

Advanced Transform Methods

Wigner-Ville Distribution

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Wigner-Ville Distribution

STFT, Wavelet transform etc:

- Compute correlations between signal and basis fns.
- Time-freq resolution determined by basis fns.

Another approach:

- Time-frequency energy density - signal's energy density in both time and freq
- (c.f. power spectrum: energy in freq only)
- e.g. Wigner-Ville distribution

An “instantaneous power spectrum”

Recall the power spectrum :

$$P(\omega) = |S(\omega)|^2 = \int_{-\infty}^{\infty} R(\tau) e^{-j\omega\tau} d\tau$$

where $R(\tau)$ is the autocorrelation function (acf)

$$R(\tau) = \int_{-\infty}^{\infty} s(t) s^*(t - \tau) dt = \int_{-\infty}^{\infty} s(t + \tau/2) s^*(t - \tau/2) dt$$

What happens if we use *instantaneous* autocorrelation :

$$R(t, \tau) = s(t + \tau/2) s^*(t - \tau/2)$$

instead of $R(\tau) = \int_{-\infty}^{\infty} R(t, \tau) dt$? We get :

$$WVD_s(t, \omega) = \int_{-\infty}^{\infty} s(t + \tau/2) s^*(t - \tau/2) e^{-j\omega\tau} d\tau$$

which is the *Wigner - Ville Distribution* (WVD).

Cross-WVD vs auto-WVD

Since we can define a cross - correlation,
we can also define a cross - Wigner - Ville distribution :

$$WVD_{s,g}(t,\omega) = \int_{-\infty}^{\infty} s(t + \tau / 2) g^*(t - \tau / 2) e^{-j\omega\tau} d\tau$$

Taking complex conjugates we find

$$WVD_{s,g}(t,\omega) = WVD_{g,s}^*(t,\omega)$$

So for the usual WVD ("auto - WVD") we have

$$WVD_s(t,\omega) = WVD_{s,s}(t,\omega) = WVD_s^*(t,\omega)$$

so the auto - WVD is always *real*.

Example: Gaussian function

Signal $s(t) = \sqrt{\frac{\alpha}{\pi}} e^{-\alpha t^2 / 2}$ (normalized to unit energy)

For WVD, we get

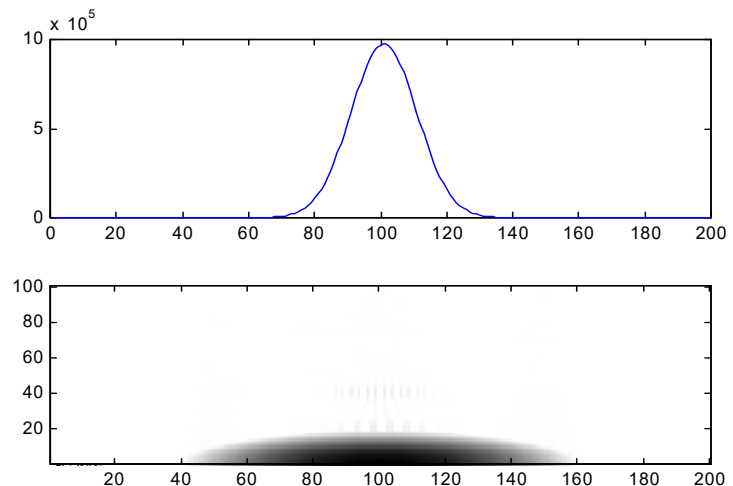
$$\begin{aligned} WVD_s(t, \omega) &= \sqrt{\frac{\alpha}{\pi}} \int_{-\infty}^{\infty} \exp \left\{ -\frac{\alpha}{2} \left[\left(t + \frac{\tau}{2} \right)^2 + \left(t - \frac{\tau}{2} \right)^2 \right] \right\} e^{-j\omega \tau} d\tau \\ &= e^{-\alpha t^2} \sqrt{\frac{\alpha}{\pi}} \int_{-\infty}^{\infty} \exp \left\{ -\frac{\alpha}{4} \tau^2 \right\} e^{-j\omega \tau} d\tau \\ &= 2 \exp \left\{ -\left[\alpha t^2 + \frac{1}{\alpha} \omega^2 \right] \right\} \end{aligned}$$

i.e. concentrated around (0,0).

α controls spread :

"time - width": $|t| < \sqrt{\frac{1}{\alpha}}$

"freq - width": $|\omega| < \sqrt{\alpha}$



Example 2: Gaussian chirplet

Signal :

$$s(t) = 4\sqrt{\frac{\alpha}{\pi}} \exp\left\{-\frac{\alpha}{2}t^2 + j\frac{\beta}{2}t^2\right\}$$

