

EBU6018

Advanced Transform Methods

Week 4.3 – Wigner-Ville Distribution (WVD)

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Lecture Outline

➤ Wigner-Ville Distribution

- ❑ Definition
- ❑ Pros and cons
- ❑ Connection to other transform methods

Wigner-Ville Distribution - Background

- So far we have looked at transforms that compute the correlation between a signal and basis functions that are functions of time and frequency (or of scale and translation). The time-frequency resolution is determined by the basis functions
 - E.g., STFT, Wavelet Transform
- An alternative approach is to compute directly:
 - ❑ Time-frequency energy density- signal's energy density in both time and frequency
 - In contrast to power spectrum: energy in frequency only
 - ❑ An example of this is the Wigner-Ville distribution

Wigner-Ville Distribution - Background

- Comparison of STFT and CWT

- Similarities:

- They are both Windowed Transforms
 - signal is multiplied by a function, and the transform is computed separately for different segments of signals.
 - can be written in inner product form

$$STFT(b, \omega) = \left\langle s(t), \gamma(t-b)e^{j\omega t} \right\rangle$$

$$CWT(b, a) = \left\langle s(t), \frac{1}{\sqrt{a}} \psi \left(\frac{t-b}{a} \right) \right\rangle$$

- Difference:

- Fixed time duration and frequency bandwidths of $\gamma(t)$
 - Variable time duration and bandwidth of $\psi(t)$

Wigner-Ville Distribution – Energy Distribution

- The purpose of energy distributions is to **distribute the energy of the signal over time and frequency**.
- The energy of a signal $s(t)$ can be found from the **squared modulus** of either **the signal** or its **Fourier Transform**:
- $E_s = \int |s(t)|^2 dt = \int |s(\omega)|^2 d\omega$
- $|s(t)|^2$ and $|s(\omega)|^2$ can be interpreted as energy densities in time and frequency respectively.

Wigner-Ville Distribution – Wiener-Khinchin Theorem

- The energy density and autocorrelation function of a signal are related by the Wiener-Khinchin Theorem.
- According to the Wiener-Khinchin theorem, the power spectrum is the **Fourier Transform** of the **Autocorrelation Function**.

$$P(\omega) = |s(\omega)|^2 = \int_{-\infty}^{\infty} R(\tau) e^{-j\omega\tau} d\tau$$

(For an example of the derivation look at:

[https://mathworld.wolfram.com/Wiener-Khinchin Theorem.html](https://mathworld.wolfram.com/Wiener-Khinchin%20Theorem.html))

Wigner-Ville Distribution – Autocorrelation

$$R(\tau) = \int s(t)s(t + \tau)dt$$

τ is the shift of the signal with respect to itself.

- In the standard autocorrelation function, **time is integrated out of the result, and $R(\tau)$ is a function of only the time lag τ .**
- There is a class of distribution called Cohen Class (of which the Wigner-Ville Distribution is a member) that uses a **variation of the autocorrelation function where time remains in the result**, this is the Instantaneous Autocorrelation Function:

$$R(t, \tau) = s(t + \tau/2)s^*(t - \tau/2) \quad (\text{If the signal is real, then } s^* = s)$$

- where τ is the time lag and $*$ represents the complex conjugate of signal s .

Wigner-Ville Distribution – Instantaneous power spectrum

Recall the power spectrum :

$$P(\omega) = |S(\omega)|^2 = \int_{-\infty}^{\infty} R(\tau) e^{-j\omega\tau} d\tau$$

Power spectrum of a signal is the Fourier Transform of its autocorrelation function

where $R(\tau)$ is the autocorrelation function (acf)

$$R(\tau) = \int_{-\infty}^{\infty} s(t) s^*(t - \tau) dt = \int_{-\infty}^{\infty} s(t + \tau/2) s^*(t - \tau/2) dt$$

What happens if we use *instantaneous* autocorrelation :

$$R(t, \tau) = s(t + \tau/2) s^*(t - \tau/2)$$

instead of $R(\tau) = \int_{-\infty}^{\infty} R(t, \tau) dt$? We get :

The WVD is the Fourier Transform of the instantaneous autocorrelation function

$$WVD_s(t, \omega) = \int_{-\infty}^{\infty} s(t + \tau/2) s^*(t - \tau/2) e^{-j\omega\tau} d\tau$$

which is the *Wigner - Ville Distribution* (WVD).

Wigner-Ville Distribution

- The Wigner-Ville Distribution is not a Windowed Transform in a similar way to the STFT
- However, the **shifted version of the same signal** in the autocorrelation function **could be considered as a window**.
- The WVD compares the information in the signal with its own information at other times and frequencies.
- The WVD has some interesting and useful properties.

Wigner-Ville Distribution

➤ Cross-WVD vs auto-WVD

Since we can define a cross-correlation,
we can also define a **cross-Wigner-Ville distribution**:

$$WVD_{s,g}(t, \omega) = \int_{-\infty}^{\infty} s(t + \tau/2) g^*(t - \tau/2) e^{-j\omega\tau} d\tau$$

Taking complex conjugates we find

$$WVD_{s,g}(t, \omega) = WVD_{g,s}^*(t, \omega)$$

So for the usual WVD (**auto-WVD**) we have

$$WVD_s(t, \omega) = WVD_{s,s}(t, \omega) = WVD_s^*(t, \omega)$$

so the auto-WVD is always real.

