

length. Note that the flow rate in one of the parallel branches is proportional to its diameter to the power 5/2 and is inversely proportional to the square root of its length and friction factor.

The analysis of piping networks, no matter how complex they are, is based on two simple principles:

1. *Conservation of mass throughout the system must be satisfied.* This is done by requiring the total flow into a junction to be equal to the total flow out of the junction for all junctions in the system. Also, the flow rate must remain constant in pipes connected in series regardless of the changes in diameters.
2. *Pressure drop (and thus head loss) between two junctions must be the same for all paths between the two junctions.* This is because pressure is a point function and it cannot have two values at a specified point. In practice this rule is used by requiring that the algebraic sum of head losses in a loop (for all loops) be equal to zero. (A head loss is taken to be positive for flow in the clockwise direction and negative for flow in the counterclockwise direction.)

Therefore, the analysis of piping networks is very similar to the analysis of electric circuits, with flow rate corresponding to electric current and pressure corresponding to electric potential. However, the situation is much more complex here since, unlike the electrical resistance, the “flow resistance” is a highly nonlinear function. Therefore, the analysis of piping networks requires the simultaneous solution of a system of nonlinear equations. The analysis of such systems is beyond the scope of this introductory text.

## Piping Systems with Pumps and Turbines

When a piping system involves a pump and/or turbine, the steady-flow energy equation on a unit-mass basis can be expressed as (see Section 5–7)

$$\frac{P_1}{\rho} + \alpha_1 \frac{V_1^2}{2} + gz_1 + w_{\text{pump}, u} = \frac{P_2}{\rho} + \alpha_2 \frac{V_2^2}{2} + gz_2 + w_{\text{turbine}, e} + gh_L \quad (8-61)$$

It can also be expressed in terms of heads as

$$\frac{P_1}{\rho g} + \alpha_1 \frac{V_1^2}{2g} + z_1 + h_{\text{pump}, u} = \frac{P_2}{\rho g} + \alpha_2 \frac{V_2^2}{2g} + z_2 + h_{\text{turbine}, e} + h_L \quad (8-62)$$

where  $h_{\text{pump}, u} = w_{\text{pump}, u}/g$  is the useful pump head delivered to the fluid,  $h_{\text{turbine}, e} = w_{\text{turbine}, e}/g$  is the turbine head extracted from the fluid,  $\alpha$  is the kinetic energy correction factor whose value is nearly 1 for most (turbulent) flows encountered in practice, and  $h_L$  is the total head loss in piping (including the minor losses if they are significant) between points 1 and 2. The pump head is zero if the piping system does not involve a pump or a fan, the turbine head is zero if the system does not involve a turbine, and both are zero if the system does not involve any mechanical work-producing or work-consuming devices.

Many practical piping systems involve a pump to move a fluid from one reservoir to another. Taking points 1 and 2 to be at the *free surfaces* of the reservoirs, the energy equation in this case reduces for the useful pump head required to (Fig. 8–44)

$$h_{\text{pump}, u} = (z_2 - z_1) + h_L \quad (8-63)$$

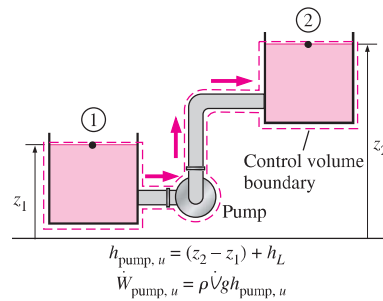


FIGURE 8-44

When a pump moves a fluid from one reservoir to another, the useful pump head requirement is equal to the elevation difference between the two reservoirs plus the head loss.

since the velocities at free surfaces are negligible and the pressures are at atmospheric pressure. Therefore, the useful pump head is equal to the elevation difference between the two reservoirs plus the head loss. If the head loss is negligible compared to  $z_2 - z_1$ , the useful pump head is simply equal to the elevation difference between the two reservoirs. In the case of  $z_1 > z_2$  (the first reservoir being at a higher elevation than the second one) with no pump, the flow is driven by gravity at a flow rate that causes a head loss equal to the elevation difference. A similar argument can be given for the turbine head for a hydroelectric power plant by replacing  $h_{pump,u}$  in Eq. 8-63 by  $-h_{turbine,e}$ .

