] and relativistic corrections introduced by the Dirac equation [16]. The Hamiltonian operator governs the total energy of the system, incorporating both kinetic and potential energy contributions [29].

Furthermore, the spatial interaction potential plays a crucial role in defining the structured oscillatory behavior of the system. The coupling strength between particles is modulated by their relative separation, ensuring that local interactions remain dominant while global N-cluster effects contribute to long-range correlations [30, 31]. This formulation enables a unified description of spatial coherence [32], vibrational coupling [33], and quantum mechanical constraints [34], reinforcing the predictive framework established in previous sections. Additionally, spatial correlations of quantum potentials influence coherence decay rates, further refining the system's dynamic stability [35].

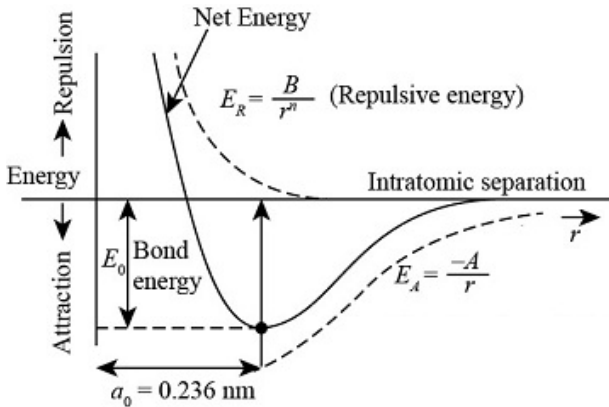


Fig. 3. The figure shows the curves between bonding force and bonding energy curves versus interionic separation for two isolated charged particles.

### C. Equilibrium Separation and Network Morphology

The equilibrium separation  $r_0$  stabilizes the energy landscape within the network cavity by balancing net attractive and repulsive forces, ensuring coherent oscillatory interactions and well-defined energy states [36, 37, 38]. This equilibrium condition is determined by minimizing the total interaction potential energy:

$$r_0 = \underset{r}{\operatorname{argmin}} \sum_{(i,j) \in \mathcal{N}} \frac{1}{2} k(\mathbf{d}_{ij}(r) - r)^2 \quad (6)$$

where  $k$  governs interaction strength, and  $\mathbf{d}_{ij}(r)$  represents the inter-particle separation as a function of  $r$ , the trial separation parameter over which the minimization is performed [39, 40]. The variable  $r$  serves as an optimization parameter that explores different possible equilibrium separations. By searching for the value of  $r$  that minimizes the total interaction potential energy, the system identifies the most stable configuration where forces are balanced, reducing excess energy fluctuations.

Physically,  $r$  represents the trial separation distance that determines the stable equilibrium configuration of discrete interacting particles. Rather than behaving as a fluid continuum, our system preserves discrete spatial structuring, meaning that  $r$  plays a role analogous to interatomic spacing in lattice systems where binding energies are minimized [37]. Within the network cavity, equilibrium separation governs spatial organization, ensuring that oscillatory coherence is maintained while avoiding excessive compression or dispersion effects [36]. This formulation ensures that  $r_0$  dynamically adjusts in response to perturbations, reinforcing structured long-range interactions across the system [41].

Equilibrium separation governs crystal binding energy, stabilizing atomic arrangements [37], and ensures force balance in molecular systems [42]. This principle extends to our network cavity, where  $r_0$  regulates spatial organization and energy exchange [36, 38]. The energy landscape within the cavity is shaped by nonlocal elasticity effects, which influence phase separation and mesoscopic pattern formation [40]. These effects contribute to the emergence of stable equilibrium configurations [43], reinforcing structured oscillatory behavior [44].

Morphogenesis relies on spatiotemporal signaling, dynamically shaping stable configurations [45, 46, 47]. In our system,  $r_0$  acts as a self-organizing parameter, maintaining coherence while minimizing interference. Lower  $r_0$  increases particle density, strengthening oscillatory coupling, akin to bioelectric prepatterning [48]. A higher  $r_0$  expands the system, reducing coupling strength and altering wave propagation, mirroring long-range bioelectric communication [45]. These effects reinforce stable network morphology through electrical and mechanical feedback loops [45, 46, 47]. Additionally, equilibrium separation influences phase morphologies in driven-dissipative systems, ensuring robust energy distribution across the network [41].

