

### Problem 1 (Purcell 9.16)

Since  $\vec{E} = E_0(\hat{x} + \hat{y}) \sin\left[\left(\frac{2\pi}{\lambda}\right)(z + ct)\right]$ , the wave travels in the  $-\hat{z}$  direction. We know that  $\vec{B}$  is perpendicular to  $-\hat{z}$  and to  $\vec{E}$ . Therefore,  $\vec{E}$  points in the  $\pm(\hat{x} - \hat{y})$  direction. Using the right hand rule,  $\vec{B}$  points in the  $(\hat{x} - \hat{y})$  direction. Since  $|\vec{B}| = \frac{1}{c}|\vec{E}|$ , we have that  $\vec{B}$  is:

$$\vec{B} = \frac{E_0(\hat{x} - \hat{y})}{c} \sin\left[\left(\frac{2\pi}{\lambda}\right)(z + ct)\right] \text{ Now, taking } E_0 = 20 \text{ V/m, we find that } B_0 = \frac{1}{c}E_0 = 6.67 \cdot 10^{-8} \text{ T.}$$

### Problem 2 (Purcell 9.20)

We are told that  $\vec{E}$  follows Eq 9.28, so  $\vec{E} = \frac{E_0 \hat{y}}{1 + (x/ct)^2}$ . We are given that  $E_0 = 100 \text{ kV/m}$  and  $l = 1 \text{ ft}$ . Since  $F = \frac{dp}{dt}$ , we can find the momentum acquired by the proton during the pulse to be:

$$p_y = \int_{-\infty}^{\infty} e E_y dt = \frac{e l E_0}{c} \left( \arctan\left(\frac{t}{t_0}\right) \right) \Big|_{-\infty}^{\infty} = \frac{e l E_0 \pi}{c}$$

Therefore, the velocity is  $v_y = \frac{p_y}{m}$ , so  $\frac{\pi e l E_0}{m c} \approx 3.1 \cdot 10^4 \text{ m/s}$ . Therefore, since we know the velocity, we can find the displacement after one microsecond to be  $v_y t = (3.1 \cdot 10^4)(10^{-6}) = 3.1 \text{ cm}$ . Therefore, the proton is 3.1 cm away from where it was a microsecond after the pulse.

### Problem 3 (Purcell 9.24)

Since the leaving signal has 1000 km in diameter, it sweeps an area of  $8 \cdot 10^{14} \text{ m}^2$ . Since we know the power of this signal, the power density is  $10^4 / 8 \cdot 10^{14}$ , so the power density is roughly  $10^{-8} \text{ W/m}^2$ . Since  $10^{-8} = \frac{E^2}{377 \Omega}$ ,  $E_{\text{ant}}$  at the receptor is roughly  $0.002 \text{ V/m}$  at the receiving end!

### Problem 4 (Purcell 9.25)

We have that the energy density is  $4 \cdot 10^{-14} \text{ J/m}^3$ . From section 9.6,  $E_{\text{rms}}^2 = \frac{U}{\epsilon_0}$ , where  $U$  is the average energy density. Therefore,  $E_{\text{rms}}^2 = 4.5 \cdot 10^{-3} \text{ V}^2/\text{m}^2$ , so  $E_{\text{rms}} = 0.067 \text{ V/m}$ . Now, if our 1 kW radio transmitter is "spherically symmetrical", then the power density at a radius  $R$  is just  $10^3 \text{ W} / 4\pi R^2$ . Therefore,  $U = \frac{1}{\epsilon_0} \cdot \frac{10^3 \text{ W}}{4\pi R^2}$ , so  $R = 2600 \text{ m}$ .

