Advanced Transform Methods

Heisenberg Uncertainty Principle

Andy Watson

Heisenberg Uncertainty Principle

$$\Delta_t \Delta_\omega \ge 1/2$$

Time Convolution

The Fourier transform of the convolution is the product of the transforms.

$$\int_{-\infty}^{\infty} \left[s_1(t) * s_2(t) \right] e^{-j\omega t} dt = S_1(\omega) S_2(\omega)$$

So the convolution of two signals is the inverse Fourier transform of the product of the transforms.

$$\int_{-\infty}^{\infty} s_1(\tau) s_2(t-\tau) d\tau = \frac{1}{2\pi} \int_{-\infty}^{\infty} S_1(\omega) S_2(\omega) e^{j\omega t} d\omega$$

Frequency Convolution

To find the inverse Fourier transform of the convolution

$$S_1(\omega) * S_2(\omega) \equiv S(\omega)$$

$$s(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \left[S_1(\omega) * S_2(\omega) \right] e^{j\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{j\omega t} \left[\int_{-\infty}^{\infty} S_1(\alpha) S_2(\omega - \alpha) d\alpha \right] d\omega$$

$$= \frac{1}{2\pi} \int_{-\infty-\infty}^{\infty} \int_{-\infty-\infty}^{\infty} e^{j(x+\alpha)t} S_1(\alpha) S_2(x) d\alpha dx = \frac{1}{2\pi} \int_{-\infty}^{\infty} S_1(\alpha) e^{j\alpha t} d\alpha \int_{-\infty}^{\infty} e^{jxt} S_2(x) dx$$

$$= 2\pi S_1(t) S_2(t)$$

Inverse Fourier Transform of the convolution of 2 Fourier transforms is 2π times the product of the signals.



The convolution of 2 Fourier transforms is 2π times the Fourier transform of the product of the signals

$$\int_{-\infty}^{\infty} S_1(\alpha) S_2(\omega - \alpha) d\alpha = 2\pi \int_{-\infty}^{\infty} S_1(t) S_2(t) e^{-j\omega t} dt$$

Parseval's Formula

From the time convolution formula,

$$\int_{-\infty}^{\infty} s_1(\tau) s_2(t-\tau) d\tau = \frac{1}{2\pi} \int_{-\infty}^{\infty} S_1(\omega) S_2(\omega) e^{j\omega t} d\omega$$

So

$$\int_{-\infty}^{\infty} s_1(t) s_2(-t) dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} S_1(\omega) S_2(\omega) d\omega$$

If $S_2^*(t) = f(-t)$ then the Hermitian property, $S_2^*(t) \leftrightarrow S_2^*(-\omega)$ tells us $S_2^*(-\omega) = F(-\omega)$

$$S_2(\omega) = F^*(\omega)$$

So

$$\int_{-\infty}^{\infty} s_1(t) f^*(t) dt = \frac{1}{2\pi} \int_{\text{EBU718U}}^{\infty} S_1(\omega) F^*(\omega) d\omega$$

Energy

Parseval's formula states,

$$\int_{-\infty}^{\infty} s(t)f^{*}(t)dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} S(\omega)F^{*}(\omega)d\omega$$

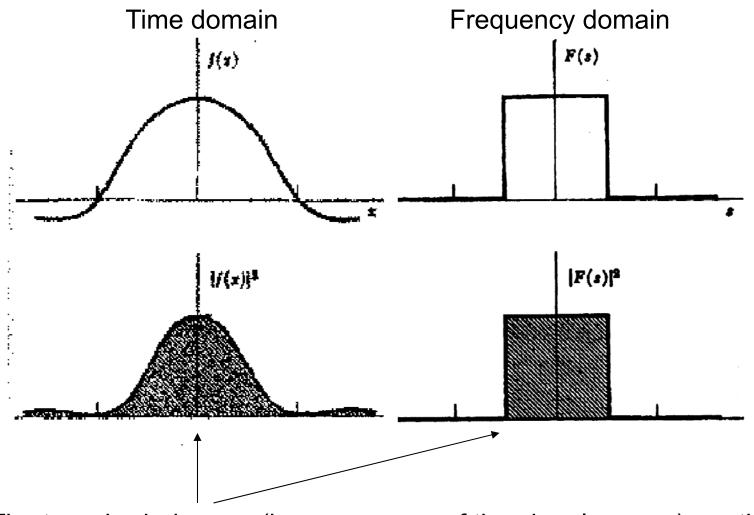
If
$$s=f$$
,
$$\int_{-\infty}^{\infty} |s(t)|^2 dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |S(\omega)|^2 d\omega$$

