EBU6018 Fourier Series (Revision)

Andy Watson

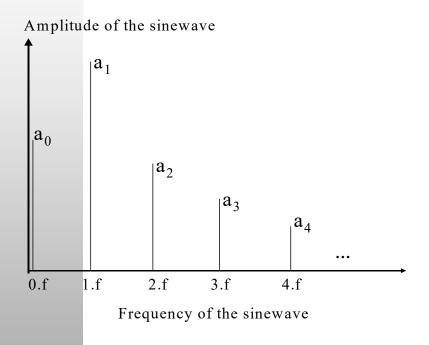
Fourier Series (FS)

- •Periodic signals can be expressed as a sum of sinusoids. The frequency spectrum can be generated by computation of the *Fourier series*.
- •The Fourier series is named after the French physicist Jean Baptiste Fourier (1768-1830), who was the first one to propose that periodic waveforms could be represented by a sum of sinusoids (or complex exponentials).

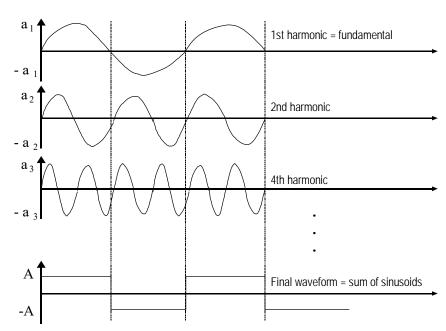
An example showing the Fourier series at work

http://www.falstad.com/fourier/

Introduction to Fourier Series

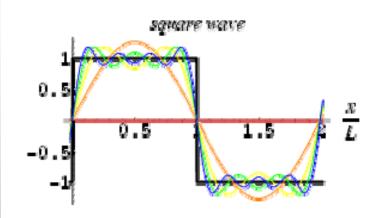


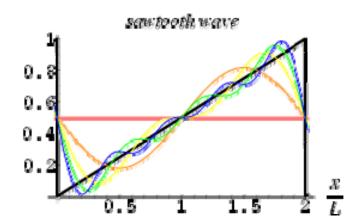
This diagram represents the frequency domain

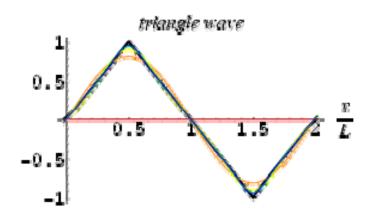


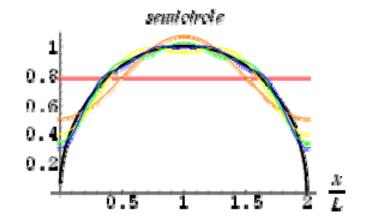
This diagram represents the time domain (but those sinewaves do not give that time signal)

Fourier Representation









A periodic signal, x(t), whose period is T, can be represented by the appropriate sum of sin and cos components:

$$x(t) = a_0 + \sum_{n=1}^{\infty} a_n \cdot \cos(n \cdot \omega \cdot t) + \sum_{n=1}^{\infty} b_n \cdot \sin(n \cdot \omega \cdot t)$$
 (1)

 a_0 is the **mean value**, or **zero frequency** term.

Integrating both sides of eqn (1), between = -T/2 and T/2:

$$\int\limits_{-T/2}^{T/2} x(t) \ dt = \int\limits_{-T/2}^{T/2} a_0 + \int\limits_{-T/2}^{T/2} \left[\begin{array}{c} \sum\limits_{n=1}^{\infty} a_n.cos(n.\omega.t) + \sum\limits_{n=1}^{\infty} b_n.sin(n.\omega.t) \end{array} \right] \ dt$$

$$\int\limits_{-T/2}^{T/2} x(t) \ dt \ = \int\limits_{-T/2}^{T/2} a_0 + \int\limits_{-T/2}^{T/2} \frac{\sum\limits_{\mathbf{n}=1}^{\infty} a_{\mathbf{n}}.cos(\mathbf{n}.\omega.t) + \sum\limits_{\mathbf{n}=1}^{\infty} b_{\mathbf{n}}.sin(\mathbf{n}.\omega.t) \] \ dt$$

$$\int\limits_{-T/2}^{T/2} \ x(t) \ dt \ = \int\limits_{-T/2}^{T/2} a_0 \ dt \ = a_0.T$$

$$a_0 = 1/T \int_{-T/2}^{T/2} x(t) dt$$

To find a formula for an it is necessary to multiply both sides of eqn(1) by $cos(m.\omega.t)$ and then integrate over the same limits:

$$\int_{-T/2}^{T/2} x(t) \cos(m.\omega.t) dt = \int_{-T/2}^{T/2} a_{0.} \cos(m.\omega.t) + \int_{-T/2}^{T/2} \sum_{n=1}^{\infty} \cos(m.\omega.t) a_{n.} \cos(n.\omega.t) + \sum_{n=1}^{\infty} \cos(m.\omega.t) b_{n.} \sin(n.\omega.t) dt$$

$$= \int_{-T/2}^{T/2} \sum_{n=1}^{\infty} \cos(m.\omega.t) a_{n.} \cos(n.\omega.t) + \sum_{n=1}^{\infty} \cos(m.\omega.t) b_{n.} \sin(n.\omega.t) dt$$
the "cos.cos" terms the "cos.sin" terms

