

Lecture-14

1. Line integral

- (a) Path independence .
- (b) Conservative fields .

2. Green's Theorem .

EE1203: Vector Calculus

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Path Independence:

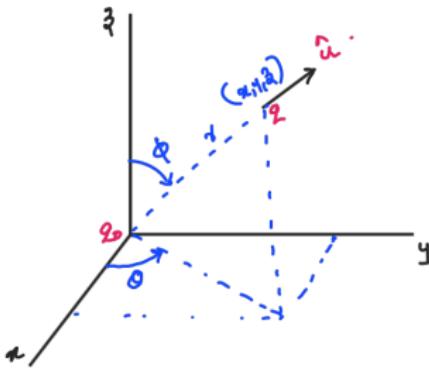
* In line integral the path of integration:

One of the ingredients that determines the very function we integrate.

* Under some conditions value of integral does not depend on the path.

Let us show this path independence in Coulomb's Force.

Let a charge q_0 be fixed at the origin;
another charge q be situated at (x, y, z) .



Coulomb Force on ' q ' is

$$\vec{F} = \frac{1}{4\pi\epsilon_0} \frac{q q_0}{r^2} \hat{u}$$

$$r = (x^2 + y^2 + z^2)^{1/2} \quad (\text{distance b/w}) \\ 2 \pi \epsilon_0$$

\hat{u} = unit vector pointing from q_0 to q

$$\hat{u} = \frac{\vec{r}}{|\vec{r}|} = \frac{\vec{r}}{r} = \frac{x\hat{i} + y\hat{j} + z\hat{k}}{r}$$

$$\vec{F} = \frac{1}{4\pi\epsilon_0} \frac{q q_0}{r^3} (x\hat{i} + y\hat{j} + z\hat{k})$$

$$\vec{F} = \langle F_x, F_y, F_z \rangle = \frac{1}{4\pi\epsilon_0} \frac{q q_0}{r^3} \langle x, y, z \rangle$$

$$\begin{aligned}\vec{F} \cdot \hat{t} \, ds &= F_x dx + F_y dy + F_z dz \\ &= \frac{q q_0}{4\pi\epsilon_0} \frac{x \, dx + y \, dy + z \, dz}{r^3}\end{aligned}$$

The trick is: use the relation

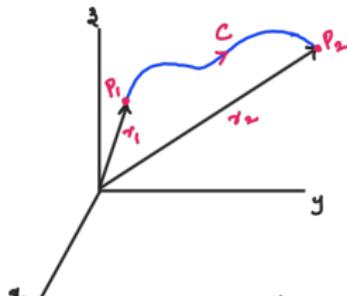
$$r^2 = x^2 + y^2 + z^2$$

Taking differentials in the equation

$$r \, dr = x \, dx + y \, dy + z \, dz$$

$$\Rightarrow \vec{F} \cdot \hat{t} \, ds = \frac{q q_0}{4\pi\epsilon_0} \frac{r \, dr}{r^3} = \frac{q q_0}{4\pi\epsilon_0} \frac{dr}{r^2}$$

If the charge 'q' moves from a point P_1 at a distance r_1 from origin, to a point P_2 at a distance r_2 over some path C connecting P_1, P_2 ; then.



$$\int_C \bar{F} \cdot \hat{t} ds = \frac{q q_0}{4\pi G_0} \int_{r_1}^{r_2} \frac{dt}{r^2} = \frac{q q_0}{4\pi G_0} \left(\frac{1}{r_1} - \frac{1}{r_2} \right)$$

⇒ not depending on the path 'C'.

i.e., $\int_C \vec{F} \cdot d\vec{s}$ with \vec{F} (Coulomb's force) is
path independent

$$\text{But } \vec{F} = 2 \vec{E}$$

$$\int_C \vec{F} \cdot \hat{\vec{t}} \, ds = 2 \int_C \vec{E} \cdot \hat{\vec{t}} \, ds$$

path independent *path independent*

The path independence solely depends on two properties of Coulomb's force;

- (1) It depends only on the distance between two charges
- (2) It acts along the line joining them.

