

ASEN 3112

Spring 2020

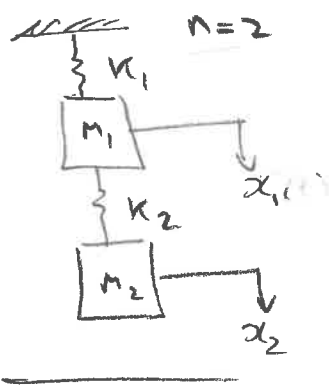
Lecture 20

Whiteboard

April 2, 2020

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MDOF Numerical Example (uncoupled)



obtain vibration characteristics
 \rightarrow natural frequencies (2 of them)
 \rightarrow mode shapes (2 of them)

$$\underline{M} = \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \text{ kg}, \quad \underline{K} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \text{ N/m}$$

(i) Form Eigenvalue problem:

$$[\underline{K} - \lambda \underline{M}] \underline{u} = \underline{0} \quad \text{or} \quad \underline{K} \underline{u} = \lambda \underline{M} \underline{u}$$

$$\begin{bmatrix} 1-2\lambda & -\lambda \\ -\lambda & 1-2\lambda \end{bmatrix} \underline{u} = \underline{0}$$

ii) Solve EVP. There are several ways to do so. If n is small, we set up the characteristic equation and use it to obtain the natural frequencies and mode shapes.

$$\begin{vmatrix} 1-2\lambda & -\lambda \\ -\lambda & 1-2\lambda \end{vmatrix} = 0$$

$$(1-2\lambda)(1-2\lambda) - (-\lambda)(-\lambda) = 0$$

$$1-2\lambda-2\lambda+4\lambda^2-\lambda^2=0$$

$$\underbrace{3\lambda^2}_{a} - \underbrace{4\lambda}_{b} + \underbrace{1}_{c} = 0$$

$$\lambda = \frac{-b \pm \sqrt{b^2 - 4ac}}{2a} = \frac{4 \pm \sqrt{(-4)(-4) + (4)(3)(1)}}{(2)(3)}$$

$$= \frac{4 \pm \sqrt{16-12}}{6} = 1, \frac{1}{3}$$

$\frac{L20}{2}$

Tradition: order in increasing size

$$\lambda_1 = \frac{1}{3} \Rightarrow \omega_{n1} = \sqrt{\frac{1}{3}} \text{ rad/s} \text{ — 1st natural frequency}$$

$$\lambda_2 = 1 \Rightarrow \omega_{n2} = \sqrt{1} \text{ rad/s} \text{ — 2nd natural frequency}$$

$$\text{Let } \underline{\tilde{u}}_i = \begin{Bmatrix} u_{i1} \\ u_{i2} \end{Bmatrix}, i=1,2$$

that is

$$\begin{array}{c} \text{1st} \\ \text{mode} \\ \text{shape} \end{array} \nearrow \underline{\tilde{u}}_1 = \begin{Bmatrix} u_{11} \\ u_{12} \end{Bmatrix}, \quad \underline{\tilde{u}}_2 = \begin{Bmatrix} u_{21} \\ u_{22} \end{Bmatrix} \nearrow \begin{array}{c} \text{2nd} \\ \text{mode} \\ \text{shape} \end{array}$$

$$\text{S EVP: } \begin{bmatrix} 1-2\lambda_i & -\lambda_i \\ -\lambda_i & 1-2\lambda_i \end{bmatrix} \begin{bmatrix} u_{i1} \\ u_{i2} \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

that is

$$\rightarrow \begin{bmatrix} 1-2\lambda_1 & -\lambda_1 \\ -\lambda_1 & 1-2\lambda_1 \end{bmatrix} \begin{bmatrix} u_{11} \\ u_{12} \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \Rightarrow \text{1st mode}$$

$$\begin{bmatrix} 1-2\lambda_2 & -\lambda_2 \\ -\lambda_2 & 1-2\lambda_2 \end{bmatrix} \begin{bmatrix} u_{21} \\ u_{22} \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \Rightarrow \text{2nd mode}$$

