

Search for top squarks and dark matter particles in opposite-charge dilepton final states at $\sqrt{s} = 13$ TeV

A. M. Sirunyan *et al.*^{*}
(CMS Collaboration)



(Received 2 November 2017; published 15 February 2018)

A search for new physics is presented in final states with two oppositely charged leptons (electrons or muons), jets identified as originating from b quarks, and missing transverse momentum (p_T^{miss}). The search uses proton-proton collision data at $\sqrt{s} = 13$ TeV amounting to 35.9 fb^{-1} of integrated luminosity collected using the CMS detector in 2016. Hypothetical signal events are efficiently separated from the dominant $t\bar{t}$ background with requirements on p_T^{miss} and transverse-mass variables. No significant deviation is observed from the expected background. Exclusion limits are set in the context of simplified supersymmetric models with pair-produced top squarks. For top squarks, decaying exclusively to a top quark and a neutralino, exclusion limits are placed at 95% confidence level on the mass of the lightest top squark up to 800 GeV and on the lightest neutralino up to 360 GeV. These results, combined with searches in the single-lepton and all-jet final states, raise the exclusion limits up to 1050 GeV for the lightest top squark and up to 500 GeV for the lightest neutralino. For top squarks undergoing a cascade decay through charginos and sleptons, the mass limits reach up to 1300 GeV for top squarks and up to 800 GeV for the lightest neutralino. The results are also interpreted in a simplified model with a dark matter (DM) particle coupled to the top quark through a scalar or pseudoscalar mediator. For light DM, mediator masses up to 100 (50) GeV are excluded for scalar (pseudoscalar) mediators. The result for the scalar mediator achieves some of the most stringent limits to date in this model.

DOI: [10.1103/PhysRevD.97.032009](https://doi.org/10.1103/PhysRevD.97.032009)

I. INTRODUCTION

The top quark couples to the Higgs boson more strongly than other fermions because of its large mass. As a result, it plays a prominent role in the so-called hierarchy problem [1,2] of the standard model (SM) of particle physics, since its dominant contribution in the loop corrections to the Higgs boson mass exposes the theory to higher energy scales present in nature. Supersymmetry (SUSY) [3–10] is a well-motivated theory beyond the SM that provides a solution to the hierarchy problem. In addition, in R -parity conserving SUSY [11], the lightest SUSY particle (LSP) is stable and can be a viable dark matter (DM) candidate, assuming it is neutral and weakly interacting. Presently, the lighter SUSY particles may have masses in the TeV range and therefore could be produced in proton-proton (pp) collisions at the CERN LHC. The scalar partners of the right- and left-handed top quarks, the top squarks \tilde{t}_R and \tilde{t}_L , can be among these particles. These two states mix into the

mass eigenstates \tilde{t}_1 and \tilde{t}_2 . The lighter one, \tilde{t}_1 , could be within the LHC energy reach to provide a natural solution to the hierarchy problem [12], which strongly motivates searches for top squark production.

In this paper, we present a search for top squark pair production in a final state with two leptons (electrons or muons), hadronic jets identified as originating from b quarks, and significant transverse momentum imbalance. The search is performed using data from pp collisions collected with the CMS detector at the LHC during 2016 at a center-of-mass energy of 13 TeV, corresponding to an integrated luminosity of 35.9 fb^{-1} . We employ an efficient background reduction strategy that suppresses the large background from SM $t\bar{t}$ events by several orders of magnitude through use of dedicated transverse-mass variables [13,14]. The predicted SM backgrounds in the various search regions are validated in data control samples orthogonal in selection to the signal regions in data.

The search is interpreted in simplified models [15–17] describing the strong production of pairs of top squarks. We consider different decay modes, following the naming convention in Ref. [18]. In the T2 $t\bar{t}$ model (Fig. 1, upper left), each top squark decays into a top quark and the lightest neutralino $\tilde{\chi}_1^0$, which is the LSP. Alternatively, we consider the T2 bW model (Fig. 1, upper right), where both top squarks decay into a b quark and an intermediate

^{*}Full author list given at the end of the article.