We can define energy in both time and frequency domains

$$E = ||s(t)||^2 = \int_{-\infty}^{\infty} |s(t)|^2 dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |S(\omega)|^2 d\omega$$

Energy is conserved

Parseval's Theorem in action



The two shaded areas (i.e., measures of the signal energy) are the same.

Mean Time and Time Duration

$$E = \int_{-\infty}^{\infty} |s(t)|^2 dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |S(\omega)|^2 d\omega$$

- The signal's gravitational centres in the time domain.

- Mean Time
$$\langle t \rangle = \frac{1}{E} \int_{-\infty}^{\infty} t |s(t)|^2 dt$$

- The signal's energy spread in the time domain
 - Time Duration Δ_t $\Delta_t^2 = \frac{1}{E} \int_{-\infty}^{\infty} (t \langle t \rangle)^2 |s(t)|^2 dt$

$$\Delta_t^2 = \frac{1}{E} \int_{-\infty}^{\infty} \left(t - \left\langle t \right\rangle \right)^2 \left| s(t) \right|^2 dt$$

Mean Frequency and Frequency Bandwidth

$$E = \int_{-\infty}^{\infty} |s(t)|^2 dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |S(\omega)|^2 d\omega$$

The signal's gravitational centre in the frequency domain.

- Mean Frequency
$$\langle \omega \rangle = \frac{1}{2\pi E} \int_{-\infty}^{\infty} \omega |S(\omega)|^2 d\omega$$

- The signal's energy spread in the frequency domain.
 - Frequency Bandwidth Δ_{ω}

$$\Delta_{\omega}^{2} = \frac{1}{2\pi E} \int_{\text{EBPJ718U}}^{\infty} \left(\omega - \left\langle \omega \right\rangle\right)^{2} \left|S(\omega)\right|^{2} d\omega$$

Duration and Bandwidth

Time Duration

$$\Delta_{t}^{2} = \frac{1}{E} \int_{-\infty}^{\infty} (t^{2} - 2t \langle t \rangle + \langle t \rangle^{2}) |s(t)|^{2} dt$$

$$= \frac{1}{E} \int_{-\infty}^{\infty} t^{2} |s(t)|^{2} dt - \frac{2 \langle t \rangle}{E} \int_{-\infty}^{\infty} |s(t)|^{2} t dt + \frac{\langle t \rangle^{2}}{E} \int_{-\infty}^{\infty} |s(t)|^{2} dt$$

$$= \frac{1}{E} \int_{-\infty}^{\infty} t^{2} |s(t)|^{2} dt - 2 \langle t \rangle \langle t \rangle + \langle t \rangle^{2} = \frac{1}{E} \int_{-\infty}^{\infty} t^{2} |s(t)|^{2} dt - \langle t \rangle^{2}$$

- Frequency Bandwidth

$$\begin{split} &\Delta_{\omega}^{\ 2} = \frac{1}{2\pi E} \int\limits_{-\infty}^{\infty} (\omega^2 - 2\omega \left\langle \omega \right\rangle + \left\langle \omega \right\rangle^2) \left| S(\omega) \right|^2 d\omega \\ &= \frac{1}{2\pi E} \int\limits_{-\infty}^{\infty} \omega^2 \left| S(\omega) \right|^2 d\omega - \frac{2\left\langle \omega \right\rangle}{2\pi E} \int\limits_{-\infty}^{\infty} \omega \left| S(\omega) \right|^2 d\omega + \frac{\left\langle \omega \right\rangle^2}{2\pi E} \int\limits_{-\infty}^{\infty} \left| S(\omega) \right|^2 d\omega \\ &= \frac{1}{2\pi E} \int\limits_{-\infty}^{\infty} \omega^2 \left| S(\omega) \right|^2 d\omega - 2\left\langle \omega \right\rangle^2 \int\limits_{-\infty}^{\infty} \left| S(\omega) \right|^2 d\omega + \frac{2}{2\pi E} \int\limits_{-\infty}^{\infty} \left| S(\omega) \right|^2 d\omega - \left\langle \omega \right\rangle^2 \right] \end{split}$$

Normalised in Time and Frequency

– Given a signal s(t), can one find a signal with the same energy, bandwidth and duration but normalised such that the mean frequency and the mean time of the signal are both set to 0?

