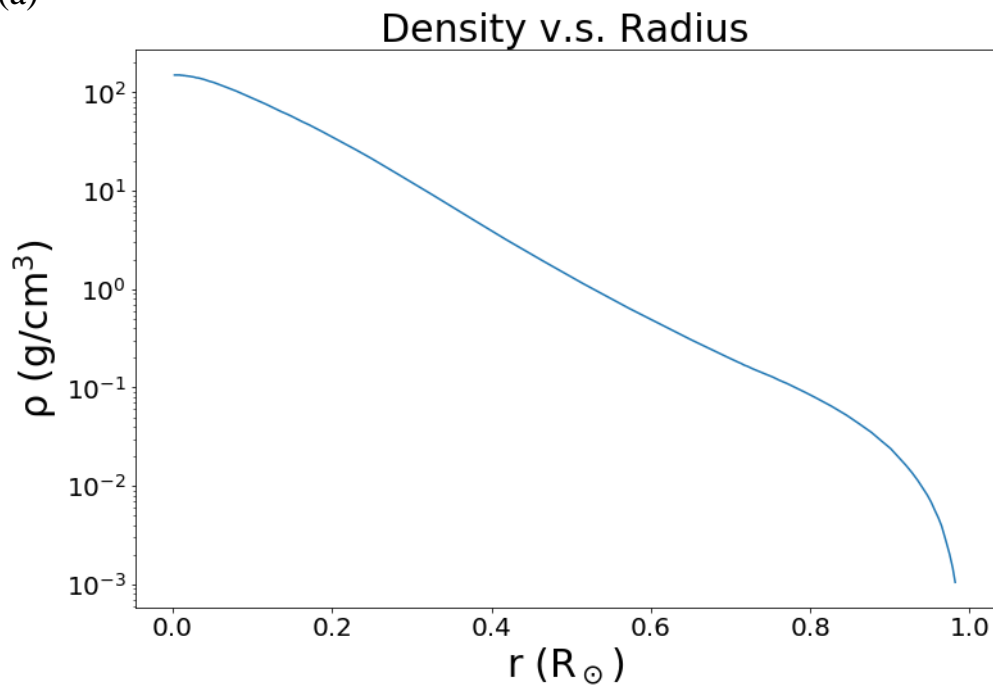
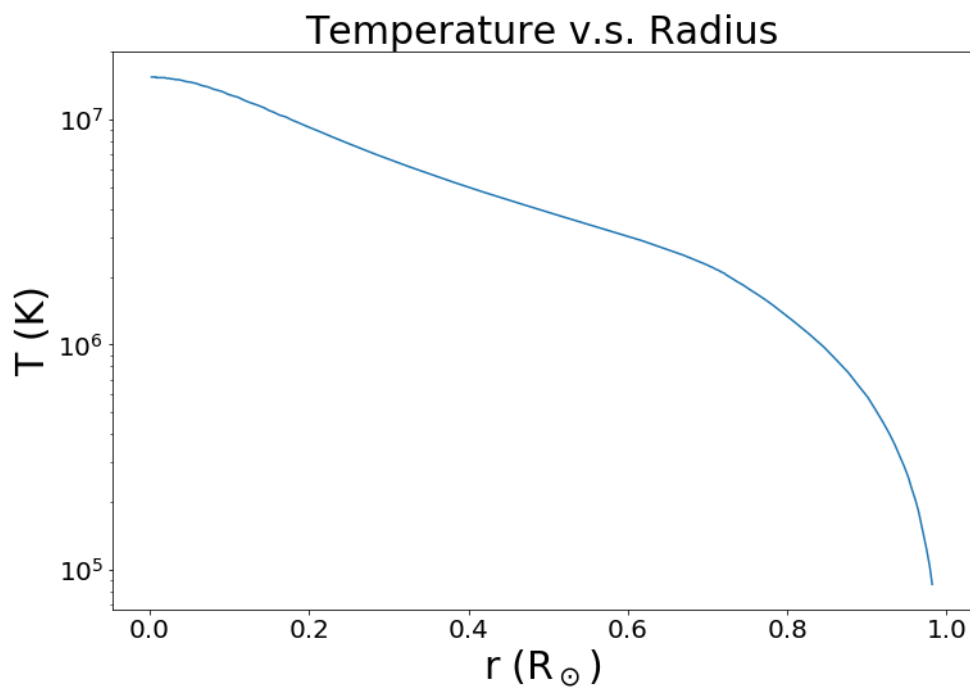