Published by the American Physical Society under the terms of the [Creative Commons Attribution 4.0 International license](https://creativecommons.org/licenses/by/4.0/). Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Funded by SCOAP³.

the computing centers and personnel of the Worldwide LHC Computing Grid for delivering so effectively the computing infrastructure essential to our analyses. Finally, we acknowledge the enduring support for the construction and operation of the LHC and the CMS detector provided by the following funding agencies: BMFWF and FWF (Austria); FNRS and FWO (Belgium); CNPq, CAPES, FAPERJ, and FAPESP (Brazil); MES (Bulgaria); CERN; CAS, MoST, and NSFC (China); COLCIENCIAS (Colombia); MSES and CSF (Croatia); RPF (Cyprus); SENESCYT (Ecuador); MoER, ERC IUT, and ERDF (Estonia); Academy of Finland, MEC, and HIP (Finland); CEA and CNRS/IN2P3 (France); BMBF, DFG, and HGF (Germany); GSRT (Greece); OTKA and NIH (Hungary); DAE and DST (India); IPM (Iran); SFI (Ireland); INFN (Italy); MSIP and NRF (Republic of Korea); LAS (Lithuania); MOE and UM (Malaysia); BUAP, CINVESTAV, CONACYT, LNS, SEP, and UASLP-FAI (Mexico); MBIE (New Zealand); PAEC (Pakistan); MSHE and NSC (Poland); FCT (Portugal); JINR (Dubna); MON, RosAtom, RAS, RFBR and RAEP (Russia); MESTD (Serbia); SEIDI, CPAN, PCTI and FEDER (Spain); Swiss Funding Agencies (Switzerland); MST (Taipei); ThEPCenter, IPST, STAR, and NSTDA (Thailand); TUBITAK and TAEK (Turkey); NASU and SFFR (Ukraine); STFC (U.K.); DOE and NSF (U.S.). Individuals have received support from the Marie-Curie

program and the European Research Council and Horizon 2020 Grant, Contract No. 675440 (European Union); the Leventis Foundation; the A. P. Sloan Foundation; the Alexander von Humboldt Foundation; the Belgian Federal Science Policy Office; the Fonds pour la Formation à la Recherche dans l'Industrie et dans l'Agriculture (FRIA-Belgium); the Agentschap voor Innovatie door Wetenschap en Technologie (IWT-Belgium); the Ministry of Education, Youth and Sports (MEYS) of the Czech Republic; the Council of Science and Industrial Research, India; the HOMING PLUS program of the Foundation for Polish Science, cofinanced from European Union, Regional Development Fund, the Mobility Plus program of the Ministry of Science and Higher Education, the National Science Center (Poland), contracts Harmonia 2014/14/M/ST2/00428, Opus 2014/13/B/ST2/02543, 2014/15/B/ST2/03998, and 2015/19/B/ST2/02861, Sonata-bis 2012/07/E/ST2/01406; the National Priorities Research Program by Qatar National Research Fund; the Programa Severo Ochoa del Principado de Asturias; the Thalís and Aristeia programs cofinanced by EU-ESF and the Greek NSRF; the Rachadapisek Sompot Fund for Postdoctoral Fellowship, Chulalongkorn University and the Chulalongkorn Academic into Its 2nd Century Project Advancement Project (Thailand); the Welch Foundation, Contract No. C-1845; and the Weston Havens Foundation (U.S.).

