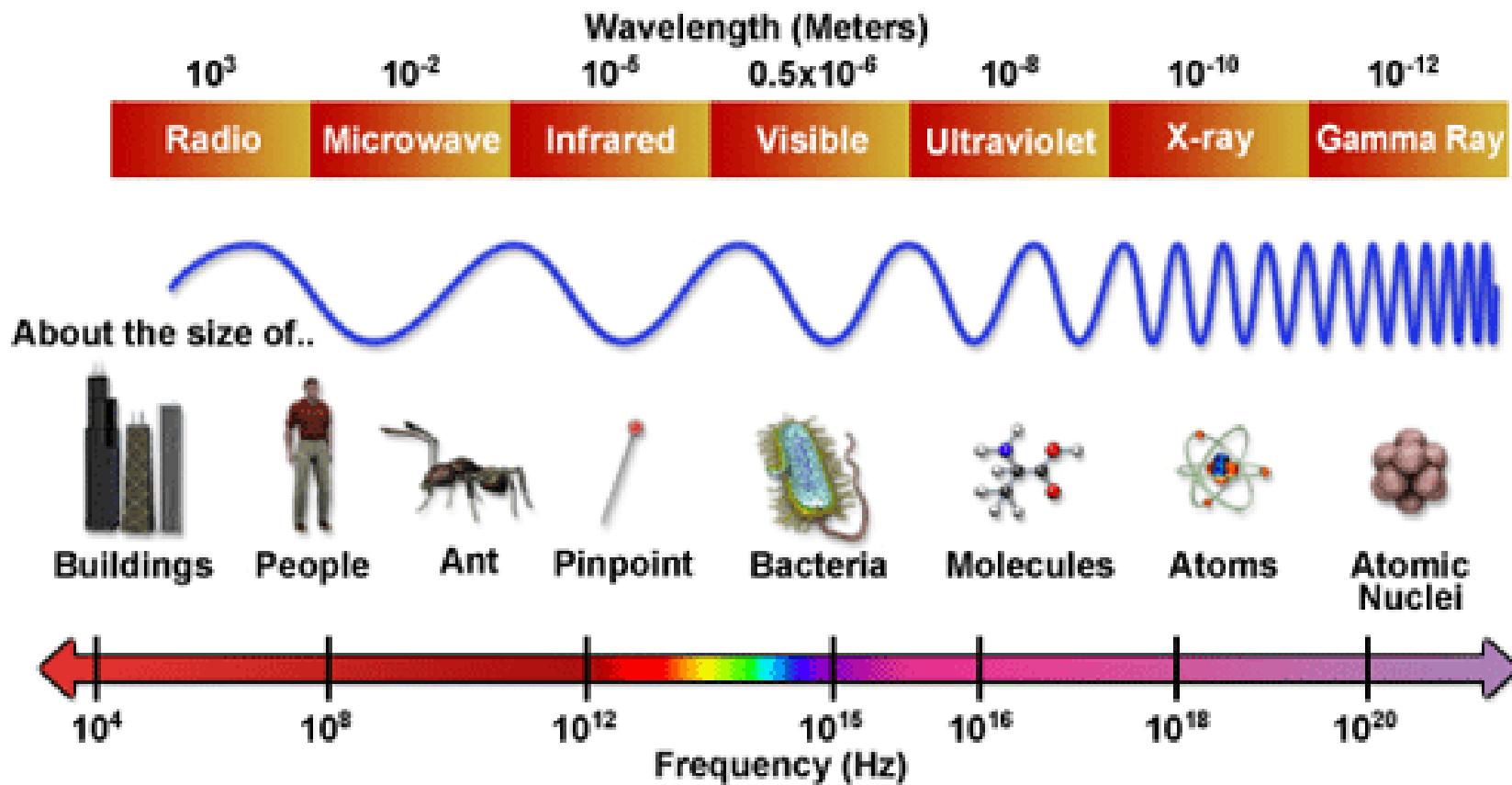


Class 17 Electromagnetic Waves

(3/21/24)