### Problem 5 (Purcell 9.29)

To find the  $\vec{B}$  field inside the capacitor, we integrate  $\nabla \times \vec{B} = \mu_0 \vec{J} + \epsilon_0 \mu_0 \frac{\partial \vec{E}}{\partial t}$  over a disk. Using Stokes's Theorem, we have:

$$\int \vec{B} \cdot d\vec{s} = \mu_0 I - \epsilon_0 \mu_0 \frac{\partial E}{\partial t} A \xrightarrow{\text{over a disk}} B = \frac{\mu_0 I}{2\pi r} - \frac{\epsilon_0 \mu_0 r}{2} \frac{\partial E}{\partial t}.$$

Since  $I = \epsilon_0 (\pi b^2) \frac{\partial E}{\partial t}$ ,  $B = \frac{\epsilon_0 \mu_0 r}{2} \frac{\partial E}{\partial t} \left( \frac{b^2}{r} - r \right)$ . For all  $r < b$  inside the capacitor,  $B > 0$ . Since  $B$  points into the page on the right side of the capacitor,  $\vec{E}$  points downwards, and  $\vec{S} = \frac{1}{\mu_0} (\vec{E} \times \vec{B})$  points away. Therefore, energy flows out away from the wire. Since  $\vec{E} \perp \vec{B}$ , we have that  $|\vec{S}| = \frac{E}{2} \frac{\partial E}{\partial t} \left( \frac{b^2}{r} - r \right)$ .

Now, since the total energy flowing out of a cylinder is  $S \cdot 2\pi r h$ , the power is given by

$$P = S \cdot 2\pi r h = \frac{d}{dt} \left( \frac{\epsilon_0 E^2 h}{2} (\pi b^2 - \pi r^2) \right). \text{ Therefore, the rate of flux of the Poynting vector is the}$$

as the energy stored in the field. We can just confirm this by taking limiting cases  $r=0$  or  $r=b$ .

### Problem 6 (Arcell 10.17)

We know that the maximum field is just  $9.5 \cdot 10^5 \text{ V/m}$ , after converting from  $\text{kV/mm}$ . Since the capacitance of the Mylar-filled capacitor is  $\frac{k\epsilon_0 A}{s}$ , and the energy stored in it is  $\frac{1}{2} C \phi^2$ , the maximum energy density is  $\frac{\text{energy}}{\text{volume}} = \frac{1}{2} k\epsilon_0 E^2 = 4.4 \cdot 10^6 \text{ J/m}^3$ . Therefore, the max. energy per kg of Mylar is  $3100 \text{ J/kg}$ . To determine how high the capacitor could operate, we take  $E = mgh$ , where  $m$  is the total mass of the capacitor, and plug in to find that  $h = 240 \text{ m}$ . The cell in problem 4.11 had about 60 times more energy than this capacitor, but the Mylar capacitor can "transfer" its energy a lot quicker!

### Problem 7 (Arcell 10.18)

The second capacitor is no more than two capacitors in series. Both capacitors have separation  $\frac{s}{2}$  and area  $A$ . The capacitance for each half, with vacuum or with dielectric, is  $C_{\text{vacuum}} = \frac{\epsilon_0 A}{\frac{s}{2}}$  and  $C_{\text{dielectric}} = k \frac{\epsilon_0 A}{\frac{s}{2}}$ . Taking  $C_0 = \frac{\epsilon_0 A}{s}$ ,  $C_v = 2C_0$  and  $C_d = 2kC_0$ . Since capacitors add in series, the total capacitance is:  $\frac{1}{C_{\text{tot}}} = \frac{1}{C_v} + \frac{1}{C_d} = \frac{1}{2k} \frac{1}{C_0}$ . The third capacitor is two capacitors in parallel. Both capacitances are now  $C_v = \frac{A\epsilon_0}{\frac{s}{2}}$  and  $C_d = \frac{kA\epsilon_0}{\frac{s}{2}}$ , so  $C_d = \frac{k}{2} C_0$  and  $C_v = \frac{1}{2} C_0$ . Therefore,  $C_{\text{total}} = \frac{1+k}{2} C_0$ .