Wigner-Ville Distribution

➤ Example: Gaussian function

Signal $s(t) = \sqrt{\frac{\alpha}{\pi}} e^{-\alpha t^2/2}$ (normalized to unit energy)

For WVD, we get

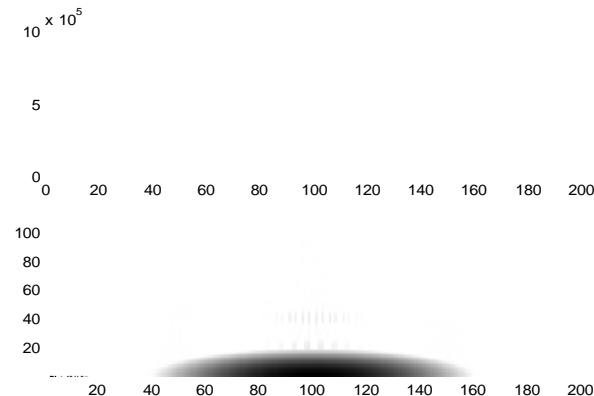
$$\begin{aligned} WVD_s(t, \omega) &= \sqrt{\frac{\alpha}{\pi}} \int_{-\infty}^{\infty} \exp \left\{ -\frac{\alpha}{2} \left[\left(t + \frac{\tau}{2} \right)^2 + \left(t - \frac{\tau}{2} \right)^2 \right] \right\} e^{-j\omega\tau} d\tau \\ &= e^{-\alpha t^2} \sqrt{\frac{\alpha}{\pi}} \int_{-\infty}^{\infty} \exp \left\{ -\frac{\alpha}{4} \tau^2 \right\} e^{-j\omega\tau} d\tau \quad \text{Gaussian in time,} \\ &= 2 \exp \left\{ -\left[\alpha t^2 + \frac{1}{\alpha} \omega^2 \right] \right\} \quad \text{Gaussian in t-f} \end{aligned}$$

i.e. concentrated around (0,0).

α controls spread:

"time-width": $|t| < \sqrt{\frac{1}{\alpha}}$

"freq-width": $|\omega| < \sqrt{\alpha}$ Similar to $\Delta t \Delta \omega = k$ from the UP



Wigner-Ville Distribution

➤ Example: Gaussian chirplet

Signal:
$$s(t) = \sqrt[4]{\frac{\alpha}{\pi}} \exp \left\{ -\frac{\alpha}{2} t^2 + j \frac{\beta}{2} t^2 \right\}$$

Power spectrum:
$$|S(\omega)|^2 = \sqrt{\frac{4\pi(\alpha^2 + \beta^2)}{\alpha}} \exp \left\{ -\frac{\alpha}{\alpha^2 + \beta^2} \omega^2 \right\}$$

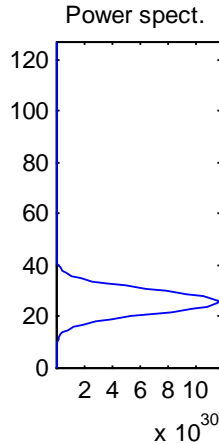
tells us which freqs $s(t)$ contains, not when. Compare:

$$\begin{aligned} WVDs(t, \omega) &= \sqrt{\frac{\alpha}{\pi}} \int_{-\infty}^{\infty} e^{-\frac{\alpha}{2} \left[\left(t + \frac{\tau}{2} \right)^2 + \left(t - \frac{\tau}{2} \right)^2 \right] + \frac{j\beta}{2} \left[\left(t + \frac{\tau}{2} \right)^2 - \left(t - \frac{\tau}{2} \right)^2 \right]} e^{-j\omega\tau} d\tau \\ &= e^{-\alpha t^2} \sqrt{\frac{\alpha}{\pi}} \int_{-\infty}^{\infty} e^{-\frac{\alpha}{4} \tau^2} e^{-j(\omega - \beta t)\tau} d\tau \\ &= 2e^{-\left[\alpha t^2 + \frac{1}{\alpha} (\omega - \beta t)^2 \right]} \end{aligned}$$

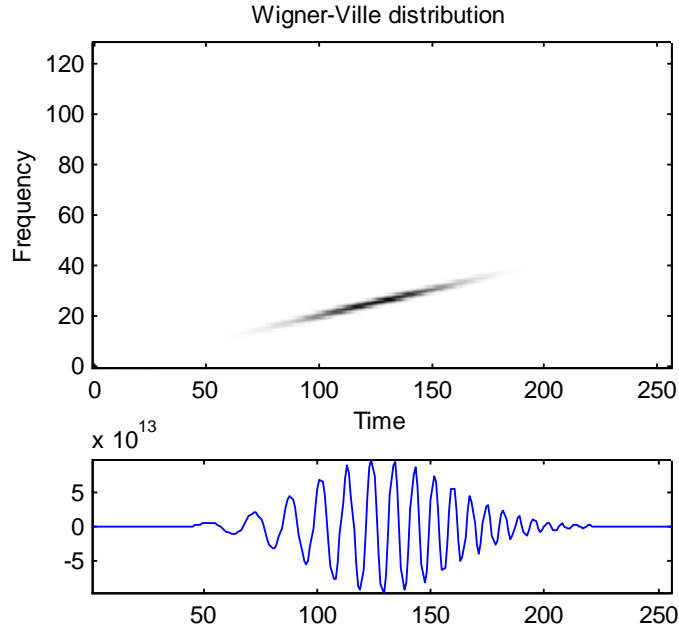
i.e. energy concentrated at $\omega = \beta t$, changing with time.

