

# Ch-6 lec - 2

## Review Spatial diffusion of $\Phi$

Consider steady state  $\Phi$  balance.  $\Phi$ s are averaged over energy - one group.

### Natural balance (steady state)

Loss rate = production rate  
 $\rightarrow$  leakage + absorption rate = source emittant fission  $\oplus$  gen. rate

$$(1) \quad \nabla \cdot J(r) + \Sigma_a(r) \Phi(r) = S^u(r) + \sigma \Sigma_f(r) \Phi(r)$$

$$\text{Fick's law: } J(r) = -D(r) \nabla \Phi(r)$$

### Diffusion Equation.

$$[\nabla \cdot D(r) \nabla \Phi(r) + \Sigma_a(r) \Phi(r)] = S^u(r) + \sigma \Sigma_f(r) \Phi(r)$$

$\rightarrow D(r)$ ?

Use 1D transport equation with isotropic scattering.

$$\frac{1}{\mu} \frac{\partial \Psi(x, u)}{\partial x} + \Sigma_f(r) \Psi(x, u) = \frac{1}{2} \Sigma_s(x) \Phi(x) + \frac{1}{2} S(x)$$

$$\Phi(x) = \int_{-1}^1 du \Psi(x, u)$$

$$\Phi_1(x) = J(x) = \int_{-1}^1 du u \Psi(x, u)$$

↗ combine

Take 0<sup>th</sup> & 1<sup>st</sup> moment of transport equation re-arrange  
to get a diffusion equation. Finally,

$$\boxed{D = \frac{1}{3\sigma_t}}$$

In general

$$\Psi(x, u) = \sum_{l=0}^{\infty} \frac{d^l u}{l!} P_l(u) \Phi_l(x)$$

$$\Phi_l(x) = \int_{-\infty}^x du P_l(u) \Psi(x, u)$$

For diffusion, because we get diffusion equation just by taking 0<sup>th</sup> & 1<sup>st</sup> moment of transport equation,  $\Psi$  only depends on  $\Phi$  &  $J$ .

$$\Rightarrow \Psi(x, u) \approx \sum_{l=0}^{l=1} \frac{d^l u}{l!} P_l(u) \Phi_l(x)$$

$$\boxed{\Psi(x, u) = \frac{1}{2} \Phi(x) + \frac{3}{2} u J(x)}$$

(2)

## 1D Diffusion Equation:

$$-\frac{\partial}{\partial x} D(x) \frac{\partial \phi(x)}{\partial x} + \epsilon_a(x) \phi(x) = s(x) + \nu \epsilon_f(x) \phi(x).$$

## Boundary conditions:

### (1) Reflecting / symmetric boundary

Net current at boundary = 0.

⇒ no  $\phi$  cross boundary.

Suppose we consider 1D system

$$J_x(\text{boundary}) = -D \frac{\partial \phi}{\partial x} \Big|_{x=\text{boundary}} = 0$$

$$\Rightarrow \frac{\partial \phi}{\partial x} \Big|_{x=\text{boundary}} = 0.$$

### (2) Vacuum boundary.

Nothing comes in thru boundaries.

$$\rightarrow | \quad \# @ \text{ coming in} = 0$$

$$| \leftarrow \quad \# @ \text{ coming in} = 0$$

In order to mathematically represent we introduce partial currents.

*represented by*

$$\text{Current} = J(x) \rightarrow J(x) = \int du u \psi(x, u)$$

Decompose into <sup>current partial currents</sup> <sub>into</sub> the & -ve directions:

$$J^+(x) \quad J^-(x)$$

$$J(x) = \int_{-1}^1 du u \psi(x, u) = \int_{-1}^0 du u \psi(x, u) + \int_0^1 du u \psi(x, u)$$

$$J(x) = J^+(x) + J^-(x)$$

$$\Rightarrow J(x) = \int_0^1 du u \psi(x, u) - \int_{-1}^0 du |u| \psi(x, u)$$

-ve sign because -ve angles from  $-1 \rightarrow 0$ .