Once the useful pump head is known, the *mechanical power that needs to be delivered by the pump to the fluid* and the *electric power consumed by the motor of the pump* for a specified flow rate are determined from

$$\dot{W}_{pump, shaft} = \frac{\rho \dot{V} g h_{pump,u}}{\eta_{pump}} \quad \text{and} \quad \dot{W}_{elect} = \frac{\rho \dot{V} g h_{pump,u}}{\eta_{pump-motor}} \quad (8-64)$$

where  $\eta_{pump-motor}$  is the *efficiency of the pump-motor combination*, which is the product of the pump and the motor efficiencies (Fig. 8-45). The pump-motor efficiency is defined as the ratio of the net mechanical energy delivered to the fluid by the pump to the electric energy consumed by the motor of the pump, and it usually ranges between 50 and 85 percent.

The head loss of a piping system increases (usually quadratically) with the flow rate. A plot of required useful pump head  $h_{pump,u}$  as a function of flow rate is called the **system (or demand) curve**. The head produced by a pump is not a constant either. Both the pump head and the pump efficiency vary with the flow rate, and pump manufacturers supply this variation in tabular or graphical form, as shown in Fig. 8-46. These experimentally determined  $h_{pump,u}$  and  $\eta_{pump,u}$  versus  $\dot{V}$  curves are called **characteristic (or supply or performance) curves**. Note that the flow rate of a pump increases as the required head decreases. The intersection point of the pump head curve with the vertical axis typically represents the *maximum head* the pump can provide, while the intersection point with the horizontal axis indicates the *maximum flow rate* (called the **free delivery**) that the pump can supply.

The *efficiency* of a pump is sufficiently high for a certain range of head and flow rate combination. Therefore, a pump that can supply the required head and flow rate is not necessarily a good choice for a piping system unless the efficiency of the pump at those conditions is sufficiently high. The pump installed in a piping system will operate at the point where the *system curve* and the *characteristic curve* intersect. This point of intersection is called the **operating point**, as shown in Fig. 8-46. The useful head

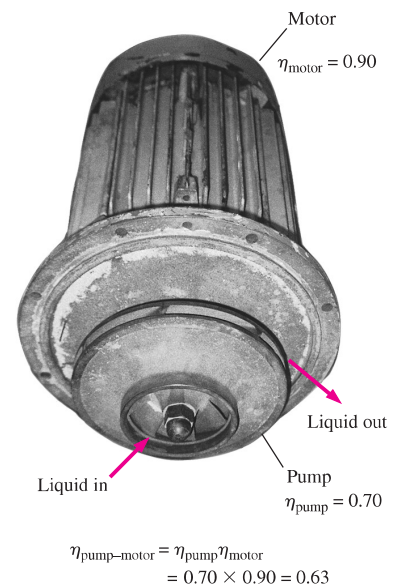
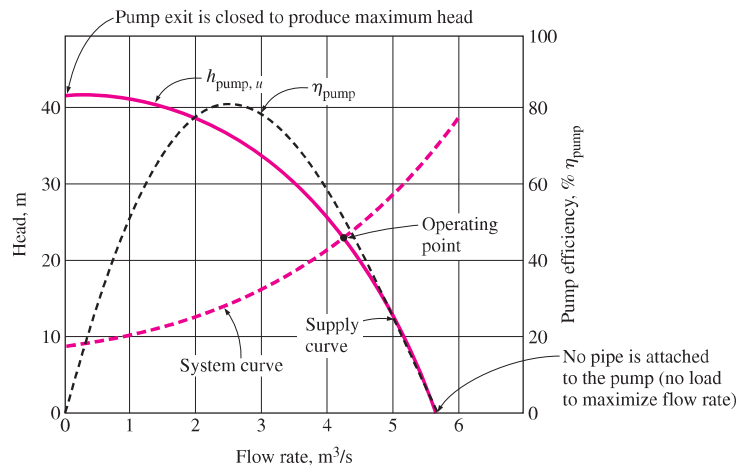


FIGURE 8-45

The efficiency of the pump-motor combination is the product of the pump and the motor efficiencies.