Wigner-Ville Distribution

➤ Illustration: chirplet



$$s(t) = \sqrt[4]{\frac{\alpha}{\pi}} \exp \left\{ -\frac{\alpha}{2} (t - t_0)^2 + j \frac{\beta}{2} t^2 \right\}$$



Wigner-Ville Distribution

➤ Time-limited or band-limited signals

- If $s(t)$ is time-limited, i.e. $s(t) = 0$ outside some interval $[t_0, t_1]$, then the WVD is also time-limited, i.e.

$$WVD_s(t, \omega) = 0 \text{ for } t \notin [t_0, t_1]$$

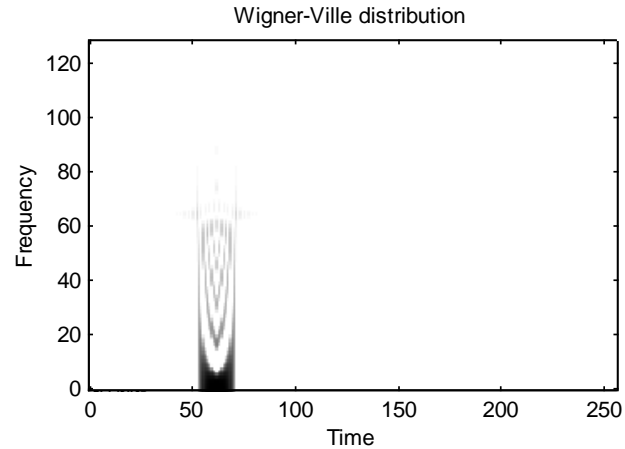
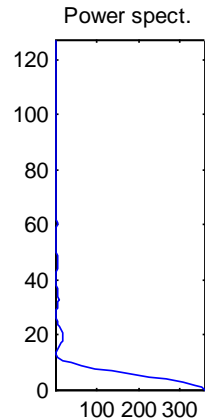
since no value for τ can make both $s(t + \tau/2)$ and $s(t - \tau/2)$ non-zero if t is outside this range.

(Either $s(t + \tau/2)$ or $s(t - \tau/2)$ will be zero)

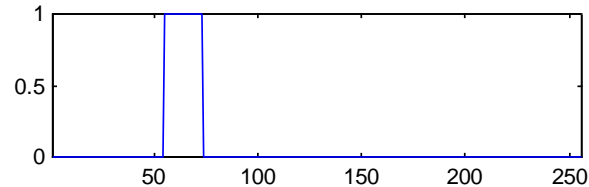
- A similar result is true for frequency band-limited signals.

Wigner-Ville Distribution

➤ Example: WVD of a Time-limited signal



The isolated pulse is time-bound
but not frequency-bound.
The FT is a sinc function.



Wigner-Ville Distribution

- WVD Properties: **Time Marginal Condition**

$$\begin{aligned}\frac{1}{2\pi} \int_{-\infty}^{\infty} WVD_s(t, \omega) d\omega &= \int_{-\infty}^{\infty} s\left(t + \frac{\tau}{2}\right) s^* \left(t - \frac{\tau}{2}\right) \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-j\omega\tau} d\omega d\tau \\ &= \int_{-\infty}^{\infty} s\left(t + \frac{\tau}{2}\right) s^* \left(t - \frac{\tau}{2}\right) \delta(\tau) d\tau \\ &= |s(t)|^2\end{aligned}$$

So integral over frequency of WVD is the signal power density at time t

Compare similar result for probability densities:

$$p_X(x) = \int_{-\infty}^{\infty} p_{X,Y}(x, y) dy$$

Wigner-Ville Distribution

➤ WVD Properties: **Frequency Marginal Condition**

$$\begin{aligned}\int_{-\infty}^{\infty} WVD_s(t, \omega) dt &= \int_{-\infty}^{\infty} s\left(t + \frac{\tau}{2}\right) s^* \left(t - \frac{\tau}{2}\right) \int_{-\infty}^{\infty} e^{-j\omega\tau} d\tau d\tau \\ &= \int_{-\infty}^{\infty} e^{-j\omega\tau} \int_{-\infty}^{\infty} s(t) s^*(t - \tau) dt d\tau \\ &= \int_{-\infty}^{\infty} e^{-j\omega\tau} R(\tau) d\tau \\ &= |S(\omega)|^2\end{aligned}$$

So integral over time of WVD is the power spectral density

We also have:

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} WVD_s(t, \omega) dt d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} |S(\omega)|^2 d\omega = \int_{-\infty}^{\infty} |s(t)|^2 dt$$

i.e. the WVD is unitary: the energy in $WVD_s(t, \omega)$ is equal to energy in original signal $s(t)$.