Homework 4

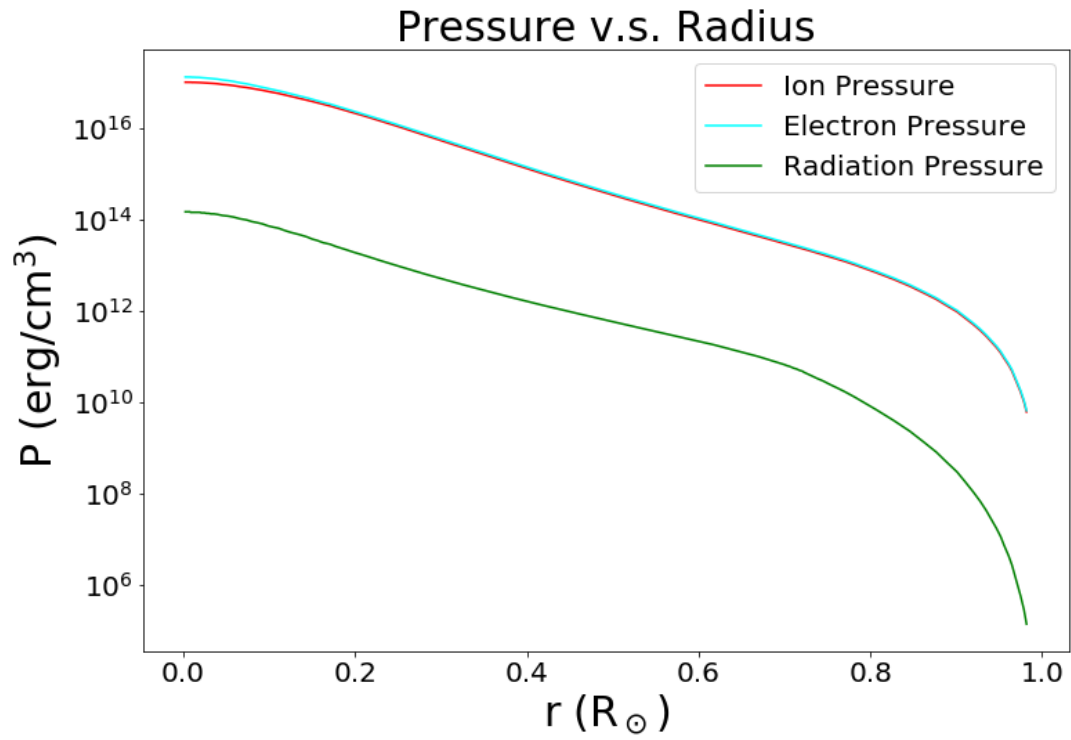
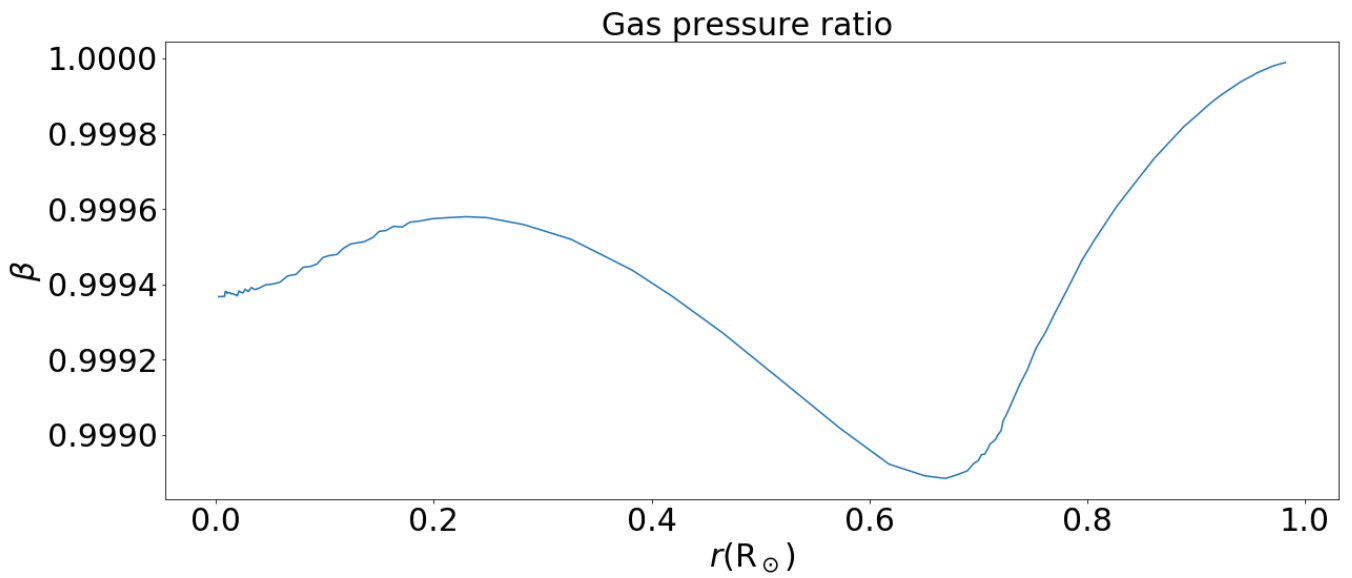
Zhiwei Pan 1901110222

