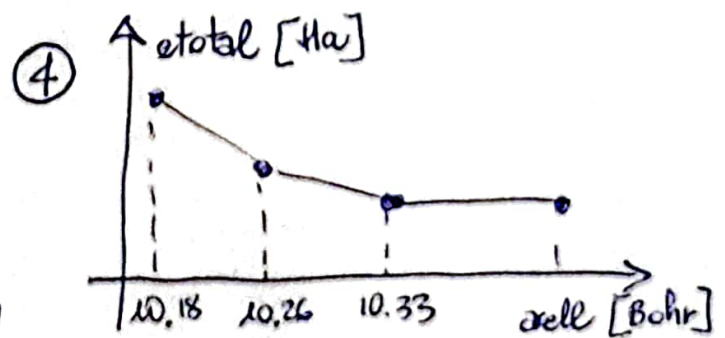
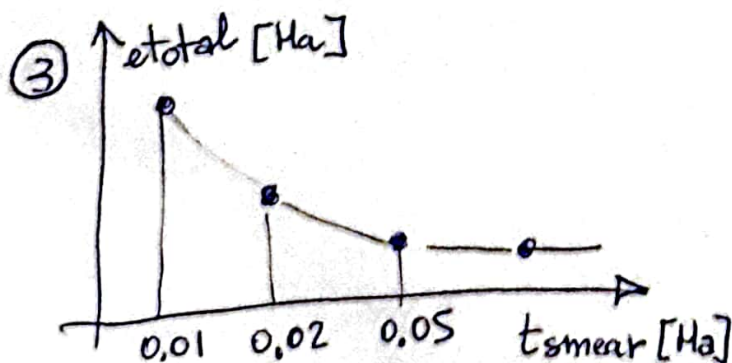
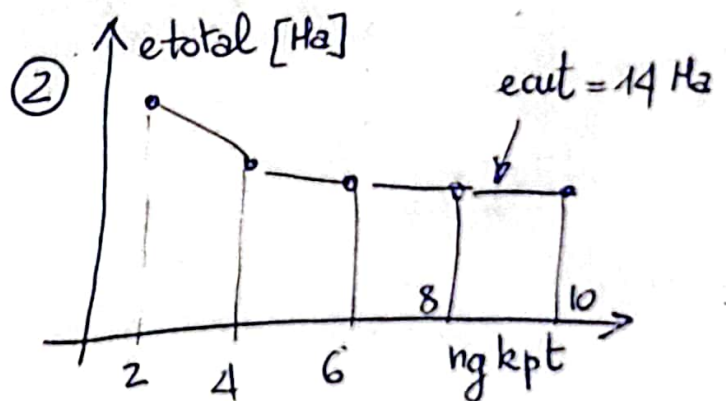
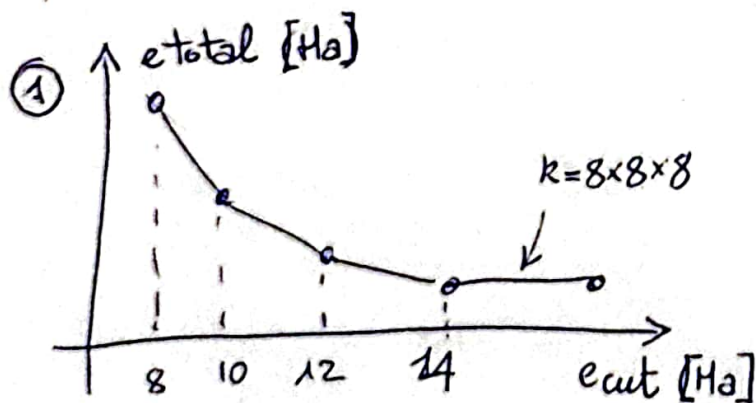


# DFT - Steps, general results and analysis

- Following and summarizing "ABINIT Hands-On" by Matteo Caguard:

- 1) Find convergence w.r.t. cut-off energy (using multi-dataset mode) for a certain grid of k-points (pagg. 25-27)
  - 2) Find convergence w.r.t. number of k-points (using multi-dataset mode) for the "optimized" cut-off energy
  - 3) (only in metals, but doped diamond is conducting)  
→ Find convergence w.r.t. smearing parameter. (pagg. 28-29)
- N.B. In principle these three convergences are not independent but...

