

## Summary Sheet for Power Series and Fourier Transform

$$y'' + p(x)y' + qy = 0 \text{ --- (1)}$$

If  $p(x)$  and  $q(x)$  of the equation (1) are analytic at  $x_0$  (that is there exists a power series in powers of  $x - x_0$ ), then the solution to (1) has the form

$$y(x) = \sum_{n=0}^{\infty} a_n (x - x_0)^n$$

### Legendre Equation:

$$(1 - x^2)y'' - 2xy' + n(n + 1)y = 0$$

Any solution of Legendre equation is called a Legendre function

Recurrence Relation:

$$a_{s+2} = -\frac{(n-s)(n+s+1)}{(s+2)(s+1)} a_s$$
$$y_1 = 1 - \frac{n(n+1)}{2!}x^2 + \frac{(n-2)n(n+1)(n+3)}{4!}x^4 - \dots$$
$$y_2 = x - \frac{(n-1)(n+2)}{3!}x^3 + \frac{(n-3)(n-1)(n+2)(n+4)}{5!}x^5 - \dots$$

### Legendre Polynomial:

$$P_n(x) = \sum_{m=0}^M \frac{(-1)^m (2n-2m)!}{2^n m! (n-m)! (n-2m)!} x^{n-2m}, M = \frac{n}{2} \text{ or } \frac{n-2}{2}$$

### Rodrigues Formula:

$$P_n(x) = \frac{1}{2^n n!} \frac{d^n}{dx^n} (x^2 - 1)^n$$

### Properties of Legendre Polynomial:

$$P_n(-x) = (-1)^n P_n(x) \quad P_n(1) = 1 \quad P_n(-1) = (-1)^n \quad P_{2n-1}(0) = 0$$

$$P'_{2n}(0) = 0 \quad (n+1)P_{n+1} - (2n+1)xP_n + nP_{n-1} \quad P'_{n+1} - P'_{n-1} = (2n+1)P_n$$

$$\int_{-1}^1 P_n P_m dx = \begin{cases} 0 & \text{if } n \neq m \\ \frac{2}{2n+1} & \text{if } n = m \end{cases}$$

### Frobenius Method

Consider the following PDE

$$y'' + \frac{b(x)}{x} y' + \frac{c(x)}{x^2} y = 0 \text{ --- (2)}$$

If  $b(x)$  and  $c(x)$  in equation (2) are analytic at  $x = 0$ . Then Eq.(2) has at least one solution that can be represented in the form

$$y(x) = x^r \sum_{n=0}^{\infty} a_n x^n \text{ --- (3)}$$

where the exponent  $r$  may be any number (real or complex). The ODE (2) has a second solution that may be similar to (3) but with a different  $r$  and different coefficients or may contain a logarithmic term. Since  $b(x)$  and  $c(x)$  are analytic, it can also be written as

$$b(x) = \sum_{n=0}^{\infty} b_n x^n, c(x) = \sum_{n=0}^{\infty} c_n x^n$$

Then the indicial equation in  $r$  is  $r(r-1) + b_0 r + c_0 = 0$ . Let  $r_1, r_2$  be two roots of the indicial equation

**Case 1:** Distinct Roots not differing by an integer ( $r_1 \neq r_2, r_1 - r_2$  is not divisible by an integer)

$$y_1(x) = x^{r_1} \sum_{n=0}^{\infty} a_n x^n, y_2(x) = x^{r_2} \sum_{n=0}^{\infty} A_n x^n$$

**Case 2:** Distinct Roots but differing by an integer ( $r_1 \neq r_2, r_1 - r_2$  is divisible by an integer)

$$y_1(x) = x^{r_1} \sum_{n=0}^{\infty} a_n x^n, y_2(x) = k y_1(x) \ln x + x^{r_2} \sum_{n=0}^{\infty} A_n x^n$$

**Case 1:** Double Roots ( $r_1 = r_2 = r$ )

$$y_1(x) = x^r \sum_{n=0}^{\infty} a_n x^n, y_2(x) = y_1(x) \ln x + x^r \sum_{n=0}^{\infty} A_n x^n$$

**Bessel's Equation:**

$$x^2 y'' + x y' + (x^2 - \nu^2) y = 0$$

**Bessel Functions of the First Kind**

$r = \nu$ :

$$a_{2n} = -\frac{1}{2^{2n}(n+\nu)} a_{2n-2} = \frac{(-1)^n}{2^{2n+\nu} n! \Gamma(n+\nu+1)} a_0$$

$$J_\nu(x) = \sum_{n=0}^{\infty} \frac{(-1)^n}{n! \Gamma(n+\nu+1)} \left(\frac{x}{2}\right)^{2n+\nu}$$

$r = -\nu$ :

$$J_{-\nu}(x) = \sum_{n=0}^{\infty} \frac{(-1)^n}{n! \Gamma(n-\nu+1)} \left(\frac{x}{2}\right)^{2n-\nu}$$