Q1

(a)

**Figure 1****Figure 2**

(b)

**Figure 3****Figure 4**

Assume ideal gas then we can calculate the pressure:

$$P_{ion} = \rho kT \left(\frac{X}{m_H} + \frac{Y}{m_{He}} \right)$$

$$P_e = \rho kT \left(\frac{X}{m_H} + \frac{2Y}{m_{He}} \right)$$

$$P_{rad} = \frac{1}{3} a T^4$$

$$P_{gas} = P_{ion} + P_e$$

$$P_{total} = P_{gas} + P_{rad}$$

Then we define: $\beta = \frac{P_{gas}}{P_{total}}$

Figure 3 and 4 show the pressure and β as a function of radius.

(c)

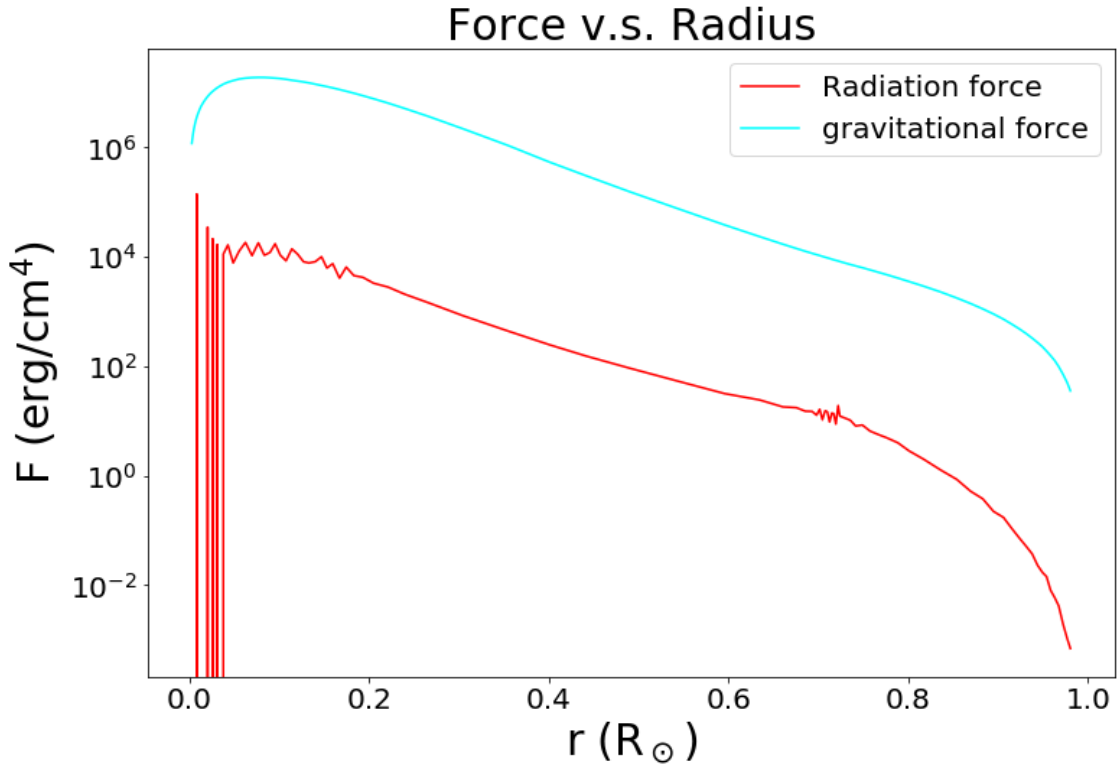


Figure 5

As figure 5 shows, radiation force is much smaller than the gravitational force so it's not important.

(d)

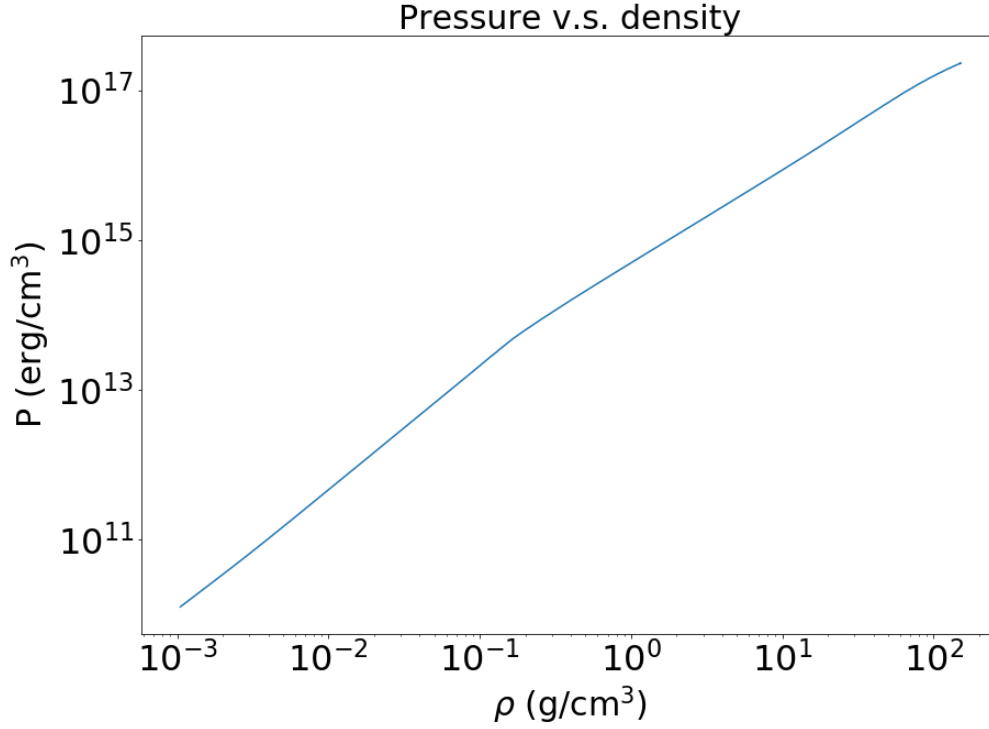


Figure 6

Figure 6 shows the relation between Pressure and density in log scale. We can find that it's very similar to linear relation, which means that: $P \sim \rho^\gamma$. Because of the similarity of such a kind of equation of state we can say that a polytrope model is very close to Standard Solar Model.