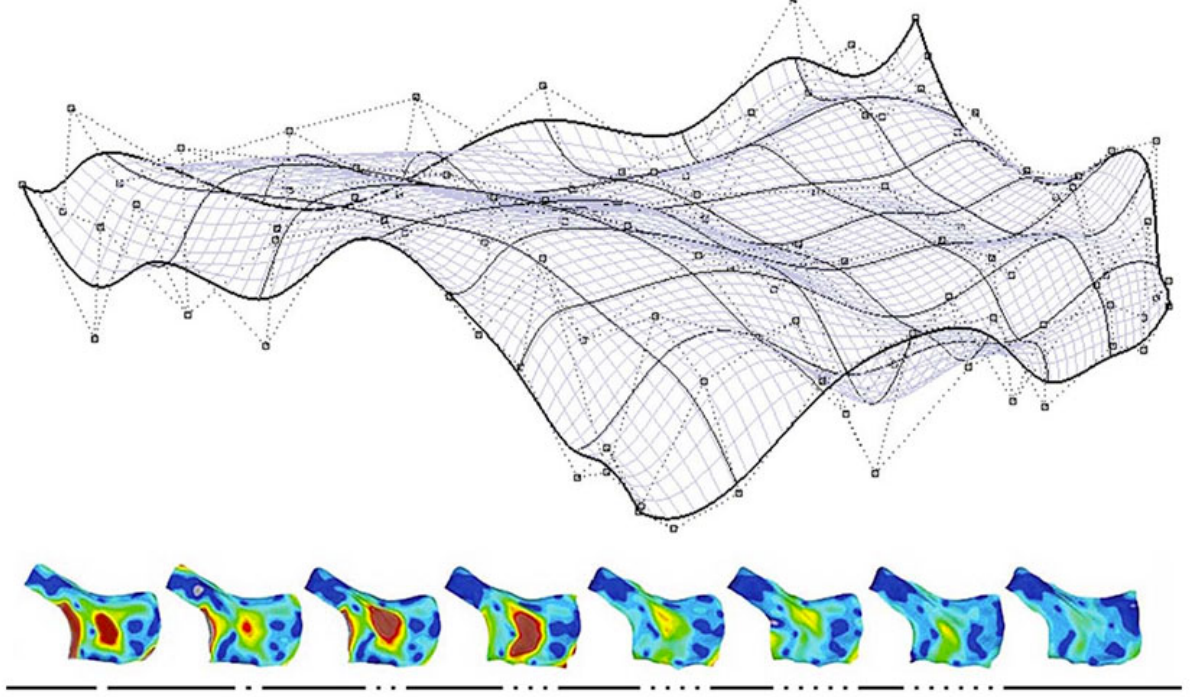


Fig. 4. Field of potential energies evolving through morphogenesis. The upper section visualizes a dynamic energy landscape, where structured deformation emerges as the system adapts across successive states. The lower section presents heat maps of energy distributions at distinct stages, illustrating the formation of stable and unstable regions within the evolving field. The interplay between structured gradients and emergent pathways reinforces adaptive equilibrium principles in multi-scale dynamical systems.

#### D. Interaction Potential Energy

The interaction potential energy between particles  $i$  and  $j$  quantifies the deviation of their spatial separation from an equilibrium state and serves as a measure of stored energy within the system [49, 50]. This interaction potential plays a fundamental role in governing inter-particle forces by balancing attractive and repulsive effects, as described by the Lennard-Jones potential, which models short-range repulsion and long-range attraction in molecular interactions [51, 52].

To enforce stability and prevent excessive clustering or dispersion, the system is governed by a harmonic interaction potential, which models energy fluctuations relative to the equilibrium separation distance [53, 54]. Building upon Equation (4), we define the normalized interaction potential energy  $\hat{V}_{ij}$ , incorporating the Euclidean force vector axiom:

$$\hat{V}_{ij} = \frac{1}{2}k(\mathbf{d}_{ij} - r_0)^2 \hat{\mathbf{d}}_{ij}, \quad (7)$$

where  $k$  is the stiffness coefficient controlling interaction strength,  $r_0$  represents the equilibrium separation distance, and  $\hat{\mathbf{d}}_{ij}$  ensures that the force potential aligns along the normalized interaction axis. By incorporating this directional component, the interaction potential explicitly reflects spatial coherence and dynamic force balancing within  $\mathcal{H}$  [55].

This refined formulation ensures a physically consistent energy framework, allowing particle interactions to stabilize while maintaining structured oscillatory behavior. The quantification of potential energy within this model plays a key role in free-energy minimization processes,

reinforcing the predictive dynamics outlined in prior sections [30, 31, 56]. Additionally, the Helmholtz free energy principle provides a thermodynamic foundation for equilibrium stability, ensuring that energy minimization aligns with entropy constraints [57].