$$J^+(x) = \int_0^1 u \psi(x, u) du$$

$$= \int_0^1 u \left[ \frac{1}{2} \phi(x) + \frac{3}{2} u J(x) \right] du$$

$$= \frac{\phi(x)}{2} \int_0^1 u du + \frac{3}{2} J(x) \int_0^1 u^2 du$$

$$J^+(x) = \frac{\phi(x)}{4} + \frac{1}{2} J(x)$$

(3)

$$\text{Now } J(x) = -D \frac{\partial}{\partial x} \Phi(x).$$

$\Rightarrow$   ~~$J^+$~~

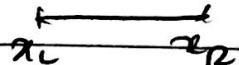
$$J^+ = \frac{1}{4} \Phi(x) - \frac{1}{2} D \frac{\partial}{\partial x} \Phi(x)$$

$$\text{Similarly, } J^-(x) = \int_{-\infty}^0 du |u| \Phi(x, u)$$

$$J^- = \frac{1}{4} \Phi(x) - \frac{1}{2} J(x)$$

$$J^- = \frac{1}{4} \Phi(x) + D \frac{\partial}{\partial x} \Phi(x)$$

for vacuum boundaries then,



$$J^+(x_1) = 0 \quad \& \quad J^-(x_2) = 0$$

$$J^+(x_1) = 0 = \frac{\Phi(x)}{4} - \frac{1}{2} D \frac{d}{dx} \Phi$$

$$\Rightarrow D \frac{\partial \Phi}{\partial x} \Big|_{x=x_1} = \frac{\Phi(x)}{2}$$

At right boundary,

$$J^-(x_2) = 0 = \frac{1}{4} \Phi(x_2) - \frac{1}{2} D \frac{d}{dx} \Phi(x) \Big|_{x=x_2}$$

right  
or derivative  
can write  
without  
value  
sign

So now, how do we solve diffusion equation?

Different methods:

④ Finite differences

Finite volume

finite element

mesh free methods - Peridynamic differential operator

- In general finite volume & discontinuous FEM  
Used in mechanics.

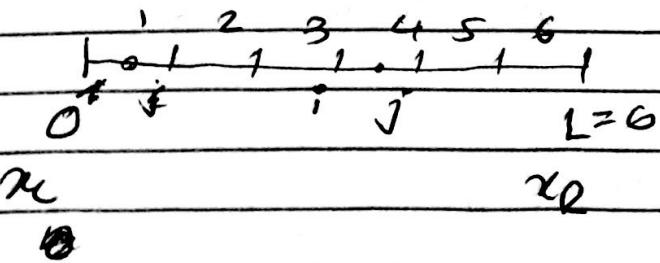
Finite element method:

In general never write your own code unless that's your job. Several free packages exist

- Ansys
- Comsol
- Moose
- OpenFOAM
- FENICS
- XFEM
- ADINA

General procedure for FEM:

- ① Divide domain into subdomains called "elements"



(4)

- ① Develop Write equation in weighted residual form  
to develop weak form of diff. eqn. over each element.
- ② Assume ~~test~~<sup>weight</sup> trial functions.
- ③ Find matrix equation for each element.
- ④ Assemble global matrix eqn by combining eqn. for each element.
- ⑤ Solve matrix eqn. using direct / indirect methods.

1

Write equation in weighted residual form.

DISCRETISATION

Each element is represented by discrete set of points called nodes with unique coordinates on coordinate system.

In the domain in fig. on prev page,

element #	node i	coord	node j	coord.
1	1	(0,0)	2	(1,0)
2	2	(1,0)	3	(2,0)
3	3	(2,0)	4	(3,0)
4	4	(3,0)	5	(4,0)
5	5	(4,0)	6	(5,0)
6	6	(5,0)	7	(2,0)

Consider generic element starting at i & endy at j

Depend on the type of element, each

## (e) WEIGHTED RESIDUAL & WEAK FORM OVER ELEMENT

diffusion eqn:  $\frac{-\partial}{\partial x} \frac{\partial \phi}{\partial x} + \epsilon_a \phi = s$

residual form:  $\frac{-\partial}{\partial x} \frac{\partial \phi}{\partial x} + \epsilon_a \phi - s = 0$

must be minimized.

