

Quantum Physics 1

Class 13

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Simple Harmonic Oscillator

Review:

c) Functional Vector Space.

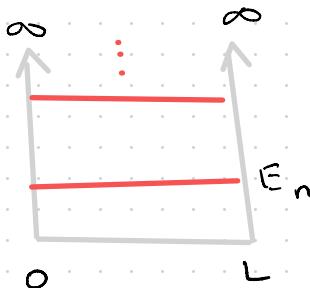
c) Energy eigenstates of
the time independent

S.E.

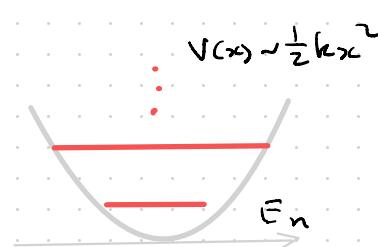
$$-\frac{\hbar^2}{2M} \frac{d^2}{dx^2} \varphi_n + V(x) \varphi_n = E_n \varphi_n$$

$\underbrace{\qquad\qquad\qquad}_{\hat{H} \in \text{Hamiltonian operator}}$

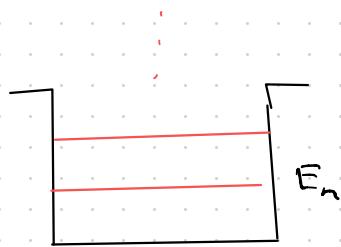
Examples of different eigenstates that manifest from different $V(x)$:



square well

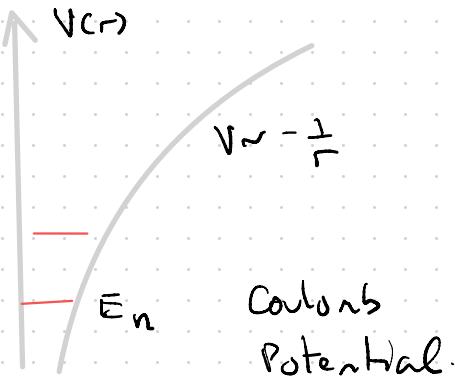


harmonic Potential Well



Finite square well

$\Psi(x), \varphi_n, \dots$
 $\hat{x}, \hat{p}, \hat{H}, \dots$



- Can write $\Psi_{(x)} = \sum c_n \varphi_n(x)$, expansion of eigenstates of the \hat{H} . Solutions of time dependent S.E. :

$$\underline{\Psi}(x, t) = \sum c_n \varphi_n(x) e^{-i E_n t / \hbar}$$

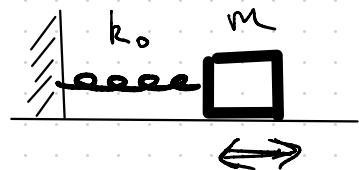
- Can solve Fourier coefficients by considering

$$\begin{aligned}
 \int \varphi_m^* \underline{\Psi}(x, t=0) dx &= \int \sum c_n \varphi_n^*(x) \varphi_n(x) dx \\
 &= \sum c_n \delta_{mn} \\
 &= \underline{\underline{c_n}}
 \end{aligned}$$