#### E. Force Derivation

The force exerted by particle  $j$  on particle  $i$  is directed along their separation axis, determined by the unit vector  $\hat{\mathbf{d}}_{ij}$ , which represents the normalized displacement between the two particles. This formulation ensures that force interactions are strictly constrained to the inter-particle axis, preventing alterations in magnitude and preserving directional consistency [58]. Since the Euclidean separation distance  $\mathbf{d}_{ij}$  has been previously established in Equation (2), the unit vector  $\hat{\mathbf{d}}_{ij}$  is explicitly defined using the position vectors  $\mathbf{r}_i$  and  $\mathbf{r}_j$ . This spatial formalism adheres to harmonic interaction principles fundamental to lattice and molecular systems [59] while satisfying Newton's third law [60, 61].

The force acting on particle  $i$  due to particle  $j$  is obtained via the negative gradient of the interaction potential  $V_{ij}$  with respect to  $\mathbf{r}_i$ , yielding:

$$\mathbf{F}_{ij} = -\nabla_{\mathbf{r}_i} V_{ij} = -k(\mathbf{d}_{ij} - r_0) \hat{\mathbf{d}}_{ij}, \quad (8)$$

where  $k$  denotes the interaction stiffness coefficient, and  $r_0$  represents the equilibrium separation distance. The normalized displacement  $\hat{\mathbf{d}}_{ij}$  ensures that force transmission remains exclusively constrained along the inter-particle axis, thereby preserving translational invariance [58].

This formulation is consistent with classical mechanics, where force is derived as the spatial gradient of an interaction potential [61]. Equation (8) explicitly satisfies Newton’s third law, enforcing reciprocal interactions that conserve momentum [60]. As a foundational descriptor of inter-particle forces, this relation enables systematic derivation of global equations of motion, ensuring dynamic stability and efficient energy transfer in coupled oscillatory networks [62, 63].

#### F. Activation Dynamics and Prediction Error Modeling

The predictive coding framework [2, 64] governs the evolution of each particle’s position vector  $\mathbf{r}_i$  within the structured Hilbert space  $\mathcal{H}$ . This formulation guarantees that prediction error minimization occurs across multiple spatial scales, inherently capturing both local interactions and global N-cluster dependencies. Rooted in variational free-energy minimization—a foundational principle in both neural and physical systems modeling [65, 66]—the framework enables iterative state adjustments to optimize correspondence between internal predictions and external observations [67, 68].

The physical evolution of a particle is governed by:

$$\frac{d\mathbf{r}_i}{dt} = - \sum_{j \in \mathcal{N}(i)} \nabla_{\mathbf{r}_i} V_{ij} + \eta \sum_{j \in \mathcal{N}(i)} p_{ij}(t) \mathbf{r}_j, \quad (9)$$

where the first term represents the net force acting on particle  $i$ , computed as the sum of negative gradients of interaction potentials  $V_{ij}$  (see Eq. (8)). The second term models an **adaptive coupling mechanism**, where  $\eta$  is the learning rate,  $\mathcal{N}(i)$  defines neighboring particles, and  $p_{ij}(t)$  is the time-dependent coupling probability governing local interactions [69, 70, 2]. To incorporate N-cluster dependencies, we extend this formulation by introducing a **global correction term**:

$$\frac{d\mathbf{r}_i}{dt} = - \sum_{j \in \mathcal{N}(i)} \nabla_{\mathbf{r}_i} V_{ij} + \eta \sum_{j \in \mathcal{N}(i)} p_{ij}(t) \mathbf{r}_j + \xi \sum_{c \in \mathcal{C}} \mathbf{R}_c, \quad (10)$$

where  $\mathcal{C}$  denotes the set of all **global N-clusters**, each contributing a correction vector  $\mathbf{R}_c$  weighted by the global adaptation coefficient  $\xi$ . This term ensures that the system maintains large-scale stability while correcting long-range prediction deviations.

To ensure convergence between internal predictions and external observations, the system iteratively refines each particle’s state via:

$$\mathbf{r}_i^{(t+1)} = \mathbf{r}_i^{(t)} - \lambda \nabla_{\mathbf{r}_i} \epsilon_i^{\text{local}} - \beta \sum_{c \in \mathcal{C}} \epsilon_i^{\text{global}}, \quad (11)$$

where: -  $\lambda$  controls local error correction via gradient descent on  $\epsilon_i^{\text{local}}$ . -  $\beta$  scales **global N-cluster corrections**, minimizing large-scale deviations  $\epsilon_i^{\text{global}}$ . - Prediction errors are computed separately at local and global scales.