- •Using the appropriate trig identities we can see that the cos.sin terms all produce $cos(A).sin(B) = \frac{1}{2} (sin(A+B) + sin(A-B))$ odd waveforms which all disappear under integration.
- •The cos.cos terms produce:

$$\cos(A).\cos(B) = \frac{1}{2} (\cos(A+B) + -\cos(A-B))$$

which will not necessarily disappear under integration:

HOWEVER, we are integrating over $-T/2 \rightarrow +T/2$ and this represents an integer number of cycles of the sinusoid, whatever the value of 'm' and 'n'. BUT when m=n, we have a non-zero term after integration:

$$\int_{-T/2}^{T/2} x(t) . cos(m.\omega.t) \ dt = \int_{-T/2}^{T/2} a_{0.} - cos(m.\omega.t) + \int_{-T/2}^{T/2} a_{n}. \ ^{1}/_{2} \ cos((0).\omega.t))$$

$$+ \int_{-T/2}^{T/2} \sum_{n=1}^{\infty} \frac{cos(m.\omega.t).a_{n}.cos(n.\omega.t) + \sum_{n=1}^{\infty} cos(m.\omega.t) b_{n}.sin(n.\omega.t)]dt$$

$$\int_{-T/2}^{T/2} x(t) \ cos(m.\omega.t) \ dt = (a_{n}./2) |t|_{-T/2}^{T/2} = a_{n}. \ T/2$$

BUT m=n, so:

$$\int_{-T/2}^{T/2} x(t) \cos(n.\omega.t) dt = a_n./2 |t|_{-T/2}^{T/2} = a_n \cdot T/2$$

$$a_{n} = 2/T \int_{-T/2}^{T/2} x(t).\cos(n.\omega.t) dt$$

And by similar reasoning:

$$b_n = 2/T \int_{-T/2}^{T/2} x(t).\sin(n.\omega.t) dt$$

Trigonometric Fourier Series – Cosine-with-phase form

The trigonometric Fourier series given by equ (1) can also be written in the cosine-with-phase form

$$x(t) = a_0 + \sum_{n=1}^{\infty} A_n \cos(n\omega_0 t + \theta_k) \qquad -\infty < t < \infty$$

$$A_n = \sqrt{a_n^2 + b_n^2}$$
 , $n = 1, 2,$

$$\theta_n = \begin{cases} \tan^{-1}(-\frac{b_n}{a_n}), & n = 1, 2, ..., when \ a_n \ge 0 \\ \pi + \tan^{-1}(-\frac{b_n}{a_n}), & n = 1, 2, ..., when \ a_n < 0 \end{cases}$$

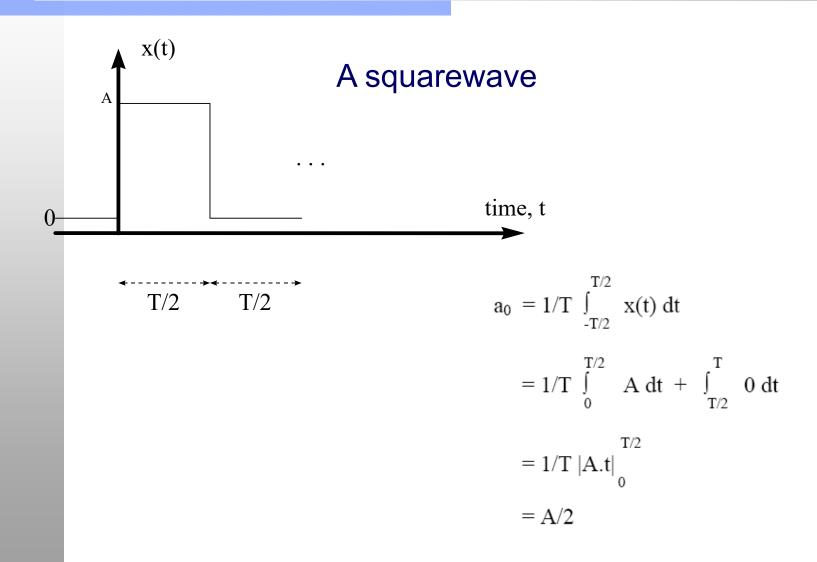
Trigonometric Fourier Series – Dirichlet conditions

