Search for WIMP Inelastic Scattering Off Xenon Nuclei With Xenon100 Data

M. Alfonsi, ¹ E. Aprile, ² K. Arisaka, ³ F. Arneodo, ⁴ C. Balan, ⁵ L. Baudis, ⁶ A. Behrens, ⁶ P. Beltrame, ³ K. Bokeloh, ⁷ E. Brown, ⁷ G. Bruno, ⁴ R. Budnik, ² J. M. R. Cardoso, ⁵ W.-T. Chen, ⁸ B. Choi, ² H. Contreras, ² J. P. Cussonneau, ⁸ M. P. Decowski, ¹ E. Duchovni, ⁹ S. Fattori, ¹⁰ A. D. Ferella, ⁶ W. Fulgione, ¹¹ F. Gao, ¹² M. Garbini, ¹³ K.-L. Giboni, ² L. W. Goetzke, ² C. Grignon, ¹⁰ E. Gross, ⁹ W. Hampel, ¹⁴ H. Kettling, ⁷ A. Kish, ⁶ J. Lamblin, ⁸ H. Landsman, ⁹ R. F. Lang, ^{15, 2} M. Le Calloch, ⁸ C. Levy, ⁷ K. E. Lim, ² Q. Lin, ¹² S. Lindemann, ¹⁴ M. Lindner, ¹⁴ J. A. M. Lopes, ⁵ K. Lung, ³ T. Marrodán Undagoitia, ^{6, *} F. V. Massoli, ¹³ Y. Mei, ^{16, 10} A. J. Melgarejo Fernandez, ² Y. Meng, ³ A. Molinario, ¹¹ E. Nativ, ⁹ K. Ni, ¹² U. Oberlack, ^{16, 10} S. E. A. Orrigo, ⁵ E. Pantic, ^{3, †} D. Pätzold, ¹⁰ R. Persiani, ¹³ G. Plante, ² N. Priel, ⁹ A. C. C. Ribeiro, ⁵ A. Rizzo, ² S. Rosendahl, ⁷ J. M. F. dos Santos, ⁵ G. Sartorelli, ¹³ J. Schreiner, ¹⁴ M. Schumann, ⁶ L. Scotto Lavina, ⁸ P. R. Scovell, ³ M. Selvi, ¹³ P. Shagin, ¹⁶ H. Simgen, ¹⁴ A. Teymourian, ³ D. Thers, ⁸ O. Vitells, ⁹ H. Wang, ³ M. Weber, ¹⁴ and C. Weinheimer, ⁷ (The XENON100 Collaboration)

¹Nikhef and the University of Amsterdam, Science park, Amsterdam, Netherlands ²Physics Department, Columbia University, New York, NY 10027, USA ³Physics & Astronomy Department, University of California, Los Angeles, USA ⁴INFN, Laboratori Nazionali del Gran Sasso, Assergi, 67100, Italy ⁵Department of Physics, University of Coimbra, R. Larga, 3004-516, Coimbra, Portugal ⁶Physics Institute, University of Zürich, Winterthurerstr. 190, CH-8057, Switzerland ⁷Institut für Kernphysik, Wilhelms-Universität Münster, 48149 Münster, Germany ⁸SUBATECH, Ecole des Mines de Nantes, CNRS/In2p3, Université de Nantes, 44307 Nantes, France ⁹Department of Particle Physics and Astrophysics, Weizmann Institute of Science, 76100 Rehovot, Israel ¹⁰Institut für Physik, Johannes Gutenberg Universität Mainz, 55099 Mainz, Germany ¹¹University of Torino and INFN-Torino, Torino, Italy ¹²Department of Physics, Shanghai Jiao Tong University, Shanghai, 200240, China ¹³ University of Bologna and INFN-Bologna, Bologna, Italy ¹⁴Max-Planck-Institut für Kernphysik, Saupfercheckweg 1, 69117 Heidelberg, Germany ¹⁵Department of Physics, Purdue University, West Lafayette, IN 47907, USA ¹⁶Department of Physics and Astronomy, Rice University, Houston, TX 77005 - 1892, USA (Dated: January 10, 2017)

Some nice abstract here.....