(e)(f)

For a polytrope of $n = 3, \beta = 0.9995853$:

$$r = \alpha z \quad \text{with} \quad \alpha = \frac{R_\odot}{z_3}$$

$$\rho = \rho_c w^3$$

$$\rho_c \approx 54.1825 \bar{\rho} = 54.1825 \frac{3M}{4\pi R^2}$$

$$K = N_3 G M^{\frac{2}{3}} \quad \text{with} \quad N_n = \frac{(4\pi)^{\frac{1}{3}}}{4} \Theta_n^{-\frac{2}{3}}$$

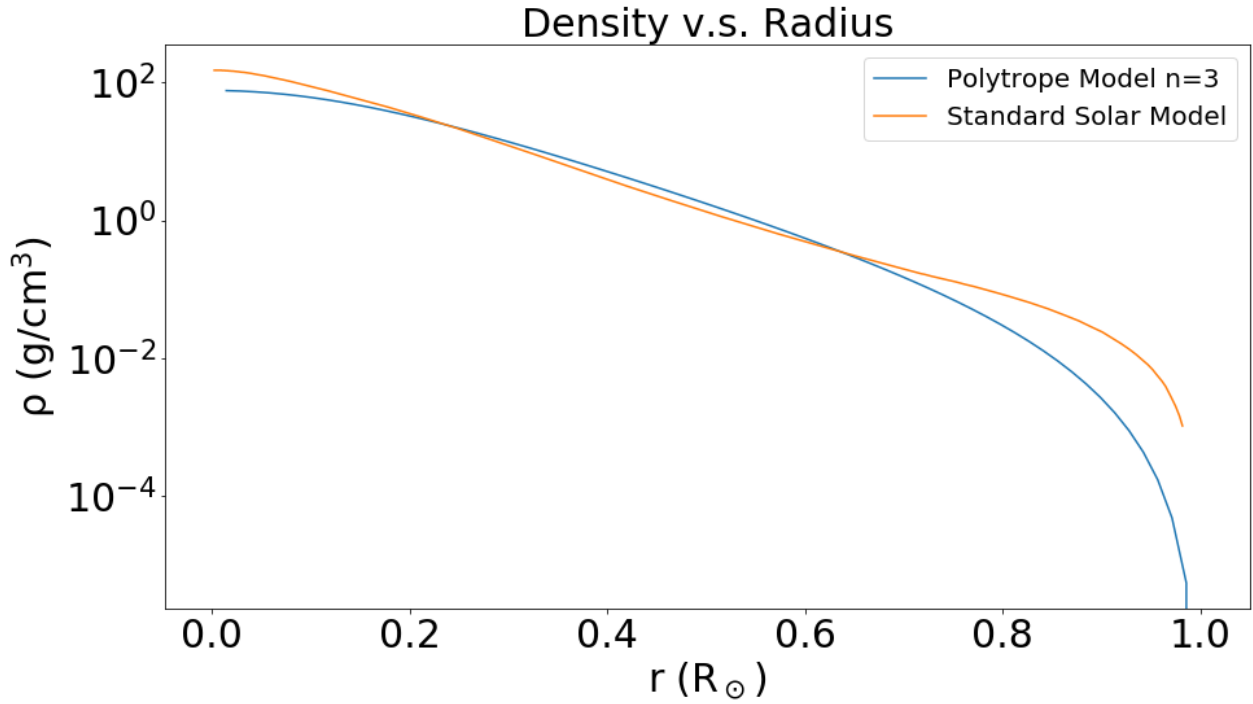
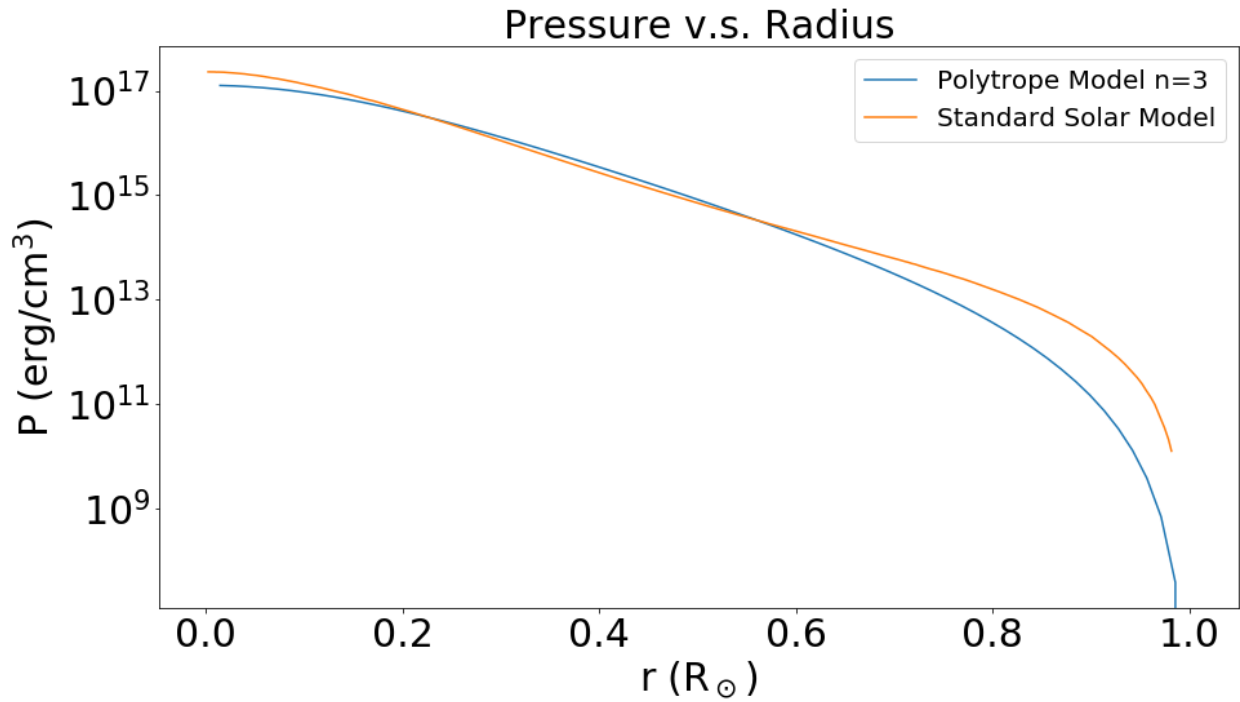
**Figure 7****Figure 8**

