

The velocity of a point with position \mathbf{r} relative to the center of mass of a cell with its main axis along the unit vector $\hat{\mathbf{a}}$ is:

$$\begin{aligned} \mathbf{v_r} &= \mathbf{v}_{\text{linear}} + \mathbf{v}_{\text{angular}} + \mathbf{v}_{\text{growth}} \\ &= \mathbf{v} + \boldsymbol{\omega} \times \mathbf{r} + \frac{1}{2} \frac{\hat{\mathbf{a}} \cdot \mathbf{r}}{l} v_{growth} \\ &= \mathbf{v} - \mathbf{r} \times \boldsymbol{\omega} + \frac{1}{2} \frac{\hat{\mathbf{a}} \cdot \mathbf{r}}{l} v_{growth} \\ &= \frac{1}{m} \mathbf{p} - \mathbf{r} \times \mathbf{I}^{-1} \mathbf{L} + \frac{1}{2} \frac{\hat{\mathbf{a}} \cdot \mathbf{r}}{l} \frac{g}{G} \\ &= \frac{1}{m} \mathbf{p} - [\mathbf{r}]_{\times} \mathbf{I}^{-1} \mathbf{L} + \frac{1}{2} \frac{\hat{\mathbf{a}} \cdot \mathbf{r}}{l} \frac{g}{G} \end{aligned}$$

This expression is linear in \mathbf{p} , \mathbf{L} , and g, so the velocity of the point after an impulse $\Delta \mathbf{p}$ is:

$$\begin{aligned} \mathbf{v_r} &= \mathbf{v_{r0}} + \Delta \mathbf{v} \\ &= \mathbf{v_{r0}} + \frac{1}{m} \Delta \mathbf{p} - [\mathbf{r}]_{\times} \mathbf{I}^{-1} \Delta \mathbf{L} \times \mathbf{r} + \frac{1}{2} \frac{\hat{\mathbf{x}} \cdot \mathbf{r}}{l} \frac{\Delta g}{G} \\ &= \mathbf{v_{r0}} + \mathbf{M_{\hat{\mathbf{a}}, \mathbf{r}}} \Delta \mathbf{p} \end{aligned}$$

 \dots where:

$$\mathbf{M}_{\hat{\mathbf{a}},\mathbf{r}} = \begin{bmatrix} \frac{1}{m}I & [\mathbf{r}]_{\times}\mathbf{I}^{-1} & \frac{\hat{\mathbf{a}}\cdot\mathbf{r}}{2lG}\hat{\mathbf{a}} \end{bmatrix}$$

 \dots and:

$$\Delta \mathbf{p} = \left[egin{array}{c} \Delta \mathbf{p} \ \Delta \mathbf{p} \ \Delta \mathbf{p} \ \Delta \mathbf{p}_z \ \Delta L_x \ \Delta L_y \ \Delta L_z \ \Delta g \end{array}
ight] = \left[egin{array}{c} \Delta p_x \ \Delta p_z \ \Delta L_x \ \Delta L_y \ \Delta L_z \ \Delta g \end{array}
ight]$$

In the above, I is the identity matrix and \mathbf{I}^{-1} is the inverse inertia tensor. No confusion should arise from the use of $\Delta \mathbf{p}$ to signify both general momentum and specifically linear momentum.

The inertia tensor for a rod of mass m, radius r, and length l centered at the origin and lying along the x axis is:

$$\mathbf{I} = \frac{1}{12} m \begin{bmatrix} 6r^2 & 0 & 0\\ 0 & 3r^2 + l^2 & 0\\ 0 & 0 & 3r^2 + l^2 \end{bmatrix}$$

Therefore the inverse inertia tensor for the same cylinder is:

$$\mathbf{I}^{-1} = \frac{12}{m} \begin{bmatrix} \frac{1}{6r^2} & 0 & 0\\ 0 & \frac{1}{3r^2 + l^2} & 0\\ 0 & 0 & \frac{1}{3r^2 + l^2} \end{bmatrix}$$