Weighted residual:  $\int_{x_i}^{x_{i+1}} w (-\frac{\partial}{\partial x} \frac{\partial \phi}{\partial x} + \epsilon_a \phi - s) dx = 0$

Weak form: To get weak form, we can integrate weighted residual eqn by parts.

$x_i^o = \text{dim}$  (for ~~forwards~~ backwards)

$$= \int_{x_i^o}^{x_{i+1}} (-\frac{\partial}{\partial x} \frac{\partial \phi}{\partial x} + \epsilon_a \phi - s) w dx$$

$$= \int_{x_i^o}^{x_{i+1}} \left( -\frac{\partial}{\partial x} \frac{\partial \phi}{\partial x} \right) w dx + \int \epsilon_a \phi w dx +$$

$$\int_{x_i^o}^{x_{i+1}} -s w dx.$$

$$\int_{x_i^o}^{x_{i+1}} -s w dx.$$

$$\Rightarrow I = I_1 + I_2 + I_3$$

3

$$I_1 = \int_{x_1}^{x_{1H}} \left( -\frac{d}{dx} \Omega \frac{dd\phi}{dx} \right) w \, dx.$$

Integrate by parts:

$$J_1 = \int_{x_1}^{x_{in}} \left( -\frac{d}{dx} D \frac{d\phi}{dx} \right) w dx$$

$$= - \left[ \int_{z_1}^{z_2} \frac{-d\omega}{dx} D \frac{d\phi}{dx} dx + \left[ \omega D \frac{d\phi}{dx} \right]_{z_1}^{z_2} \right]$$

$$I_i = D_i \int_{x_i}^{x_{i+1}} \frac{dw}{dx} \frac{dq}{dx} dx - \left[ w D_i \frac{dq}{dx} \right]_{x_i}^{x_{i+1}}$$

assume  $T$  constant over each element;

$I_{i,1}$        $\uparrow$   
 $I_{i,2}$

$$I = I_1 + I_2 + I_3$$

## Weighted residual

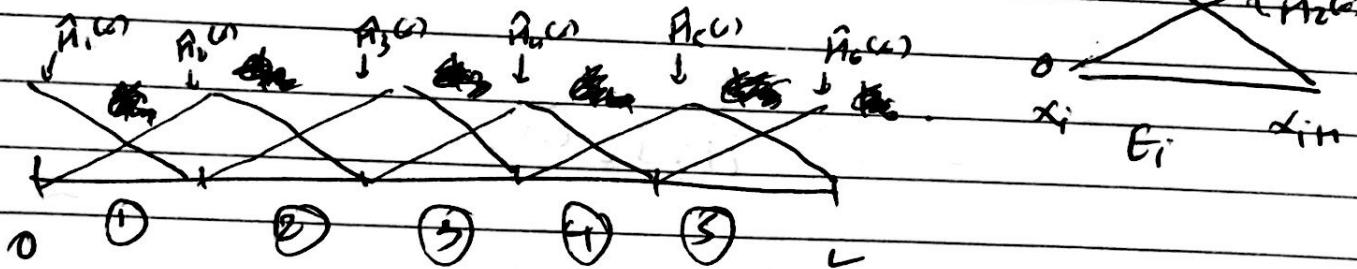
$\Rightarrow$  Equation is

$$I_1 + I_2 + I_3 = 0$$

$$\Rightarrow \left[ D_i \int_{x_i}^{x_{i+1}} \frac{\partial w}{\partial x} \frac{\partial \psi}{\partial x} dx - \left[ \omega D_i \frac{\partial \psi}{\partial x} \right]_{x_i}^{x_{i+1}} \right] - \text{weak form as } \frac{d^L \psi}{dx^2} \text{ reduced to } d\psi/dx$$

### ③ WEIGHT & TRIAL F

Here, we choose linear elements (hat functions)



over element,

$$\Phi = H_1(x)\Phi_1 + H_2(x)\Phi_{1n} \leftarrow \text{Trial function}$$

Galerkin Method: Weight function has the same form as trial function.