Courtesy Yunus Çengel

**FIGURE 8-46**

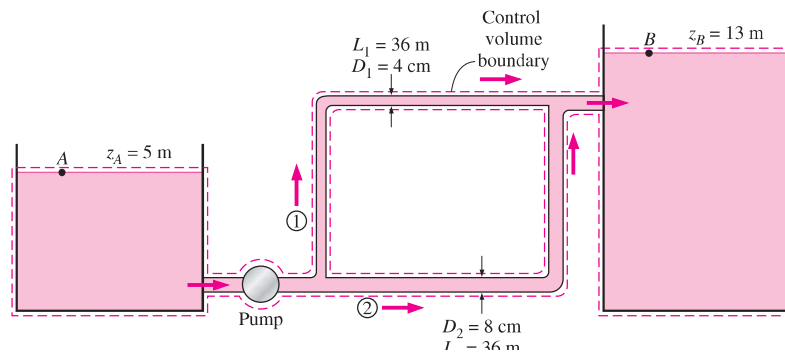
Characteristic pump curves for centrifugal pumps, the system curve for a piping system, and the operating point.

produced by the pump at this point matches the head requirements of the system at that flow rate. Also, the efficiency of the pump during operation is the value corresponding to that flow rate.

#### EXAMPLE 8-7 Pumping Water through Two Parallel Pipes

Water at 20°C is to be pumped from a reservoir ( $z_A = 5$  m) to another reservoir at a higher elevation ( $z_B = 13$  m) through two 36-m-long pipes connected in parallel, as shown in Fig. 8-47. The pipes are made of commercial steel, and the diameters of the two pipes are 4 and 8 cm. Water is to be pumped by a 70 percent efficient motor–pump combination that draws 8 kW of electric power during operation. The minor losses and the head loss in pipes that connect the parallel pipes to the two reservoirs are considered to be negligible. Determine the total flow rate between the reservoirs and the flow rate through each of the parallel pipes.

**SOLUTION** The pumping power input to a piping system with two parallel pipes is given. The flow rates are to be determined.

**FIGURE 8-47**

The piping system discussed in Example 8-7.

**Assumptions** 1 The flow is steady and incompressible. 2 The entrance effects are negligible, and thus the flow is fully developed. 3 The elevations of the reservoirs remain constant. 4 The minor losses and the head loss in pipes other than the parallel pipes are said to be negligible. 5 Flows through both pipes are turbulent (to be verified).

**Properties** The density and dynamic viscosity of water at 20°C are  $\rho = 998 \text{ kg/m}^3$  and  $\mu = 1.002 \times 10^{-3} \text{ kg/m} \cdot \text{s}$ . The roughness of commercial steel pipe is  $\varepsilon = 0.000045 \text{ m}$ .

**Analysis** This problem cannot be solved directly since the velocities (or flow rates) in the pipes are not known. Therefore, we would normally use a trial-and-error approach here. However, nowadays equation solvers such as EES are widely available, and thus we will simply set up the equations to be solved by an equation solver. The useful head supplied by the pump to the fluid is determined from

$$\dot{W}_{\text{elect}} = \frac{\rho \dot{V} g h_{\text{pump}, u}}{\eta_{\text{pump-motor}}} \rightarrow 8000 \text{ W} = \frac{(998 \text{ kg/m}^3) \dot{V} (9.81 \text{ m/s}^2) h_{\text{pump}, u}}{0.70} \quad (1)$$

We choose points *A* and *B* at the free surfaces of the two reservoirs. Noting that the fluid at both points is open to the atmosphere (and thus  $P_A = P_B = P_{\text{atm}}$ ) and that the fluid velocities at both points are zero ( $V_A = V_B = 0$ ), the energy equation for a control volume between these two points simplifies to

$$\frac{P_A}{\rho g} + \alpha_A \frac{V_A^2}{2g} + z_A + h_{\text{pump}, u} = \frac{P_B}{\rho g} + \alpha_B \frac{V_B^2}{2g} + z_B + h_L \rightarrow h_{\text{pump}, u} = (z_B - z_A) + h_L$$

or

$$h_{\text{pump}, u} = (13 - 5) + h_L \quad (2)$$

where

$$h_L = h_{L,1} = h_{L,2} \quad (3)(4)$$

We designate the 4-cm-diameter pipe by 1 and the 8-cm-diameter pipe by 2. The average velocity, the Reynolds number, the friction factor, and the head loss in each pipe are expressed as

$$V_1 = \frac{\dot{V}_1}{A_{c,1}} = \frac{\dot{V}_1}{\pi D_1^2/4} \rightarrow V_1 = \frac{\dot{V}_1}{\pi (0.04 \text{ m})^2/4} \quad (5)$$