Wigner-Ville Distribution

➤ Time-shift & Freq-moduln. invariant

- If the WVD of $s(t)$ is $WVD_s(t, \omega)$,
then the WVD of time-shifted signal $s_0(t) = s(t - t_0)$
is a time-shifted WVD: $WVD_{s_0}(t, \omega) = WVD_s(t - t_0, \omega)$
- Further, the WVD of frequency-modulated signal
 $s_1(t) = s(t)e^{j\omega_1 t}$ is a frequency-shifted WVD:
 $WVD_{s_1}(t, \omega) = WVD_s(t, \omega - \omega_1)$

(Both follow immediately from the formulas for WVD)

Wigner-Ville Distribution

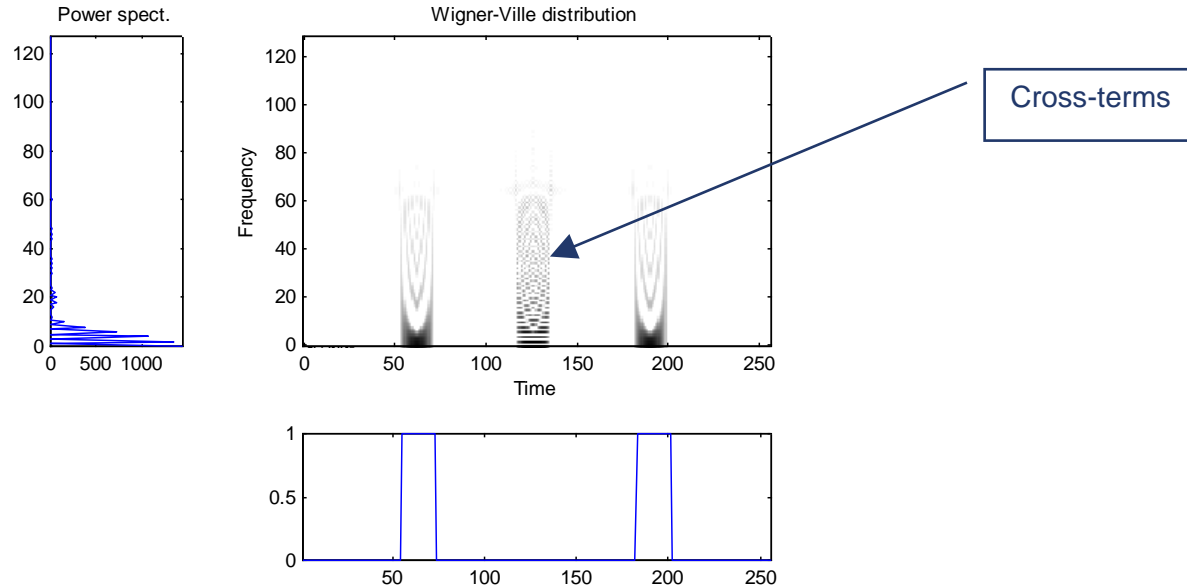
➤ WVD of multiple signals: Cross-terms

- Wigner-Ville Distribution has many useful properties, and better resolution than STFT spectrogram. BUT
- Applications are limited due to *cross-term interference*.
- Consider composite signal $s(t) = s_1(t) + s_2(t)$. Then
$$WVD_s(t, \omega) = WVD_{s_1}(t, \omega) + WVD_{s_2}(t, \omega) + 2 \operatorname{Re}\{WVD_{s_1, s_2}(t, \omega)\}$$
i.e. not only the sum of WVDs, but also *the cross-term* $WVD_{s_1, s_2}(t, \omega)$
- Also, cross-term is included at double the magnitude of the auto-terms, so often obscures useful patterns.