Petrov-Galerkin when weight fn & trial fn do not have the same form.

$$\text{We will do Galerkin} \Rightarrow w = H_1(x)\omega_1 + H_2(x)\omega_2$$

④ Find matrix equation for each element.

first thing we do is write  $\phi$  &  $\omega$  as matrices.

$$\begin{aligned}\phi &= H_1(x)\phi_i + H_2(x)\phi_{iH} \\ &= \begin{bmatrix} H_1(x) & H_2(x) \end{bmatrix} \begin{bmatrix} \phi_i \\ \phi_{iH} \end{bmatrix} \\ &= [H][\phi]\end{aligned}$$

$$\omega = H_1(x)\omega_i + H_2(x)\omega_{iH}$$

~~$$\begin{aligned}\omega &= \begin{bmatrix} H_1(x) & H_2(x) \end{bmatrix} \begin{bmatrix} \omega_i \\ \omega_{iH} \end{bmatrix} \\ &= [H][\omega]\end{aligned}$$~~

Now  $\omega$  is a scalar.

$$\begin{aligned}\Rightarrow \omega &= \omega^T \\ \Rightarrow \omega &= \omega^T = \cancel{\begin{bmatrix} H_1(x) & H_2(x) \end{bmatrix}}^{\cancel{\text{scalar}}} \cdot \cancel{\begin{bmatrix} \omega_i \\ \omega_{iH} \end{bmatrix}}^{\cancel{\text{vector}}} \\ &= \cancel{\begin{bmatrix} H_1(x) & H_2(x) \end{bmatrix}}^{\cancel{\text{matrix}}} [\omega_i \ \omega_{iH}] \begin{bmatrix} H_1(x) \\ H_2(x) \end{bmatrix}\end{aligned}$$

$$\text{Now, let's look at } I = I_1 + I_2 + I_3$$

Now, the equation becomes

$$I_1 + I_2 + I_3 = 0$$

(7)

$$\Rightarrow D_i \int_{x_i}^{x_{iH}} \frac{d\omega}{dx} \frac{\partial \phi}{\partial x} dx - \left[ \omega D_i \frac{d\phi}{dx} \right]_{x_i}^{x_{iH}} + \sum_{ai} \int_{x_i}^{x_{iH}} \omega \phi dx$$

$$= \int_{x_i}^{x_{iH}} S_i \omega dx.$$

~~$\int_{x_i}^{x_{iH}} D_i \int \frac{d[\tilde{\omega}]^T [H]^T}{dx} \frac{d[H][\phi]}{dx} dx$~~

~~$\int_{x_i}^{x_{iH}} \left[ D_i [\tilde{\omega}]^T [H]^T \frac{\partial \phi}{\partial x} \right]_{x_i}^{x_{iH}}$~~

~~$\int_{x_i}^{x_{iH}} [\tilde{\omega}]^T [H]^T [H][\phi] dx$~~

~~$= S_i \int_{x_i}^{x_{iH}} [\tilde{\omega}]^T [H]^T dx.$~~

Boundary condition  
 will have  
 off & keep it  
 so  
 like that  
 for now.

Now note that  $[\tilde{\omega}]^T$  is arbitrary and does not depend on  $x$ . Moreover all terms have it. we cancel  $[\tilde{\omega}]^T$  from all terms. We get:

$$D_i \int_{x_i}^{x_{iH}} \frac{d[H]^T}{dx} \frac{d[H]}{dx} dx [\phi] - \left[ D_i [H]^T \frac{\partial \phi}{\partial x} \right]_{x_i}^{x_{iH}}$$

$$+ \sum_{ai} \int_{x_i}^{x_{iH}} [H]^T [H] d\phi [\phi] = S_i \int_{x_i}^{x_{iH}} [H]^T dx.$$