Suppose \mathbf{T}_i transforms from world coordinates to local coordinates for cell i—i.e. it transforms the x axis to the main axis of that cell. Then the inverse inertia tensor for cell i in world coordinates is given by:

$$\mathbf{I}_i^{-1} = \mathbf{T}_i \mathbf{I}_{\text{cyl}}^{-1} \mathbf{T}_i^{-1}$$

So we now have:

$$\mathbf{M}_{\hat{\mathbf{a}},\mathbf{r}} = \begin{bmatrix} \frac{1}{m}I & [\mathbf{r}]_{\times}\mathbf{T}_{i}\mathbf{I}_{\mathrm{cyl}}^{-1}\mathbf{T}_{i}^{-1} & \frac{\hat{\mathbf{a}}\cdot\mathbf{r}}{2lG}\hat{\mathbf{a}} \end{bmatrix}$$

We assume that colliding capsules — those that overlap and are moving toward each other (the dot product of the velocity vectors of the contact points on the two capsules is negative) — collide inelastically. In this case, the relative velocity of the two contact points along the normal $\hat{\bf n}$ after the impulse provided by the collision must be zero. So:

$$\begin{split} \hat{\mathbf{n}}^{\mathrm{T}} \left(\mathbf{v}_{\mathbf{r}_{a}} - \mathbf{v}_{\mathbf{r}_{b}} \right) \cdot &= 0 \\ \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{a}} \mathbf{p}_{a} - \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{b}} \mathbf{p}_{b} &= 0 \\ \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{a}} \left(\mathbf{p}_{a_{0}} + \Delta \mathbf{p}_{a} \right) - \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{b}} \left(\mathbf{p}_{b_{0}} + \Delta \mathbf{p}_{b} \right) &= 0 \\ \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{a}} \Delta \mathbf{p}_{a} - \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{b}} \Delta \mathbf{p}_{b} &= \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{b}} \mathbf{p}_{b_{0}} - \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{a}} \mathbf{p}_{a_{0}} \\ \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{a}} \Delta \mathbf{p}_{a} - \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{b}} \Delta \mathbf{p}_{b} &= \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{v}_{b_{0}} - \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{v}_{a_{0}} \\ \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{a}} \Delta \mathbf{p}_{a} - \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{b}} \Delta \mathbf{p}_{b} &= \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{v}_{\mathrm{rel}_{a,b}} \end{split}$$

Write:

$$\begin{split} \mathbf{N}_{\hat{\mathbf{a}},\mathbf{r}} &= \left[\begin{array}{ccc} \hat{\mathbf{n}}^{\mathrm{T}} \frac{1}{m} I & \hat{\mathbf{n}}^{\mathrm{T}} [\mathbf{r}]_{\times} \mathbf{T}_{i} \mathbf{I}_{\mathrm{cyl}}^{-1} \mathbf{T}_{i}^{-1} & \hat{\mathbf{n}}^{\mathrm{T}} \frac{\hat{\mathbf{a}} \cdot \mathbf{r}}{2lG} \hat{\mathbf{a}} \end{array} \right] \\ &= \left[\begin{array}{ccc} \frac{1}{m} \hat{\mathbf{n}} & \hat{\mathbf{n}}^{\mathrm{T}} [\mathbf{r}]_{\times} \mathbf{T}_{i} \mathbf{I}_{\mathrm{cyl}}^{-1} \mathbf{T}_{i}^{-1} & \frac{1}{2lG} \left(\hat{\mathbf{a}} \cdot \mathbf{r} \right) \left(\hat{\mathbf{a}} \cdot \hat{\mathbf{n}} \right) \end{array} \right] \end{split}$$

This is a 1×7 matrix which when multiplied by the 7×1 matrix $\Delta \mathbf{p}$ gives a change in velocity along the normal $\hat{\mathbf{n}}$.