Wigner-Ville Distribution

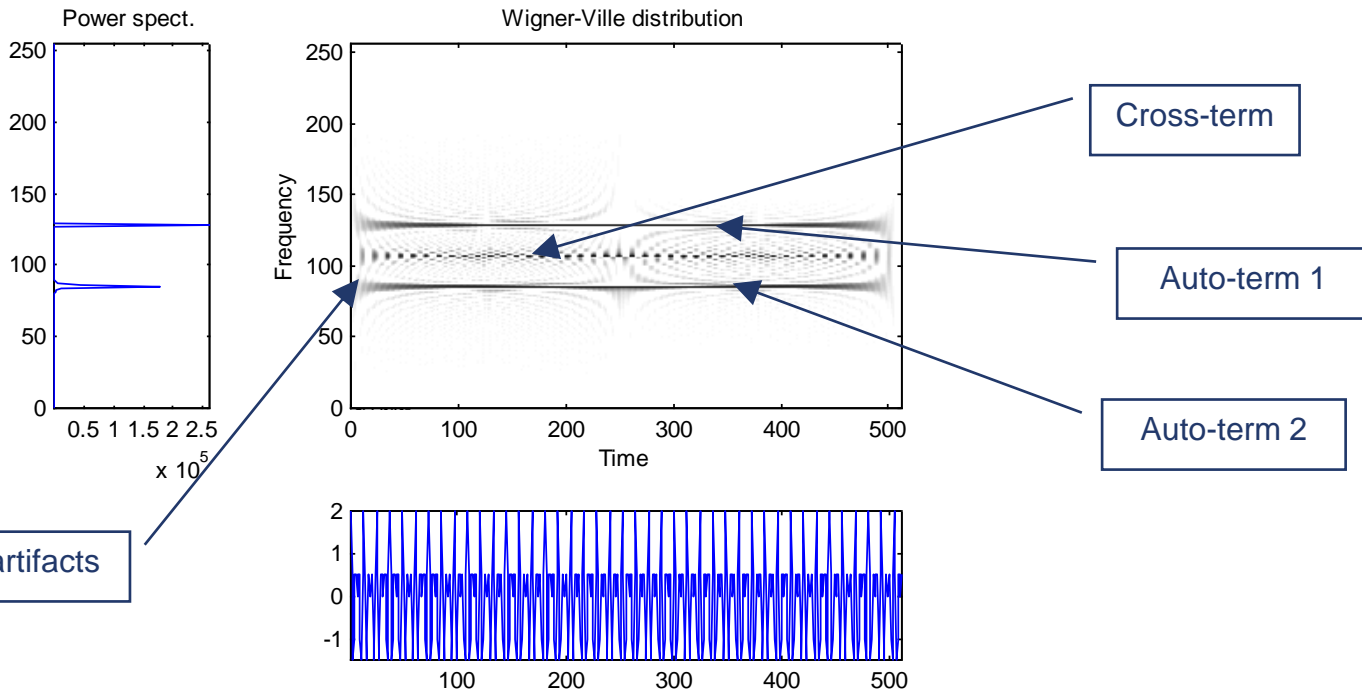
➤ Example of cross-terms

- ❑ WVD gives cross-terms for all but simplest signals, e.g.:



Wigner-Ville Distribution

➤ Example: sum of two sinusoids



Wigner-Ville Distribution

➤ Example: sum of two sinusoids

- If $s(t) = \exp(j\omega_0 t)$ then

$$WVD_s(t, \omega) = \int_{-\infty}^{\infty} \exp\left\{j\omega_0 \left(t + \frac{\tau}{2} - t + \frac{\tau}{2}\right)\right\} e^{-j\omega\tau} d\tau = 2\pi\delta(\omega - \omega_0)$$

i.e. the WVD is a "ridge" along frequency ω_0 .

- Now let $s(t) = \exp(j\omega_1 t) + \exp(j\omega_2 t)$. The power spectrum is

$$|S(\omega)|^2 = 2\pi\delta(\omega_1) + 2\pi\delta(\omega_2)$$

while the WVD is

$$WVD_s(t, \omega) = 2\pi\delta(\omega - \omega_1) + 2\pi\delta(\omega - \omega_2) + 4\pi\delta(\omega - \omega_\mu) \cos(\omega_d t)$$

where $\omega_\mu = \frac{\omega_1 + \omega_2}{2}$ and $\omega_d = \omega_1 - \omega_2$.

Wigner-Ville Distribution

➤ Example: cross-term (cont)

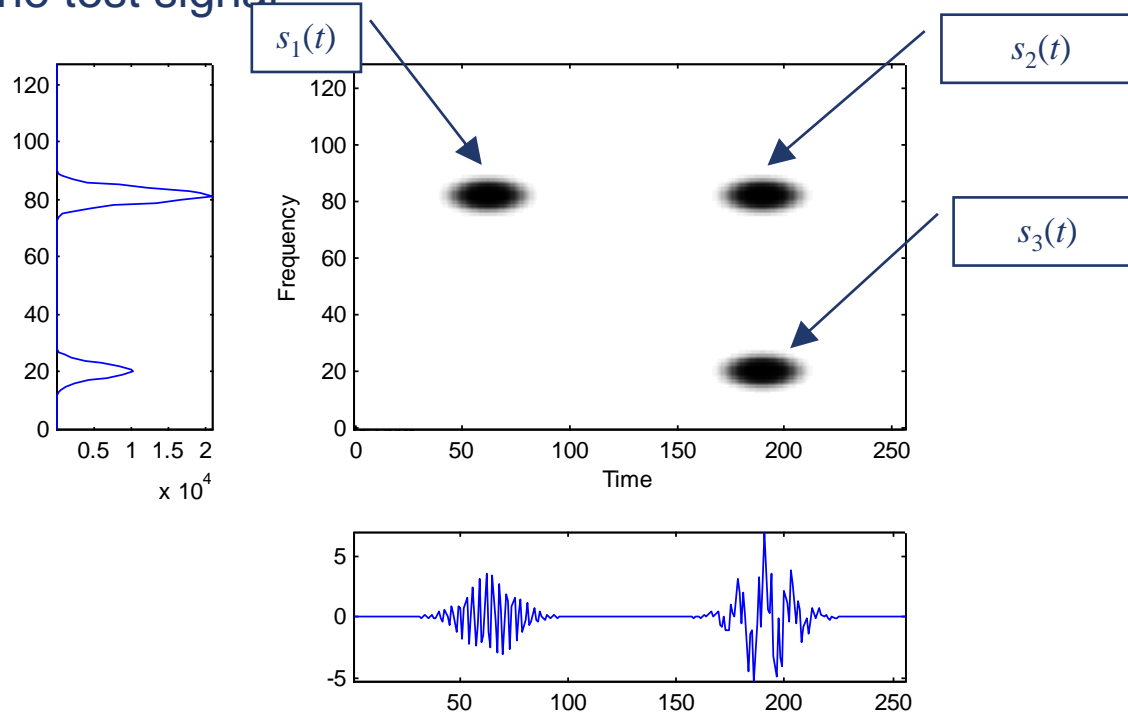
Get a large cross-term $4\pi\delta(\omega - \omega_\mu) \cos(\omega_d t)$ which varies as $\cos((\omega_1 - \omega_2)t)$ at ω_d , mid-way between the auto-terms. While the auto-terms are +ve, this oscillates +ve & \rightarrow -ve
Average of the cross-term is zero:

$$\int_{-\infty}^{\infty} 4\pi\delta(\omega - \omega_\mu) \cos(\omega_d t) dt = 0 \quad \omega_d \neq 0$$

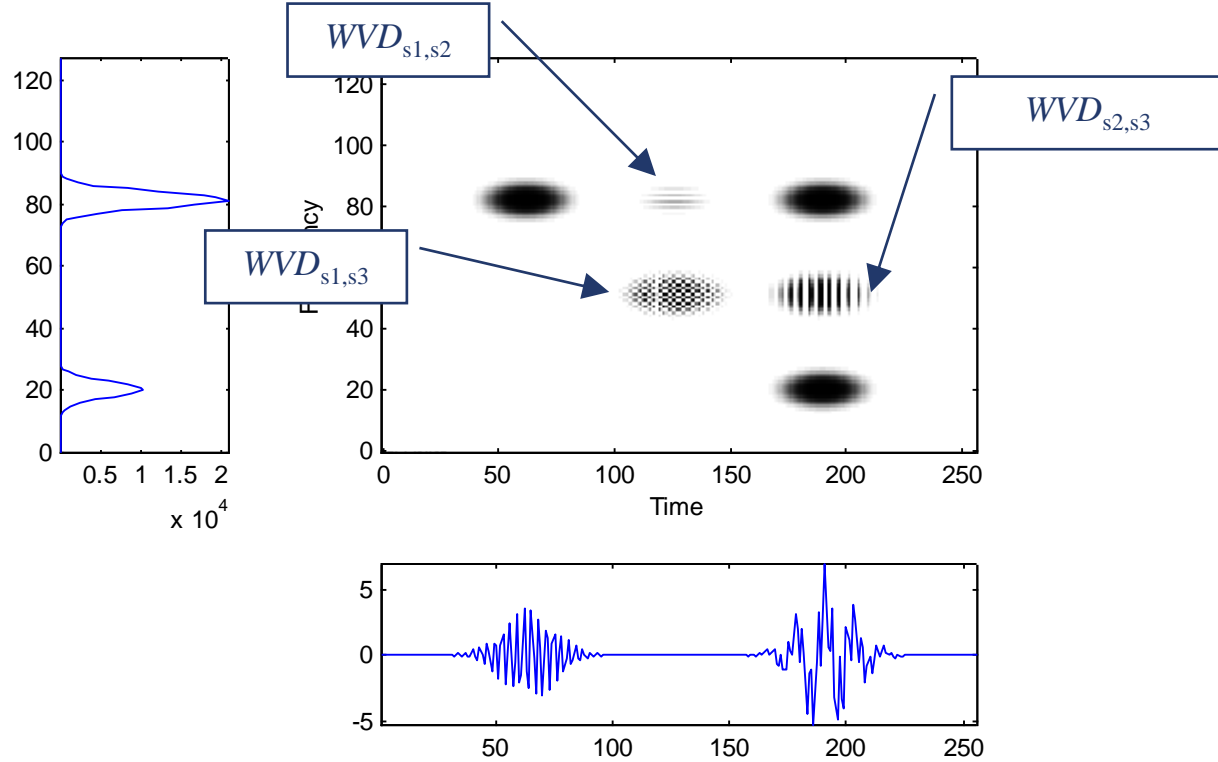
This suggests we may be able to remove these by smoothing (see later.)

