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The Mixed A_μ - b Propagator

The anomaly term is $\sim \frac{N}{4\pi} F_{03} \cdot \text{Im } \frac{\tilde{S}^b}{S^b} = -2i \frac{N}{4\pi} F_{03} \cdot \text{Arg } \tilde{S}^b$

When the phase of S^b is small, $2i \text{Arg } \tilde{S}^b \approx \frac{\tilde{S}^b - \tilde{S}^b}{2S^b}$.

Together with the kinetic terms, this gives,

$$\frac{4\pi}{N} \mathcal{L} \rightarrow \sum_{S^b \neq 0} \left[\frac{1}{4} F_{\mu\nu}^2 + \frac{1}{2} (\partial_\mu S^b)^2 - F_{03} \frac{(\tilde{S}^b - S^b)}{2S^b} [\lambda S^b] \right]$$

Neglecting loops, we put S^b to its vacuum value λ

$$\frac{4\pi}{N} \mathcal{L}_{\text{quadratic}} = \frac{1}{4} \frac{F_{\mu\nu}^2}{\lambda^2} + \frac{1}{2} \frac{(\partial_\mu S^b)^2}{\lambda^2} - \lambda F_{03} \cdot 2i \text{Im } \tilde{S}^b$$

Redefining $A_\mu \rightarrow \lambda A_\mu$, $S^b \rightarrow \lambda S^b$, we find

$$\frac{4\pi}{N} \mathcal{L} = \frac{1}{4} F_{\mu\nu}^2 + \frac{1}{2} (\partial_\mu S^b)^2 - 2i\lambda F_{03} \cdot \tilde{S}^b,$$

$$\text{where } \tilde{S}^b \equiv \text{Im } S^b$$

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Now omit the real part,

$$\frac{4\epsilon}{N} \mathcal{L}_0 = \frac{1}{4} F_{\mu\nu}^2 + \frac{1}{2} (\partial_\mu A_\nu)^2 - 2i\lambda F_0 z \cdot \vec{A},$$

We also don't really need $\frac{4\epsilon}{N}$ for these calculations,

$$S_{\text{quadratic}} = \frac{1}{4} F_{\mu\nu}^2 + \frac{1}{2} (\partial_\mu A_\nu)^2 - 2i\lambda \epsilon_{\mu\nu} \partial_\mu A_\nu \cdot \vec{A} + \frac{1}{2\alpha} (\partial_\mu A_\nu)^2$$

$$\text{where } \epsilon_{\mu\nu} \partial_\mu A_\nu = \partial_0 A_3 - \partial_3 A_0 = F_0$$

We've also added the gauge-fixing term.

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Overall,

$$\mathcal{L}_{\text{quadratic}} = \frac{1}{4} \epsilon_{\mu\nu}^2 + \frac{1}{2\alpha} (\partial_\mu A^\mu)^2 + \frac{1}{2} (\partial_\mu \delta_a)^2 - 2i\lambda \epsilon_{\mu\nu} \partial_\mu A^\nu \cdot \delta_a =$$

$$\approx -\frac{1}{2} \left\{ \Delta_\mu \right\} \underbrace{\epsilon_{\mu\nu}^2}_{2i\lambda \epsilon_{\mu\nu} \partial_\mu \delta_a^{\mu\nu}} \left\{ \begin{array}{c} \Box(\mathcal{B}_L^{\mu\nu} + \frac{1}{2} \mathcal{B}_R^{\mu\nu}) \\ -2i\lambda \epsilon_{\mu\nu} \partial_\mu \mathcal{B}_R^{\mu\nu} \end{array} \right\} \left\{ \begin{array}{c} \partial^\mu \\ \partial^\nu \end{array} \right\} \left\{ \begin{array}{c} A^\mu \\ A^\nu \end{array} \right\} \left\{ \begin{array}{c} \delta_a^{\mu\nu} \\ \sim \mathcal{E}_R \end{array} \right\}$$

- Pay attention to the overall factor of $-\frac{1}{2}$

Although we work on the Euclidean space, we

denote $\partial_\mu \delta^\mu \equiv \square$

As customary, here $\mathcal{B}_R^{\mu\nu} \approx \frac{1}{17} \partial_\mu \delta_\nu$

$$\mathcal{B}_L^{\mu\nu} \approx \delta^{\mu\nu} - \mathcal{B}_R^{\mu\nu}$$

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So the kinetic matrix is

The propagator matrix must have the shape

A diagram consisting of several mathematical symbols arranged horizontally. From left to right, it includes: a left brace (bracelet), a right brace (bracelet), a left curly brace (bracelet), a right curly brace (bracelet), a left arrow, a right arrow, and a right brace (bracelet).