Fourier believed that any periodic signal could be expressed as a sum of sinusoids. However, this turned out not to be the case, although virtually all periodic signals arising in engineering do have a Fourier series representation. In particular, a periodic signal x(t) has a Fourier series if it satisfies the following *Dirichlet conditions*:

1. x(t) is absolutely integrable over any period; that is

$$\int_{0}^{a+T} |x(t)| dt < \infty \quad \text{for any } a$$

- 2. x(t) has only a finite number of maxima and minima over any period.
- 3. x(t) has only a finite number of discontinuities over any period.



$$a_n = 2/T \int\limits_{-T/2}^{T/2} x(t).cos(n.\omega.t) \ dt = 2/T \int\limits_{0}^{T/2} \ A. \ cos(n.\omega.t) \ dt \ + \int\limits_{T/2}^{T} 0 \ dt$$

=
$$2A/T \mid \sin(n.\omega.t) / (n.\omega) \mid_{0}^{T/2}$$

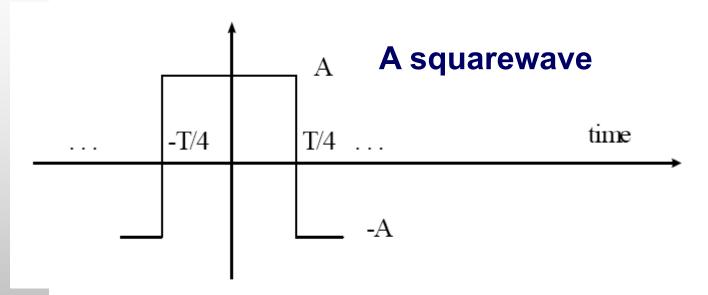
$$= A/n\pi \left[\sin(n.\pi) \right] = 0$$

$$b_n = 2/T \int_{-T/2}^{T/2} x(t).\sin(n.\omega.t) dt$$

=
$$2/T \int_{0}^{T/2} A. \sin(n.\omega.t) dt + \int_{T/2}^{T} 0 dt$$

=
$$2A/T \mid -\cos(n.\omega.t) / (n.\omega) \mid_{0}^{T/2}$$

$$= A/n\pi [1 - cos(n.\pi)]$$



$$a0 = 0$$
 by inspection

$$a_n = 2/T \int\limits_{-T/4}^{3T/4} x(t).cos(n.\omega.t) \ dt$$

$$= 2/T \int_{-T/4}^{T/4} A. \cos(n.\omega.t) dt + \int_{-T/4}^{3T/4} - A. \cos(n.\omega.t) dt$$

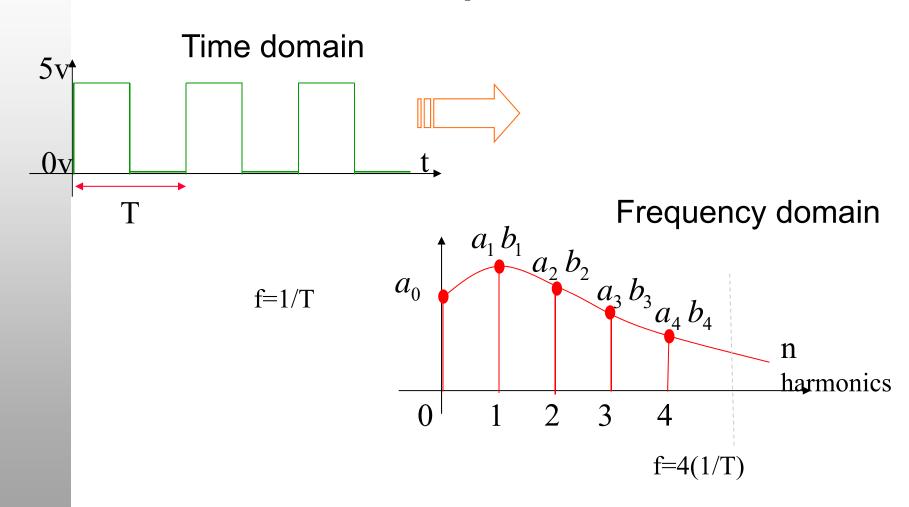
$$= 2A/T | \sin(n\omega t) / n\omega |_{-T/4}^{T/4} - 2A/T | \sin(n\omega t) / n\omega |_{T/4}^{3T/4}$$

$$= 2A/nT\omega \left[\sin(n\omega T/4) - \sin(n\omega (-T)/4) - \sin(3n\omega T/4) + \sin(n\omega T/4) \right]$$
but $\omega T = (2\pi f).(1/f) = 2\pi$, $\sin(-A) = -\sin(A)$ and $\sin(3n2\pi/4) = -\sin(n\pi/2)$ therefore:
$$= 2A/n2\pi \left[\sin(n2\pi/4) - \sin(-n2\pi/4) - \sin(3n2\pi/4) + \sin(n2\pi/4) \right]$$