Wigner-Ville Distribution

➤ Example: 3-tone test signal

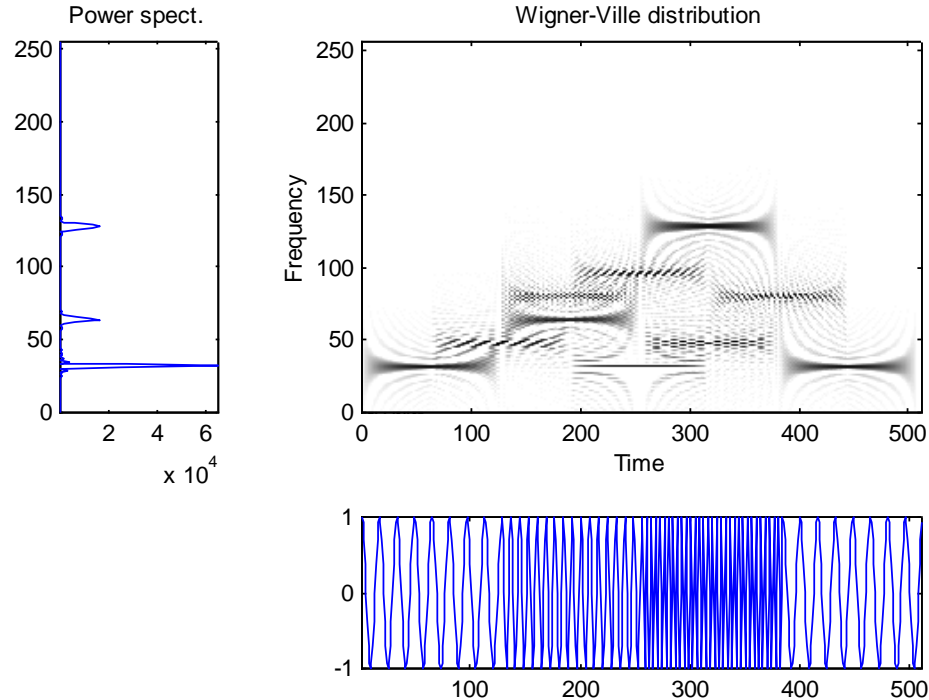


Wigner-Ville Distribution



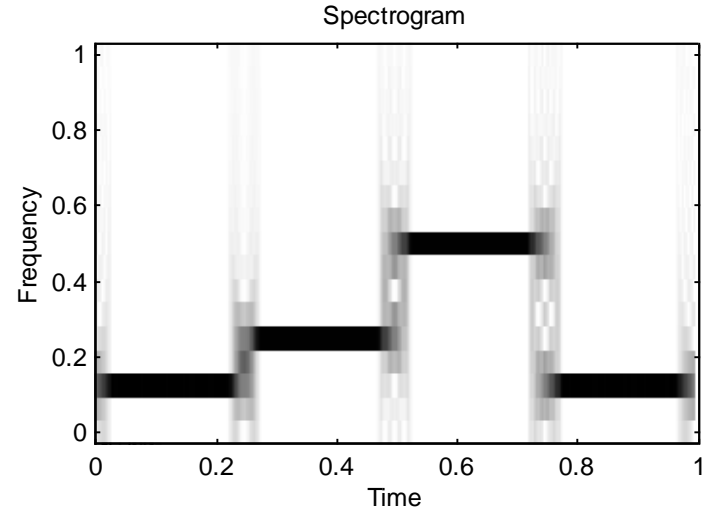
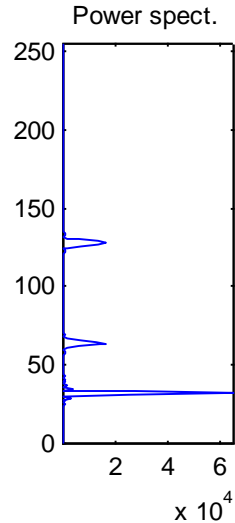
Wigner-Ville Distribution

- Another example: Frequency pulses

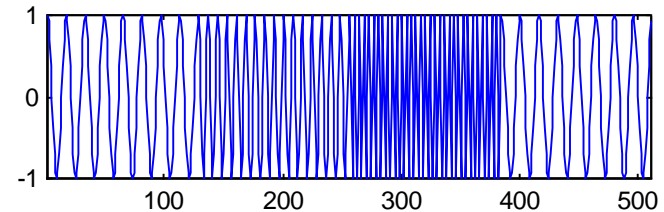


Wigner-Ville Distribution

- STFT spectrogram



Less resolution,
but no cross-terms



Wigner-Ville Distribution - Smoothing

- So, there is a **trade-off** between the beneficial properties of the WVD and the interference caused by the cross terms.
- To try to reduce the interference problem, we could use a windowed version of the WVD, called the Pseudo-WVD (PWVD):

$$PWVD_s(t, \omega) = \frac{1}{2\pi} \int s\left(t + \frac{\tau}{2}\right) s^*\left(t - \frac{\tau}{2}\right) h(\tau) e^{-j\omega\tau} d\tau$$

$h(\tau)$ is the window function.

- This is equivalent to frequency smoothing (low-pass filtering) the WVD

Wigner-Ville Distribution

- Since the cross-terms WVD are usually strongly oscillating, try removing them by using 2D low-pass filtering, to give a "smoothed Wigner-Ville distribution" (SWVD):

$$SWVD_s(t, \omega) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \varphi(x, y) WVD_s(t - x, \omega - y) dx dy$$

where $\varphi(x, y)$ is a 2D low-pass filter.