Electromagnetic Spectrum



UNITED STATES FREQUENCY ALLOCATIONS

THE RADIO SPECTRUM



ACTIVITY CODE

FEDERAL GOVERNMENT TECHNICAL/NON-FEDERAL SHARED

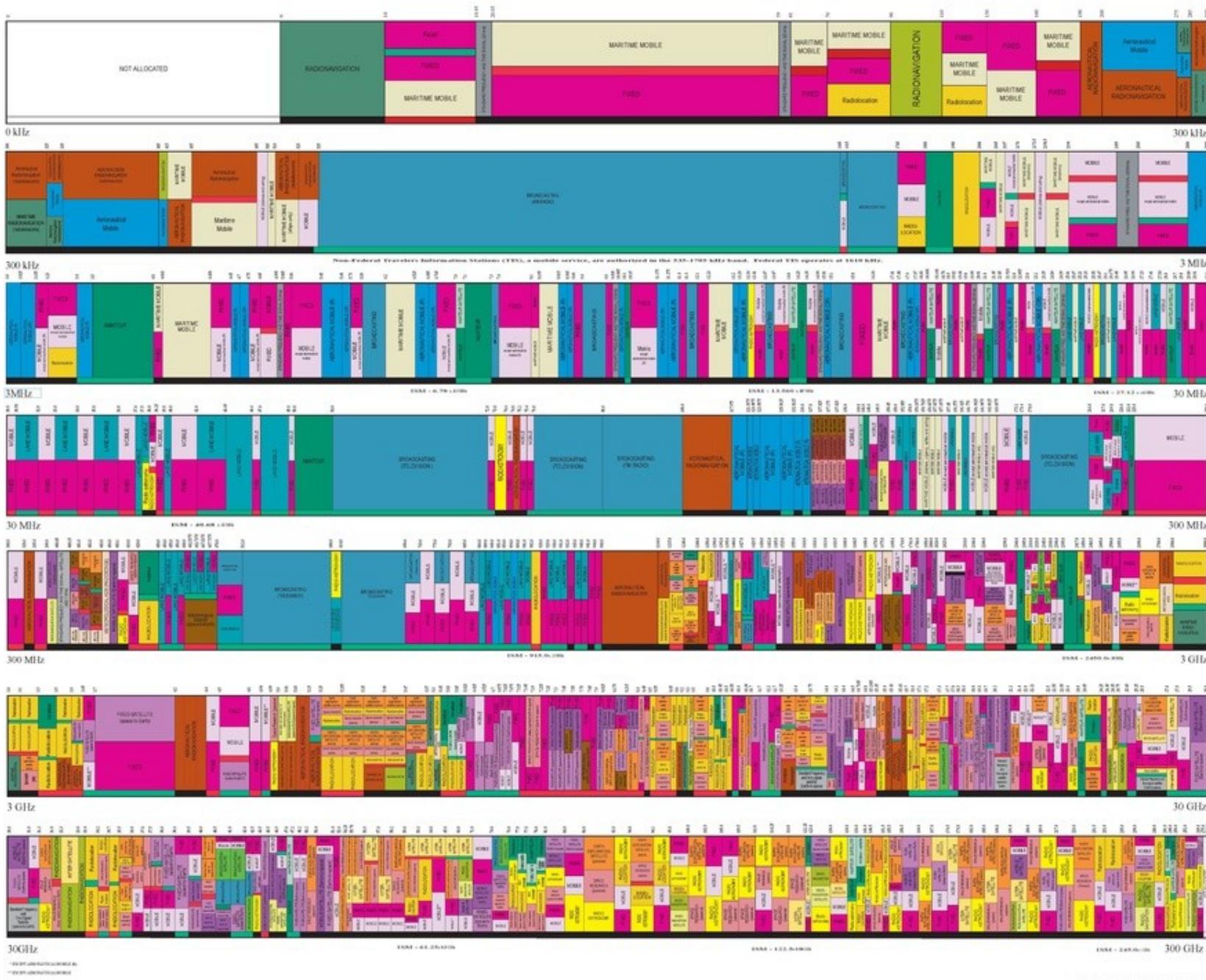
NONFEDERAL USAGE CODE

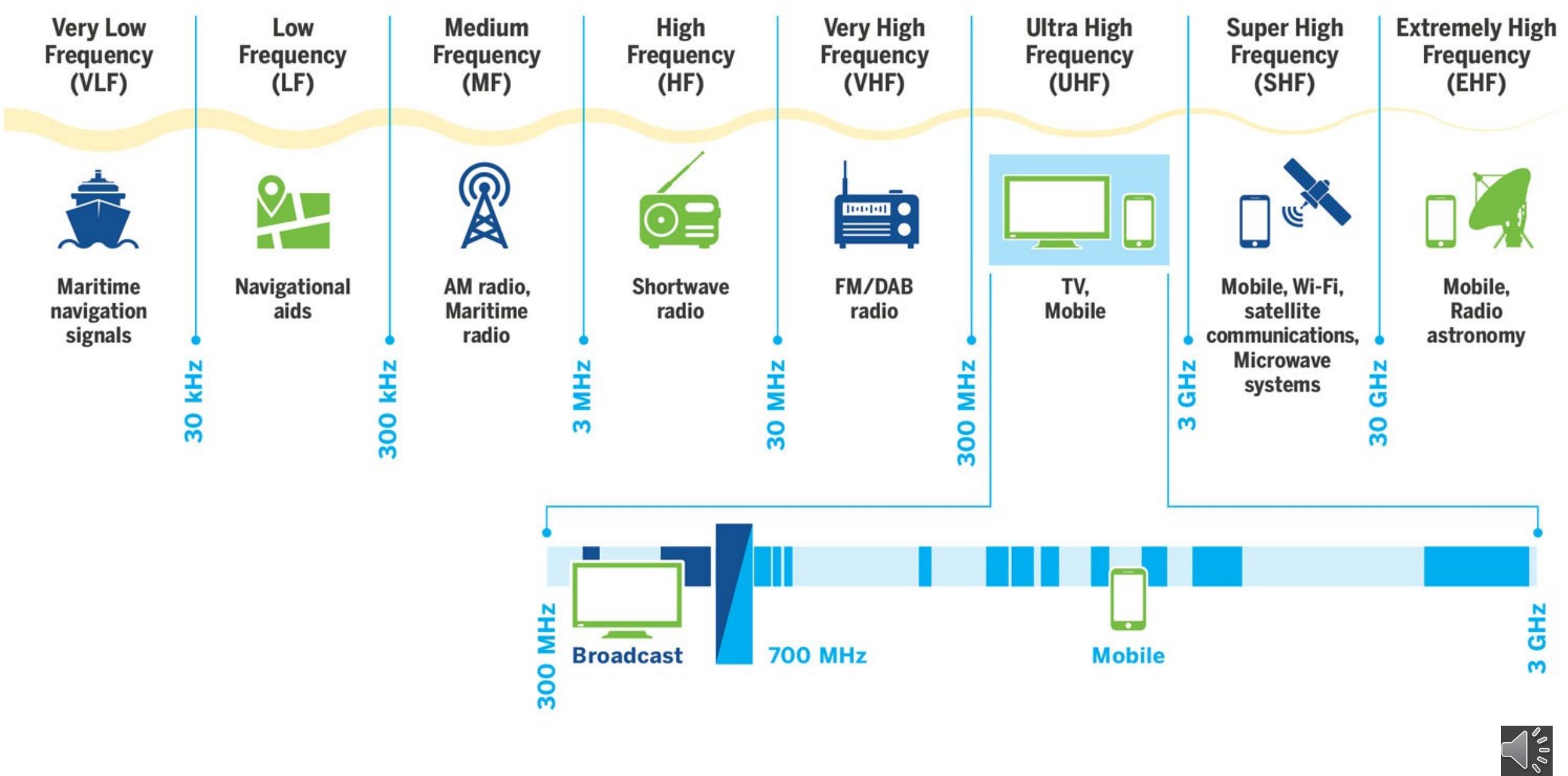
ALLOCATION USAGE DESIGNATION