$$V_2 = \frac{\dot{V}_2}{A_{c,2}} = \frac{\dot{V}_2}{\pi D_2^2/4} \rightarrow V_2 = \frac{\dot{V}_2}{\pi (0.08 \text{ m})^2/4} \quad (6)$$

$$\text{Re}_1 = \frac{\rho V_1 D_1}{\mu} \rightarrow \text{Re}_1 = \frac{(998 \text{ kg/m}^3) V_1 (0.04 \text{ m})}{1.002 \times 10^{-3} \text{ kg/m} \cdot \text{s}} \quad (7)$$

$$\text{Re}_2 = \frac{\rho V_2 D_2}{\mu} \rightarrow \text{Re}_2 = \frac{(998 \text{ kg/m}^3) V_2 (0.08 \text{ m})}{1.002 \times 10^{-3} \text{ kg/m} \cdot \text{s}} \quad (8)$$

$$\begin{aligned} \frac{1}{\sqrt{f_1}} &= -2.0 \log \left( \frac{\varepsilon/D_1}{3.7} + \frac{2.51}{\text{Re}_1 \sqrt{f_1}} \right) \\ \rightarrow \frac{1}{\sqrt{f_1}} &= -2.0 \log \left( \frac{0.000045}{3.7 \times 0.04} + \frac{2.51}{\text{Re}_1 \sqrt{f_1}} \right) \end{aligned} \quad (9)$$

$$\frac{1}{\sqrt{f_2}} = -2.0 \log \left( \frac{\varepsilon/D_2}{3.7} + \frac{2.51}{\text{Re}_2 \sqrt{f_2}} \right)$$

$$\rightarrow \frac{1}{\sqrt{f_2}} = -2.0 \log \left( \frac{0.000045}{3.7 \times 0.08} + \frac{2.51}{\text{Re}_2 \sqrt{f_2}} \right) \quad (10)$$

$$h_{L,1} = f_1 \frac{L_1}{D_1} \frac{V_1^2}{2g} \rightarrow h_{L,1} = f_1 \frac{36 \text{ m}}{0.04 \text{ m}} \frac{V_1^2}{2(9.81 \text{ m/s}^2)} \quad (11)$$

$$h_{L,2} = f_2 \frac{L_2}{D_2} \frac{V_2^2}{2g} \rightarrow h_{L,2} = f_2 \frac{36 \text{ m}}{0.08 \text{ m}} \frac{V_2^2}{2(9.81 \text{ m/s}^2)} \quad (12)$$

$$\dot{V} = \dot{V}_1 + \dot{V}_2 \quad (13)$$

This is a system of 13 equations in 13 unknowns, and their simultaneous solution by an equation solver gives

$$\dot{V} = 0.0300 \text{ m}^3/\text{s}, \quad \dot{V}_1 = 0.00415 \text{ m}^3/\text{s}, \quad \dot{V}_2 = 0.0259 \text{ m}^3/\text{s}$$

$$V_1 = 3.30 \text{ m/s}, \quad V_2 = 5.15 \text{ m/s}, \quad h_L = h_{L,1} = h_{L,2} = 11.1 \text{ m}, \quad h_{\text{pump}} = 19.1 \text{ m}$$

$$\text{Re}_1 = 131,600, \quad \text{Re}_2 = 410,000, \quad f_1 = 0.0221, \quad f_2 = 0.0182$$

Note that  $\text{Re} > 4000$  for both pipes, and thus the assumption of turbulent flow is verified.

**Discussion** The two parallel pipes are identical, except the diameter of the first pipe is half the diameter of the second one. But only 14 percent of the water flows through the first pipe. This shows the strong dependence of the flow rate (and the head loss) on diameter. Also, it can be shown that if the free surfaces of the two reservoirs were at the same elevation (and thus  $z_A = z_B$ ), the flow rate would increase by 20 percent from 0.0300 to 0.0361  $\text{m}^3/\text{s}$ . Alternately, if the reservoirs were as given but the irreversible head losses were negligible, the flow rate would become 0.0715  $\text{m}^3/\text{s}$  (an increase of 138 percent).

### EXAMPLE 8–8 Gravity-Driven Water Flow in a Pipe

Water at 10°C flows from a large reservoir to a smaller one through a 5-cm-diameter cast iron piping system, as shown in Fig. 8–48. Determine the elevation  $z_1$  for a flow rate of 6 L/s.