Figure 7 and 8 show the density and pressure distributions of these two models. Figure 9 demonstrates the pressure versus density. We can see that they are very similar so we can conclude that this polytrope model with $n=3$ is very accurate compared with Standard Solar Model.

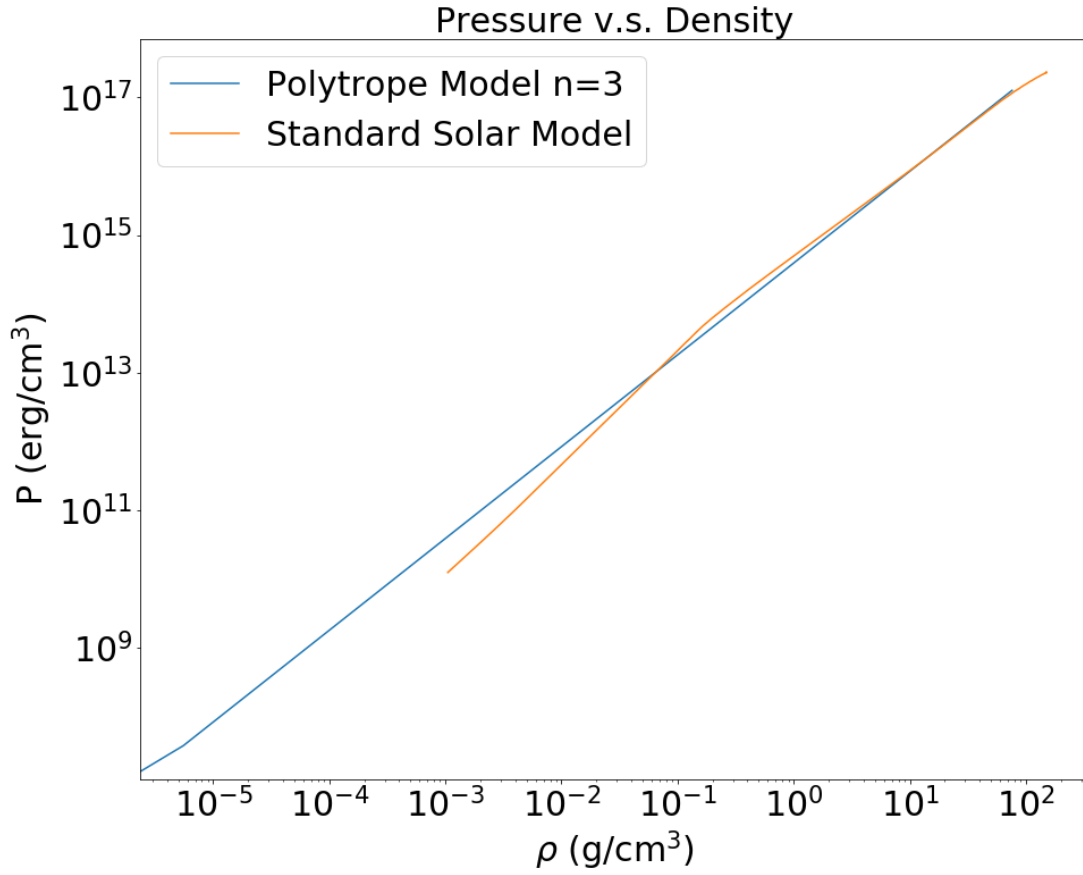


Figure 9

Q2

(a)
$$F_{conv} = \rho c_p T \left(\frac{l_c}{H_p} \right)^2 \sqrt{\frac{1}{2} g H_p} (\nabla - \nabla_{ad})^{\frac{3}{2}}$$

And we have:

$$F_{conv} = \frac{F(r)}{4\pi r^2} = \frac{L_\odot}{\pi R_\odot^2}$$

$$g = G \frac{m(r)}{r^2} = \frac{2GM_\odot}{R_\odot^2}$$

$$H_p = \frac{P}{\rho g} = \frac{kT}{\mu m_H g}$$

$$T = 10^7 \text{ K}, \rho = 1 \text{ g/cm}^3, X = 0.733, Y = 0.253, Z = 0.014$$

$$\therefore \frac{\delta T}{T} = \frac{l_c}{H_p} (\nabla - \nabla_{ad}) \approx 3.8 \times 10^{-8}$$

(b)

$$v_c \approx \sqrt{\frac{1}{2} l_c g \frac{\delta T}{T}} = 1.91 \times 10^3 \text{ cm/s}$$

$$c_s = \left(\frac{kT}{\mu m_H} \right)^{0.5} \approx 3.7 \times 10^7 \text{ cm/s}$$

$$\frac{v_c}{c_s} \approx 5 \times 10^{-5}$$

So the velocity of the convective element is much smaller than sound speed.

(c)
$$\beta = \frac{\rho v_c^2}{P} \sim \frac{\rho v_c^2}{nkT} \sim \left(\frac{v_c}{c_s} \right)^2 \sim 10^{-9}$$

So the convection doesn't alter significantly the hydrostatic structure of the region.

(d)

The crossing time is:

$$t_c = \frac{l_c}{v_c} \approx 1.8 \times 10^6 \text{ s} \approx 5.8 \times 10^{-2} \text{ yr}$$

The thermal timescale is:

$$t_{KH} = \frac{GM_\odot^2}{2R_\odot L_\odot} \sim 10^7 \text{ yr}$$

The nuclear timescale is:

$$t_{nuc} = \phi f_{nuc} \frac{M_\odot c^2}{L_\odot} \sim 10^{10} \text{ yr}$$

So the times are:

$$\frac{t_{th}}{t_c} \sim 10^8, \frac{t_{nuc}}{t_c} \sim 10^{11}$$

So that over a thermal timescale, and certainly over a nuclear timescale, convective elements will cross the region many times and makes convective region inside a star mixed homogeneously.