## Bessel Functions of the Second Kind

$$Y_\nu(x) = \frac{\cos \nu \pi J_\nu(x) - J_{-\nu}(x)}{\sin \nu \pi}$$

**Properties of Bessel Functions:**  $n$  is an integer

$$[x^\nu J_\nu(x)]' = x^\nu J_{\nu-1}(x) \quad [x^{-\nu} J_\nu(x)]' = -x^\nu J_{\nu+1}(x)$$

$$J_{\nu-1}(x) + J_{\nu+1}(x) = \frac{2\nu}{x} J_\nu(x) \quad J_{\nu-1}(x) - J_{\nu+1}(x) = 2J'_\nu(x)$$

$$J_{\frac{1}{2}}(x) = \sqrt{\frac{2}{\pi x}} \sin x \quad J_{-\frac{1}{2}}(x) = \sqrt{\frac{2}{\pi x}} \cos x$$

$$J_{-n}(x) = (-1)^n J_n(x) \quad J_n(-x) = (-1)^n J_n(x) \quad \lim_{x \rightarrow 0} Y_n(x) = -\infty$$

$$J_n(0) = \begin{cases} 0 & \text{if } n > 0 \\ 1 & \text{if } n = 0 \end{cases} \quad \int_0^1 x J_n(\lambda x) J_n(\mu x) dx = \begin{cases} 0 & \text{if } \lambda \neq \mu > 0 \\ \frac{1}{2} J_{n+1}^2(\lambda) & \text{if } \lambda = \mu \end{cases}$$

$$\cos x = J_0(x) + 2 \sum_{n=1}^{\infty} (-1)^n J_{2n}(x) \quad \sin x = 2 \sum_{n=1}^{\infty} (-1)^n J_{2n+1}(x)$$

## Fourier Series

Suppose  $f(x)$  is a given function of period  $2\pi$

$$f(x) = a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx)$$

$$a_0 = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) dx, \quad a_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos nx dx, \quad b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin nx dx$$

## Representation by a Fourier Series

Let  $f(x)$  be periodic with period  $2\pi$  and piecewise continuous in the interval  $[-\pi, \pi]$ . Furthermore, let  $f(x)$  have a left-hand derivative and a righthand derivative at each point of that interval. Then the Fourier series of  $f(x)$  converges. Its sum is  $f(x)$ , except at points  $x_0$  where  $f(x)$  is discontinuous. There sum of the series is the average of the left- and right-hand limits of  $f(x)$  at  $x_0$ .

## Fourier Series (period $2L$ )

$$f(x) = a_0 + \sum_{n=1}^{\infty} \left( a_n \cos \frac{n\pi}{L} x + b_n \sin \frac{n\pi}{L} x \right)$$

$$a_0 = \frac{1}{2L} \int_{-L}^L f(x) dx, \quad a_n = \frac{1}{L} \int_{-L}^L f(x) \cos n \frac{n\pi}{L} x dx, \quad b_n = \frac{1}{L} \int_{-L}^L f(x) \sin n \frac{n\pi}{L} x dx$$

If  $f$  is even and  $L = \pi$

$$f(x) = a_0 + \sum_{n=1}^{\infty} a_n \cos nx, \quad a_0 = \frac{1}{\pi} \int_0^{\pi} f(x) dx, \quad a_n = \frac{2}{\pi} \int_0^{\pi} f(x) \cos nx dx$$

If  $f$  is odd and  $L = \pi$

$$f(x) = \sum_{n=0}^{\infty} b_n \sin nx, \quad b_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin nx \, dx$$

### **Sturm-Liouville Problems:**

$$[p(x)y']' + [q(x) + \lambda r(x)]y = 0, \quad x \in [a, b]$$

$$k_1 y(a) + k_2 y'(a) = 0, \quad l_1 y(b) + l_2 y'(b) = 0$$

Here  $\lambda$  is a parameter,  $k_i, l_i$  are real constants.

### **Orthogonality of Eigenfunction:**

Suppose that the functions  $p, q, r, p'$  in the Sturm-Liouville problem are real valued and continuous and  $r(x) > 0$  on  $[a, b]$ . Let  $y_m, y_n$  be eigenfunctions of the Sturm-Liouville problem, that correspond to different eigenvalues  $\lambda_m, \lambda_n$  respectively. Then  $y_m, y_n$  are orthogonal on that interval w.r.to the weight function  $r$ ,

$$(y_m, y_n) = \int_a^b r(x) y_m(x) y_n(x) dx = 0, \quad m \neq n$$

### **Generalized Fourier Series**

Let  $y_0, y_1, y_2, \dots$  be orthogonal w.r.to a weight function  $r(x)$  on  $[a, b]$ . Let  $f(x)$  be a function that can be represented by a convergent series

$$f(x) = \sum_{n=0}^{\infty} a_n y_n(x)$$

This is called orthogonal series, orthogonal expansion or generalized Fourier Series. If  $y_n$ 's are eigenfunctions of the Sturm-Liouville problem, then it is called as an eigenfunction expansion.

### **Fourier Integral**

$$f(x) = \int_0^{\infty} [A(w) \cos wx + B(w) \sin wx] dw$$

$$A(w) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(v) \cos wv \, dv, \quad B(w) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(v) \sin wv \, dv$$

### **Fourier Cosine Integral**

$$f(x) = \int_0^{\infty} A(w) \cos wx \, dw, \quad A(w) = \frac{2}{\pi} \int_0^{\infty} f(v) \cos wv \, dv$$

### **Fourier Sine Integral**

$$f(x) = \int_0^{\infty} A(w) \sin wx \, dw, \quad A(w) = \frac{2}{\pi} \int_0^{\infty} f(v) \sin wv \, dv$$