-
- [1] Edward Witten, Dynamical breaking of supersymmetry, *Nucl. Phys.* **B188**, 513 (1981).
 - [2] S. Dimopoulos and H. Georgi, Softly broken supersymmetry and SU(5), *Nucl. Phys.* **B193**, 150 (1981).
 - [3] P. Ramond, Dual theory for free fermions, *Phys. Rev. D* **3**, 2415 (1971).
 - [4] Y. A. Gol'fand and E. P. Likhtman, Extension of the algebra of Poincaré group generators and violation of P invariance, *JETP Lett.* **13**, 323 (1971).
 - [5] A. Neveu and J. H. Schwarz, Factorizable dual model of pions, *Nucl. Phys.* **B31**, 86 (1971).
 - [6] D. V. Volkov and V. P. Akulov, Possible universal neutrino interaction, *JETP Lett.* **16**, 438 (1972).
 - [7] J. Wess and B. Zumino, A lagrangian model invariant under supergauge transformations, *Phys. Lett. B* **49**, 52 (1974).
 - [8] J. Wess and B. Zumino, Supergauge transformations in four dimensions, *Nucl. Phys.* **B70**, 39 (1974).
 - [9] P. Fayet, Supergauge invariant extension of the Higgs mechanism and a model for the electron and its neutrino, *Nucl. Phys.* **B90**, 104 (1975).
 - [10] H. P. Nilles, Supersymmetry, supergravity and particle physics, *Phys. Rep.* **110**, 1 (1984).
 - [11] G. R. Farrar and P. Fayet, Phenomenology of the production, decay, and detection of new hadronic states associated with supersymmetry, *Phys. Lett. B* **76**, 575 (1978).
 - [12] M. Papucci, J. T. Ruderman, and A. Weiler, Natural SUSY endures, *J. High Energy Phys.* **09** (2012) 035.
 - [13] J. Smith, W. L. van Neerven, and J. A. M. Vermaseren, The Transverse Mass and Width of the W Boson, *Phys. Rev. Lett.* **50**, 1738 (1983).
 - [14] C. G. Lester and D. J. Summers, Measuring masses of semi-invisibly decaying particles pair produced at hadron colliders, *Phys. Lett. B* **463**, 99 (1999).
 - [15] J. Alwall, P. Schuster, and N. Toro, Simplified models for a first characterization of new physics at the LHC, *Phys. Rev. D* **79**, 075020 (2009).
 - [16] J. Alwall, M.-P. Le, M. Lisanti, and J. G. Wacker, Model-independent jets plus missing energy searches, *Phys. Rev. D* **79**, 015005 (2009).
 - [17] D. Alves (LHC New Physics Working Group), Simplified models for LHC new physics searches, *J. Phys. G* **39**, 105005 (2012).
 - [18] CMS Collaboration, Interpretation of searches for supersymmetry with simplified models, *Phys. Rev. D* **88**, 052017 (2013).