- 4) Perform structural relaxation (pagg. 31-34)



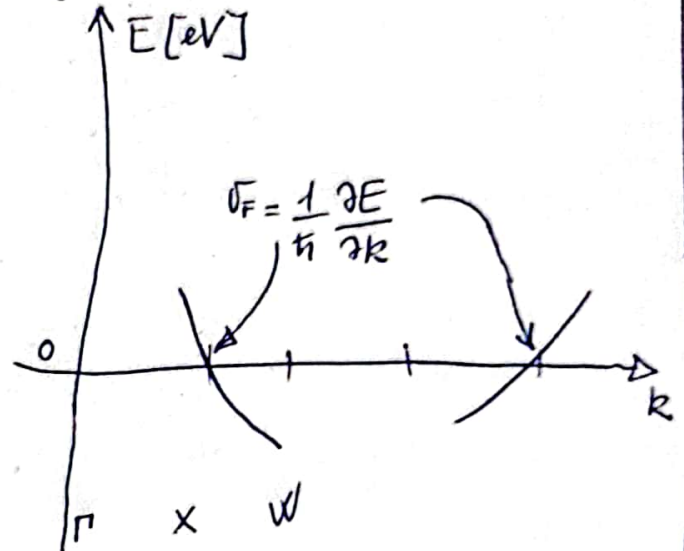
## 5) ELECTRONIC PROPERTIES

- Band structure (eventually adding Spin-Orbit Coupling <sup>2/</sup> (SOC) in heavy metals) → (pagg. 35-42)

WHAT YOU CAN GET FROM BAND STRUCTURE?

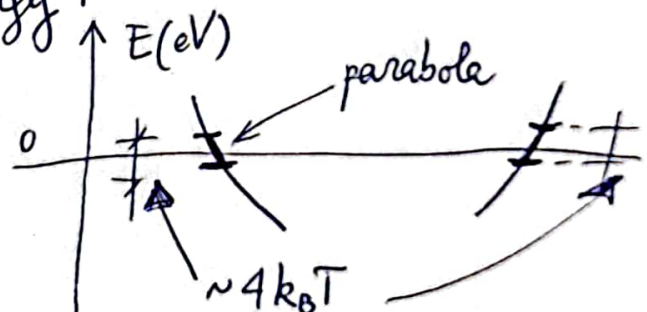
- Comparison ~~from~~ with results from literature
- Calculation of the Fermi velocity in any point where a band crosses the Fermi energy

$$v_F = \frac{1}{\hbar} \left[ \frac{\partial E(k)}{\partial k} \right]_{k=k_F}$$



- Calculation of the effective mass in any point where a band crosses the Fermi energy:

- fit by a quadratic (parabolic) behaviour a small region (of extension  $\sim 4k_B T$ ) around the  $E_F$  of the band



- Then you can obtain the effective mass in that  $k$  point by:

$$m^*(k_F) = \frac{\hbar^2}{\left( d^2 E / dk^2 \right)_{k=k_F}}$$

- The total (average) values of  $v_F$  and  $m^*$  on the Fermi surface, that can be used in other analyses on experimental values can be obtained, in first approximation, by a simple average of the different values found in the different  $k_F$  points



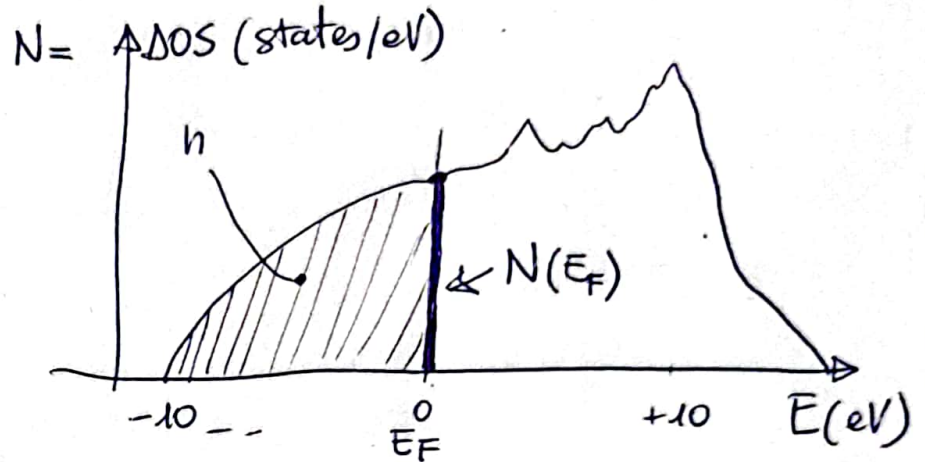
• Density of States → (page 41-47)