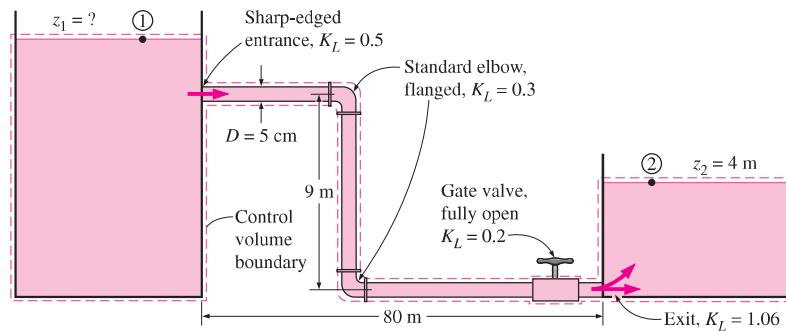


FIGURE 8–48

The piping system discussed in Example 8–8.

**SOLUTION** The flow rate through a piping system connecting two reservoirs is given. The elevation of the source is to be determined.

**Assumptions** 1 The flow is steady and incompressible. 2 The elevations of the reservoirs remain constant. 3 There are no pumps or turbines in the line.

**Properties** The density and dynamic viscosity of water at 10°C are  $\rho = 999.7 \text{ kg/m}^3$  and  $\mu = 1.307 \times 10^{-3} \text{ kg/m} \cdot \text{s}$ . The roughness of cast iron pipe is  $\varepsilon = 0.00026 \text{ m}$ .

**Analysis** The piping system involves 89 m of piping, a sharp-edged entrance ( $K_L = 0.5$ ), two standard flanged elbows ( $K_L = 0.3$  each), a fully open gate valve ( $K_L = 0.2$ ), and a submerged exit ( $K_L = 1.06$ ). We choose points 1 and 2 at the free surfaces of the two reservoirs. Noting that the fluid at both points is open to the atmosphere (and thus  $P_1 = P_2 = P_{\text{atm}}$ ) and that the fluid velocities at both points are zero ( $V_1 = V_2 = 0$ ), the energy equation for a control volume between these two points simplifies to

$$\frac{P_1}{\rho g} + \alpha_1 \frac{V_1^2}{2g} + z_1 = \frac{P_2}{\rho g} + \alpha_2 \frac{V_2^2}{2g} + z_2 + h_L \quad \rightarrow \quad z_1 = z_2 + h_L$$

where

$$h_L = h_{L, \text{total}} = h_{L, \text{major}} + h_{L, \text{minor}} = \left( f \frac{L}{D} + \sum K_L \right) \frac{V^2}{2g}$$

since the diameter of the piping system is constant. The average velocity in the pipe and the Reynolds number are

$$V = \frac{\dot{V}}{A_c} = \frac{\dot{V}}{\pi D^2/4} = \frac{0.006 \text{ m}^3/\text{s}}{\pi(0.05 \text{ m})^2/4} = 3.06 \text{ m/s}$$

$$\text{Re} = \frac{\rho V D}{\mu} = \frac{(999.7 \text{ kg/m}^3)(3.06 \text{ m/s})(0.05 \text{ m})}{1.307 \times 10^{-3} \text{ kg/m} \cdot \text{s}} = 117,000$$

The flow is turbulent since  $\text{Re} > 4000$ . Noting that  $\varepsilon/D = 0.00026/0.05 = 0.0052$ , the friction factor can be determined from the Colebrook equation (or the Moody chart),

$$\frac{1}{\sqrt{f}} = -2.0 \log \left( \frac{\varepsilon/D}{3.7} + \frac{2.51}{\text{Re} \sqrt{f}} \right) \quad \rightarrow \quad \frac{1}{\sqrt{f}} = -2.0 \log \left( \frac{0.0052}{3.7} + \frac{2.51}{117,000 \sqrt{f}} \right)$$

It gives  $f = 0.0315$ . The sum of the loss coefficients is

$$\begin{aligned} \sum K_L &= K_{L, \text{entrance}} + 2K_{L, \text{elbow}} + K_{L, \text{valve}} + K_{L, \text{exit}} \\ &= 0.5 + 2 \times 0.3 + 0.2 + 1.06 = 2.36 \end{aligned}$$

Then the total head loss and the elevation of the source become

$$h_L = \left( f \frac{L}{D} + \sum K_L \right) \frac{V^2}{2g} = \left( 0.0315 \frac{89 \text{ m}}{0.05 \text{ m}} + 2.36 \right) \frac{(3.06 \text{ m/s})^2}{2(9.81 \text{ m/s}^2)} = 27.9 \text{ m}$$

$$z_1 = z_2 + h_L = 4 + 27.9 = \mathbf{31.9 \text{ m}}$$

Therefore, the free surface of the first reservoir must be 31.9 m above the ground level to ensure water flow between the two reservoirs at the specified rate.