E. Palencia Cortezon,¹¹⁰ S. Sanchez Cruz,¹¹⁰ P. Vischia,¹¹⁰ J. M. Vizan Garcia,¹¹⁰ I. J. Cabrillo,¹¹¹ A. Calderon,¹¹¹ B. Chazin Quero,¹¹¹ E. Curras,¹¹¹ J. Duarte Campderros,¹¹¹ M. Fernandez,¹¹¹ J. Garcia-Ferrero,¹¹¹ G. Gomez,¹¹¹ A. Lopez Virto,¹¹¹ J. Marco,¹¹¹ C. Martinez Rivero,¹¹¹ P. Martinez Ruiz del Arbol,¹¹¹ F. Matorras,¹¹¹ J. Piedra Gomez,¹¹¹ T. Rodrigo,¹¹¹ A. Ruiz-Jimeno,¹¹¹ L. Scodellaro,¹¹¹ N. Trevisani,¹¹¹ I. Vila,¹¹¹ R. Vilar Cortabitarte,¹¹¹ D. Abbaneo,¹¹² B. Akgun,¹¹² E. Auffray,¹¹² P. Baillon,¹¹² A. H. Ball,¹¹² D. Barney,¹¹² J. Bendavid,¹¹² M. Bianco,¹¹² P. Bloch,¹¹² A. Bocci,¹¹² C. Botta,¹¹² T. Camporesi,¹¹² R. Castello,¹¹² M. Cepeda,¹¹² G. Cerminara,¹¹² E. Chapon,¹¹² Y. Chen,¹¹² D. d'Enterria,¹¹² A. Dabrowski,¹¹² V. Daponte,¹¹² A. David,¹¹² M. De Gruttola,¹¹² A. De Roeck,¹¹² N. Deelen,¹¹² M. Dobson,¹¹² T. du Pree,¹¹² M. Dünser,¹¹² N. Dupont,¹¹² A. Elliott-Peisert,¹¹² P. Everaerts,¹¹² F. Fallavollita,¹¹² G. Franzoni,¹¹² J. Fulcher,¹¹² W. Funk,¹¹² D. Gigi,¹¹² A. Gilbert,¹¹² K. Gill,¹¹² F. Glege,¹¹² D. Gulhan,¹¹² P. Harris,¹¹² J. Hegeman,¹¹² V. Innocente,¹¹² A. Jafari,¹¹² P. Janot,¹¹² O. Karacheban,^{112,r} J. Kieseler,¹¹² V. Knünz,¹¹² A. Kornmayer,¹¹² M. J. Kortelainen,¹¹² M. Krammer,^{112,b} C. Lange,¹¹² P. Lecoq,¹¹² C. Lourenço,¹¹² M. T. Lucchini,¹¹² L. Malgeri,¹¹² M. Mannelli,¹¹² A. Martelli,¹¹² F. Meijers,¹¹² J. A. Merlin,¹¹² S. Mersi,¹¹² E. Meschi,¹¹² P. Milenovic,^{112,r} F. Moortgat,¹¹² M. Mulders,¹¹² H. Neugebauer,¹¹² J. Ngadiuba,¹¹² S. Orfanelli,¹¹² L. Orsini,¹¹² L. Pape,¹¹² E. Perez,¹¹² M. Peruzzi,¹¹² A. Petrilli,¹¹² G. Petrucciani,¹¹² A. Pfeiffer,¹¹² M. Pierini,¹¹² D. Rabadý,¹¹² A. Racz,¹¹² T. Reis,¹¹² G. Rolandi,^{112,ss} M. Rovere,¹¹² H. Sakulin,¹¹² C. Schäfer,¹¹² C. Schwick,¹¹² M. Seidel,¹¹² M. Selvaggi,¹¹² A. Sharma,¹¹² P. Silva,¹¹² P. Sphicas,^{112,t} A. Stakia,¹¹² J. Steggemann,¹¹² M. Stoye,¹¹² M. Tosi,¹¹² D. Treille,¹¹² A. Triossi,¹¹² A. Tsirou,¹¹² V. Veckalns,^{112,uu} M. Verweij,¹¹² W. D. Zeuner,¹¹² W. Bertl,^{113,a} L. Caminada,^{113,vv} K. Deiters,¹¹³ W. Erdmann,¹¹³ R. Horisberger,¹¹³ Q. Ingram,¹¹³ H. C. Kaestli,¹¹³ D. Kotlinski,¹¹³ U. Langenegger,¹¹³ T. Rohe,¹¹³ S. A. Wiederkehr,¹¹³ M. Backhaus,¹¹⁴ L. Bäni,¹¹⁴ P. Berger,¹¹⁴ L. Bianchini,¹¹⁴ B. Casal,¹¹⁴ G. Dissertori,¹¹⁴ M. Dittmar,¹¹⁴ M. Donegà,¹¹⁴ C. Dorfer,¹¹⁴ C. Grab,¹¹⁴ C. Heidegger,¹¹⁴ D. Hits,¹¹⁴ J. Hoss,¹¹⁴ G. Kasieczka,¹¹⁴ T. Klijnsma,¹¹⁴ W. Lustermann,¹¹⁴ B. Mangano,¹¹⁴ M. Marionneau,¹¹⁴ M. T. Meinhard,¹¹⁴ D. Meister,¹¹⁴ F. Micheli,¹¹⁴ P. Musella,¹¹⁴ F. Nessi-Tedaldi,¹¹⁴ F. Pandolfi,¹¹⁴ J. Pata,¹¹⁴ F. Pauss,¹¹⁴ G. Perrin,¹¹⁴ L. Perrozzi,¹¹⁴ M. Quittnat,¹¹⁴ M. Reichmann,¹¹⁴ D. A. Sanz Becerra,¹¹⁴ M. Schönenberger,¹¹⁴ L. Shchutska,¹¹⁴ V. R. Tavolaro,¹¹⁴ K. Theofilatos,¹¹⁴ M. L. Vesterbacka Olsson,¹¹⁴ R. Wallny,¹¹⁴ D. H. Zhu,¹¹⁴ T. K. Aarrestad,¹¹⁵ C. Amsler,^{115,ww} M. F. Canelli,¹¹⁵ A. De Cosa,¹¹⁵ R. Del Burgo,¹¹⁵ S. Donato,¹¹⁵ C. Galloni,¹¹⁵ T. Hreus,¹¹⁵ B. Kilminster,¹¹⁵ D. Pinna,¹¹⁵ G. Rauco,¹¹⁵ P. Robmann,¹¹⁵ D. Salerno,¹¹⁵ K. Schweiger,¹¹⁵ C. Seitz,¹¹⁵ Y. Takahashi,¹¹⁵ A. Zucchetta,¹¹⁵ V. Candelise,¹¹⁶ T. H. Doan,¹¹⁶ Sh. Jain,¹¹⁶ R. Khurana,¹¹⁶ C. M. Kuo,¹¹⁶ W. Lin,¹¹⁶ A. Pozdnyakov,¹¹⁶ S. S. Yu,¹¹⁶ Arun Kumar,¹¹⁷ P. Chang,¹¹⁷ Y. Chao,¹¹⁷ K. F. Chen,¹¹⁷ P. H. Chen,¹¹⁷ F. Fiori,¹¹⁷ W.-S. Hou,¹¹⁷ Y. Hsiung,¹¹⁷ Y. F. Liu,¹¹⁷ R.-S. Lu,¹¹⁷ E. Paganis,¹¹⁷ A. Psallidas,¹¹⁷ A. Steen,¹¹⁷ J. f. Tsai,¹¹⁷ B. Asavapibhop,¹¹⁸ K. Kovitanggoon,¹¹⁸ G. Singh,¹¹⁸ N. Srimanobhas,¹¹⁸ M. N. Bakirci,^{119,xx} A. Bat,¹¹⁹ F. Boran,¹¹⁹ S. Cerci,^{119,yy} S. Damarseckin,¹¹⁹ Z. S. Demiroglu,¹¹⁹ C. Dozen,¹¹⁹ S. Girgis,¹¹⁹ G. Gokbulut,¹¹⁹ Y. Guler,¹¹⁹ I. Hos,^{119,zz} E. E. Kangal,^{119,aaa} O. Kara,¹¹⁹ U. Kiminsu,¹¹⁹ M. Oglakci,¹¹⁹ G. Onengut,^{119,bbb} K. Ozdemir,^{119,ccc} S. Ozturk,^{119,xx} A. Polatoz,¹¹⁹ U. G. Tok,¹¹⁹ H. Topakli,^{119,xx} S. Turkcapar,¹¹⁹ I. S. Zorbakir,¹¹⁹ C. Zorbilmez,¹¹⁹ B. Bilin,¹²⁰ G. Karapinar,^{120,ddd} K. Ocalan,^{120,eee} M. Yalvac,¹²⁰ M. Zeyrek,¹²⁰ E. Gülmez,¹²¹ M. Kaya,^{121,fff} O. Kaya,^{121,ggg} S. Tekten,¹²¹ E. A. Yetkin,^{121,hhh} M. N. Agaras,¹²² S. Atay,¹²² A. Cakir,¹²² K. Cankocak,¹²² I. Köseoglu,¹²² B. Grynyov,¹²³ L. Levchuk,¹²⁴ F. Ball,¹²⁵ L. Beck,¹²⁵ J. J. Brooke,¹²⁵ D. Burns,¹²⁵ E. Clement,¹²⁵ D. Cussans,¹²⁵ O. Davignon,¹²⁵ H. Flacher,¹²⁵ J. Goldstein,¹²⁵ G. P. Heath,¹²⁵ H. F. Heath,¹²⁵ L. Kreczko,¹²⁵ D. M. Newbold,^{125,iii} S. Paramesvaran,¹²⁵ T. Sakuma,¹²⁵ S. Seif El Nasr-storey,¹²⁵ D. Smith,¹²⁵ V. J. Smith,¹²⁵ K. W. Bell,¹²⁶ A. Belyaev,^{126,jjj} C. Brew,¹²⁶ R. M. Brown,¹²⁶ L. Calligaris,¹²⁶ D. Cieri,¹²⁶ D. J. A. Cockerill,¹²⁶ J. A. Coughlan,¹²⁶ K. Harder,¹²⁶ S. Harper,¹²⁶ J. Linacre,¹²⁶ E. Olaiya,¹²⁶ D. Petyt,¹²⁶ C. H. Shepherd-Themistocleous,¹²⁶ A. Thea,¹²⁶ I. R. Tomalin,¹²⁶ T. Williams,¹²⁶ G. Auzinger,¹²⁷ R. Bainbridge,¹²⁷ J. Borg,¹²⁷ S. Breeze,¹²⁷ O. Buchmuller,¹²⁷ A. Bundock,¹²⁷ S. Casasso,¹²⁷ M. Citron,¹²⁷ D. Colling,¹²⁷ L. Corpe,¹²⁷ P. Dauncey,¹²⁷ G. Davies,¹²⁷ A. De Wit,¹²⁷ M. Della Negra,¹²⁷ R. Di Maria,¹²⁷ A. Elwood,¹²⁷ Y. Haddad,¹²⁷ G. Hall,¹²⁷ G. Iles,¹²⁷ T. James,¹²⁷ R. Lane,¹²⁷ C. Laner,¹²⁷ L. Lyons,¹²⁷ A.-M. Magnan,¹²⁷ S. Malik,¹²⁷ L. Mastrolorenzo,¹²⁷ T. Matsushita,¹²⁷ J. Nash,¹²⁷ A. Nikitenko,^{127,h} V. Palladino,¹²⁷ M. Pesaresi,¹²⁷ D. M. Raymond,¹²⁷ A. Richards,¹²⁷ A. Rose,¹²⁷ E. Scott,¹²⁷ C. Seez,¹²⁷ A. Shtipliyski,¹²⁷ S. Summers,¹²⁷ A. Tapper,¹²⁷ K. Uchida,¹²⁷ M. Vazquez Acosta,^{127,kkk} T. Virdee,^{127,o} N. Wardle,¹²⁷ D. Winterbottom,¹²⁷ J. Wright,¹²⁷ S. C. Zenz,¹²⁷ J. E. Cole,¹²⁸ P. R. Hobson,¹²⁸ A. Khan,¹²⁸ P. Kyberd,¹²⁸ I. D. Reid,¹²⁸ L. Teodorescu,¹²⁸ M. Turner,¹²⁸ S. Zahid,¹²⁸ A. Borzou,¹²⁹ K. Call,¹²⁹ J. Dittmann,¹²⁹ K. Hatakeyama,¹²⁹ H. Liu,¹²⁹ N. Pastika,¹²⁹ C. Smith,¹²⁹ R. Bartek,¹³⁰ A. Dominguez,¹³⁰ A. Buccilli,¹³¹ S. I. Cooper,¹³¹ C. Henderson,¹³¹ P. Rumerio,¹³¹ C. West,¹³¹ D. Arcaro,¹³² A. Avetisyan,¹³² T. Bose,¹³² D. Gastler,¹³² D. Rankin,¹³² C. Richardson,¹³² J. Rohlf,¹³² L. Sulak,¹³² D. Zou,¹³² G. Benelli,¹³³ D. Cutts,¹³³ A. Garabedian,¹³³ M. Hadley,¹³³