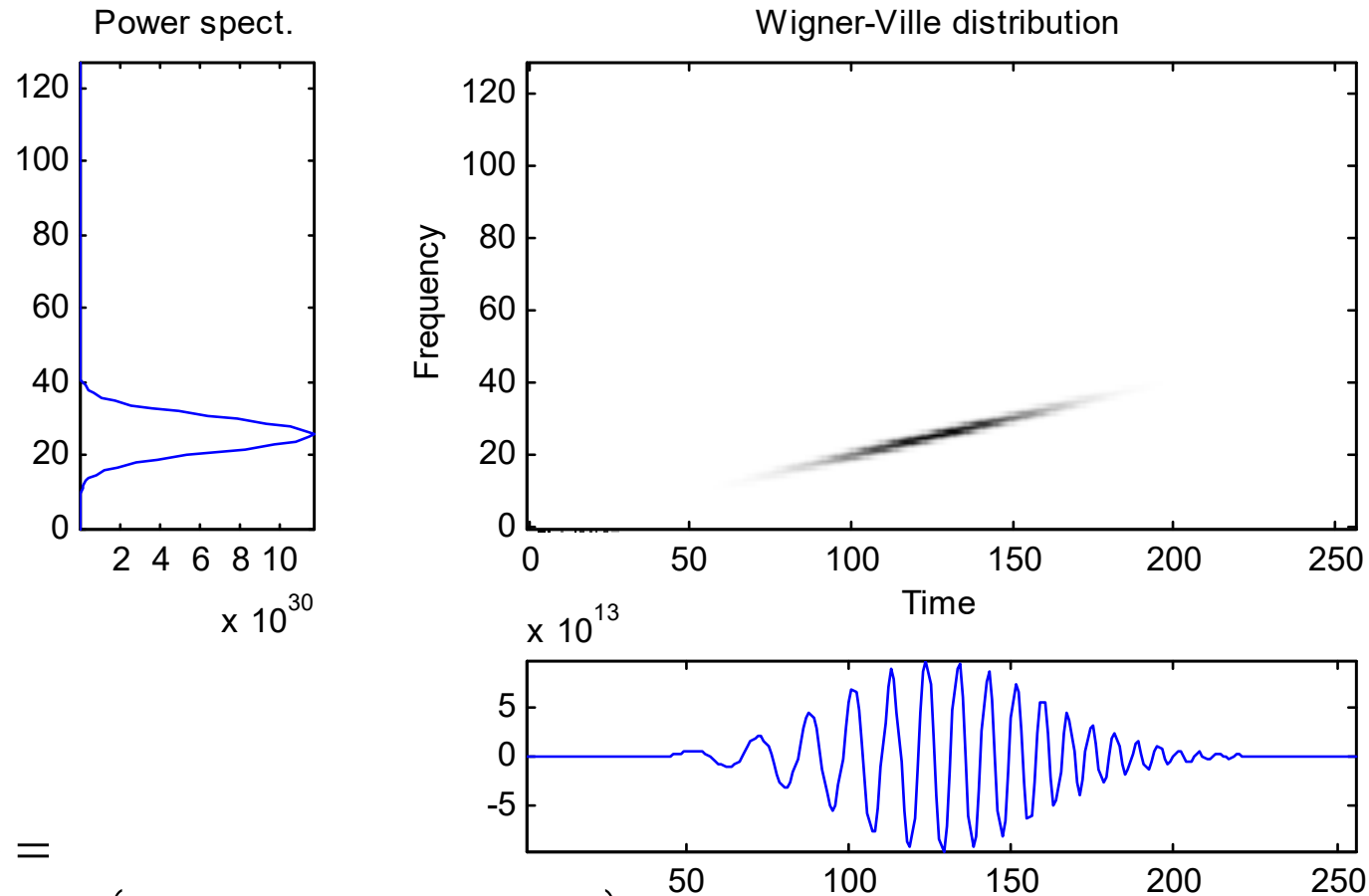
Power spectrum : $|S(\omega)|^2 = \sqrt{\frac{4\pi(\alpha^2 + \beta^2)}{\alpha}} \exp\left\{-\frac{\alpha}{\alpha^2 + \beta^2}\omega^2\right\}$

tells us *which* freqs $s(t)$ contains, not *when*. Compare :

$$\begin{aligned} WVD_s(t, \omega) &= \sqrt{\frac{\alpha}{\pi}} \int_{-\infty}^{\infty} e^{-\frac{\alpha}{2}\left[\left(t+\frac{\tau}{2}\right)^2 + \left(t-\frac{\tau}{2}\right)^2\right] + \frac{j\beta}{2}\left[\left(t+\frac{\tau}{2}\right)^2 - \left(t-\frac{\tau}{2}\right)^2\right]} e^{-j\omega\tau} d\tau \\ &= e^{-\alpha t^2} \sqrt{\frac{\alpha}{\pi}} \int_{-\infty}^{\infty} e^{-\frac{\alpha}{4}\tau^2} e^{-j(\omega - \beta t)\tau} d\tau \\ &= 2e^{-\left[\alpha t^2 + \frac{1}{\alpha}(\omega - \beta t)^2\right]} \end{aligned}$$

i.e. energy concentrated at $\omega = \beta t$, changing with time.

Illustration: chirplet



$s(t) =$

$$4\sqrt{\frac{\alpha}{\pi}} \exp\left\{-\frac{\alpha}{2}(t-t_0)^2 + j\frac{\beta}{2}t^2\right\}$$

Time-limited or band-limited signals

If $s(t)$ is time - limited, i.e. $s(t) = 0$ outside some interval $[t_0, t_1]$, then the WVD is also time - limited, i.e.