**Discussion** Note that  $fL/D = 56.1$  in this case, which is about 24 times the total minor loss coefficient. Therefore, ignoring the sources of minor losses in this case would result in about 4 percent error.

It can be shown that the total head loss would be 35.9 m (instead of 27.9 m) if the valve were three-fourths closed, and it would drop to 24.8 m if the pipe between the two reservoirs were straight at the ground level (thus

eliminating the elbows and the vertical section of the pipe). The head loss could be reduced further (from 24.8 to 24.6 m) by rounding the entrance. The head loss can be reduced significantly (from 27.9 to 16.0 m) by replacing the cast iron pipes by smooth pipes such as those made of plastic.

### EXAMPLE 8–9 Effect of Flushing on Flow Rate from a Shower

The bathroom plumbing of a building consists of 1.5-cm-diameter copper pipes with threaded connectors, as shown in Fig. 8–49. (a) If the gage pressure at the inlet of the system is 200 kPa during a shower and the toilet reservoir is full (no flow in that branch), determine the flow rate of water through the shower head. (b) Determine the effect of flushing of the toilet on the flow rate through the shower head. Take the loss coefficients of the shower head and the reservoir to be 12 and 14, respectively.

**SOLUTION** The cold-water plumbing system of a bathroom is given. The flow rate through the shower and the effect of flushing the toilet on the flow rate are to be determined.

**Assumptions** 1 The flow is steady and incompressible. 2 The flow is turbulent and fully developed. 3 The reservoir is open to the atmosphere. 4 The velocity heads are negligible.

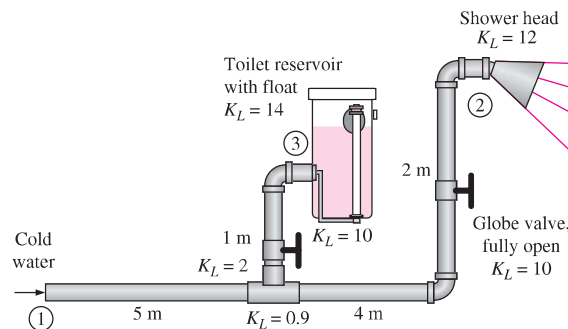
**Properties** The properties of water at 20°C are  $\rho = 998 \text{ kg/m}^3$ ,  $\mu = 1.002 \times 10^{-3} \text{ kg/m} \cdot \text{s}$ , and  $\nu = \mu/\rho = 1.004 \times 10^{-6} \text{ m}^2/\text{s}$ . The roughness of copper pipes is  $\varepsilon = 1.5 \times 10^{-6} \text{ m}$ .

**Analysis** This is a problem of the second type since it involves the determination of the flow rate for a specified pipe diameter and pressure drop. The solution involves an iterative approach since the flow rate (and thus the flow velocity) is not known.

(a) The piping system of the shower alone involves 11 m of piping, a tee with line flow ( $K_L = 0.9$ ), two standard elbows ( $K_L = 0.9$  each), a fully open globe valve ( $K_L = 10$ ), and a shower head ( $K_L = 12$ ). Therefore,  $\Sigma K_L = 0.9 + 2 \times 0.9 + 10 + 12 = 24.7$ . Noting that the shower head is open to the atmosphere, and the velocity heads are negligible, the energy equation for a control volume between points 1 and 2 simplifies to

$$\frac{P_1}{\rho g} + \alpha_1 \frac{V_1^2}{2g} + z_1 + h_{\text{pump}, u} = \frac{P_2}{\rho g} + \alpha_2 \frac{V_2^2}{2g} + z_2 + h_{\text{turbine}, e} + h_L$$

$$\rightarrow \frac{P_{1, \text{gage}}}{\rho g} = (z_2 - z_1) + h_L$$



**FIGURE 8–49**  
Schematic for Example 8–9.

Therefore, the head loss is

$$h_L = \frac{200,000 \text{ N/m}^2}{(998 \text{ kg/m}^3)(9.81 \text{ m/s}^2)} - 2 \text{ m} = 18.4 \text{ m}$$