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Use 1st Equation:

$$(1 - 2\lambda_i) u_{i1} - \lambda_i u_{i2} = 0 \quad i=1,2$$

$$\frac{(1 - 2\lambda_i)}{\lambda_i} = \frac{u_{i2}}{u_{i1}} \quad *$$

1st mode shape:

$$\lambda_1 = \frac{1}{3} \Rightarrow \text{From } (*) \quad \frac{(1 - 2(\frac{1}{3}))}{\frac{1}{3}} = \frac{u_{12}}{u_{11}}$$

$$1 = \frac{u_{12}}{u_{11}}$$

$$\underline{\underline{\tilde{u}_1 = \begin{Bmatrix} 1 \\ 1 \end{Bmatrix}}}$$

1st
mode
shape

2nd mode shape:

$$\lambda_2 = 1 \Rightarrow \text{From } (*) \quad \frac{(1 - 2(1))}{1} = \frac{u_{22}}{u_{21}}$$
$$-1 = \frac{u_{22}}{u_{21}}$$

$$\underline{\underline{\tilde{u}_2 = \begin{Bmatrix} 1 \\ -1 \end{Bmatrix}}}$$

Eigenvalues:
(Vibration characteristics)

$$\underbrace{\left[\frac{1}{3}, \begin{Bmatrix} 1 \\ 1 \end{Bmatrix} \right]}_{\text{1st mode}}, \underbrace{\left[1, \begin{Bmatrix} 1 \\ -1 \end{Bmatrix} \right]}_{\text{2nd mode}}$$

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Solving for "response" of MDOF system

— Undamped

— Unforced (Free)

(no damping in MDOF systems in this section)

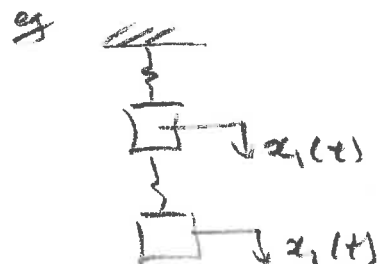
→ By Modal Analysis
(Expansion Theorem)

Recall

$$\underline{\underline{x(t)}} = \begin{Bmatrix} x_1(t) \\ x_2(t) \end{Bmatrix}$$

$$\lambda_1 \rightarrow \omega_{n1} ; \underline{\underline{U}}_1$$

$$\lambda_2 \rightarrow \omega_{n2} ; \underline{\underline{U}}_2$$



$$\underline{\underline{x(t)}} = \beta_1(t) \underline{\underline{U}}_1 + \beta_2(t) \underline{\underline{U}}_2$$

In matrix notation

expansion of
solution $\underline{\underline{x(t)}}$
in "modal space"

$$\underline{\underline{x(t)}} = \underline{\underline{U}} \underline{\underline{\beta}}$$

Expansion

$$\begin{Bmatrix} x_1(t) \\ x_2(t) \end{Bmatrix} = \begin{bmatrix} \underline{\underline{U}}_1 & \underline{\underline{U}}_2 \end{bmatrix} \begin{Bmatrix} \beta_1(t) \\ \beta_2(t) \end{Bmatrix}$$

$\underline{\underline{U}}$
Modal
matrix

vector of
generalized
coordinates
(principle coordinates)

$$\begin{Bmatrix} x_1(t) \\ x_2(t) \end{Bmatrix} = \begin{bmatrix} U_{11} & U_{21} \\ U_{12} & U_{22} \end{bmatrix} \begin{Bmatrix} \beta_1(t) \\ \beta_2(t) \end{Bmatrix}$$

S

Physical
space

$$x_1(t) = U_{11} \beta_1(t) + U_{21} \beta_2(t) \quad \text{--- 1st equation}$$

motion of
mass 1

$$x_2(t) = U_{21} \beta_1(t) + U_{22} \beta_2(t) \quad \text{--- 2nd equation}$$

motion of
mass 2

Back to EOM

$$\underline{M} \ddot{\underline{x}} + \underline{K} \underline{x} = \underline{0} \quad \text{IC}$$

$$\underline{x}_0 = \underline{x}(t=0)$$

$$\dot{\underline{x}}_0 = \underline{V}_0 = \dot{\underline{x}}(t=0)$$

Replace \underline{x} with $\underline{U} \underline{\beta}$

$$\underline{M} \underline{U} \ddot{\underline{\beta}} + \underline{K} \underline{U} \underline{\beta} = \underline{0}$$

Premultiply by \underline{U}^T

$$\underline{U}^T \underline{M} \underline{U} \ddot{\underline{\beta}} + \underline{U}^T \underline{K} \underline{U} \underline{\beta} = \underline{0}$$

We show will show later that this process "uncouples"
the two equations!

eg $n=2$

$$D_{11}^M \ddot{\beta}_1(t) + D_{11}^K \beta_1(t) = 0 \quad \text{--- Eq 1}$$

in nodal

$$D_{22}^M \ddot{\beta}_2(t) + D_{22}^K \beta_2(t) = 0 \quad \text{--- Eq 2}$$

space

Two uncoupled ODE's.