I. INTRODUCTION

Astrophysical evidence indicates that the dominant mass fraction of our Universe consists of some yet unsknown form of dark matter. Well motivated models predict Dark Matter in the form of Weakly Interacting Massive Particles (WIMPs), hypothesis which is currently being tested by several direct and indirect detection experiment. some citation missing.

Most of direct detection searches focuses on elastic scattering of dark matter particles off nuclei. In this analysis instead an inelastic scattering process is explored:
we consider the ¹²⁹Xe isotope being excited to a lowlying state with subsequent prompt de-excitation via the emission of a photon. This isotope is an excellent target since its abundance in natural xenon is of 26.4% and a relatively low energy is necessary to excite its 3/2+ state above the 1/2+ spin ground state. These type of processes were previously studied in detail for liquid xenon detectors in [1]. Inelastic WIMP-nucleus scattering in

10

11

12

13

14

15

16

17

18

19

20

21

22

23

24

25

26

27

28

29

 $_{49}$ xenon is complementary to elastic scattering for spin- $_{50}$ dependent interactions, the former dominating the inte- $_{51}$ grated rate above $\simeq 10~{\rm keV}$ of energy deposition. Fur- $_{52}$ thermore, in case of dark matter detection, this channel $_{53}$ can be employed to asses whether the nature of the fun- $_{54}$ damental interaction is spin-dependent or not.

II. XENON100 DETECTOR

The Xenon100 experiment is a dual phase liquid xenon TPC. For a given interaction in the liquid target this type of detector produces two separated signals, one proportional to the prompt scintillation (S1) the other to ionization (S2).

To add: describe and explain cS1 and cS2 definitions, some sentences about detector stability, maybe science run data used and calibration campaign goes here, maybe Ly and Y measurements used goes here. Note that the corrected S2 observed by the bottom PMT array is used in this analysis.

^{*} marrodan@physik.uzh.ch

[†] pantic@physics.ucla.edu

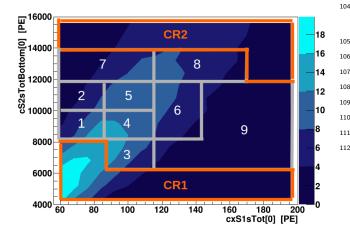


FIG. 1. Signal region and control region.

DATA ANALYSIS

The inelastic scattering of a WIMP with the nucleus of ¹²⁹Xe produces an energy deposit via nuclear recoil with subsequent emission of a 39.6 keV de-excitation photon. The largest fraction of the energy released in the event is via electronic recoil due to the emitted photon, this represents an unusual signature for this kind of detector and brings the possible signal to overlay a phase space region with large backgrounds. The chosen region of interest (ROI) for this analysis surrounds the 39.6 keV xenon line in the cS1-cS2 plane and is further divided into subregions, as shown in Figure 1 and 2.

Events are asked, other than falling in the defined region of interest, to fulfill several selection criteria which can be summarized as: selection aimed to reduce noise impact and including energy (S1) and S2 thresholds, veto of events with energy release in the detector's outer volume, and additionally, events must be of single scatter 92 nature and fall into a predefined fiducial volume. This 93 analysis follows the selection criteria described in detail 94 in [2] for Run-II, with only few exceptions reported in what follows. The selection on S2 width as a function of drift time has been optimized on a sample of events selected from the 39.6 keV line and set to a 95% acceptance on these. Events are required to be single scatter by applying a threshold on the second largest S2 peak 100 size, for this analysis this threshold has been optimized 101 to 160 PE and set constant as function of S2 signal size. 102 Finally the chosen fiducial volume corresponds to 34 Kg 103 of liquid xenon.

Signal Simulation

The detector response to inelastic scattering of WIMPs off ¹²⁹Xe nucleus was simulated using an empirical model. The total deposited energy is divided into two independent contributions: one related to the 39.6 keV deexcitation photon and the other relative to the simulta-110 neous nuclear recoil of the xenon atom. The number of photons and charge yield detected is simulated separately 112 for each contribution and then added together.