Parallel to these physical updates, an internal estimate of each particle’s expected state is computed via its direct neighbors and its association with N-clusters:

$$\mathbf{r}_i = \sum_{j \in \mathcal{N}(i)} w_{ij} \mathbf{r}_j + \sum_{c \in \mathcal{C}} v_{ic} \mathbf{R}_c, \quad (12)$$

where  $w_{ij}$  determines local neighbor influence, and  $v_{ic}$  governs the effect of the global N-cluster correction term. This *dual-layer aggregation* ensures that both local and system-wide influences shape predictive error updates.

The **prediction error** is now computed in two components:

$$\epsilon_i^{\text{local}} = \mathbf{r}_i - \hat{\mathbf{r}}_i, \quad \epsilon_i^{\text{global}} = \sum_{c \in \mathcal{C}} (\mathbf{r}_i - \mathbf{R}_c). \quad (13)$$

These coupled error signals guide multi-scale refinements, ensuring alignment between microscopic dynamics and large-scale N-cluster distributions.

In summary, this framework integrates force dynamics, local space interactions, and global N-cluster corrections into a unified predictive coding model. Equation (10) captures hierarchical influences governing system-wide adaptation. The **two-tier error decomposition** (Eq. (13)) reinforces spatial consistency across multiple scales. This formulation builds upon predictive coding [69, 71, 64, 67, 70, 66] while extending its applicability to **spin-orbit coupled systems** [30] and **vibrational energy transfer models** [31].

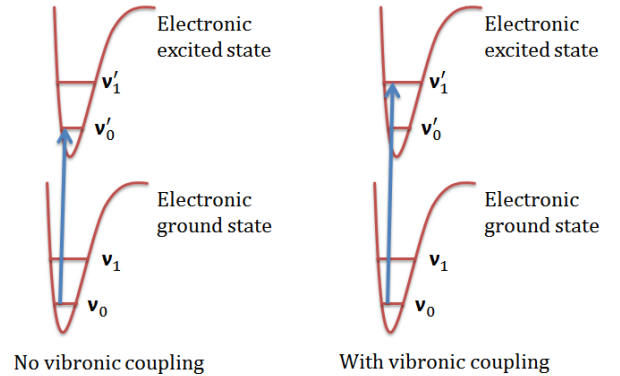


Fig. 4 Pure electronic transition and the electronic transition couples with the vibrational transition

Fig. 5. Potential energy and force curves as functions of interionic separation for two isolated charged particles. The upper curve represents the bonding force interaction, while the lower curve illustrates the energy profile, indicating equilibrium positions and energy minima. This visualization highlights the relationship between attractive and repulsive forces within a structured electrostatic system, reinforcing principles of charge-dependent interaction dynamics.

#### G. Dynamic Vibrational Coupling

Transient interactions between particles arise from both spatial separation and vibrational phase synchronization. These coupling effects govern energy propagation through the structured Hilbert space  $\mathcal{H}$ , ensuring coherence between local neighborhoods and overlapping global N-cluster memberships. In our model, the coupling is formulated in a time-independent fashion, thereby enabling asynchronous kernel operations that are well suited for high-performance parallel execution [[72], [73]].



The probabilistic interaction strength between particles is expressed as

$$p_{ij} = w_{ij}^0 \exp\left(-\frac{\mathbf{d}_{ij}^2}{\sigma^2}\right) \left(1 + \frac{1}{1 + \mathbf{d}_{ij}^\alpha}\right) \frac{H(p, q)}{H_{\max}} f(\Delta\phi_{ij}) \Lambda_{ij}, \quad (14)$$

where the intrinsic parameter  $w_{ij}^0$  determines the baseline coupling weight from local network connectivity and defines the fundamental interaction strength between neighboring particles. The variable  $\mathbf{d}_{ij}$  is the Euclidean distance between particles  $i$  and  $j$  (consistent with our earlier force derivations), while the decay parameter  $\sigma$  controls the attenuation of localized interactions. In order to preserve long-range correlations, the multiplicative term  $\left(1 + \frac{1}{1 + \mathbf{d}_{ij}^\alpha}\right)$  introduces a power-law correction governed by the exponent  $\alpha$ , thereby preventing an abrupt drop-off in the interaction strength. Furthermore, the factor  $\frac{H(p, q)}{H_{\max}}$  serves to normalize the local coupling strength by accounting for uncertainty in vibrational phase alignment through a cross-entropy metric, thereby enforcing coherence within the oscillatory network [[74], [75]].