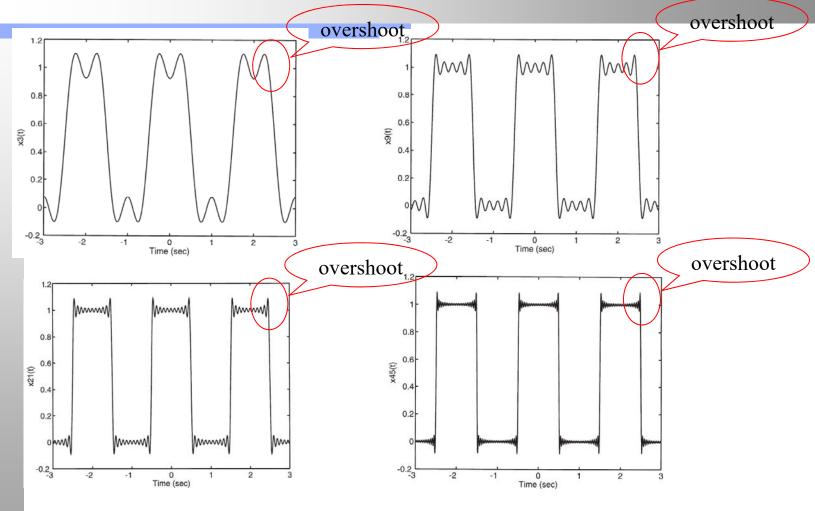
$$= 4A/n\pi \left[\sin(n\pi/2) \right]$$

$$b_n = 0 \quad \text{by inspection}$$

Example: transform the signal to line spectra



Gibbs Phenomenon



Gibbs Phenomenon

- The overshoot at the corners is still present even in the limit as N approaches to infinity. This characteristic was first discovered by Josiah Willard Gibbs (1893-1903), and this overshoot is referred to as the *Gibbs phenomenon*.
- Now let x(t) be an arbitrary periodic signal. As a consequence of the Gibbs phenomenon, the Fourier series representation of x(t) is not actually equal to the true value of x(t) at any points where x(t) is discontinuous.
- If x(t) is discontinuous at $t = t_1$, the Fourier series representation is off by approximately 9% at t_1^+ and t_1^- .

The exponential form of the Fourier Series

Let's recall the original form of Fourier series:

$$x(t) = a_0 + \sum_{n=1}^{\infty} a_n . cos(n.\omega.t) + \sum_{n=1}^{\infty} b_n . sin(n.\omega.t)$$

 In order to reduce the amount of 'writing out' the Fourier series, an exponential form can be expressed as:

$$\begin{split} a_n.cos(n.\omega.t) &= (a_n/2) \cdot \left[e^{jn\omega t} + e^{-jn\omega t}\right] \\ b_n.sin(n.\omega.t) &= (b_n/2j) \cdot \left[e^{jn\omega t} - e^{-jn\omega t}\right] \\ a_n.cos(n.\omega.t) &+ b_n.sin(n.\omega.t) = (a_n/2) \left[e^{jn\omega t} + e^{-jn\omega t}\right] + (b_n/2j) \left[e^{jn\omega t} - e^{-jn\omega t}\right] \\ &= X_n \cdot e^{jn\omega t} + X_{-n} \cdot e^{-jn\omega t} \end{split}$$
 where: $X_n = \frac{1}{2} (a_n - j b_n) \quad n \neq 0$ $X_{-n} = \frac{1}{2} (a_n + j b_n) \quad n \neq 0$

So the original Fourier Series can be written out as: $x(t) = \sum_{n=-\infty} X_n$. $e^{jn\omega t}$

Where we have defined: $X_0 = a_0$

Summary of the Fourier Series

Three forms

Original (sine and cosine components)

$$x(t) = a_0 + \sum_{n=1}^{\infty} a_n . cos(n.\omega.t) + \sum_{n=1}^{\infty} b_n . sin(n.\omega.t)$$

Cosine-with-phase form

$$x(t) = a_0 + \sum_{n=1}^{\infty} A_n \cos(n\omega_0 t + \theta_k) \qquad -\infty < t < \infty$$

Exponential form

$$x(t) = \sum_{n=-\infty}^{\infty} X_n. e^{jn\omega t}$$

- Dirichlet conditions
- Gibbs Phenomenon