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Here's our guess for the first column of the propagator,

$$\left\{ \begin{array}{l} x \beta_L^{v_k} + y \beta_R^{v_k} \\ z \cdot \epsilon_{aL} \cdot \delta x \delta L^k \end{array} \right\}$$

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First row times first column:

$$x \square \beta_{\perp}^{\mu k} + \sum_{\alpha} \beta_{\parallel}^{\mu k} -$$

$$- 2i\lambda \epsilon_{\rho b} \delta_{\rho} \delta_{\mu}^k \cdot z \cdot \epsilon_{\alpha \lambda} \delta_{\alpha} \delta_{\lambda}^k =$$

$$\begin{array}{c} \vdots \\ \vdots \\ \vdots \\ \vdots \\ \vdots \end{array} - - - - - \quad \begin{array}{c} \vdots \\ \vdots \\ \vdots \\ \vdots \\ \vdots \end{array}$$

$$\epsilon_{\rho b} \epsilon_{\alpha \lambda} = \delta_{\rho}^{\alpha} \delta_{\lambda}^b$$

$$= x \square \beta_{\perp}^{\mu k} + \sum_{\alpha} \beta_{\parallel}^{\mu k} -$$

$$- 2i\lambda \delta_{\rho}^{\alpha} \delta_{\mu}^k \cdot z \cdot \delta_{\rho} \delta_{\mu}^k \cdot \delta_{\alpha} \delta_{\lambda}^k =$$

$$= x \square \beta_{\perp}^{\mu k} + \sum_{\alpha} \beta_{\parallel}^{\mu k} -$$

$$- 2i\lambda z \delta_{\rho} \delta_{\mu}^k \cdot \delta_{\rho} \delta_{\mu}^k =$$

$$\overbrace{\square \delta^{\mu k} - \delta^{\mu k}} = \square \beta_{\perp}^{\mu k}$$

$$= (x \square - 2i\lambda z \square) \beta_{\perp}^{\mu k} + \sum_{\alpha} \beta_{\parallel}^{\mu k} \equiv \delta^{\mu k}$$

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We have

$$x_n - 2\lambda z_n \approx 1 \quad \left. \begin{array}{l} \\ \end{array} \right\}$$

$$y_n = 1$$

$$\underline{x} \approx \underline{y_n} + \underline{2\lambda z}$$

$$x \approx y_n + 2\lambda z$$

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Second row times First column:

$$2i\lambda \epsilon_{eb} \delta^{\nu\xi} \cdot (\overset{\nu\xi}{x \beta_2 + y \beta_{11}}) + \\ + \square \cdot \epsilon_{\alpha\lambda} \delta^{\nu\xi}$$

| Now $\delta_{\alpha\rho} \delta_{\beta\eta} \cdot \delta_\nu \delta_\xi = 0 \leftarrow \text{annihilates the parallel parts}$

| So, out of bracket $(x \dots + y \dots)$ only

$x \cdot \delta^{\nu\xi}$ will contribute

$$\approx 2i\lambda x \cdot \epsilon_{eb} \delta^{\nu\xi} + \square z \cdot \epsilon_{\alpha\lambda} \delta^{\nu\xi} \equiv 0.$$

$$\Rightarrow 2i\lambda x + \square z = 0$$

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$$\text{From p. 7, } x = \frac{1}{\omega} + 2i\lambda z$$

$$\Rightarrow x\omega = 1 + 2i\lambda z\omega \quad (9.1)$$

Coupled now with $2i\lambda x + \omega z = 0$ (what we just found)

$$\Rightarrow \omega z = -2i\lambda x$$

Plug it into (9.1)

$$x\omega = 1 + 2i\lambda(-2i\lambda x)$$

$$\text{or } x\omega = 1 + 4\lambda^2 - \omega x$$

$$\Rightarrow x = \frac{1}{\omega - 4\lambda^2}$$

$$\Rightarrow z = -\frac{2i\lambda x}{\omega} = -\frac{2i\lambda}{\omega(\omega - 4\lambda^2)}$$

$$y = \omega/\omega \quad (\text{see p. 7})$$

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Hypothesis about the second column:

$$\left. \begin{array}{l} \text{1 - z. } \mathbb{E}[\epsilon] \neq \delta_2 \\ \text{2. } u/n \end{array} \right\}$$

z should turn out to
be the same as on p. 9

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First row times second column:

$$\Leftrightarrow \left(\beta_1^{\mu\nu} + \gamma_\alpha \beta_{11}^{\mu\nu} \right) \cdot (-z) \epsilon_{\mu\nu} \delta_\alpha \delta_\lambda^\nu = 2i \lambda \epsilon_{\mu\nu} \delta_\alpha \delta_\lambda^\mu \cdot u_\alpha$$