WHAT YOU CAN GET FROM DOS ?

a) Comparison with literature

b) Density of states at Fermi energy :

$N(E_F)$



c) Density of charge carriers n :

$$n = \int_{E_{\min}}^{E_F} N(E) dE$$

$n$  is very important for calculations and analysis of the transport properties (see Erik)

d) Comparison with  $n$  from literature ... but attention

→ most of the values of  $n$  present in the books are calculated in free electron model approximation

→ difference ? ...

• Fermi surface → (page 46-47)

a) Comparison with literature



## 6) PHONON PROPERTIES

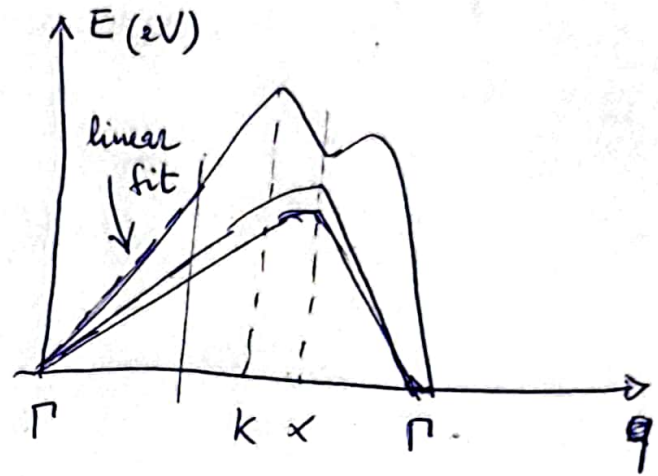
4/

- Dispersion relations (phonon bands) → (pagg. 53-58)

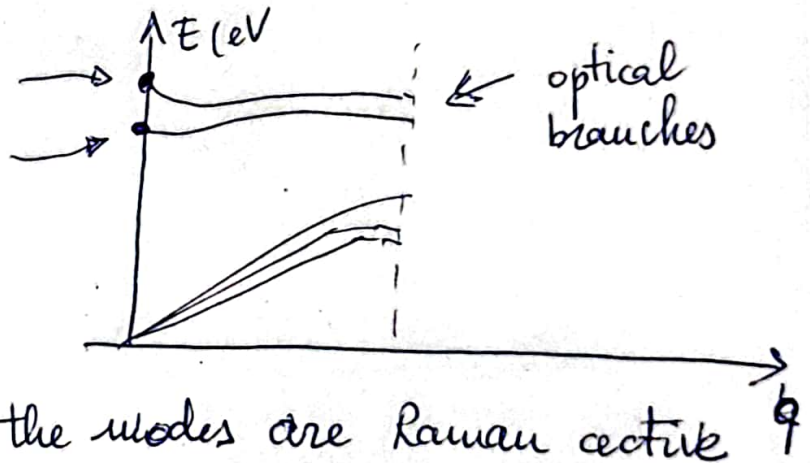
### WHAT YOU CAN GET FROM DISPERSION RELATIONS

a) Comparison with literature

b) Speed of sound in directions of high symmetry  
→ by a linear fit of the acoustic branches at energies close to zero i.e. at  $q$  vectors close to  $\Gamma$



c) The energy values of optical branches at  $q=0$  can give info on the energies of Raman peaks but