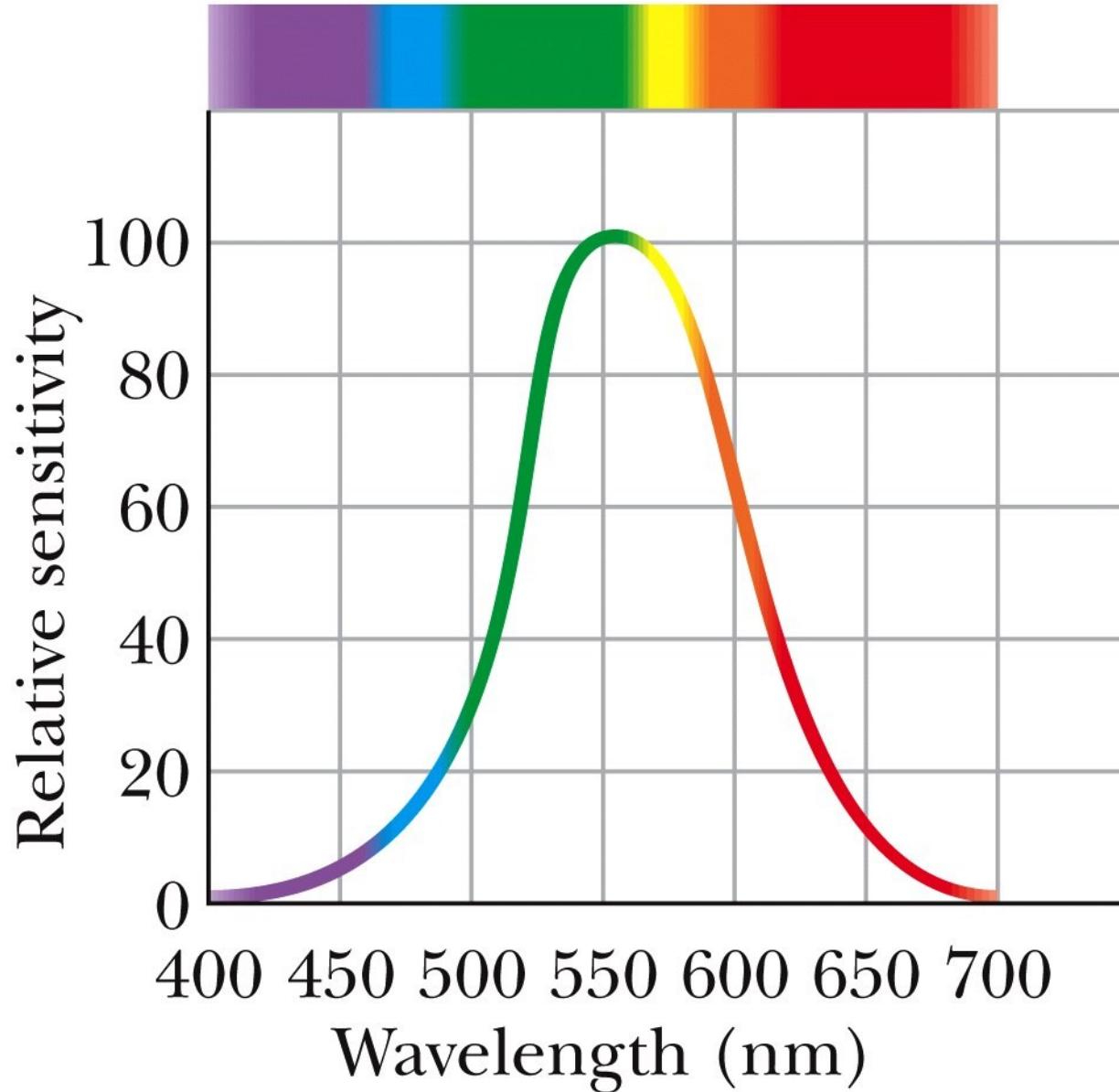
USAGE	EXAMPLE	DESCRIPTION
Federal	FIXED	Capital Letter
Nonfederal	Mobile	For Capital with lower case letter

This chart is a graphic representation of the Radio Frequency Allocation and Shared Spectrum Usage in the United States. It is not a complete catalog of all allocations and usage changes made in the Radio Frequency Allocation. Therefore, for complete information, users should consult the Federal Communications Commission's (FCC) database for current radio frequency assignments.