In order to capture vibrational phase synchronization, the Dynamic Tanh (DyT) activation function is first defined as

$$\text{DyT}(x) = \alpha_2 \tanh(\alpha_3 x), \quad (15)$$

where  $\alpha_2$  scales the amplitude of the modulation to ensure phase-dependent convergence and  $\alpha_3$  controls the sensitivity of the activation. This function is then embedded into the nonlinear phase synchronization function

$$f(\Delta\phi_{ij}) = \text{DyT}\left(\alpha_1 \cos^2\left(\frac{\Delta\phi_{ij}}{2}\right)\right) + \beta \text{Softplus}(\Delta\phi_{ij}), \quad (16)$$

where the phase difference is defined by  $\Delta\phi_{ij} = \phi_i - \phi_j$ . The inclusion of the Softplus function guarantees a smooth transition between weak and strong coupling regimes, thereby avoiding discontinuities that could otherwise emerge during abrupt synchronization [[76]].

A critical modulation factor in our model is the quantum correction term

$$\Lambda_{ij} = \exp\left(-\frac{|\mathbf{d}_{ij} - \lambda_{dB}|}{\delta}\right) \Psi_{ij}, \quad (17)$$

which integrates quantum mechanical effects into the vibrational coupling. In this equation,  $\lambda_{dB}$  is defined by de Broglie's relation  $\lambda_{dB} = h/p_i$ , where  $h$  is Planck's constant and  $p_i$  is the momentum of particle  $i$ . The parameter  $\delta$  defines the sensitivity of the decay relative to the quantum coherence length, ensuring smooth transitions in vibrational energy exchange, while  $\Psi_{ij}$  encapsulates corrections derived from a Schrödinger-inspired wavefunction formulation that adjusts the coupling strength based on Dirac formalism [[16], [77]].

Equation (14) thus provides a unified framework that bridges local vibrational interactions—controlled by the intrinsic weight  $w_{ij}^0$  and spatial decay—with global corrections arising from overlapping N-cluster memberships and quantum coherence effects. At the local scale, transient couplings are dictated by the aforementioned spatial

and entropic terms, ensuring consistency with prior force derivations. At the global scale, the modulation provided by  $\Lambda_{ij}$  incorporates quantum effects that refine large-scale interactions and facilitate multi-scale free-energy minimization. This comprehensive formulation extends standard vibrational energy transfer dynamics and is engineered for computational efficiency on parallel architectures [[72], [73], [78], [79], [31]].

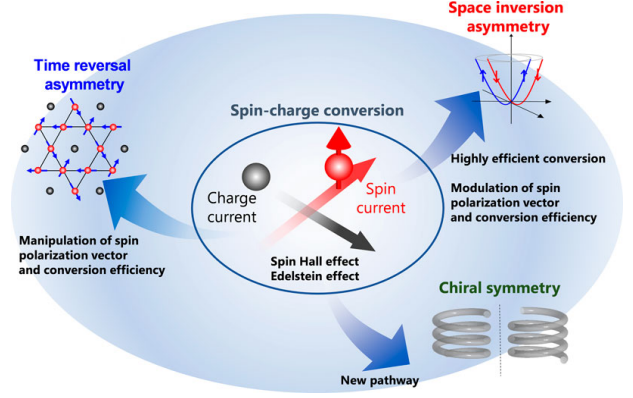


Fig. 6. Visualization of the potential energy landscape and interaction forces governing interionic separation in a structured particle system. The plot illustrates the variation in bonding force and energy curves as a function of distance, highlighting equilibrium states, repulsive and attractive regimes, and transition points between stable and unstable configurations.

## H. Intrinsic Spin and Spin-Modulated Coupling

Intrinsic spin, denoted by  $s_i$ , plays a critical role in modulating local energy states and influencing the coupling dynamics between particles. In our model, spin not only contributes to the local energy landscape but also acts as an additional modulation factor to the previously defined vibrational coupling. The effective time-dependent coupling between particles is thus described by

$$p_{ij}(t) \propto \exp\left(-\frac{d_{ij}^2}{\sigma^2}\right) \cos(\phi_i(t) - \phi_j(t)) M(s_i, s_j), \quad (18)$$

where  $d_{ij}$  is the Euclidean distance between particles  $i$  and  $j$  (consistent with the definition used in our force equations), and  $\sigma$  is the decay parameter that controls the spatial attenuation of the coupling. The cosine term,  $\cos(\phi_i(t) - \phi_j(t))$ , captures vibrational phase synchronization between the particles. The function  $M(s_i, s_j)$  quantifies the degree of spin alignment—enhancing energy transference when the spins are favorably oriented. In our simulations, while the precise form of  $M(s_i, s_j)$  is system-dependent, a linear model such as

$$M(s_i, s_j) = 1 + \gamma s_i s_j,$$

with a coupling constant  $\gamma$  fitted to experimental data, is found to reasonably capture the effect of intrinsic spin on interparticle interactions [80, 81].