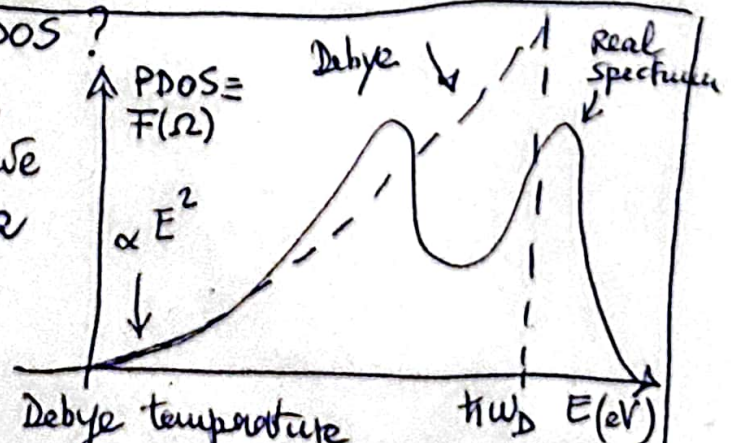
attention → not all the modes are Raman active

→ There are specific commands in ABINIT for the calculation of Raman and IR active-modes energies  
↳ comparison with Raman spectra in doped diamond

- Phonon DOS → (pag. 58)

### WHAT YOU CAN GET FROM PDOS?

a) By fitting the low-energy part ( $E \leq E_{\text{max}}/10$ ) by a parabolic curve and then extending this parabola up to the value that produce the same integral of the true PDOS → Debye energy or Debye temperature  
→ comparison with literature





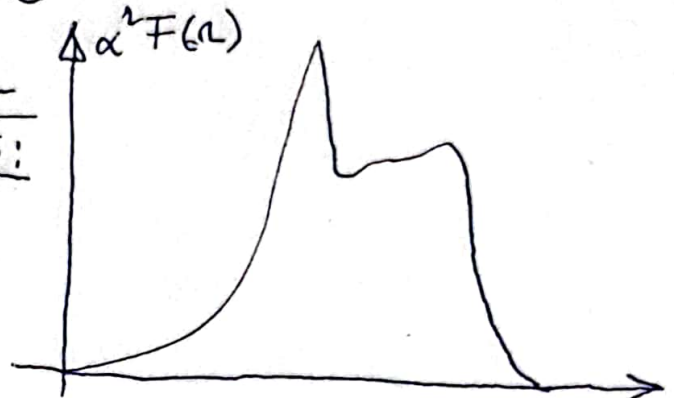
# • Electron-phonon spectral function $\alpha^2 F(\Omega) \rightarrow$ (page 59-67) <sup>5/</sup>

WHAT YOU CAN GET FROM  $\alpha^2 F(\Omega)$ ?

a) comparison with literature

b) The dynamical electron-phonon coupling constant:

$$\lambda = 2 \int_0^{\Omega_{\max}} \frac{\alpha^2 F(\Omega)}{\Omega} d\Omega$$



- Eventually you can compare this  $\lambda$  with the transport one  $\lambda_{tr}$  that you can obtain from the fit of the linear  $T$  dependence of resistivity at high  $T$  (see page 42-45 of AEP lecture Notes - Chapt. 5)

$$\rho = \rho_0 + \frac{2\pi k_B}{e^2 \hbar} \frac{m_{\text{eff}}}{n} \lambda_{tr} T \quad \text{where } m_{\text{eff}} \text{ can be obtained from DFT band structure and DOS}$$

Remember that  $\lambda$  and  $\lambda_{tr}$  are similar, but not identical because  $\alpha^2 F(\Omega)_{tr}$  is different from  $\alpha^2 F(\Omega)$  (due to the contribution of small-angle scattering in transport properties).  $\lambda$  and  $\lambda_{tr}$  are dimensionless and of order of 1.

c) Neglecting (as a first approximation) the previous difference, you can use  $\alpha^2 F(\Omega)$  (obtained from DFT) for the estimation of the measured resistivity by using eq. (1) at page 42 of Chapt. 5 of the notes.

d) If the metal is superconducting (as Nb), from  $\lambda$  you can obtain the critical temperature  $T_c$  by using the Allen-Dynes formula (page 13 and 14 of Chapt. 6 of the AEP notes) and by fixing (from literature) the value of the repulsive Coulomb pseudo-potential  $\mu^*$ .

This is automatically done by the routine of Abinit that calculates  $T_c$