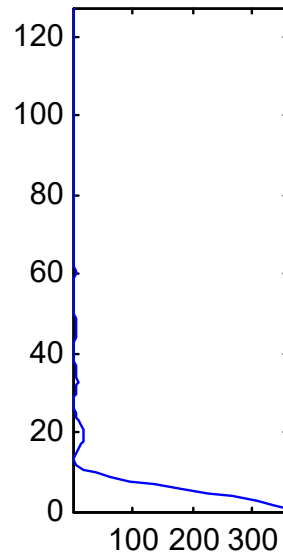
$$WVD_s(t, \omega) = 0 \text{ for } t \notin [t_0, t_1]$$

since no value for τ can make both $s(t + \tau / 2)$ and $s(t - \tau / 2)$ non - zero if t is outside this range.

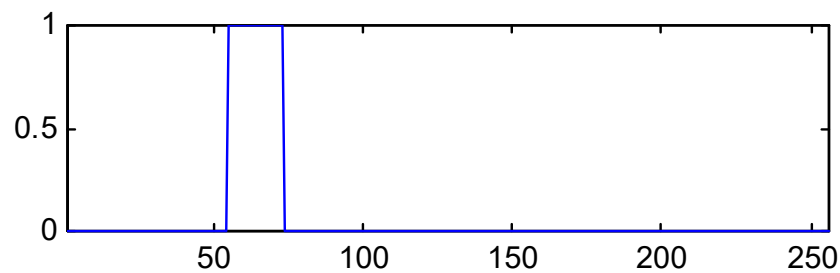
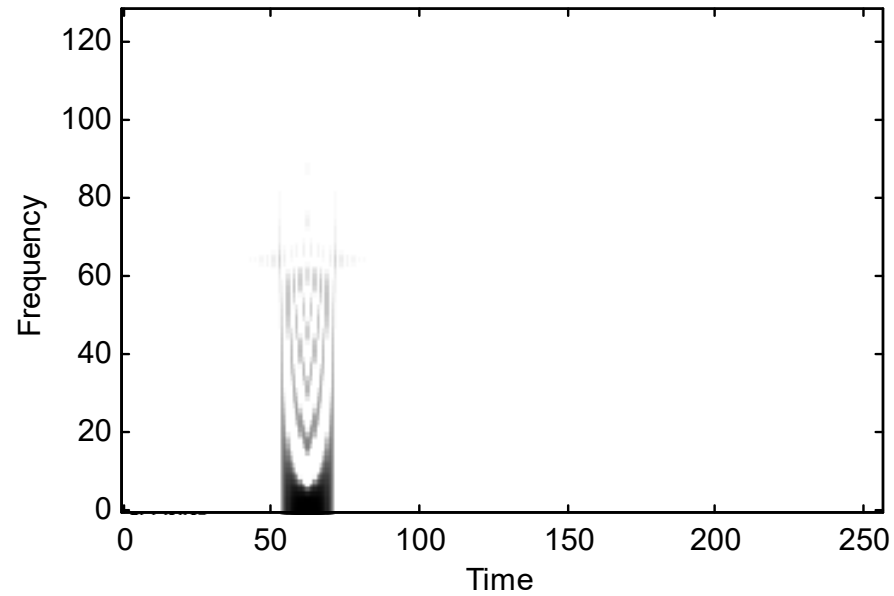
A similar result is true for freq band - limited signals.

Example: WVD of a Time-limited signal

Power spect.



Wigner-Ville distribution



WVD Properties: Time Marginal Condition

$$\begin{aligned}\frac{1}{2\pi} \int_{-\infty}^{\infty} WVD_s(t, \omega) d\omega &= \int_{-\infty}^{\infty} s\left(t + \frac{\tau}{2}\right) s^*\left(t - \frac{\tau}{2}\right) \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-j\omega\tau} d\omega d\tau \\ &= \int_{-\infty}^{\infty} s\left(t + \frac{\tau}{2}\right) s^*\left(t - \frac{\tau}{2}\right) \delta(\tau) d\tau \\ &= |s(t)|^2\end{aligned}$$

So integral over frequency of WVD is the signal power density at time t

Compare similar result for probability densities :

$$p_X(x) = \int_{-\infty}^{\infty} p_{X,Y}(x, y) dy$$

Frequency Marginal Condition

$$\begin{aligned}\int_{-\infty}^{\infty} WVD_s(t, \omega) dt &= \int_{-\infty}^{\infty} s\left(t + \frac{\tau}{2}\right) s^*\left(t - \frac{\tau}{2}\right) \int_{-\infty}^{\infty} e^{-j\omega\tau} d\tau dt \\ &= \int_{-\infty}^{\infty} e^{-j\omega\tau} \int_{-\infty}^{\infty} s(t) s^*(t - \tau) dt d\tau \\ &= \int_{-\infty}^{\infty} e^{-j\omega\tau} R(\tau) d\tau \\ &= |S(\omega)|^2\end{aligned}$$

So integral over time of WVD is the power spectral density

We also have :

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} WVD_s(t, \omega) dt d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} |S(\omega)|^2 d\omega = \int_{-\infty}^{\infty} |s(t)|^2 dt$$

i.e. the WVD is unitary : the energy in $WVD_s(t, \omega)$ is equal to energy in original signal $s(t)$.

Time-shift & Freq-moduln. invariant

If the WVD of $s(t)$ is $WVD_s(t, \omega)$,
then the WVD of time - shifted signal $s_0(t) = s(t - t_0)$
is a time - shifted WVD : $WVD_{s_0}(t, \omega) = WVD_s(t - t_0, \omega)$

Further, the WVD of frequency - modulated signal

$s_1(t) = s(t)e^{j\omega_1 t}$ is a frequency - shifted WVD :
 $WVD_{s_1}(t, \omega) = WVD_s(t, \omega - \omega_1)$

(Both follow immediately from the formulas for WVD)

WVD of multiple signals: Cross-terms

Wigner-Ville Distribution has many useful properties, and better resolution than STFT spectrogram. BUT Applications are limited due to *cross-term interference*.