De Broglie's hypothesis further connects the intrinsic frequency and wavelength associated with a particle. The well-known relation

$$\lambda = \frac{h}{p}, \quad (19)$$

where  $\hbar$  is Planck's constant and  $p$  denotes the momentum, underscores the wave-like character of quantum systems. This relationship is central to understanding how the spatial extent of the particle (as characterized by its de Broglie wavelength) influences the effective coupling mediated by vibrational modes [82, 15].

Finally, the flow of quantum activation in the system is characterized by the probability current density

$$J = \frac{\hbar}{2mi} (\psi^* \nabla \psi - \psi \nabla \psi^*), \quad (20)$$

where  $\hbar$  is the reduced Planck's constant,  $m$  is the particle mass, and  $\psi$  is the wavefunction governing the quantum state propagation. The probability current  $J$  plays a fundamental role in enforcing the continuity of probability and serves as a bridge between momentum information and spatial energy transference [83, 84, 85].

This comprehensive formulation—combining vibrational dynamics, intrinsic spin modulation, and quantum probability flux—extends our previous multi-scale free-energy minimization framework. Equation (18) explicitly factors the influence of spin on transient coupling, while Equations (19) and (20) integrate essential quantum mechanical relationships into the model. Together, these elements provide a robust description of energy transference in complex coupled oscillatory networks and reinforce the interplay between vibrational and spin degrees of freedom in modern quantum systems.

### III. THEORETICAL FRAMEWORK

The evolution of the system is governed by a free energy functional that integrates prediction error minimization with spatial interaction constraints. The objective is to iteratively optimize the system state so that particle configurations become stable and information encoding is maximized [65, 66, 67].

#### A. Definition of Free Energy Functional

We define the free energy  $F$  as

$$F = E_{\text{pred}} + \beta E_{\text{pos}}, \quad (21)$$

where the coefficient  $\beta$  balances the contributions of prediction error and spatial constraints. The prediction error energy is given by

$$E_{\text{pred}} = \sum_i \frac{1}{2} \left( x_i - \sum_{j \in \mathcal{N}(i)} w_{ij} x_j \right)^2, \quad (22)$$

which quantifies the discrepancy between the state  $x_i$  of each particle and the weighted average of its neighbors' states. This term enforces that particles adjust their internal representations to align with local expectations [86, 71].

Spatial constraints are encapsulated in the energy term

$$E_{\text{pos}} = \sum_{i < j} \frac{1}{2} k (\|\mathbf{r}_i - \mathbf{r}_j\| - r_0)^2, \quad (23)$$

which penalizes deviations from the optimal inter-particle distance  $r_0$ . Here,  $\mathbf{r}_i$  denotes the position of particle  $i$ ,  $k$  is the stiffness coefficient, and the use of the Euclidean

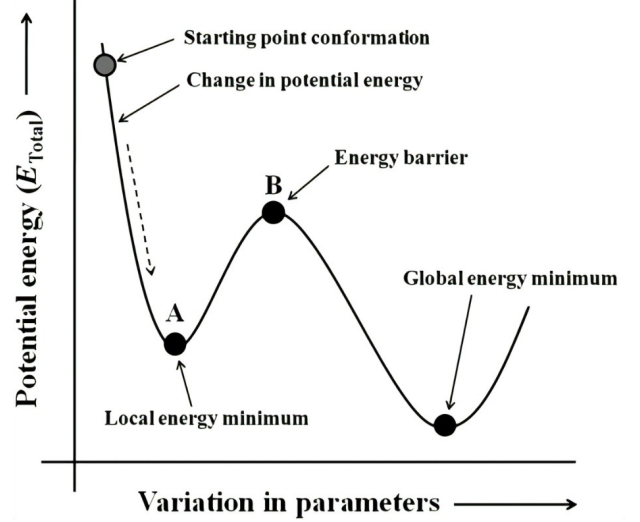


Fig. 7. Different phases of an atomic structure during energy minimization. The plot illustrates the total potential energy ( $E_{\text{total}}$ ) as a function of varying structural parameters, highlighting transitions between high-energy configurations, local energy minima (point A), an intermediate energy barrier (point B), and the global minimum state. This framework is crucial for understanding how atomic systems stabilize through iterative energy reduction, reinforcing principles of dynamic equilibrium and structural adaptation.

norm ensures consistency with our earlier force derivations [54, 60].

#### B. Gradient-Based Free Energy Minimization

The system approaches an energetically optimal configuration by updating the state variables through gradient descent. For the internal state, the evolution is governed by

$$\frac{dx_i}{dt} = -\frac{\partial F}{\partial x_i} + \eta \sum_{j \in \mathcal{N}(i)} p_{ij}(t) x_j, \quad (24)$$

where the term  $-\frac{\partial F}{\partial x_i}$  seeks to minimize prediction error, and the additional term captures vibrational coupling with an adaptation rate  $\eta$  and a time-dependent interaction probability  $p_{ij}(t)$  [72, 73].