Also,

$$h_L = \left( f \frac{L}{D} + \sum K_L \right) \frac{V^2}{2g} \quad \rightarrow \quad 18.4 = \left( f \frac{11 \text{ m}}{0.015 \text{ m}} + 24.7 \right) \frac{V^2}{2(9.81 \text{ m/s}^2)}$$

since the diameter of the piping system is constant. The average velocity in the pipe, the Reynolds number, and the friction factor are

$$\begin{aligned} V &= \frac{\dot{V}}{A_c} = \frac{\dot{V}}{\pi D^2/4} \quad \rightarrow \quad V = \frac{\dot{V}}{\pi(0.015 \text{ m})^2/4} \\ \text{Re} &= \frac{VD}{\nu} \quad \rightarrow \quad \text{Re} = \frac{V(0.015 \text{ m})}{1.004 \times 10^{-6} \text{ m}^2/\text{s}} \\ \frac{1}{\sqrt{f}} &= -2.0 \log \left( \frac{\epsilon/D}{3.7} + \frac{2.51}{\text{Re}\sqrt{f}} \right) \\ \rightarrow \quad \frac{1}{\sqrt{f}} &= -2.0 \log \left( \frac{1.5 \times 10^{-6} \text{ m}}{3.7(0.015 \text{ m})} + \frac{2.51}{\text{Re}\sqrt{f}} \right) \end{aligned}$$

This is a set of four equations with four unknowns, and solving them with an equation solver such as EES gives

$$\dot{V} = 0.00053 \text{ m}^3/\text{s}, \quad f = 0.0218, \quad V = 2.98 \text{ m/s}, \quad \text{and} \quad \text{Re} = 44,550$$

Therefore, the flow rate of water through the shower head is **0.53 L/s**.

(b) When the toilet is flushed, the float moves and opens the valve. The discharged water starts to refill the reservoir, resulting in parallel flow after the tee connection. The head loss and minor loss coefficients for the shower branch were determined in (a) to be  $h_{L,2} = 18.4 \text{ m}$  and  $\sum K_{L,2} = 24.7$ , respectively. The corresponding quantities for the reservoir branch can be determined similarly to be

$$\begin{aligned} h_{L,3} &= \frac{200,000 \text{ N/m}^2}{(998 \text{ kg/m}^3)(9.81 \text{ m/s}^2)} - 1 \text{ m} = 19.4 \text{ m} \\ \sum K_{L,3} &= 2 + 10 + 0.9 + 14 = 26.9 \end{aligned}$$

The relevant equations in this case are

$$\begin{aligned} \dot{V}_1 &= \dot{V}_2 + \dot{V}_3 \\ h_{L,2} &= f_1 \frac{5 \text{ m}}{0.015 \text{ m}} \frac{V_1^2}{2(9.81 \text{ m/s}^2)} + \left( f_2 \frac{6 \text{ m}}{0.015 \text{ m}} + 24.7 \right) \frac{V_2^2}{2(9.81 \text{ m/s}^2)} = 18.4 \\ h_{L,3} &= f_1 \frac{5 \text{ m}}{0.015 \text{ m}} \frac{V_1^2}{2(9.81 \text{ m/s}^2)} + \left( f_3 \frac{1 \text{ m}}{0.015 \text{ m}} + 26.9 \right) \frac{V_3^2}{2(9.81 \text{ m/s}^2)} = 19.4 \\ V_1 &= \frac{\dot{V}_1}{\pi(0.015 \text{ m})^2/4}, \quad V_2 = \frac{\dot{V}_2}{\pi(0.015 \text{ m})^2/4}, \quad V_3 = \frac{\dot{V}_3}{\pi(0.015 \text{ m})^2/4} \\ \text{Re}_1 &= \frac{V_1(0.015 \text{ m})}{1.004 \times 10^{-6} \text{ m}^2/\text{s}}, \quad \text{Re}_2 = \frac{V_2(0.015 \text{ m})}{1.004 \times 10^{-6} \text{ m}^2/\text{s}}, \quad \text{Re}_3 = \frac{V_3(0.015 \text{ m})}{1.004 \times 10^{-6} \text{ m}^2/\text{s}} \end{aligned}$$