In block matrix form, we can write:

$$\left[\begin{array}{cc} \hat{\mathbf{n}}^{\mathrm{T}}\mathbf{M}_{\mathbf{r}_{a}} & -\hat{\mathbf{n}}^{\mathrm{T}}\mathbf{M}_{\mathbf{r}_{b}} \end{array}\right] \left[\begin{array}{c} \Delta\mathbf{p}_{a} \\ \Delta\mathbf{p}_{b} \end{array}\right] = \hat{\mathbf{n}}^{\mathrm{T}}\mathbf{v}_{\mathrm{rel}_{a,b}}$$

More generally, we will have a system of equations with one such equation for each pair of colliding cells. We build a larger block matrix with a row for each constraint and a column for each cell in the system. If the ith constraint involves cells a and b (in that order), then:

$$\mathbf{M}_{ij} = \left\{ \begin{array}{ll} \hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{a}} & \text{if } j = a \\ -\hat{\mathbf{n}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{b}} & \text{if } j = b \\ 0 & \text{otherwise} \end{array} \right\}$$

A sytem with four cells and four collisions might have the equation:

$$\begin{bmatrix} \mathbf{M_{r_{1a}}} & -\mathbf{M_{r_{1b}}} & 0 & 0 \\ 0 & \mathbf{M_{r_{2b}}} & -\mathbf{M_{r_{2c}}} & 0 \\ \mathbf{M_{r_{3a}}} & 0 & 0 & -\mathbf{M_{r_{3d}}} \\ 0 & 0 & \mathbf{M_{r_{4c}}} & -\mathbf{M_{r_{4d}}} \end{bmatrix} \begin{bmatrix} \Delta \mathbf{p}_a \\ \Delta \mathbf{p}_b \\ \Delta \mathbf{p}_c \\ \Delta \mathbf{p}_d \end{bmatrix} = \begin{bmatrix} \mathbf{v}_{\mathrm{rel}_{a,b}} \\ \mathbf{v}_{\mathrm{rel}_{b,c}} \\ \mathbf{v}_{\mathrm{rel}_{a,d}} \\ \mathbf{v}_{\mathrm{rel}_{c,d}} \end{bmatrix}$$

We can transform this into a square matrix by multiplying by \mathbf{M}^{T} .

$$\mathbf{M}^{\mathrm{T}}\mathbf{M}\Delta\mathbf{p} = \mathbf{M}^{\mathrm{T}}\mathbf{v}_{\mathrm{rel}}$$

Consider $\mathbf{A} = \mathbf{M}^{\mathrm{T}} \mathbf{M}$. We have:

$$\mathbf{A}_{ij} = \sum_k \mathbf{M}_{ik}^{\mathrm{T}} \mathbf{M}_{kj}$$

 \dots where $k \in \text{constraints}$.

If $i \neq j$ and there is no constraint between cells i and j then $\mathbf{M}^{\mathrm{T}}{}_{ik}\mathbf{M}_{kj}$ will be zero. If there is a constraint between i and j, then:

$$\mathbf{A}_{ij} = -\mathbf{M}_{\mathbf{r}_{ik}}^{\mathrm{T}}\mathbf{M}_{\mathbf{r}_{jk}}$$

... where \mathbf{r}_{ik} is the vector from the center of cell i to the point on the surface of that cell associated with contact k.