-Yes:

$$r(t) = e^{-jt\langle\omega\rangle} s(t + \langle t\rangle)$$

 That is, a shift left in time and a shift left in frequency. The shift in frequency is obtained by modulation in time.

Uncertainty Principle

To show:

if $\sqrt{t}s(t) \to 0$ as $|t| \to \infty$ (i.e. s(t) decays fast enough) and the signal has unit energy:

$$E = \int_{-\infty}^{\infty} |s(t)|^2 dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |S(\omega)|^2 d\omega = 1$$

then

$$\Delta_t \Delta_\omega \ge \frac{1}{2}$$

where equality holds when s(t) is a Gaussian, $s(t) = Ae^{-\alpha t^2}$

We will assume that $\langle t \rangle = 0$ and $\langle \omega \rangle = 0$. (Normalised)

{ NB Will do for simple case of real s(t) = s * (t) }

Squared time width is given by: $\Delta_t^2 = \int_{-\infty}^{\infty} t^2 |s(t)|^2 dt$

and sq. frequency width by: $\Delta_{\omega}^{2} = \frac{1}{2\pi} \int_{-\infty}^{\infty} \omega^{2} |S(\omega)|^{2} d\omega$

with their product as:

$$\Delta_t^2 \Delta_\omega^2 = \int_{-\infty}^\infty t^2 |s(t)|^2 dt \frac{1}{2\pi} \int_{-\infty}^\infty \omega^2 |S(\omega)|^2 d\omega$$

From the time derivative property, we know

if
$$h(t) = \frac{d}{dt}s(t)$$
 then $H(\omega) = j\omega S(\omega)$

So using this in the frequency width we get

$$\Delta_{\omega}^{2} = \frac{1}{2\pi} \int_{-\infty}^{\infty} \omega^{2} |S(\omega)|^{2} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} |j\omega S(\omega)|^{2} d\omega$$
$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} |H(\omega)|^{2} d\omega$$

So we have

$$\Delta_{\omega}^{2} = \frac{1}{2\pi} \int_{-\infty}^{\infty} |H(\omega)|^{2} d\omega$$

But from Parseval we have the conservation of energy:

$$\int_{-\infty}^{\infty} |h(t)|^2 dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |H(\omega)|^2 d\omega$$
so $\Delta_{\omega}^2 = \int_{-\infty}^{\infty} |h(t)|^2 dt = \int_{-\infty}^{\infty} \left| \frac{d}{dt} s(t) \right|^2 dt$ since $h(t) = \frac{d}{dt} s(t)$

Therefore for the original product we get

$$\Delta_t^2 \Delta_\omega^2 = \int_{-\infty}^\infty t^2 |s(t)|^2 dt \int_{-\infty}^\infty \left| \frac{d}{dt} s(t) \right|^2 dt$$

So we have

$$\Delta_t^2 \Delta_\omega^2 = \int_{-\infty}^\infty t^2 |s(t)|^2 dt \int_{-\infty}^\infty \left| \frac{d}{dt} s(t) \right|^2 dt$$

The Schwarz Inequality states: $\|\psi_1\|^2 \|\psi_2\|^2 \ge |\langle \psi_1, \psi_2 \rangle|^2$ i.e.

$$\int_{-\infty}^{\infty} |f(t)|^2 dt \cdot \int_{-\infty}^{\infty} |g(t)|^2 dt \ge \left| \int_{-\infty}^{\infty} f(t)g(t) dt \right|^2$$

which we can use with f(t) = ts(t) and $g(t) = \frac{d}{dt}s(t)$ to get

$$\Delta_t^2 \Delta_\omega^2 = \int_{-\infty}^\infty t^2 |s(t)|^2 dt \int_{-\infty}^\infty \left| \frac{d}{dt} s(t) \right|^2 dt \qquad \ge \qquad \left| \int_{-\infty}^\infty t s(t) \frac{d}{dt} s(t) dt \right|^2$$