- Example: 2D Gaussian

$$\varphi(x, y) = e^{-\alpha t^2 - \beta \omega^2} \quad \alpha, \beta > 0$$

- We have a trade-off:

more smoothing \rightarrow less cross-terms, BUT

more smoothing \rightarrow reduced resolution

i.e. not as good

Wigner-Ville Distribution

- STFT Spectrogram from WVD
 - The STFT spectrogram is a smoothed WVD, with the WVD of the analysis function $\gamma(t)$ doing the smoothing:

$$|STFT_s(t, \omega)|^2 = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} WVD_{\gamma}(x, y) WVD_s(t - x, \omega - y) dx dy$$

Convolution of the WVD of $s(t)$ and
the WVD of the STFT window function

Wigner-Ville Distribution

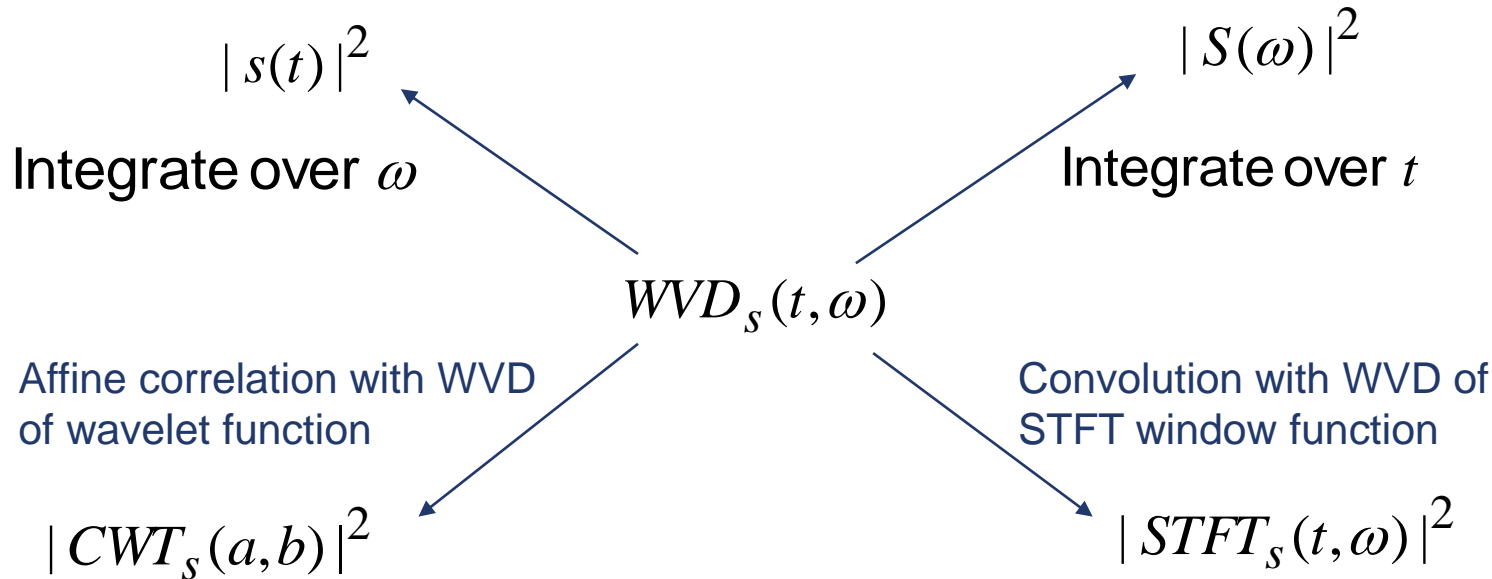
- Wavelet Scalogram from WVD
 - The scalogram (square of the wavelet transform) can be written in terms of the WVD:

$$\begin{aligned} SCAL_s(a, b) &= |CWT_s(a, b)|^2 \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} WVD_s(x, y) WVD_{\psi} \left(\frac{x-b}{a}, ay \right) dx dy \end{aligned}$$

where $WVD_s(x, y)$ is the WVD of the signal $s(t)$ and $WVD_{\psi} \left(\frac{x-b}{a}, ay \right)$ is the WVD of the mother wavelet $\psi(t)$.

- This operation is known as **affine correlation**.
 - Affine means “related to” and in mathematics can refer to a transformation that maps from one space to another while preserving shape.

Wigner-Ville Distribution - From WVD to ...?



Both the STFT spectrogram and the WT scalogram are smoothed versions of the WVD, explaining why the WVD has the best time-frequency resolution. It is still constrained by the Uncertainty Principle.

Wigner-Ville Distribution - Summary

- **Kind of Decomposition**

Time-Frequency

- **Analyzing Function**

Uses the signal itself. Motivated by time-frequency energy density (c.f. a probability density).

- **Variable**

Time and Frequency. Has high resolution in both time and frequency.

- **Suited for**

Simple (i.e. not composite) signals (non-stationary), e.g. linear chirp, gaussian pulse

- **Notes**

More “complex” (composite) signals lead to undesired “cross-terms”. Can be suppressed with smoothing, but lose high resolution in the process.



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