Consider composite signal $s(t) = s_1(t) + s_2(t)$. Then

$$WVD_s(t, \omega) = WVD_{s_1}(t, \omega) + WVD_{s_2}(t, \omega) + 2 \operatorname{Re}\{WVD_{s_1, s_2}(t, \omega)\}$$

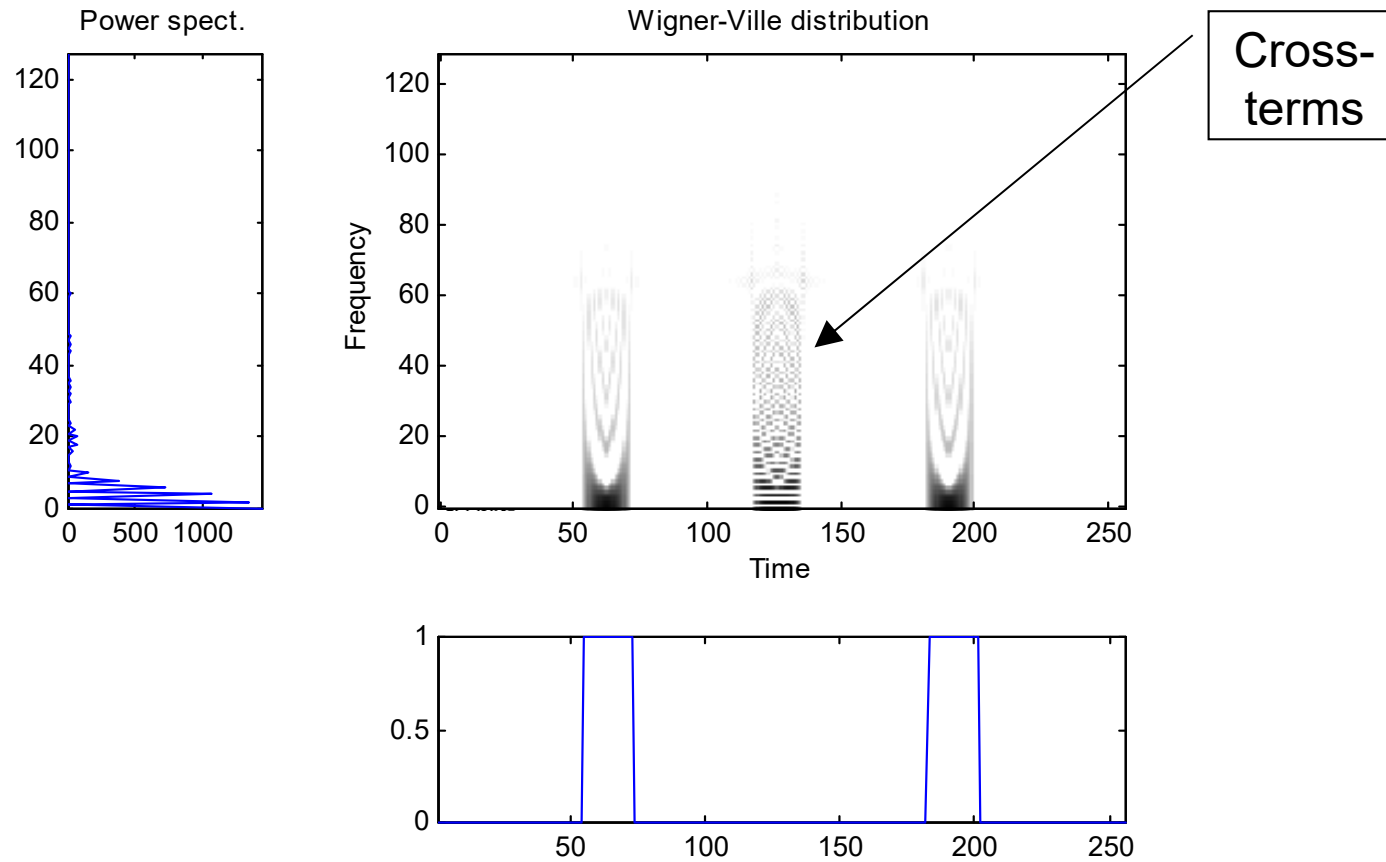
i.e. not only the sum of WVDs, but also the cross - term