If i = j, then:

$$\mathbf{A}_{ii} = \sum_k \mathbf{M}_{\mathbf{r}_{ik}}^{\mathrm{T}} \mathbf{M}_{\mathbf{r}_{ik}}$$

 \ldots where k ranges over all constraints involving cell i. In general:

$$\mathbf{M_{s}^{T}M_{r}} = \begin{bmatrix} \frac{1}{M_{1}M_{2}} & 0 & 0 & 0 & \frac{\mathbf{s}_{z}}{I_{2}} & -\frac{\mathbf{s}_{y}}{I_{2}} & \frac{\mathbf{a}_{\mathbf{s}_{x}}}{2l_{2}G_{2}} \\ 0 & \frac{1}{M_{1}M_{2}} & 0 & -\frac{\mathbf{s}_{z}}{I_{2}} & 0 & \frac{\mathbf{s}_{x}}{I_{2}} & \frac{\mathbf{a}_{\mathbf{s}_{y}}}{2l_{2}G_{2}} \\ 0 & 0 & \frac{1}{M_{1}M_{2}} & \frac{\mathbf{s}_{y}}{I_{2}} & -\frac{\mathbf{s}_{x}}{I_{2}} & 0 & \frac{\mathbf{a}_{\mathbf{s}_{z}}}{2l_{2}G_{2}} \\ 0 & -\frac{\mathbf{r}_{z}}{I_{1}} & \frac{\mathbf{r}_{y}}{I_{1}} & \frac{\mathbf{r}_{y}\mathbf{s}_{y}+\mathbf{r}_{z}\mathbf{s}_{z}}{I_{1}I_{2}} & -\frac{\mathbf{r}_{y}\mathbf{s}_{x}}{I_{1}I_{2}} & -\frac{\mathbf{r}_{z}\mathbf{s}_{x}}{I_{1}I_{2}} & \frac{\mathbf{r}_{y}\mathbf{a}_{\mathbf{s}_{z}}}{2l_{1}l_{2}G_{2}} \\ \frac{\mathbf{r}_{z}}{I_{1}} & 0 & -\frac{\mathbf{r}_{x}}{I_{1}} & -\frac{\mathbf{r}_{x}\mathbf{s}_{y}}{I_{1}I_{2}} & \frac{\mathbf{r}_{x}\mathbf{s}_{x}+\mathbf{r}_{z}\mathbf{s}_{z}}{I_{1}I_{2}} & -\frac{\mathbf{r}_{z}\mathbf{s}_{y}}{I_{1}I_{2}} & \frac{\mathbf{r}_{z}\mathbf{a}_{\mathbf{s}_{x}}}{2l_{1}l_{2}G_{2}} \\ -\frac{\mathbf{r}_{y}}{I_{1}} & \frac{\mathbf{r}_{x}}{I_{1}} & 0 & -\frac{\mathbf{r}_{x}\mathbf{s}_{z}}{I_{1}I_{2}} & -\frac{\mathbf{r}_{y}\mathbf{s}_{z}}{I_{1}I_{2}} & \frac{\mathbf{r}_{y}\mathbf{s}_{y}+\mathbf{r}_{z}\mathbf{s}_{z}}{I_{1}I_{2}G_{2}} & \frac{\mathbf{r}_{x}\mathbf{a}_{\mathbf{s}_{y}}}{2l_{1}l_{2}G_{2}} \\ \frac{\mathbf{a}_{\mathbf{r}_{x}}}{2l_{1}G_{1}} & \frac{\mathbf{a}_{\mathbf{r}_{y}}}{2l_{1}G_{1}} & \frac{\mathbf{s}_{y}\mathbf{a}_{\mathbf{r}_{z}}-\mathbf{s}_{z}\mathbf{a}_{\mathbf{r}_{y}}}{2l_{2}l_{1}G_{1}} & \frac{\mathbf{s}_{x}\mathbf{a}_{\mathbf{r}_{y}}-\mathbf{s}_{y}\mathbf{a}_{\mathbf{r}_{x}}}{2l_{2}l_{1}G_{1}} & \frac{\mathbf{a}_{\mathbf{r}^{*}}\mathbf{a}_{\mathbf{s}_{z}}}{4l_{1}l_{2}G_{1}G_{2}} \end{bmatrix}$$