Any \vec{F} with these two properties \Rightarrow Central Force.

$\int_C \vec{F} \cdot \hat{t} ds$ is independent of path for any central force.

* Many other functions which are not central forces also have path independent line integral.

Theorem: $\int_C \vec{F} \cdot \hat{t} ds$ is independent of path, then

$$\int_{C_1} \vec{F} \cdot \hat{t} ds = \int_{C_2} \vec{F} \cdot \hat{t} ds. \quad C_1, C_2 : \text{two different arbitrary paths connecting } P_1 \text{ & } P_2$$



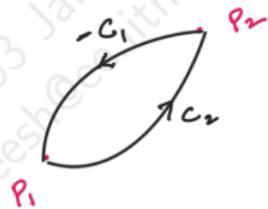
If we integrate from P_2 to P_1 through C_1

$$\Rightarrow \int_{-C_1} \bar{F} \cdot \hat{t} ds = - \int_{C_1} \bar{F} \cdot \hat{t} ds ; \text{ Thus}$$

$$\int_{C_2} \bar{F} \cdot \hat{t} ds = - \int_{-C_1} \bar{F} \cdot \hat{t} ds$$

$$\text{or } \int_{-C_1 + C_2} \bar{F} \cdot \hat{t} ds = 0.$$

$-C_1 + C_2$ is just the closed loop from P_1 to P_2 & back.



Thus if $\int \bar{F} \cdot \hat{t} ds$ is independent of path then,

$$\boxed{\oint \bar{F} \cdot \hat{t} ds = 0.}$$

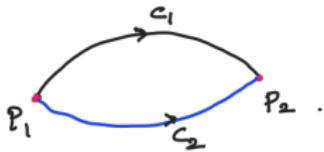
\oint \Rightarrow Line integral around a closed path.

If \vec{E} is an electrostatic field,

$\oint \vec{E} \cdot \hat{T} \, ds = 0 \Rightarrow$ The circulation of the electrostatic field is zero.
 \Rightarrow Circulation law:

Theorem: In a region R' the line integral $\int_C \vec{F} \cdot \hat{T} \, ds$ is independent of path between any two points in R if and only if $\oint_C \vec{F} \cdot \hat{T} \, ds = 0$, for every closed curve C which is contained in R .

Proof: Suppose $\oint_C \vec{F} \cdot \hat{T} \, ds = 0$ for every closed curve C which is contained in R . P_1, P_2 be two distinct points in R and C_1 be a curve from P_1 to P_2 & C_2 be another from P_1 to P_2 in R .



Then $C = C_1 \cup -C_2$ is closed curve in \mathbb{R} (P_1 to P_1)

and so $\oint_C \vec{F} \cdot \hat{t} ds = 0$

$$\begin{aligned} 0 &= \oint_C \vec{F} \cdot \hat{t} ds \\ &= \int_{C_1} \vec{F} \cdot \hat{t} ds + \int_{-C_2} \vec{F} \cdot \hat{t} ds \\ &= \int_{C_1} \vec{F} \cdot \hat{t} ds - \int_{C_2} \vec{F} \cdot \hat{t} ds. \end{aligned}$$

$$\Rightarrow \boxed{\int_{C_1} \vec{F} \cdot \hat{t} ds = \int_{C_2} \vec{F} \cdot \hat{t} ds.}$$

* The theorem does not give a practical way to verify the path independence, as it is almost impossible to check the line integrals around all the closed curves in a region.

Practical method to verify path independence.