↑
z again, only $\delta^{\mu\nu}$ survives z

$$z - z \square \cdot \epsilon_{\mu\nu} \delta_\alpha \delta_\lambda^\nu = 2i \lambda \epsilon_{\mu\nu} \delta_\alpha \delta_\lambda^\mu u_\alpha \approx 0$$

$$\Rightarrow -z \square - 2i \lambda u_\alpha / \alpha \approx 0$$

$$z = - \frac{2i \lambda u}{\alpha^2} \quad (11.1)$$

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Second row times second column!

$$2i\lambda \epsilon_{pb} \delta_p \delta_b^v \cdot (-z) \epsilon_{ax} \delta_a \delta_x^v + u =$$

$$= -2i\lambda z \cdot \delta_p \overset{[x \quad \lambda]}{\delta_b} \delta_p \delta_b^v \cdot \delta_a \delta_x^v + u =$$

$$= -2i\lambda z \cdot \delta_p \overset{x}{\delta_b} \cdot \delta_a \delta_b^v \cdot \delta_a \delta_x^v + u =$$

apply then

$$= -2i\lambda z \cdot \delta_a \delta_b^v \cdot \delta_p \delta_x^v + u =$$

$$= -2i\lambda z \cdot [\delta^{av} \square - \delta_p \delta_b \delta^{pb}] + u =$$

$$= -2i\lambda z \cdot \square + u = 1$$

$$\underline{\underline{u - 2i\lambda z \cdot \square = 1}}$$

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Equations (11.1) and (12.1):

$$z = -\frac{2i\lambda u}{\Omega^2}$$



$$u - 2i\lambda z \cdot \Omega = 1$$

$$u = 1 + 2i\lambda z \cdot \Omega$$

$$\Rightarrow z = -\frac{2i\lambda}{\Omega^2} (1 + 2i\lambda z \cdot \Omega) = -\frac{2i\lambda}{\Omega^2} + (4\lambda^2) \frac{z}{\Omega}$$

$$z (1 + \cancel{4\lambda^2/\Omega}) = -\frac{2i\lambda}{\Omega^2}$$

or,
$$z = -\frac{2i\lambda}{\Omega(\Omega - 4\lambda^2)}$$

Again,

was anticipated.

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Then

$$u = -\frac{\Omega^2}{2i\lambda} z_2 = -\frac{\Omega^2}{2i\lambda} \cdot (-1) \frac{2i\lambda}{\Omega(\Omega - 4\lambda^2)} =$$

$$z_2 = \frac{-\frac{\Omega}{\Omega - 4\lambda^2}}{\overbrace{\Omega - 4\lambda^2}^{\longrightarrow}}$$

Together,



$$z_2 = -\frac{2i\lambda}{\Omega(\Omega - 4\lambda^2)}$$

$$u = -\frac{\Omega}{\Omega - 4\lambda^2} \Rightarrow \text{element}(2,2) \text{ is}$$

$$u_2 = \frac{1}{\Omega - 4\lambda^2}$$

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The first column is (p. 9)

The second column is (p. 54)

$$\left\{ \begin{array}{l} \frac{\beta_2 v_k}{\Omega - 4\lambda^2} + \frac{\alpha}{\Omega} \beta_{11} v_k \\ - \frac{2i\lambda}{\Omega - 4\lambda^2} \text{Gal} \frac{\delta \omega \delta \lambda}{\Omega} \end{array} \right\}$$

$$\left\{ \begin{array}{l} \frac{2i\lambda}{\Omega - 4\lambda^2} \text{Gal} \frac{\delta \omega \delta \lambda}{\Omega} \\ \frac{1}{\Omega - 4\lambda^2} \end{array} \right\}$$

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Altogether, the propagator matrix is

$$\leftarrow \xi \rightarrow$$

$$\begin{pmatrix} & & & \\ & \frac{\beta_L^{\nu\xi}}{L - 4\lambda^2} + \alpha \frac{\beta_R^{\nu\xi}}{L} & & \\ & & \frac{i\lambda}{L - 4\lambda^2} \cdot \cancel{\epsilon} \cancel{\lambda} \frac{\delta_{\nu\nu} \delta_{\lambda\lambda}}{L} & \\ & & & \\ & - \frac{i\lambda}{L - 4\lambda^2} \cancel{\epsilon} \cancel{\lambda} \frac{\delta_{\nu\nu} \delta_{\lambda\lambda}}{L} & & \frac{1}{L - 4\lambda^2} \\ & & & \end{pmatrix}$$