 \dots and:

$$\mathbf{M_r^T} \mathbf{M_r} = \begin{bmatrix} \frac{1}{M_1^2} & 0 & 0 & 0 & \frac{\mathbf{r_z}}{I_1} & -\frac{\mathbf{r_y}}{I_1} & \frac{\mathbf{a_{r_x}}}{2l_1G_1} \\ 0 & \frac{1}{M_1^2} & 0 & -\frac{\mathbf{r_z}}{I_1} & 0 & \frac{\mathbf{r_z}}{I_1} & \frac{\mathbf{a_{r_y}}}{I_1} \\ 0 & 0 & \frac{1}{M_1^2} & \frac{\mathbf{r_y}}{I_1} & -\frac{\mathbf{r_x}}{I_1} & 0 & \frac{\mathbf{a_{r_z}}}{2l_1G_1} \\ 0 & -\frac{\mathbf{r_z}}{I_1} & \frac{\mathbf{r_y}}{I_1} & \frac{\mathbf{r_y^2 + r_z^2}}{I_1^2} & -\frac{\mathbf{r_y r_x}}{I_1^2} & -\frac{\mathbf{r_z r_x}}{I_1^2} & \frac{\mathbf{r_y a_{rz}}}{2l_1l_1G_1} \\ \frac{\mathbf{r_z}}{I_1} & 0 & -\frac{\mathbf{r_x}}{I_1} & -\frac{\mathbf{r_x r_y}}{I_1^2} & \frac{\mathbf{r_x^2 + r_z}}{I_1^2} & -\frac{\mathbf{r_z r_x}}{I_1^2} & \frac{\mathbf{r_z a_{s_x}}}{2l_1l_1G_1} \\ -\frac{\mathbf{r_y}}{I_1} & \frac{\mathbf{r_x}}{I_1} & 0 & -\frac{\mathbf{r_x r_y}}{I_1^2} & -\frac{\mathbf{r_y r_z}}{I_1^2} & \frac{\mathbf{r_y^2 + r_z}}{I_1^2} & \frac{\mathbf{r_x a_{r_y}}}{I_1^2} \\ \frac{\mathbf{a_{r_x}}}{2l_1G_1} & \frac{\mathbf{a_{r_y}}}{2l_1G_1} & \frac{\mathbf{a_{r_z}}}{2l_1l_1G_1} & \frac{\mathbf{r_y a_{r_z} - r_z a_{r_y}}}{2l_1l_1G_1} & \frac{\mathbf{r_z a_{r_x} - r_x a_{r_z}}}{2l_1l_1G_1} & \frac{\mathbf{r_x a_{r_y} - r_y a_{r_x}}}{2l_1l_1G_1} & \frac{\mathbf{a_{r_x} a_{r_x} - r_x a_{r_y}}}{2l_1l_1G_1} & \frac{\mathbf{a_{r_x} a_{r_x} - r_x a_{r_x}}}{2l_1l_1G_1} & \frac{\mathbf{a_{r_x} a_{r_x} - r_x a_{r_x}}}{2l_1l_1G_1} & \frac{\mathbf{a_{r_x} a_{r_x} - r_x a_{r_x}}}{2l_1l_1G_1} & \frac{\mathbf{a_{r_x} a_{r_x}$$

 \dots where:

$$\mathbf{a_r} = (\hat{\mathbf{x}}_1 \cdot \mathbf{r}) \, \mathbf{r}$$
$$\mathbf{a_s} = (\hat{\mathbf{x}}_2 \cdot \mathbf{s}) \, \mathbf{s}$$

So we are left with the equation:

$$\mathbf{A}\Delta\mathbf{p} = \Delta\mathbf{v}$$

 \dots where:

$$\mathbf{A}_{ij} = \left\{ \begin{array}{ll} \mathbf{M}_i^{\mathrm{T}} \mathbf{M}_i & \text{ if } i = j \\ -\mathbf{M}_i^{\mathrm{T}} \mathbf{M}_j & \text{ if cell } i \text{ and cell } j \text{ are colliding} \end{array} \right\}$$