U.S. DEPARTMENT OF COMMERCE
National Telecommunications and Information Administration
Office of Spectrum Management
JANUARY 2016







As human beings we are outfitted with sensors (eyes) which detect electromagnetic waves between 400nm and 800nm. This is the reason behind calling this spectral range the “Visible Light”.



9.2 Electromagnetic waves in vacuum

EM waves travel through vacuum, corresponding to set of Maxwell's eqns:

$$\vec{\nabla} \cdot \vec{E} = 0, \quad \vec{\nabla} \cdot \vec{B} = 0, \quad \vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}, \quad \vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} = \mu_0 \epsilon_0 \vec{E}$$

derivation of wave equation:

$$\vec{\nabla} \times \vec{\nabla} \times \vec{E} = -\vec{\nabla} \times \frac{\partial \vec{B}}{\partial t} = -\mu_0 \epsilon_0 \vec{E}$$

$$\vec{\nabla} \times \vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \vec{\nabla} \times \vec{E} = -\mu_0 \epsilon_0 \vec{B}$$

$$\vec{\nabla}(\vec{\nabla} \cdot \vec{E}) - \vec{\nabla}^2 \vec{E} = -\mu_0 \epsilon_0 \vec{E} \Rightarrow \vec{\nabla}^2 \vec{E} = \mu_0 \epsilon_0 \vec{E}$$

$$\vec{\nabla}(\vec{\nabla} \cdot \vec{B}) - \vec{\nabla}^2 \vec{B} = -\mu_0 \epsilon_0 \vec{B} \Rightarrow \vec{\nabla}^2 \vec{B} = \mu_0 \epsilon_0 \vec{B}$$