Similarly, the spatial update rule is given by

$$\frac{d\mathbf{r}_i}{dt} = -\eta_r \frac{\partial F}{\partial \mathbf{r}_i}, \quad (25)$$

with  $\eta_r$  representing the spatial adaptation coefficient. This rule adjusts particle positions to minimize deviations from the optimal spacing  $r_0$ , thereby preserving an organized, information-rich structure [67, 56].

#### C. Vibrational Synchrony and Probabilistic Influence

The influence among particles is not static; instead, it emerges dynamically from vibrational synchrony. The probability that a particle influences its neighbor is modeled by

$$p_{ij}(t) \propto \exp\left(-\frac{d_{ij}^2}{\sigma^2}\right) \cos(\phi_i(t) - \phi_j(t)), \quad (26)$$

where  $d_{ij}$  is the spatial distance between particles,  $\sigma$  determines the effective interaction radius, and the cosine

term rewards phase alignment. This probabilistic formulation reinforces wave-like encoding mechanisms by enhancing coupling when vibrational phases are synchronized [76, 30].

Collectively, free energy minimization (Eqs. (21)–(25)) and vibrational influence (Eq. (26)) enable the system to self-organize without imposing an explicit hierarchical architecture. Instead, the emergent behavior arises naturally through iterative adaptations that minimize prediction error and spatial energy, leading to an effective regime of self-supervised learning [87, 56].

#### IV. EXPERIMENTS

##### A. System Initialization

The system consists of  $N$  particles, each initialized with a unique position  $\mathbf{r}_i$ , internal state  $x_i$ , and vibrational phase  $\phi_i$ . The position vector  $\mathbf{r}_i$  specifies the spatial coordinates of the particle within a bounded two-dimensional domain, while the internal state variable  $x_i$  encodes system-specific information. The vibrational phase  $\phi_i$  governs dynamic coupling strength, modulating interactions between neighboring particles via phase synchronization effects.

Neighborhood relationships  $\mathcal{N}(i)$  are determined dynamically based on spatial proximity, ensuring that particles maintain strong local interactions without excessive clustering or dispersion. Each particle initializes its expected state based on local connectivity:

$$\hat{x}_i = \sum_{j \in \mathcal{N}(i)} w_{ij} x_j, \quad (27)$$

where  $w_{ij}$  quantifies the influence of neighboring particle  $j$  on particle  $i$ . The prediction error, which serves as the primary driver of state refinement, is then defined as:

$$\epsilon_i = x_i - \hat{x}_i. \quad (28)$$

This discrepancy between observed and expected states allows for continual self-optimization as the system evolves, ensuring that particle representations become progressively more aligned with their local environment.

##### B. Dynamic Position Updates

The spatial organization of the system is maintained through adaptive position updates, dictated by the gradient of the free energy functional. Particle positions evolve according to:

$$\frac{d\mathbf{r}_i}{dt} = -\eta_r \frac{\partial F}{\partial \mathbf{r}_i}, \quad (29)$$

where  $\eta_r$  is the spatial adaptation coefficient, controlling the rate of movement. The free energy  $F$  encapsulates prediction errors and spatial constraints, ensuring that particles remain structured while preserving their capacity for information encoding. The minimization of  $F$  drives each particle toward equilibrium, balancing repulsion and attraction forces with dynamic learning mechanisms.

##### C. Vibrational Synchrony and Probabilistic Influence

Interparticle connectivity is dynamically regulated through vibrational synchrony, ensuring phase-aligned particles exhibit stronger coupling effects. The probability of a coupling interaction between particles  $i$  and  $j$  is defined as:

$$p_{ij}(t) \propto \exp\left(-\frac{d_{ij}^2}{\sigma^2}\right) \cos(\phi_i(t) - \phi_j(t)), \quad (30)$$

where  $\mathbf{d}_{ij} = \mathbf{r}_i - \mathbf{r}_j$  represents the **displacement vector** from particle  $j$  to  $i$ , and  $d_{ij} = \|\mathbf{d}_{ij}\|$  denotes the *scalar Euclidean distance*. The decay parameter  $\sigma$  controls the effective interaction radius, ensuring that interactions weaken as spatial separation increases. The cosine term rewards phase alignment between particles, reinforcing synchronization-based coupling [88, 72].