Recall: Classical Harmonic Oscillator

$$V = \frac{1}{2} k_0 x^2, \quad k_0 \equiv \text{spring constant}$$

$$F = -\frac{\partial V}{\partial x} = -k_0 x = m \frac{d^2x}{dt^2}$$

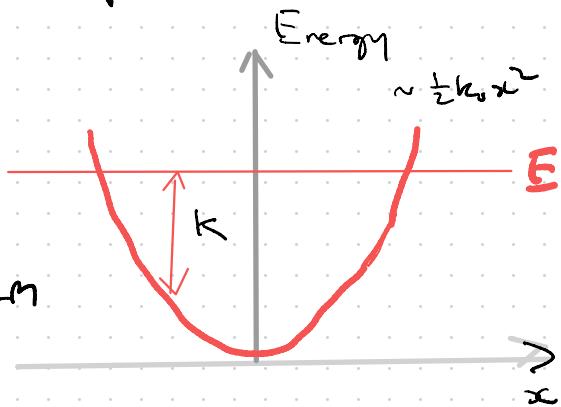


with solutions: $e^{i\omega t}$; $\omega \equiv \sqrt{k_0/m}$

$$\nabla H = \frac{\hat{p}^2}{2m} + V(x) \Rightarrow$$

$$k = \frac{\hat{p}^2}{2m}$$

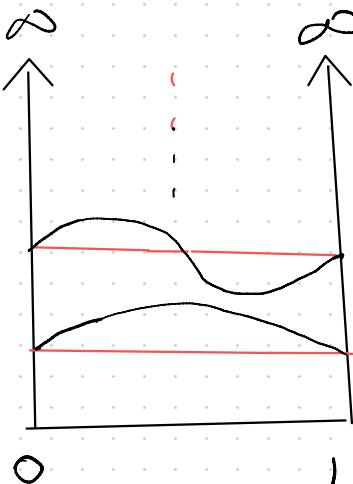
$$\therefore E = \frac{\hat{p}^2}{2m} + V(x)$$



Quantum Case:

$$\hat{H} = \frac{\hat{p}^2}{2m} + V(\hat{x}) \Rightarrow -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_{00}}{\partial x^2} + V(x) \psi_{00} = E \psi_{00}$$

In-class 13.1, 13.2



∞ -square well

$n=2$

$n=1$

wave penetration

Quantum Harmonic oscillator
(QHO)

Solu:

$$e^{ikx} \text{ for } \left[\frac{p^2}{2m} + V(x) \right] \psi = E \psi$$

$\curvearrowleft = 0 \text{ inside, } \infty \text{ outside}$

$$\text{For QHO } V(x) \sim \frac{1}{2} k x^2$$

Eigenfunctions for the QHO

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi(x) + \frac{1}{2} k_0 x^2 \psi(x) = E \psi(x)$$

$\sim x^2$

with $\psi_n(x) = A_n H_n\left(\sqrt{\frac{m\omega}{\hbar}}x\right) e^{-m\omega x^2/2\hbar}$

\nearrow
Hermite Polynomials

$$H_n(y) = (-1)^n e^{y^2} \frac{d^n}{dy^n} e^{-y^2}$$

$$E_n = (n + \frac{1}{2}) \hbar \omega ; \quad n = 0, 1, 2, \dots$$

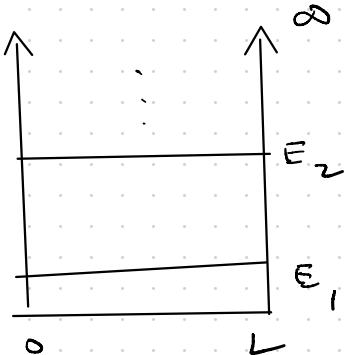
NB : $n=0$ Ground state
 $\omega \hbar t \quad E_0 = \frac{1}{2} \hbar \omega$

$$\Delta E_{n, n+1} = \hbar \omega \quad ; \quad E_{n+1} - E_n = [(n+1) + \frac{1}{2}] \hbar \omega - [n + 1] \hbar \omega \\ = \hbar \omega$$

In-class 13.3, 13.4

Compare the Infinite Square Well:

$$\begin{aligned}\Delta E_{n+1, n} &= E_{n+1} - E_n \\ &= \frac{\hbar^2 (n+1)^2 \pi^2}{2mL^2} - \frac{\hbar^2 n^2 \pi^2}{2mL^2} \\ &= \frac{2n\pi^2 \hbar^2}{2mL^2} + \frac{\pi^2 \hbar^2}{2mL^2}\end{aligned}$$