The distribution in the cS1-cS2 plane of an ER induced by the de-excitation photon is simulated assuming a two dimensional normal pdf, f(cS1, cS2), described except a normalization factor in equation 1:

$$f(cS1, cS2) = exp\left(-\frac{1}{2(1-\rho^2)} \left[\frac{(cS1 - \mu_{cS1})^2}{\sigma_{cs1}^2} + \frac{(cS2 - \mu_{cS2})^2}{\sigma_{cs2}^2} - \frac{2\rho(cS1 - \mu_{cs1})(cS2 - \mu_{cs2})}{\sigma_{cs1}\sigma_{cs2}} \right] \right)$$
(1)

where μ_{cS1} and μ_{cS2} represents the average observed cS1 This analysis is performed using XENON100 Run- 114 and cS2 given a 39.6 keVee, σ_{cs1} and σ_{cs2} are the stan- $_{69}$ II science data, which corresponds to an exposure of $_{115}$ dard deviation in cS1 and cS2 respectively, while ho stands 224.6 live days. The detector response to electronic re- 116 for the correlation between cS1 and cS2. The detector recoil (ER) has been characterized using 60 Co and 232 Th 117 lated light yield L_y at 39.6 keV, necessary to evaluate the radioactive sources, while response to inelastic nuclear 118 average number of photon detected (μ_{cS1}), is obtained recoil (NR) scattering was calibrated using an ²⁴¹AmBe ¹¹⁹ as the result of a NEST model [3–5] fit to data collected 120 with several lines. The same model is used to predict the 121 charge yield at 39.6 keV which is then scaled according to the detector's secondary scintillation gain Y, determined from detector response to single electrons [6]. The detector resolution at 39.6 keV in cS1 and cS2 has been measured to be respectively 15.8% and 14.7% and used 126 to extract the standard deviations σ_{cs1} , σ_{cs2} . The corre-127 lation parameter is assumed to be independent of energy (at least in the considered range) and measured using the 164 keV xenon activated line by ¹²⁴AmBe calibration 130 data, this line is chosen since allows to disentangle effi-131 ciently contribution from nuclear recoil. The measured ₁₃₂ correlation is $\rho = -0.45 \pm 0.10$.

> The cS1 and cS2 distributions from NR contribution 134 are predicted starting from the nuclear recoil energy spec-135 trum of WIMPs inelastic interaction [1]. The average 136 cS1 and cS2 are given by equations 2 and 3 respectively, where \mathcal{L}_{eff} is the liquid xenon NR relative scintillation 138 efficiency, while $S_{ee}=0.58$ and $S_{nr}=0.95$ describe the 139 scintillation quenching due to the electric field [7]. The 140 parameterization and uncertainties of \mathcal{L}_{eff} as a function of E_{nr} are based on existing direct measurements [8]. The 142 light yield at 122 keVee originate from the same NEST 143 model fit as described above. For the cS2 the parameter-144 ization of $Q_Y(E_{nr})$ is taken from [9]. Finally all detec-145 tor related resolution effects are introduced following the 146 prescriptions described in [2].

$$cS1_{nr} = E_{nr} \mathcal{L}_{eff}(E_{nr}) L_y \frac{S_{nr}}{S_{ee}}$$
 (2)

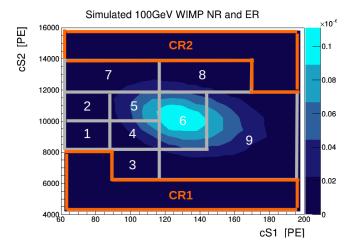


FIG. 2. Signal region and control region, for WIMP of mass 100 GeV.

$$cS2_{nr} = E_{nr} Q_Y(E_{nr}) Y (3$$

The pdf of the ER and NR contributions are then convoluted together to obtain the overall pdf of the signal. A 2D (cS1 versus cS2) acceptance map is applied to the sig-150 nal pdf to reproduce data selection effects. Acceptances are computed separately for each selection criteria using ¹²⁴AmBe calibration sample, selections as the outer volume veto and the single scatter interaction represent an exception and a dedicated computation has been performed in these cases. The selections acceptance average in the region of interest to about 0.80 ± 0.05 . Figure 1 shows an example of full simulated signal model for a WIMP of 100 GeV mass.

The signal simulation procedure has been validated re-160 producing the 39.6 keV xenon line from interaction due 194 161 to ¹²⁴AmBe source and has been compared to data. For 162 this comparison the proper ¹²⁴AmBe nuclear recoil and ₁₉₅ cal uncertainties.