Theorem: Let $\vec{F}(x,y) = P(x,y)\hat{i} + Q(x,y)\hat{j}$ be a vector field in some region R with P & Q continuously differentiable functions on R . Let C be the smooth curve in R parametrized by $x = x(s)$; $y = y(s)$; $a \leq s \leq b$. Suppose there exist a real-valued function $F(x,y)$ such that

$\nabla F = \vec{F}$ on R , Then

$$\int_C \vec{F} \cdot \hat{t} ds = F(B) - F(A)$$

where $A = (x(a), y(a))$; $B = (x(b), y(b))$ are

end-points of C . Thus the line integral is independent of the path between its endpoints, since it depends only on the values of F at these points.

Proof:

$$\int_C \vec{f} \cdot \hat{t} ds = \int_a^b \left[P(x(s), y(s)) \frac{dx}{ds} + Q(x(s), y(s)) \frac{dy}{ds} \right] ds$$

Given $\nabla F = \vec{P}$

$$\Rightarrow \vec{f} = \frac{\partial F}{\partial x} i + \frac{\partial F}{\partial y} j \Rightarrow \frac{\partial F}{\partial x} = P; \quad \frac{\partial F}{\partial y} = Q.$$

$$\Rightarrow \int_C \vec{f} \cdot \hat{t} ds = \int_a^b \left(\underbrace{\left(\frac{\partial F}{\partial x} \frac{dx}{ds} + \frac{\partial F}{\partial y} \frac{dy}{ds} \right)}_{F'(x(s), y(s))} \right) ds = \frac{dF}{ds} .$$

Recall chain rule: $z = f(x, y); \quad x = x(s); \quad y = y(s)$

$$\begin{aligned} \frac{dz}{ds} &= \frac{\partial z}{\partial x} \frac{dx}{ds} + \frac{\partial z}{\partial y} \frac{dy}{ds} \\ &= f'(x(s), y(s)) \end{aligned}$$

$$\begin{aligned}
 \Rightarrow \int_C \vec{f} \cdot \hat{t} ds &= \int_a^b F(x(s), y(s)) ds. \\
 &= F(x(s), y(s)) \Big|_{\substack{s=b \\ s=a}} = F(x(b), y(b)) - F(x(a), y(a)) \\
 &= F(B) - F(A) \quad \rightarrow \text{Fundamental theorem of calculus.}
 \end{aligned}$$

* A real valued function $F(x, y)$ such that $\nabla F(x, y) = \vec{f}(x, y)$ is called potential of \vec{f} .

* A conservative field which has potential.

$\vec{f} = \nabla F (\nabla \times \vec{f} = 0)$? } We connect these two
 also $\oint \vec{f} \cdot \hat{t} ds = 0$ ✓ using Green's theorem!



Green's Theorem

If we have a closed curve;



$$\oint_C (\bar{F} \cdot \hat{t}) ds = ?$$

(1) Direct method.

(2) Using Green's theorem.

Green's Theorem: If C is a closed curve, enclosing the region R , counterclockwise, and \vec{F} is a vector field defined & differentiable in region R enclosed, then the line integral (work done)

$$\oint_C \mathbf{f}(\mathbf{F}, \hat{\mathbf{t}}) ds = \iint_R \text{curl } \mathbf{F} dA$$

If $\vec{F}(x,y) = P(x,y)i + Q(x,y)j$ be vector field defined in Region R and C
then:

$$\oint_C (\vec{F} \cdot \hat{t}) ds = \iint_R |\nabla \times \vec{F}| dA = \iint_R \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dA.$$

* Green's theorem : Only for closed curve.

If $\text{curl } \vec{F} = 0 \Rightarrow \oint_C (\vec{F} \cdot \hat{t}) ds = \iint_R \text{curl } \vec{F} dA = 0$.

(Proof for conservative field \vec{F})

\Rightarrow If \vec{F} is conservative : $\text{curl } \vec{F} = 0$.
 $\oint_C (\vec{F} \cdot \hat{t}) ds = 0$.

* If \vec{F} is defined everywhere in the plane & $\nabla \times \vec{F} = 0$
everywhere $\Rightarrow \vec{F}$ is conservative.