{Remember using real s(t)} EBU718U

So far we have $\Delta_t^2 \Delta_\omega^2 \ge \left(\int_{-\infty}^\infty t s(t) \frac{d}{dt} s(t) dt \right)^2$ {s(t) etc are real}

Now, try differentiating $s(t)^2$

$$\frac{d}{dt}s(t)^2 = 2s(t)\frac{d}{dt}s(t)$$

so inserting this we get

$$\Delta_t^2 \Delta_\omega^2 \ge \left(\int_{-\infty}^\infty t s(t) \frac{d}{dt} s(t) dt \right)^2$$
$$= \left(\frac{1}{2} \int_{-\infty}^\infty t \frac{d}{dt} s(t)^2 dt \right)^2$$

So far we have
$$\Delta_t^2 \Delta_\omega^2 \ge \left(\frac{1}{2} \int_{-\infty}^\infty t \frac{d}{dt} s(t)^2 dt\right)^2$$

Now remember integration by parts:

$$\int_{a}^{b} u dv = [uv]_{a}^{b} - \int_{a}^{b} v du$$
and use $u = t$, $v = s^{2}(t)$, $du = dt$, $dv = \frac{d}{dt}s^{2}(t)dt$
Then
$$\frac{1}{2} \int_{-\infty}^{\infty} t \frac{d}{dt} s(t)^{2} dt = \frac{1}{2} [ts(t)^{2}]_{-\infty}^{\infty} - \frac{1}{2} \int_{-\infty}^{\infty} s(t)^{2} dt$$

But we specified that
$$\sqrt{t}s(t) \to 0$$
 as $|t| \to \infty$, so $[ts(t)^2]_{-\infty}^{\infty} = \infty s(\infty)^2 - (-\infty)s(-\infty)^2 = 0 - 0 = 0$

So we now have

$$\Delta_t^2 \Delta_\omega^2 \ge \left(-\frac{1}{2} \int_{-\infty}^\infty s^2(t) dt \right)^2$$
EBU718U

So far we have

$$\Delta_t^2 \Delta_\omega^2 \ge \left(-\frac{1}{2} \int_{-\infty}^\infty s^2(t) dt \right)^2$$

But $\int_{-\infty}^{\infty} s^2(t)dt = 1$ since it is the energy of the signal,

so
$$\Delta_t^2 \Delta_\omega^2 \ge \left(-\frac{1}{2} \times 1\right)^2$$
 i.e. $\Delta_t \Delta_\omega \ge \frac{1}{2}$

Which was what we wanted.

Finally, Schwartz' inequality is exact (an equality) when the two functions ψ_1 and ψ_2 are colinear, $\psi_1 = c \psi_2$ With our functions $\psi_1 = ts(t)$ and $\psi_2 = \frac{d}{dt} s(t)$

this happens for Gaussian $s(t) = Ae^{-\alpha t^2}$,

$$\psi_2 = \frac{d}{dt}s(t) = -2\alpha t A e^{-\alpha t^2} = -2\alpha t s(t) = -2\alpha \psi_1$$

So
$$\Delta_t \Delta_\omega = \frac{1}{2}$$
 for a Gaussia θ_{U718U}

Uncertainty Principle

- We can easily find signals normalised in both time and frequency simultaneously, but it is not possible to *localise* a function in time and frequency simultaneously.
- Different definitions of the Fourier Transform yield different versions of the Uncertainty Principle.
- * The Time-Bandwidth Product $\Delta_t \Delta_\omega$ is a measure of the pulse complexity
- A signal fixed in time has infinite bandwidth
- A signal fixed in frequency has infinite duration
- Discovered by Heisenberg and Applied to quantum mechanics