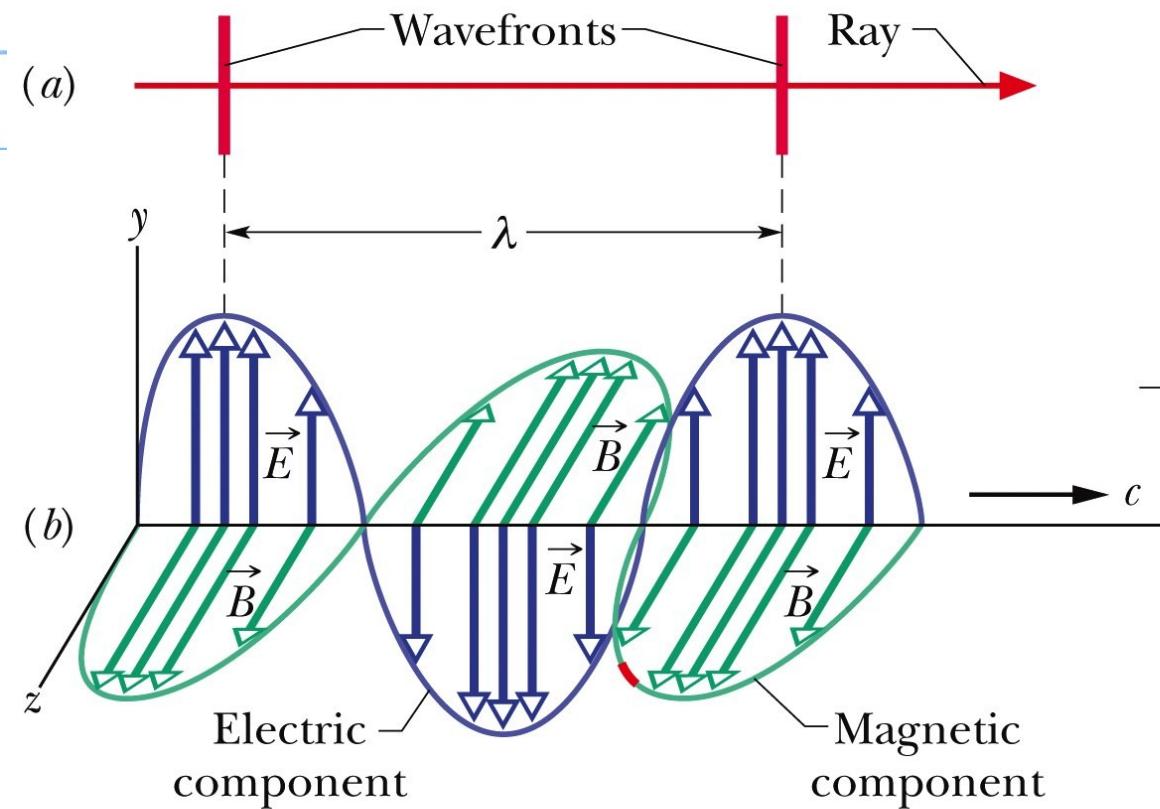
with $\mu_0 \epsilon_0 = \frac{1}{C^2}$

general solutions $\vec{E}(z, t) = \vec{E}_0 e^{i(kz - \omega t)}$

$$\vec{B}(z, t) = \vec{B}_0 e^{i(kz - \omega t)}$$

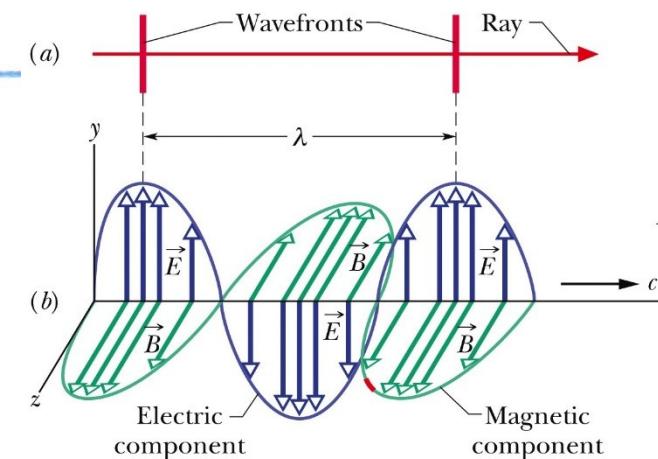
monochromatic (oscillators at a single frequency)

plane wave



em waves are transverse waves follows from

$$\vec{\nabla} \cdot \vec{E} = 0, \quad \vec{\nabla} \cdot \vec{B} = 0 \quad (\text{example})$$



Transverse em wave

$$\vec{E} = E_0 \cos(\omega z - \omega t) \hat{x} + E_0 \cos(\omega z - \omega t) \hat{y}$$

$$\vec{\nabla} \cdot \vec{E} = \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0$$

Because E_x, E_y are independent of z , i.e. $\frac{\partial E_x}{\partial z} = 0, \frac{\partial E_y}{\partial z} = 0$,

and $E_z = 0$, i.e. $\frac{\partial E_z}{\partial z} = 0$.



\vec{E} , \vec{B} , \vec{k} are perpendicular to each other follows

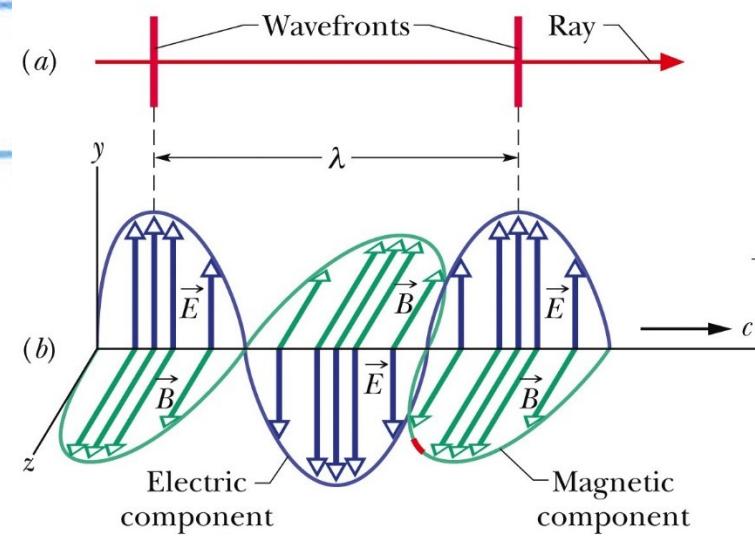
from Faraday's law $\vec{\nabla} \times \vec{E} = -\frac{\partial}{\partial t} \vec{B}$ or $\vec{B}_0 = \frac{1}{\omega} (\vec{k} \times \vec{E}_0)$

for amplitudes $B_0 = \frac{1}{\omega} E_0 = \frac{1}{c} E_0$ (example)

$$\vec{E}(z, t) = E_0 \cos(\theta z - \omega t + \delta) \hat{x}$$

$$\vec{B}(z, t) = \frac{1}{c} E_0 \cos(\theta z - \omega t + \delta) \hat{y}$$

$$\text{and } \vec{k} = \frac{c}{\lambda} \hat{z}$$



\vec{B} is \perp to \vec{E} follows from Faraday's Law.