* Important : If \vec{F} not defined in some point in region 'R' then
Green's theorem does not hold!

Proof: Consider a simple region R where boundary curve

$$C = C_1 \cup C_2 ; \quad \text{Also,}$$

C_1 = the curve $y = y_1(x)$ from point x_1 to x_2

C_2 = the curve $y = y_2(x)$ from x_2 to x_1

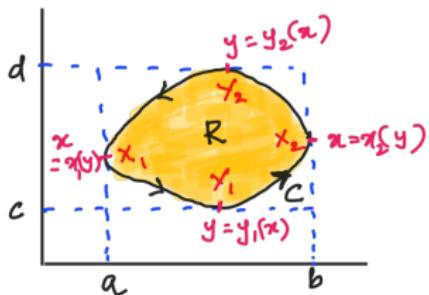
where x_1, x_2 ; points on C farthest on
left & right;

Also,

C_1 = the curve $x = y_1(y)$ from point y_2 to y_1

C_2 = the curve $a = \pi_2(y)$ from y_1 to y_2

where γ_1, γ_2 ; lowest & highest points on 'C'.



Green's theorems:

$$\int_C P(x,y) dx + Q(x,y) dy = \iint_R \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dA$$

Let us see when $\Omega = 0$; Let us prove the following;

$$\Rightarrow \oint_C p(x,y) dx = \iint_R -\frac{\partial p}{\partial y} dA \quad \dots \quad (1)$$

* Integrate $P(x,y)$ around C using $C = C_1 \cup C_2$;

$$\oint_C P(x,y) dx = \int_{C_1} P(x,y) dx + \int_{C_2} P(x,y) dx.$$

$$= \int_a^b P(x, y_1(x)) dx + \int_b^a P(x, y_2(x)) dx. = \int_a^b P(x, y_1(x)) dx - \int_a^b P(x, y_2(x)) dx.$$

$$= - \int_a^b [P(x, y_2(x)) - P(x, y_1(x))] dx. = - \int_a^b \left(P(x, y) \Big|_{y=y_1(x)}^{y=y_2(x)} \right) dx.$$

$$= - \int_{x=a}^{x=b} \int_{y=y_1(x)}^{y=y_2(x)} \frac{\partial P(x,y)}{\partial y} dy dx. = - \iint_R \frac{\partial P}{\partial y} dA. \quad (\text{By Fundamental Theorems of Calculus})$$

Similarly integrate $Q(x,y)$ around 'C'. $C = C_1 \cup C_2$

$$\oint Q(x,y) dy = \int_{C_1} Q(x,y) dy + \int_{C_2} Q(x,y) dy.$$

$$= \int_d^c Q(x_1(y), y) dy + \int_c^d Q(x_2(y), y) dy. = - \int_c^d Q(x_1(y), y) dy + \int_d^c Q(x_2(y), y) dy.$$

$$\begin{aligned}
 &= \int_C^d [Q(x,y), y] - Q(x,y)] dy = \int_C^d Q(x,y) \left\{ \begin{array}{l} x_2(y) \\ x_1(y) \end{array} \right\} dy. \\
 &= \int_C^d \int_{x=x_1(y)}^{x_2(y)} \frac{\partial Q(x,y)}{\partial x} dx dy. = \iint_R \frac{\partial Q}{\partial x} dA.
 \end{aligned}$$

$$\begin{aligned}
 \therefore \oint_C (\vec{F} \cdot \hat{t}) ds &= \oint_C P(x,y) dx + \oint_C Q(x,y) dy \\
 &= - \iint_R \frac{\partial P}{\partial y} dA + \iint_R \frac{\partial Q}{\partial x} dA. \\
 &= \iint_R \left(\underbrace{\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y}}_{\nabla \times \vec{F}} \right) dA. \quad (Q.E.D) \text{ (nicely concluded)}
 \end{aligned}$$