$$WVD_{s_1, s_2}(t, \omega)$$

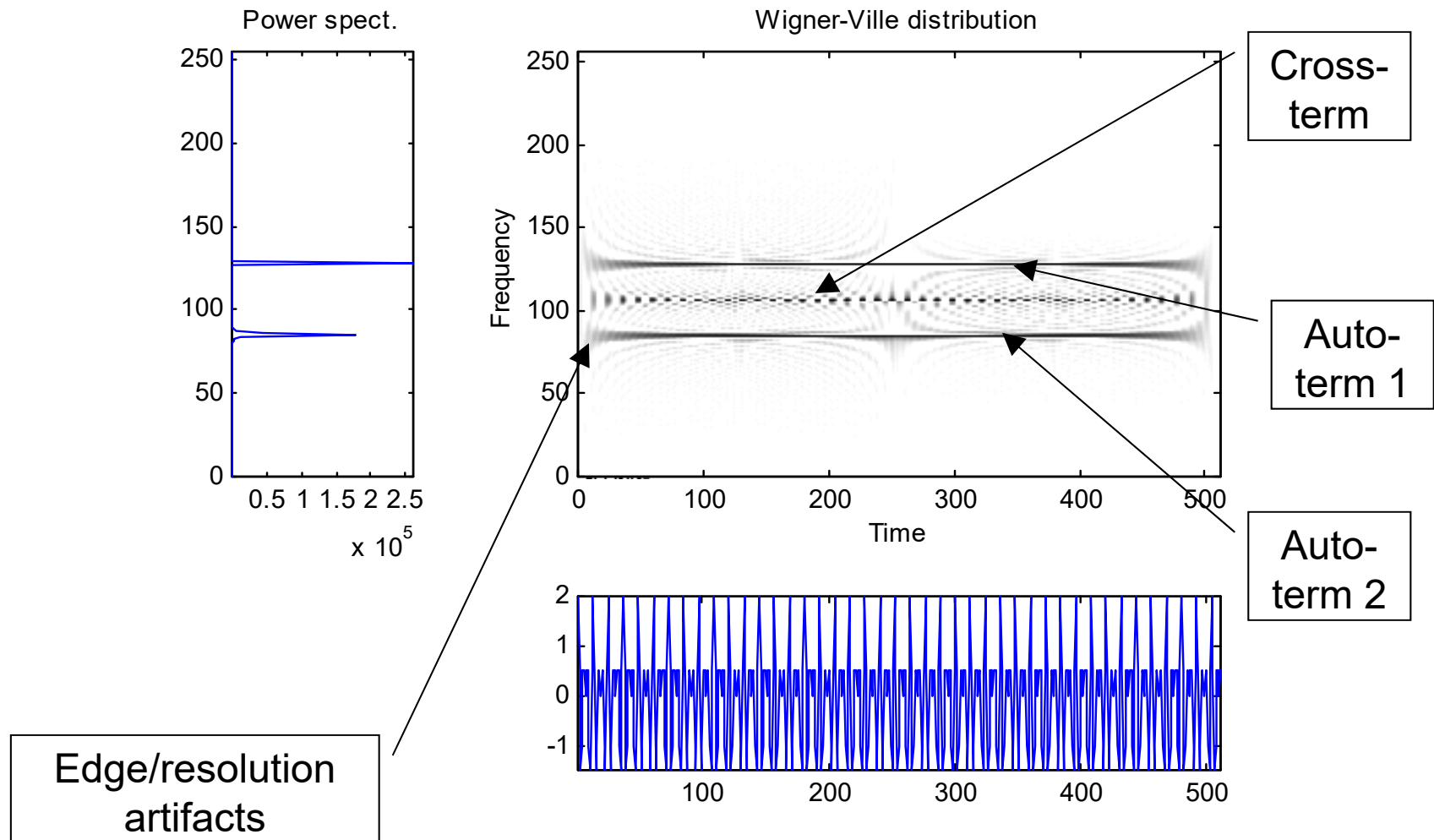
Also, cross - term is included at double the magnitude of the auto - terms, so often obscures useful patterns.

Example of cross-terms

WVD gives cross-terms for all but simplest signals, e.g.:



Example: sum of two sinusoids



Example: sum of sinusoids (cont)

If $s(t) = \exp(j\omega_0 t)$ then

$$WVD_s(t, \omega) = \int_{-\infty}^{\infty} \exp\left\{j\omega_0\left(t + \frac{\tau}{2} - t + \frac{\tau}{2}\right)\right\} e^{-j\omega\tau} d\tau = 2\pi\delta(\omega - \omega_0)$$

i.e. the WVD is a "ridge" along frequency ω_0 .

Now let $s(t) = \exp(j\omega_1 t) + \exp(j\omega_2 t)$. The power spectrum is

$$|S(\omega)|^2 = 2\pi\delta(\omega - \omega_1) + 2\pi\delta(\omega - \omega_2)$$

while the WVD is

$$WVD_s(t, \omega) = 2\pi\delta(\omega - \omega_1) + 2\pi\delta(\omega - \omega_2) + 4\pi\delta(\omega - \omega_\mu) \cos(\omega_d t)$$

where $\omega_\mu = \frac{\omega_1 + \omega_2}{2}$ and $\omega_d = \omega_1 - \omega_2$.

Example (cont): cross-term

Get a large cross - term $4\pi\delta(\omega - \omega_\mu)\cos(\omega_d t)$

which varies as $\cos((\omega_1 - \omega_2)t)$ at ω_d ,

mid - way between the auto - terms.