Now let's substitute  $[H]$ ,  $[H]'$  &  $\hat{[\Phi]}$  in  
the weak form equation from last page.

$$I_{1,1} = D_i \int_{x_i}^{x_{iH}} \frac{d[H]}{dx} \frac{d[H]}{dx} dx (\hat{\Phi})$$

$$= D_i \int_{x_i}^{x_{iH}} \begin{pmatrix} \frac{d}{dx} H_1(x) \\ \frac{d}{dx} H_2(x) \end{pmatrix} \begin{bmatrix} \frac{d}{dx} H_1(x) & \frac{d}{dx} H_2(x) \end{bmatrix} dx \begin{bmatrix} \Phi_i \\ \Phi_{iH} \end{bmatrix}$$

$$\frac{d H_1(x)}{dx} = \frac{d}{dx} \frac{(x_{iH} - x)}{x_{iH} - x_i} = -\frac{1}{h_i}$$

$$\frac{d H_2(x)}{dx} = \frac{d}{dx} \frac{(x - x_i)}{x_{iH} - x_i} = \frac{1}{h_i}$$

$$\Rightarrow I_{1,1} = D_i \int_{x_i}^{x_{iH}} \begin{bmatrix} -1/h_i \\ 1/h_i \end{bmatrix} \begin{bmatrix} 1/h_i & 1/h_i \end{bmatrix} dx \begin{bmatrix} \Phi_i \\ \Phi_{iH} \end{bmatrix}$$

$$= \frac{D_i}{h_i^2} \int_{x_i}^{x_{iH}} \begin{bmatrix} -1 & 1 \\ 1 & 1 \end{bmatrix} \begin{bmatrix} 1 & 1 \\ -1 & 1 \end{bmatrix} dx \begin{bmatrix} \Phi_i \\ \Phi_{iH} \end{bmatrix}$$

$$= \frac{D_i}{h_i^2} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \int_{x_i}^{x_{iH}} \begin{bmatrix} 1 & 1 \\ -1 & 1 \end{bmatrix} d x \begin{bmatrix} \Phi_i \\ \Phi_{iH} \end{bmatrix}$$

(2)

$$= \frac{D_i}{h_i} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} \phi_i \\ \phi_{in} \end{bmatrix} //$$

$$I_{1,2} = \left[ \otimes D_i [H]^T \frac{\partial \phi}{\partial x} \right]_{x_i}^{x_{in}} \leftarrow \text{more on this later.}$$

$$I_2 = \sum_i \int_{x_i}^{x_{in}} [H]^T [H] dx [\tilde{\phi}]$$

*(using simplicial evaluation)*

$$= \sum_i \int_{x_i}^{x_{in}} \begin{bmatrix} H_1(x) \\ H_2(x) \end{bmatrix} [H_1(x) \quad H_2(x)] dx \begin{bmatrix} \phi_i \\ \phi_{in} \end{bmatrix}$$

$$= \sum_i \begin{bmatrix} h_{i/3} & h_{i/6} \\ h_{i/6} & h_{i/3} \end{bmatrix} \begin{bmatrix} \phi_i \\ \phi_{in} \end{bmatrix} //$$

$$I_3 = s_i \int_{x_i}^{x_{in}} [B_i]^T dx = s_i \int_{x_i}^{x_{in}} \begin{bmatrix} H_1(x) \\ H_2(x) \end{bmatrix} dx$$

*(using evaluation)*

$$= s_i \int_{x_i}^{x_{in}} \begin{bmatrix} \frac{x_{i+1}-x}{h_i} \\ \frac{x-x_i}{h_i} \end{bmatrix} dx$$

$$\approx s_i \begin{bmatrix} h_{i/2} \\ h_{i/2} \end{bmatrix} //$$

Now lets come to  $I_{1,2}$

$$I_{1,2} = \left[ D_i [H_1]^\top \frac{\partial \phi}{\partial x} \right]_{x_i}^{x_{i+1}}$$

When we are at interfaces / interior of slab, interface condition states that both  $\phi$  &  $\frac{\partial \phi}{\partial x}$  must be continuous.