В. **Background Model**

166

167 168 region of interest is due to eviromental and material 204 that uncertainties belonging to different sub-regions are radioactivity, its composition is mainly represented by 205 considered independent from each other. Compton scattering photons. Background contribution 206 radiogenic neutrons is expected to be negligible.

calibration campaign, which are assumed to well repre- 210 pothesis. sent the background density distribution in the cS1-cS2 211 176 plane. The calibration sample yields about 22'000 events 212 related to detector response are parameterized using the ₁₇₇ in the ROI, these are then scaled to data according to ₂₁₃ respective uncertainties on the measure of L_y , \mathcal{L}_{eff} , Y, τ_{178} a measured scale factor τ_{bkq} . This scale factor, which is τ_{174} τ_{174} τ_{174} and τ_{174} τ_{174}

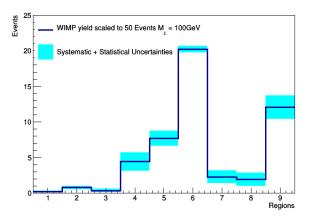


FIG. 3. Signal region, uncertainties for WIMP of mass 100 GeV.

179 merely the ratio between the data and calibration sample 180 yields, is measured in the two control regions shown in 181 Figure 1 and labelled CR1 and CR2. The two control re-182 gions give compatible results and the computed average is $\tau_{bkg} = 0.034 \pm 0.002$, where the reported uncertainty 184 is of statistical nature only.

The distribution of the calibration sample has been 186 compared to the data of the science run in the two con-187 trol regions, agreement is found within statistical uncer-188 tainties. Furthermore, ⁶⁰Co calibration data have been 189 compared in the region of interest to data from ²³²Th 190 calibration campaign, the largest deviation between the 191 two shapes is within 4%. An additional systematic uncer-192 tainty of 4% has been applied to the expected background 193 yield of each sub-region of the ROI.

Systematic Uncertainties

Uncertainties on the total prediction of background acceptances were simulated. The simulated events were 196 events arise from the uncertainty on the measure of the found in agreement with calibration data within statisti- 197 normalization factor, τ_{bkg} , and amount to 6%, contri-198 bution of radiogenic neutrons are neglected. Systematic uncertainty on the shape of the predicted background dis-200 tribution are assessed by the maximal discrepancy in the 201 ROI between the ⁶⁰Co and ²³²Th calibration samples, 202 a 4% systematic additional to statistical uncertainty is The main expected background contribution in the 203 assigned to the expected yield of each sub-region. Note

Uncertainty on the total yield of signal arising from due to the activation of the xenon 39.6 keV line from 207 selections acceptance uncertainties are found to be very 208 weakly dependent on the WIMP mass, an overall 6% ac-The background is modeled using data from the ⁶⁰Co ²⁰⁰ ceptance uncertainty is then applied to all WIMPs hy-

Uncertainties on the energy scale and, more generally,

 215 certainties mainly affect the pdf of the signal model in 216 the ROI, and very weakly the total signal yield. They are 217 taken into account by simulating several signal pseudo- 218 sample for each WIMP mass, the pseudo-samples are 219 produced varying the model parameters respectively of 220 ± 1 standard deviation. For each sub-region is then com- 221 puted an overall uncertainty by adding in quadrature the 222 variations of each pseudo-sample with respect to nominal. 223 Figure 3 is an example of such a systematic uncertainty 224 computation for a WIMP of 100 GeV mass.

All the uncertainties discussed here are parameterized within a binned profiled likelihood function using the framework [10, 11]. All parameters related to systematic uncertainties are assumed to be normally distributed.

IV. RESULTS

229

Using an exposure of 34 kg of liquid xenon and 231 224.6 live days of data a yield of 764 events is observed in 232 the region of interest, this is compatible with the expectation of 756 \pm 5 (stat.) \pm 55 (syst.) events from the back-234 ground only hypothesis. Figure 4 shows how these events 235 are distributed in the region of interest, the bottom panel 236 shows the ratio between data and expected background, 237 where the gray and orange shaded areas represent re-238 spectively statistical and systematic uncertainty on the 239 background expectation.