Special Properties of the Gaussian

• Simple form $g(t) = e^{-at^2}$

$$g(t) = e^{-at^2}$$

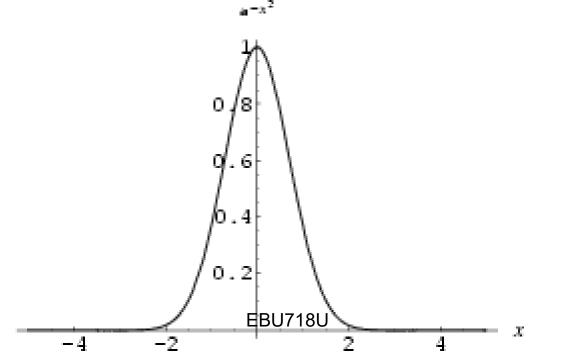
General form

$$g(t) = ae^{jbt}e^{-c(t-d)^2}$$

Normalised form

$$g(t) = ae^{jbt}e^{-c(t-d)^{2}}$$

$$g(t) = \sqrt[4]{\frac{\alpha}{\pi}}e^{-\frac{\alpha}{2}t^{2}}$$



Uncertainty Principle and Gaussian functions

$$\frac{d}{dt}e^{-at^2} = -2ate^{-at^2}$$

$$E^2 \Delta_t^2 \Delta_\omega^2 = \left(\int_{-\infty}^\infty t^2 \left| e^{-at^2} \right|^2 dt\right) \left(\int_{-\infty}^\infty \left| \frac{d}{dt} e^{-at^2} \right|^2 dt\right)$$

$$= \left(\int_{-\infty}^\infty t^2 \left(e^{-at^2} \right)^2 dt\right) \left(\int_{-\infty}^\infty \left(-2ate^{-at^2} \right)^2 dt\right)$$

$$= \left(\int_{-\infty}^\infty -2at^2 e^{-at^2} e^{-at^2} dt\right) \left(\int_{-\infty}^\infty -2at^2 e^{-at^2} e^{-at^2} dt\right)$$

$$= \left(\int_{-\infty}^\infty t e^{-at^2} \frac{d}{dt} e^{-at^2} dt\right) \left(\int_{-\infty}^\infty t e^{-at^2} \frac{d}{dt} e^{-at^2} dt\right)$$

In fact, the Gaussian is the *only* function that gives equality in the uncertainty relationship.

A Gaussian transforms to a Gaussian

$$\int_{-\infty}^{\infty} e^{-at^2} e^{-j\omega t} dt$$

$$= \int_{-\infty}^{\infty} \cos(\omega t) e^{-at^2} dt - j \int_{-\infty}^{\infty} \sin(\omega t) e^{-at^2} dt$$

$$= \int_{-\infty}^{\infty} \cos(\omega t) e^{-at^2} dt = \sqrt{\frac{\pi}{a}} e^{-\omega^2/4a}$$

The narrower a Gaussian is in one domain, the broader it is in the other domain.

Gaussian and Convolution- The Central Limit Theorem

The Central Limit Theorem says:

The convolution of the convolution of the convolution etc. of any signal approaches a Gaussian.

Mathematically,

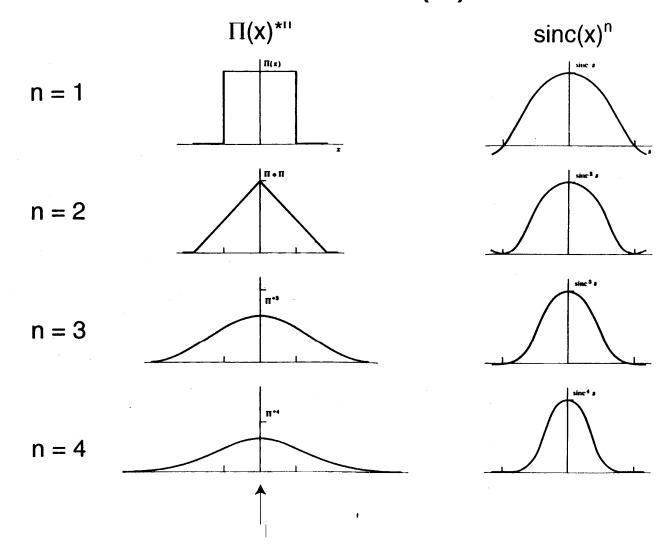
$$f(x) * f(x) * f(x) * f(x) * ... * f(x) \rightarrow e^{(-x/a)2}$$

or:

$$f(x)^{*n} \rightarrow exp[(-x/a)^2]$$

The Central Limit Theorem is why everything has a Gaussian distribution.

The Central Limit Theorem for a square function, $\Pi(x)$



Note that P(x) the ady looks like a Gaussian!