Consider $\vec{E} = E_0 \cos(kz - \omega t) \hat{x}$

Calculate:

$$\vec{\nabla} \times \vec{E} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & E_y & 0 \end{vmatrix} = \left(\frac{\partial}{\partial y} 0 - \frac{\partial E_x}{\partial z} \right) \hat{x} + \left(\frac{\partial E_x}{\partial z} - \frac{\partial 0}{\partial x} \right) \hat{y} +$$
$$= 0 \quad \frac{\partial E_x}{\partial z} = 0 \quad = 0$$
$$\left(\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \right) \hat{z}$$
$$= 0 \quad = 0$$

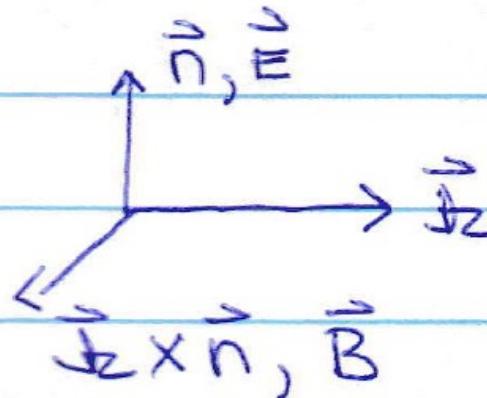
$$\vec{\nabla} \times \vec{E} = \frac{\partial E_x}{\partial z} \hat{y} = - \frac{\partial}{\partial t} \vec{B} \quad \downarrow \vec{B} \text{ in } \hat{y} \text{ direction}$$



\vec{n} : vector in direction of \vec{E}

\vec{k}_z : vector in direction of propagation

$\vec{k}_z \times \vec{n}$: vector perpendicular to \vec{n} and \vec{k}_z



general description

$$\vec{E}(\vec{r}, t) = E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} \hat{n}$$

$$\vec{B}(\vec{r}, t) = \frac{1}{c} E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} (\hat{k} \times \hat{n}) = \frac{1}{c} \hat{k} \times \vec{E}$$

$$\text{with } \hat{n} \cdot \hat{k} = 0$$



Polarization of an em wave is defined with regard to the direction of the em electric field vector \vec{E}

in general $\vec{E} = \vec{E}_x + \vec{E}_y = E_0 \cos(\omega z - \omega t + \phi) \hat{x} + E_0 \cos(\omega z - \omega t) \hat{y}$

$\phi = 0$ linear polarization, $\phi = \frac{\pi}{4}$ circular polarization

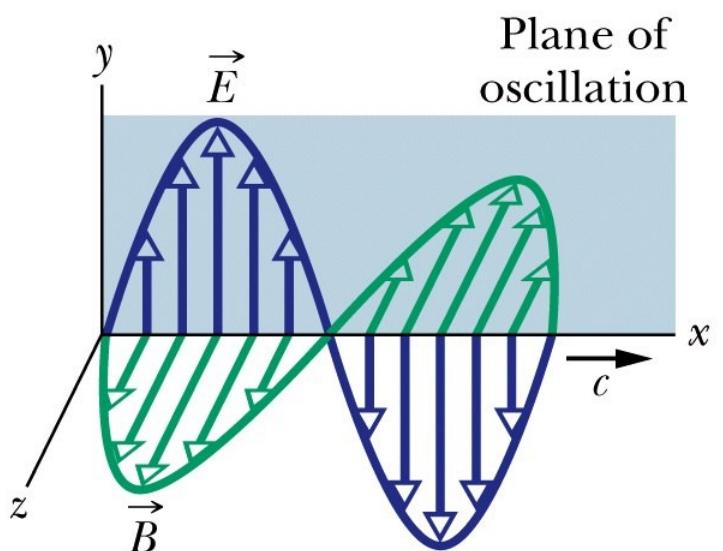
and for ϕ other than $0, \frac{\pi}{4}$, elliptical polarization