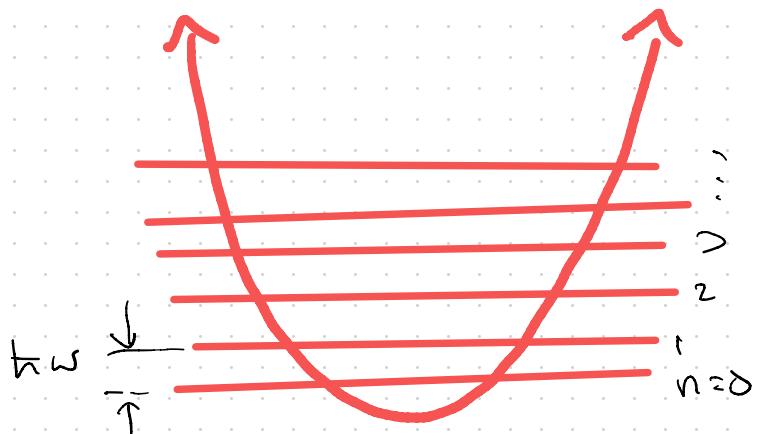


$\propto n$ (dependent on $n!$)

\therefore if $\Delta E_{n+1, n} \uparrow$ then $n \uparrow$

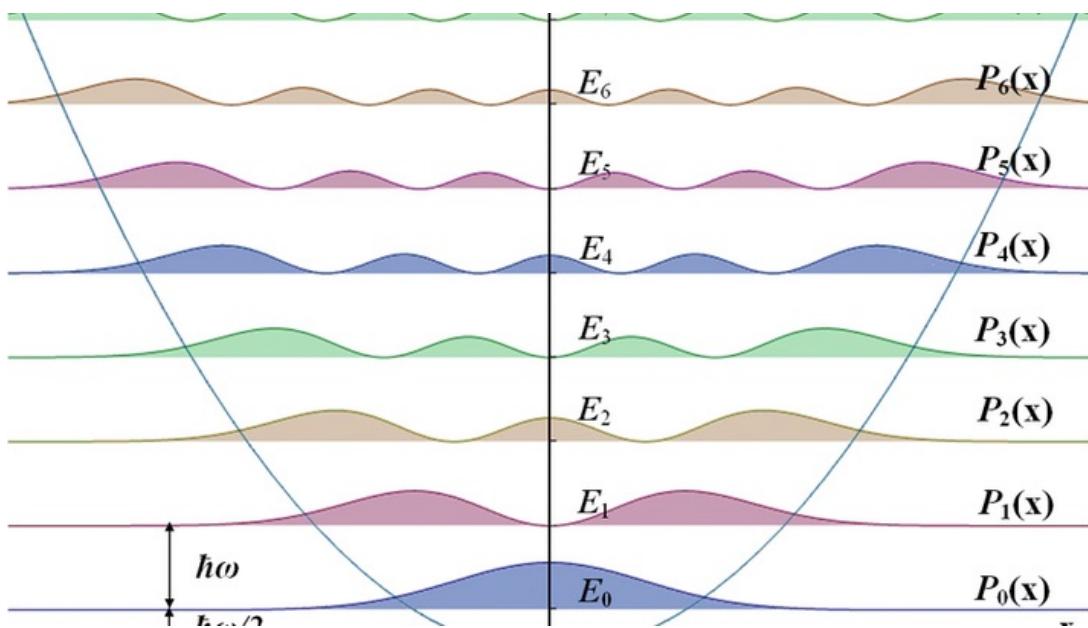
But for QHO: $\Delta E = \hbar\omega$;

equally spaced and independent of n .



Applications :

1. Phonons
2. Molecular vibrations.
3. Any system where can make parabolic approximation at small E, V .



The Science Of Agents Of S.H.I.E.L.D.: What Is Quantum Harmonic Oscillation?

<https://www.forbes.com/sites/chadorzel/2015/10/13/the-science-of-agents-of-s-h-i-e-l-d-what-is-quantum-harmonic-oscillation/#7c2b81c3e7e3>