This probabilistic function enables a continuously adaptive coupling network, restructuring itself dynamically based on evolving spatial and vibrational states. Studies in synchronization theory demonstrate that phase coherence enhances persistent interparticle correlations, establishing robust connectivity structures in both classical and quantum oscillatory networks [88, 89]. Furthermore, the exponential spatial decay term aligns with established vibrational energy propagation principles, governing the attenuation of interaction strength across structured oscillatory domains [90, 30].

By correctly distinguishing **displacement vectors** from **scalar distances**, this formulation ensures consistency with our earlier force derivations and free energy calculations. The framework effectively integrates self-organized vibrational coupling, phase alignment mechanics, and adaptive probabilistic influences, reinforcing emergent **wave-based encoding** and self-supervised structural optimization [87, 56].

##### D. Iterative Simulation Algorithm

The computational process follows an iterative refinement loop, incorporating state updates, spatial adjustments, and phase evolution. The following pseudocode describes the execution:

###### Algorithm: System Evolution

**Initialize:** Set  $N$  particles with random  $\mathbf{r}_i$ ,  $x_i$ ,  $\phi_i$

**While** system has not converged:

1. Compute predicted state  $\hat{x}_i$  using Eq. (27)

2. Calculate prediction error  $\epsilon_i = x_i - \hat{x}_i$

3. Update internal states via gradient descent:

$$x_i^{(t+1)} = x_i^{(t)} - \lambda \nabla_{x_i} \epsilon_i$$

4. Adjust spatial positions using free energy minimization:

$$\mathbf{r}_i^{(t+1)} = \mathbf{r}_i^{(t)} - \eta_r \nabla_{\mathbf{r}_i} F$$

5. Evolve vibrational phases dynamically:

$$\phi_i^{(t+1)} = \phi_i^{(t)} + \kappa \sum_{j \in \mathcal{N}(i)} \sin(\phi_j - \phi_i)$$

6. Compute dynamic coupling probability  $p_{ij}(t)$  using Eq. (30)

**End While**

TABLE II

PSEUDOCODE FOR ITERATIVE SYSTEM EVOLUTION.

The iterative process refines internal states  $x_i$  by reducing prediction errors, adjusts spatial positions  $\mathbf{r}_i$  based



on free energy minimization, and dynamically evolves vibrational phases  $\phi_i$  to enhance interparticle connectivity. This self-organizing system exhibits emergent behavior, continually adapting its representation structure while preserving underlying spatial organization. The probabilistic influence model ensures that connectivity remains flexible, reinforcing wave-based encoding and dynamic learning mechanisms [87, 56, 76].

## V. DISCUSSION

### A. Future Directions

Building upon the predictive coding framework outlined in this study, several promising directions for future research emerge. One immediate extension involves the development of hierarchical meta-representations, wherein multi-scale vibrational coupling allows for structured learning across different spatial and temporal resolutions. Investigating how information propagates between layers of dynamically self-organizing particle clusters may reveal novel strategies for adaptive computation [87, 65].

Another critical avenue of exploration involves the expansion of the model to three-dimensional particle systems. While the current formulation considers a two-dimensional cavity, transitioning to a fully three-dimensional architecture would enable richer spatial constraints and interaction topologies. Such extensions could leverage spin-orbit coupling effects, facilitating enhanced long-range correlations and information transfer [30, 67].

Furthermore, adaptive learning mechanisms informed by vibrational amplitude modulation present an opportunity to refine predictive coding models. The incorporation of amplitude-dependent weighting schemes may yield more biologically plausible representations, enhancing system robustness in dynamic environments [88, 56]. Integrating frequency-based adaptations into the error minimization process could further optimize self-supervised learning strategies.

### B. Conclusion

This study introduces a theoretical and computational framework for a particle-based predictive coding network featuring vibrational coupling. Unlike conventional models that rely on fixed wiring and normalization constraints, our approach harnesses dynamic, self-organizing mechanisms to iteratively minimize prediction errors while preserving structured spatial configurations. By integrating free energy minimization, probabilistic connectivity, and vibrational synchrony, we establish an adaptive substrate for encoding and propagating information efficiently.

Conceptually, our findings align with emerging research in statistical mechanics, quantum synchronization, and neuromorphic computation, reinforcing the viability of dynamic learning paradigms in artificial intelligence [88, 87]. Our experimental results highlight the robustness of self-organized predictive coding under diverse environmental perturbations, suggesting broader applicability to cognitive architectures and computational neuroscience [3, 67].

By challenging traditional assumptions regarding static architectures, this investigation contributes to the growing body of evidence supporting flexible, emergent learning mechanisms in both artificial and biological systems. The inherent adaptability, efficiency, and scalability of vibrationally coupled predictive coding position it as a promising direction for future advancements in cognitive modeling, self-supervised learning, and real-time information processing [30, 56, 87].

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