

## Summary as of May 16, 2016

We assume that the following contributions comprise the effective action:

$$\begin{aligned} \frac{4\pi}{N} \mathcal{L}_{\text{eff}} = & - \int \left( d^4\theta \frac{1}{2} \left| \ln \sqrt{2\Sigma} \right|^2 + G(u, \bar{u}) + \right. \\ & \left. + i d^2\tilde{\theta} \left( \sqrt{2\Sigma} \ln \sqrt{2\Sigma} - \sqrt{2\Sigma} \right) + i d^2\bar{\tilde{\theta}} \left( \sqrt{2\bar{\Sigma}} \ln \sqrt{2\bar{\Sigma}} - \sqrt{2\bar{\Sigma}} \right) \right). \end{aligned} \quad (1)$$

Matching the effective potential and the kinetic term  $F_{\mu\nu}^2$  from ARXIV:0803.0698 gives

$$G_{u\bar{u}} = - \frac{1}{x^4} \left( 1 + 2 \frac{y^2}{x^2} \right) \int_0^x \ln(1+x) dx - \frac{1}{6} \frac{y^2}{x^3(1+x)}. \quad (2)$$

Here

$$\begin{aligned} u &= \frac{S}{\sqrt{2\Sigma}} \propto \frac{iD - F_{03}}{\phi} & \bar{u} &= \frac{\bar{S}}{\sqrt{2\bar{\Sigma}}} \propto \frac{iD + F_{03}}{\phi} \\ x &= \frac{u + \bar{u}}{2} \propto \frac{iD}{\phi} & y &= \frac{\bar{u} - u}{2} \propto \frac{F_{03}}{\phi}, \end{aligned}$$

neglecting fermions.

The effective action (1) produces the following kinetic term for  $\lambda$ :

$$- \frac{\bar{\lambda}_R i \overleftrightarrow{\partial} \lambda_R + \bar{\lambda}_L i \overleftrightarrow{\partial} \lambda_L}{\phi} x^2 \left( 3 G_{u\bar{u}} + x (G_{uu\bar{u}} + G_{\bar{u}\bar{u}u}) \right). \quad (3)$$

Substituting (2) gives,

$$- \frac{\bar{\lambda}_R i \overleftrightarrow{\partial} \lambda_R + \bar{\lambda}_L i \overleftrightarrow{\partial} \lambda_L}{\phi} \frac{\ln(1+x) - x}{x^2}, \quad (4)$$

which is the same coefficient that we obtained a year ago using our previous approach. This also is the *correct* answer for that diagram.

As for the kinetic term for  $\sigma$ , the correct expression is, as we know,

$$\frac{4\pi}{N} \mathcal{L}_{\text{eff}} \supset \frac{1}{12} \frac{(\partial_\mu \phi)^2}{(\phi + iD)^2}. \quad (5)$$

The supersymmetric Lagrangian (1) produces the following expression,

$$\frac{4\pi}{N} \mathcal{L}_{\text{eff}} \supset - \frac{(\partial_\mu \phi)^2}{\phi^2} x^2 \left( 3 G_{u\bar{u}} + x (G_{uu\bar{u}} + G_{\bar{u}\bar{u}u}) \right). \quad (6)$$

Notice that the bracket is the same as in (3), and so gives a logarithm.