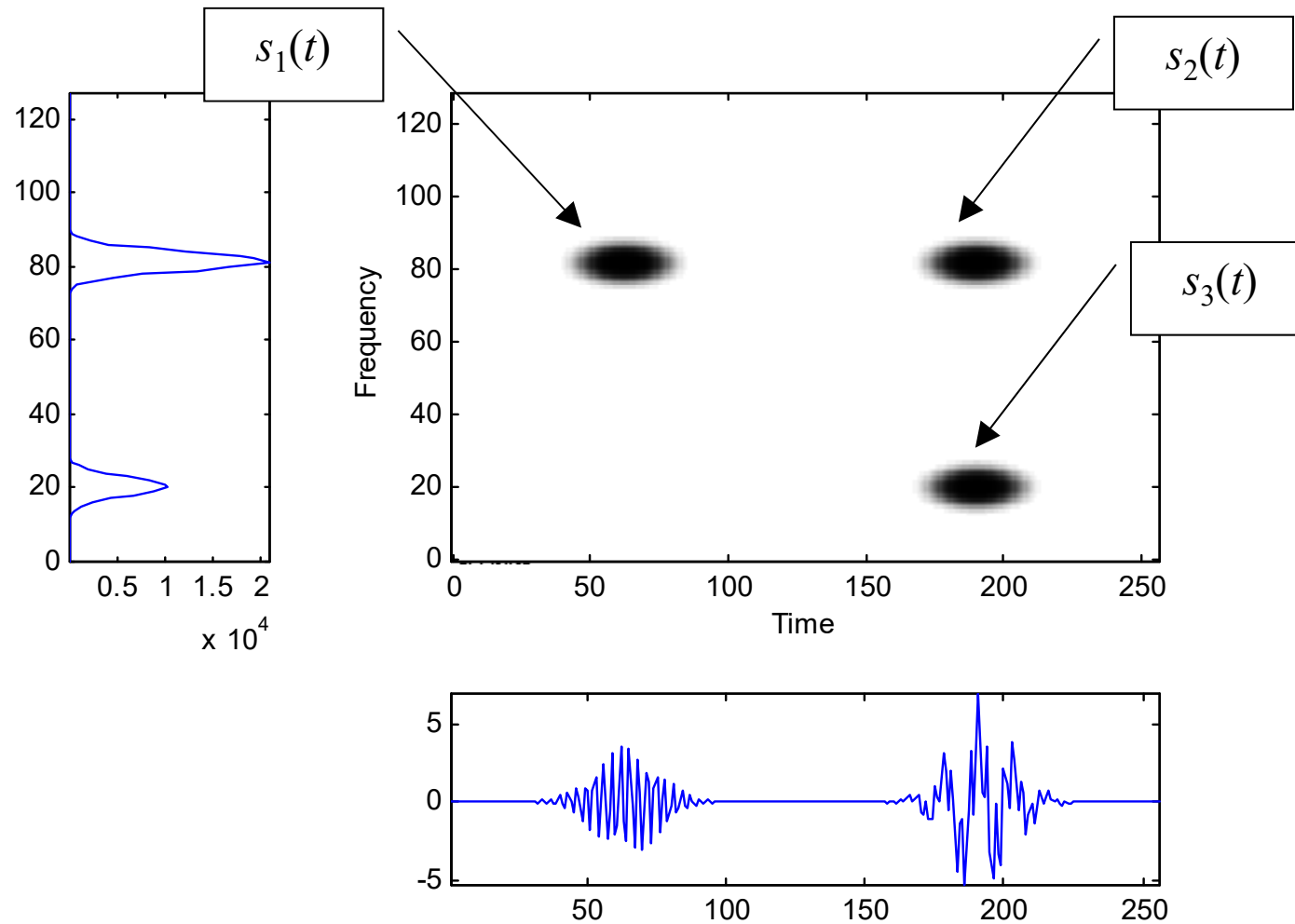
While the auto - terms are + ve, this oscillates + ve & - ve

Average of the cross - term is zero :