Q3

(a)

The value I use is: $m_e = 0.511 \text{ MeV}$

$$m_{^1H} = 938.272 \text{ MeV}, m_{^2H} = 1875.613 \text{ MeV}$$

$$m_{^3He} = 2808.392 \text{ MeV}, m_{^4He} = 3727.378 \text{ MeV}$$

Then I calculate the energy released in each reaction:

Table 1: PP1 energy released

Reaction of PPI	Energy released (MeV)
$^1H + ^1H \rightarrow ^2H + e^+ + \nu$	E1=1.442
$^2H + ^1H \rightarrow ^3He + \gamma$	E2=5.493
$^3He + ^3He \rightarrow ^4He + 2 ^1H$	E3=12.861

(b)

The total energy released in the PPI reaction is:

$$Q_{PP1} = 2E_1 + 2E_2 + 2E_3 + E_4 = 26.73 \text{ MeV}$$

(c)

The average energy of neutrino is:

$$\langle E_{\nu 1} \rangle = 0.265 \text{ MeV}, \langle E_{\nu 2} \rangle = 0.814 \text{ MeV}, \langle E_{\nu 3} \rangle = 6.71 \text{ MeV}$$

Then the energy of the three pp-chain release is:

$$Q_{PP1} = 26.2 \text{ MeV}, Q_{PP2} = 25.66 \text{ MeV}, Q_{PP3} = 19.76 \text{ MeV}$$

So PP1 release the most energy.

The energy released by unit mass is:

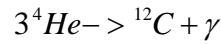
$$\frac{Q_{PP1}}{m_{He}} = 6.30 \times 10^{18} \text{ erg/g}$$

$$\frac{Q_{PP2}}{m_{He}} = 6.18 \times 10^{18} \text{ erg/g}$$

$$\frac{Q_{PP3}}{m_{He}} = 4.76 \times 10^{18} \text{ erg/g}$$

(d)

The triple- α reaction is:



The energy released is:

$$Q_{3\alpha} = (3m_{He} - m_C)c^2 = 7.275 \text{ MeV}$$

The energy release per unit mass is:

$$\frac{Q_{3\alpha}}{m_c} = 5.9 \times 10^{17} \text{ erg/g}$$

(e)

The mass of the core is:

$$m = 0.2M_{\odot} = 3.978 \times 10^{32} \text{ g}$$

The total energy release is:

$$E_{PP1} = 2.51 \times 10^{51} \text{ erg}, E_{3\alpha} = 2.35 \times 10^{50} \text{ erg}$$

So the time is:\

$$t_{PP1} = \frac{E_{PP1}}{L_{\odot}} = 6.54 \times 10^{17} \text{ s} = 2.07 \times 10^{10} \text{ yr}$$

$$t_{3\alpha} = \frac{E_{3\alpha}}{L_{\odot}} = 6.13 \times 10^{16} \text{ s} = 1.94 \times 10^9 \text{ yr}$$

Q4

(a)

the condition is :

$$F_{rot} = m\omega^2 R \ll F_{gra} = G \frac{Mm}{R^2}$$

$$\Leftrightarrow \delta = \frac{\omega^2 R^3}{GM} \ll 1$$

(b)

Table 2: rotation parameters

Spectral type	M/M_{\odot}	R/R_{\odot}	V_{rot}	δ
O5	60	12	190	0.038
B2.5	7	3.23	200	0.097
B5	4	2.35	210	0.136
A0	2.5	1.81	190	0.137
A7	1.8	1.55	160	0.116
F0	1.6	1.48	95	0.044
F5	1.4	1.36	25	3.2×10^{-3}
G0	1.2	1.13	12	7.1×10^{-4}
G2	1	1	2.3	2.8×10^{-5}

As **table 2** shows, δ of A0 is the largest so the largest departure for spherical symmetry occurs in early A stars.

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I thank my roommate Yuxuan Pang, for we did so many discussions. I also thank Jiayi Tang and Tai Zhou.

Reference:

Onno pols 《STELLAR STRUCTURE AND EVOLUTION》
Wiki