This result is interpreted via a binned profiled like-lihood approach by means of the test statistic \tilde{q} and its asymptotic distributions described in [12]. Assuming an isothermal WIMP halo with a local density of $\rho_{\chi}=0.3~GeV/cm^3$, a local circular velocity of $v_0=220~{\rm km/s}$, and a galactic escape velocity of $v_{esc}=544~{\rm km/s}$, other assumptions..., a 90% CL_s [13] confidence level limit is computed on the spin dependent inelastic WIMP-nucleon cross section, σ_{inel} , as a function of the WIMP mass, m_{χ} , and shown in Figure 5. The expected median sensitivity is reported with its relative one (green area) and two (yellow area) standard deviation uncertainty. A limit is set on σ_{inel} to $3.3\times 10^{-38}~cm^2$ at 90% CL_s confidence level for a WIMP of mass 100 GeV. This limit is compared with decide which other experiment to plot.

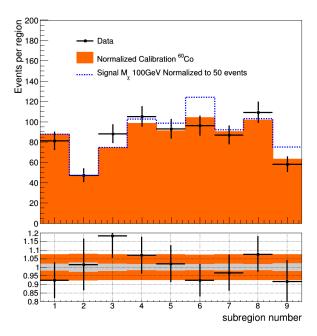


FIG. 4. Results, comparison between data and expected background.

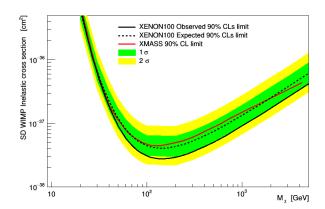


FIG. 5. Observed and expected limits.

- [1] L. Baudis, G. Kessler, P. Klos, R. F. Lang, J. Men- 278 255 ndez, S. Reichard, and A. Schwenk. Signatures of Dark 279 256 Matter Scattering Inelastically Off Nuclei. Phys. Rev., 280 257 D88(11):115014, 2013. 258
- [2] E. Aprile et al. XENON100 Dark Matter Results 282 259 from a Combination of 477 Live Days. Phys. Rev., 283 260 261 D94(12):122001, 2016.

262

263

264

265

267

270

271

272

273

- M Szydagis, A Fyhrie, D Thorngren, and M Tripathi. En- 285 hancement of nest capabilities for simulating low-energy 286 recoils in liquid xenon. Journal of Instrumentation, 287 8(10):C10003, 2013.
- 266 tions. IEEE Trans. Nucl. Sci., 53:270, 2006.
- S. Agostinelli et al. GEANT4: A Simulation toolkit. 291 268 Nucl. Instrum. Meth., A506:250-303, 2003. 269
 - E Aprile et al. Observation and applications of single- 293 electron charge signals in the xenon100 experiment. 294 Journal of Physics G: Nuclear and Particle Physics, 295 41(3):035201, 2014.
- [7]E. Aprile, C. E. Dahl, L. DeViveiros, R. Gaitskell, K. L. 297 274 Giboni, J. Kwong, P. Majewski, Kaixuan Ni, T. Shutt, 298 275 and M. Yamashita. Simultaneous measurement of ioniza- 299 276 tion and scintillation from nuclear recoils in liquid xenon 277

- as target for a dark matter experiment. Phys. Rev. Lett., 97:081302, 2006.
- [8] E. Aprile et al. Dark Matter Results from 100 Live Days of XENON100 Data. Phys. Rev. Lett., 107:131302, 2011.
- E. Aprile et al. Response of the XENON100 Dark Matter Detector to Nuclear Recoils. Phys. Rev., D88:012006, 2013.
- [10] Lorenzo Moneta, Kevin Belasco, Kyle S. Cranmer, S. Kreiss, Alfio Lazzaro, Danilo Piparo, Gregory Schott, Wouter Verkerke, and Matthias Wolf. The RooStats Project. PoS, ACAT2010:057, 2010.
- John Allison et al. Geant4 developments and applica- 289 [11] Wouter Verkerke and David P. Kirkby. The RooFit toolkit for data modeling. eConf, C0303241:MOLT007, 290 2003. [,186(2003)].
 - 292 [12] Glen Cowan, Kyle Cranmer, Eilam Gross, and Ofer Vitells. Asymptotic formulae for likelihood-based tests of new physics. Eur. Phys. J., C71:1554, 2011. [Erratum: Eur. Phys. J.C73,2501(2013)].
 - ²⁹⁶ [13] Alexander L. Read. Modified frequentist analysis of search results (The CL(s) method). In Workshop on confidence limits, CERN, Geneva, Switzerland, 17-18 Jan 2000: Proceedings, pages 81-101, 2000.