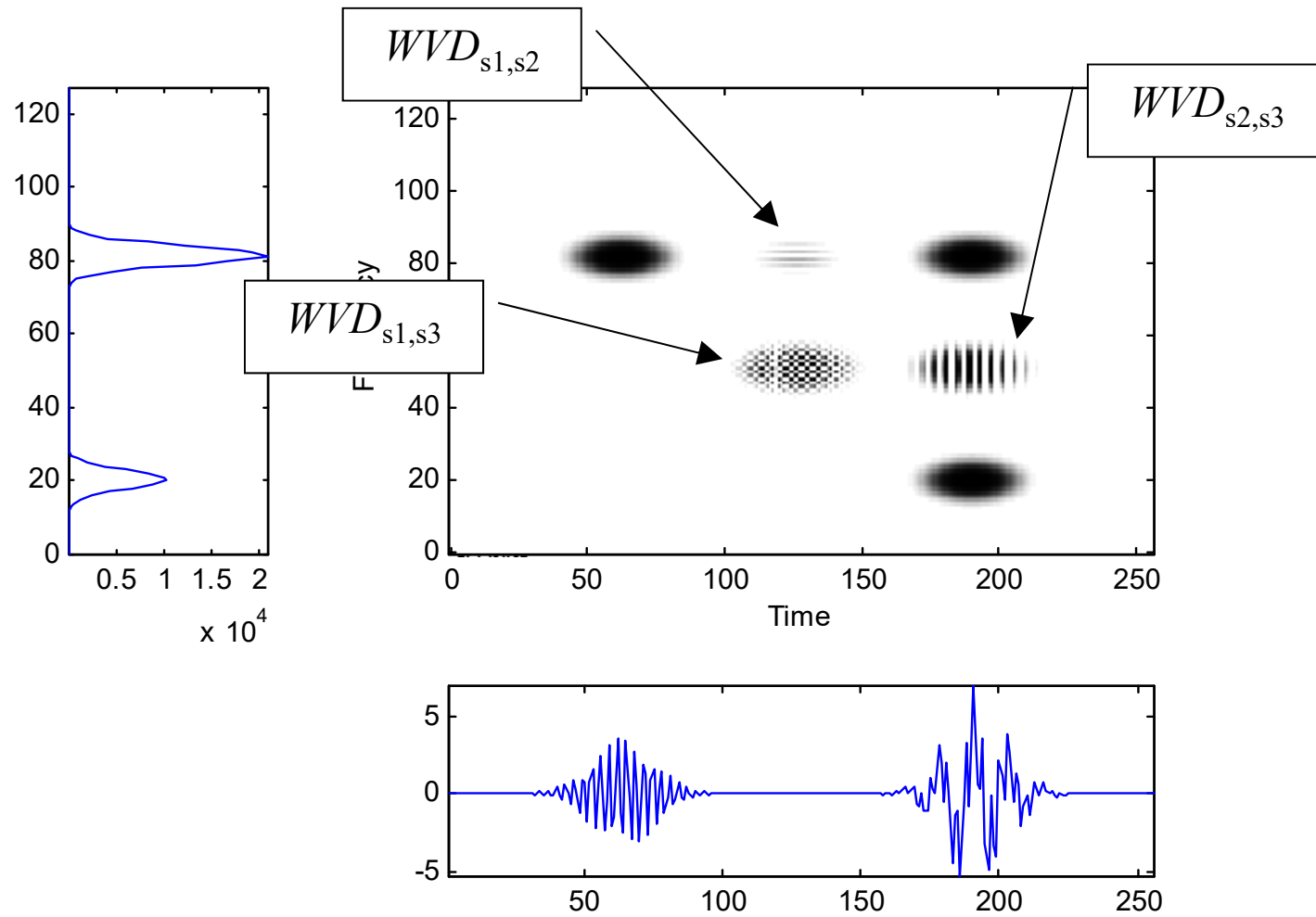
$$\int_{-\infty}^{\infty} 4\pi\delta(\omega - \omega_\mu)\cos(\omega_d t)dt = 0 \quad \omega_d \neq 0$$

This suggests we may be able to remove these
by *smoothing* (see later.)

Example: 3-tone test signal

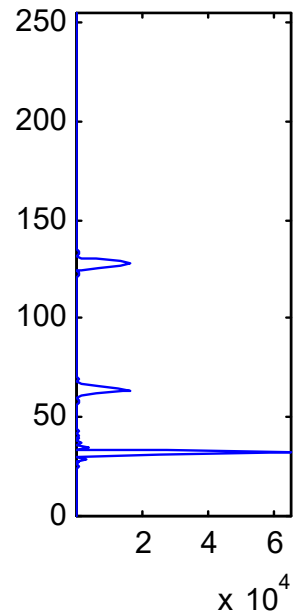


Wigner-Ville Distribution

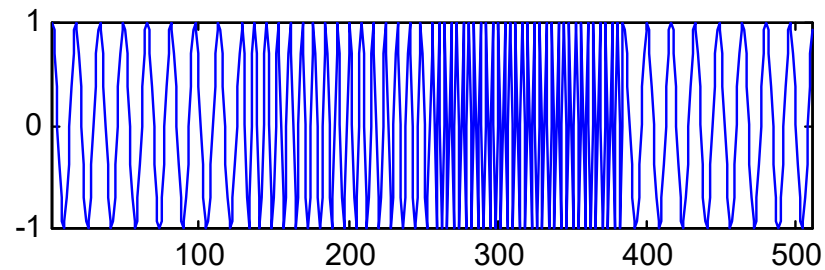
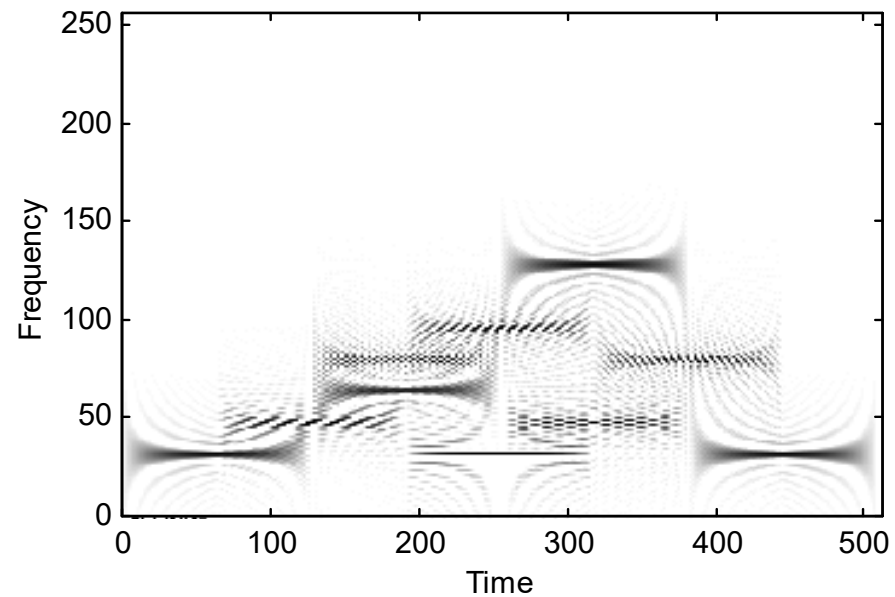