Solve as two separate 1DOF problems to obtain

Modal
space

$$\beta_1(t) = \underline{\quad}$$

$$\beta_2(t) = \underline{\quad}$$

Once you have $\beta_1(t), \beta_2(t)$, retransform back to \underline{x}

$$\underline{x} = \underline{U} \underline{\beta}$$

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Normalization

$$\text{Let } \underline{\tilde{v}}_i = \alpha_i \underline{\tilde{u}}_i \quad i=1,2$$

$$\text{Form } \underline{\tilde{v}}_i^T \underline{M} \underline{\tilde{v}}_i = \alpha_i^2 \underline{\tilde{u}}_i^T \underline{M} \underline{\tilde{u}}_i = 1 \quad (*)$$

Choose α_i such that (*) is satisfied

$$\text{Thus } \alpha_i = \frac{1}{\sqrt{\underline{\tilde{u}}_i^T \underline{M} \underline{\tilde{u}}_i}}$$

This way, we enforce $\underline{\tilde{v}}_i^T \underline{M} \underline{\tilde{v}}_i = 1$

We write normalized modal matrix

$$\underline{V} = \left[\begin{array}{c} \underline{\tilde{v}}_1 \\ \underline{\tilde{v}}_2 \end{array} \right] = \begin{bmatrix} v_{11} & v_{21} \\ v_{12} & v_{22} \end{bmatrix}$$

$$\underline{V}^T \underline{M} \underline{V} = \underline{I} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$\underline{V}^T \underline{K} \underline{V} = \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix}$$

This is an intermediate, but very useful.

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Apply Modal with normalization:

$$\underline{x} = \beta_1(t) \underline{V}_1 + \beta_2 \underline{V}_2$$

$\left. \begin{matrix} \underline{V}_1 \\ \underline{V}_2 \end{matrix} \right\}$ normalized
mode
shapes

$$\underline{x} = \underline{V} \underline{\beta}$$

Back to EOM:

$$\underline{M} \ddot{\underline{x}} + \underline{K} \underline{x} = \underline{0}$$

Replace $\underline{x} = \underline{V} \underline{\beta}$

$$\underline{M} \underline{V} \ddot{\underline{\beta}} + \underline{K} \underline{V} \underline{\beta} = \underline{0}$$

Pre-multiply by \underline{V}^T :

$$\underbrace{\underline{V}^T \underline{M} \underline{V}}_{\underline{I}} \ddot{\underline{\beta}} + \underbrace{\underline{V}^T \underline{K} \underline{V}}_{\begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix}} \underline{\beta} = \underline{0}$$

Modal EOM:

$$\ddot{\beta}_j + \lambda_j \beta_j = 0 \quad j=1,2$$

$$\ddot{\beta}_1 + \lambda_1 \beta_1 = 0$$

→ EOM #1 in modal space

$$\ddot{\beta}_2 + \lambda_2 \beta_2 = 0$$

→ EOM # 2 in modal space

Analogous
to 1DOF
 $\ddot{x} + \omega_n^2 x = 0$

Effectively we have converted our 2DOF coupled ODE's to two 1DOF uncoupled ODEs

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Initial Conditions also need to be converted to modal space:

$$\underline{\tilde{x}}_0 = \underline{V} \underline{\tilde{\beta}}_0, \quad \underline{\dot{\tilde{x}}}_0 = \underline{V} \underline{\dot{\tilde{\beta}}}_0$$

Recall $\underline{V}^T \underline{M} \underline{V} = \underline{I}$

Post multiply by \underline{V}^{-1}

$$\underline{V}^T \underline{M} \underline{V} \underline{V}^{-1} = \underline{I} \underline{V}^{-1}$$

$$\underline{V}^T \underline{M} = \underline{V}^{-1}$$

$$\underline{V}^T \underline{\tilde{x}}_0 = \underline{V}^{-1} \underline{V} \underline{\tilde{\beta}}_0$$

$$\Rightarrow \underline{\tilde{\beta}}_0 = \underline{V}^T \underline{M} \underline{\tilde{x}}_0$$

similarly

$$\Rightarrow \underline{\dot{\tilde{\beta}}}_0 = \underline{V}^T \underline{M} \underline{\dot{\tilde{x}}}_0$$

Find set in modal space

$$\ddot{\tilde{\beta}}_1 + \lambda_1 \tilde{\beta}_1 = 0$$

$$\ddot{\tilde{\beta}}_2 + \lambda_2 \tilde{\beta}_2 = 0$$

IC'

$$\tilde{\beta}_0 = \underline{V}^T \underline{M} \underline{x}_0$$

$$\dot{\tilde{\beta}}_0 = \underline{V}^T \underline{M} \underline{\dot{x}}_0$$

→ solve for $\tilde{\beta}_1(t), \tilde{\beta}_2(t)$

→ Transform back to $\underline{x} = \underline{V} \underline{\tilde{\beta}}$

DONE