That means at  $x_{i+1}$  &  $x_i$  (except at slab boundaries, where  $\phi$  conditions will be prescribed),

$$\begin{aligned} \cancel{\phi(x_-) = \phi(x_+)} \\ \Rightarrow \cancel{\frac{\partial \phi}{\partial x}} \end{aligned} \quad \phi(x_-) = \phi(x_+)$$

This will be taken care of naturally when we do global matrix assembly later.

$$I_{1,2} = \left[ \left[ D_i H_1(x) \frac{\partial \phi}{\partial x} \right]_{x_{i+1}} - \left[ D_i H_1(x) \frac{\partial \phi}{\partial x} \right]_{x_i} \right. \\ \left. - \left[ D_i H_2(x) \frac{\partial \phi}{\partial x} \right]_{x_{i+1}} - \left[ H_2(x) D_i \frac{\partial \phi}{\partial x} \right]_{x_i} \right]$$

Now, we know that  $H_1(x_{i+1}) = 0$  &  $H_2(x_i) = 0$   
moreover,  $H_1(x_i) = 1$  &  $H_2(x_{i+1}) = 1$

2)

$$I_{1,2} = \left[ -D_i \frac{d\phi}{dx} \Big|_{x_i} + D_i \frac{d\phi}{dx} \Big|_{x_{i+1}} \right]$$

(9)

Specifically for boundary elements.

Then we put  $I_1, I_2$  &  $I_3$  together to obtain a matrix equation for a generic element:

we have

$$I_1 + I_2 = I_3.$$

$$\therefore \begin{bmatrix} D_i/h_i + \epsilon_{ai} h_{i3} & -D_i/h_i + \epsilon_{ai} h_{i6} \\ -D_i/h_i + \epsilon_{ai} h_{i6} & D_i/h_i + \epsilon_{ai} h_{i3} \end{bmatrix} \begin{bmatrix} d_i \\ d_{iH} \end{bmatrix}$$

$$- \begin{bmatrix} -D_i \frac{d\phi}{dx} \Big|_{x_i} \\ D_i \frac{d\phi}{dx} \Big|_{x_{iH}} \end{bmatrix} = s_i \begin{bmatrix} h_{i2} \\ h_{i2} \end{bmatrix}$$

~~1~~

Now, let us look at ~~Boundary~~ specific boundary conditions.

Reflecting boundary:

$\frac{\partial \phi}{\partial x} = 0 \Rightarrow$  for left reflecting boundary.  
we have

$$I_{1,2} = \begin{bmatrix} 0 \\ D_i \frac{d\phi}{dx} \Big|_{x_{iH}} \end{bmatrix}.$$

Similarly for Right reflecting boundary,

$$J_{1,2} = \begin{bmatrix} -D_i \frac{d\phi}{dx} \Big|_{x_i} \\ 0 \end{bmatrix}$$

for vacuum boundary condition.

$$J_x^+ (x) = \frac{1}{4} \phi(x) + \frac{1}{2} D \frac{d\phi}{dx}$$

for left boundary

$$J^+ = 0 \Rightarrow \frac{1}{4} \phi(x) = \frac{1}{2} \frac{d\phi}{dx}$$

$$\Rightarrow +D \frac{d\phi}{dx} = \frac{1}{2} \phi(x)$$

$$\Rightarrow J_{1,2}^{\text{left, vac}} = \begin{bmatrix} -\frac{\phi(x)}{2} \Big|_{x_1} \\ D \frac{d\phi}{dx} \Big|_{x_{i+1}} \end{bmatrix}$$

for right boundary,

$$-D \frac{d\phi}{dx} = \frac{\phi(x)}{2} \Rightarrow J_{1,2} = \begin{bmatrix} -D \frac{d\phi}{dx} \Big|_{x_1} \\ -\frac{\phi(x)}{2} \Big|_{x_{i+1}} \end{bmatrix}$$