Another example: Frequency pulses

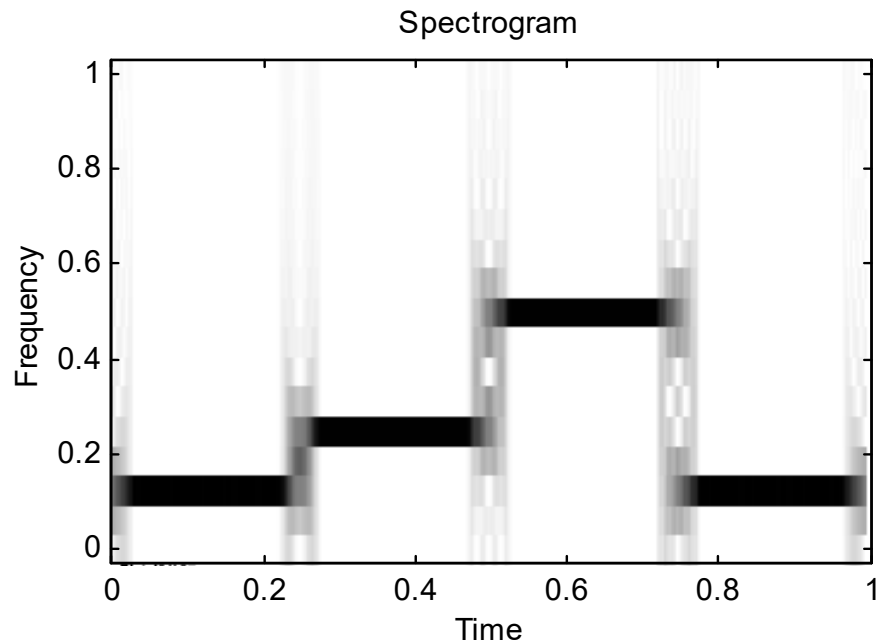
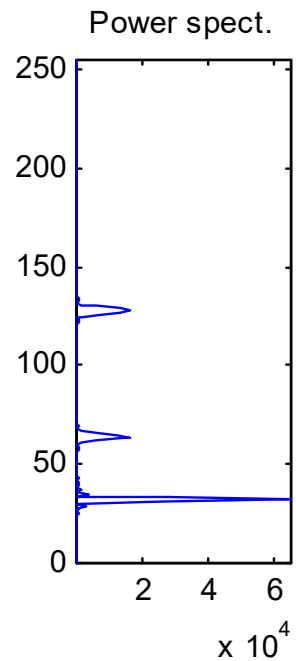
Power spect.



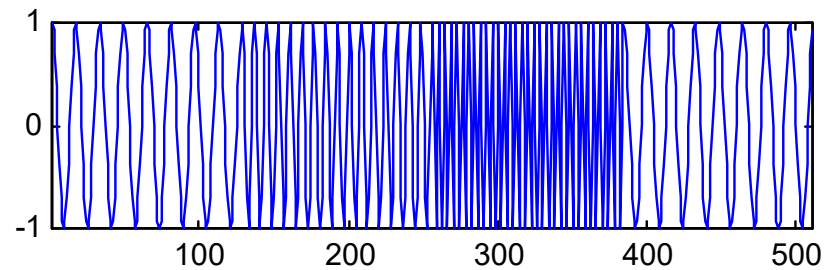
Wigner-Ville distribution



Compare spectrogram



Less resolution,
but no cross-terms



Wigner-Ville Distribution

Since the cross - terms WVD are usually strongly oscillating, try removing them by using 2D low - pass filtering, to give a "smoothed Wigner - Ville distribution" (SWVD) :

$$SWVD_s(t, \omega) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \phi(x, y) WVD_s(t - x, \omega - y) dx dy$$

where $\phi(x, y)$ is a 2D low - pass filter.

Example : 2D Gaussian

$$\phi(x, y) = e^{-\alpha t^2 - \beta \omega^2} \quad \alpha, \beta > 0$$

We have a trade - off :

more smoothing \rightarrow less cross - terms, BUT
more smoothing \rightarrow reduced resolution

STFT Spectrogram from WVD

In general, WVD can have - ve values.

But if $\alpha\beta \geq 1$ in the Gaussian 2D filter,
then the smoothed WVD will be non - negative.

Special case : when $\alpha\beta = 1$ in $\phi(x, y) = e^{-\alpha t^2 - \beta \omega^2}$
then $\phi(x, y)$ is actually a WVD of a Gaussian function.

The STFT spectrogram is a smoothed WVD, with
the WVD of the analysis function $\gamma(t)$ doing the smoothing :

$$|STFT_s(t, \omega)|^2 = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} WVD_{\gamma}(x, y) WVD_s(t - x, \omega - y) dx dy$$

(Quan p163)

Wavelet Scalogram from WVD

The scalogram (square of the wavelet transform) can be written in terms of the WVD :

$$\begin{aligned} SCAL_s(a,b) &= |CWT_s(a,b)|^2 \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} WVD_s(x,y) WVD_{\psi}\left(\frac{x-b}{a}, ay\right) dx dy \end{aligned}$$

where $WVD_s(x,y)$ is the WVD of the signal $s(t)$ and $WVD_{\psi}\left(\frac{x-b}{a}, ay\right)$ is the WVD of the mother wavelet $\psi(t)$.

This operation is known as *affine correlation*.

From WVD to ...?