derived from eliminating t from wave equation

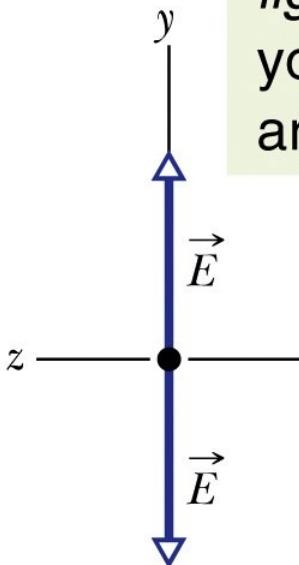
for components E_x, E_y , i.e $E_x = E_{0x} \cos(\omega t + \phi)$ and

$$E_y = E_{0y} \cos \omega t$$





(a)

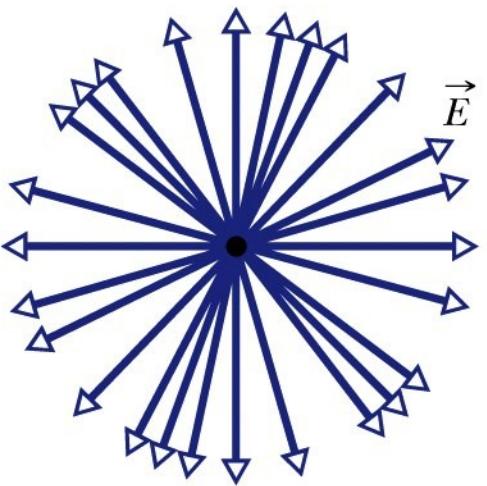


(b)

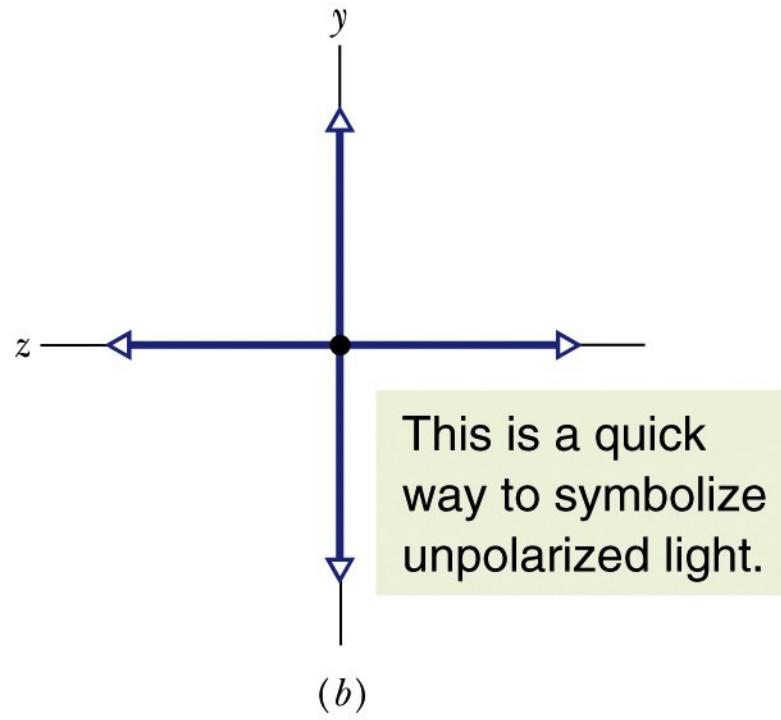
*Vertically polarized
light headed toward
you—the electric fields
are all vertical.*



Unpolarized light
headed toward
you—the electric
fields are in all
directions in the
plane.



(a)



This is a quick
way to symbolize
unpolarized light.

(b)



Energy & momentum in electro magnetic waves

em energy density $\epsilon_0 = \frac{1}{2} (\epsilon_0 E^2 + \frac{1}{\mu_0} B^2)$

for monochromatic plane wave $B^2 = \frac{1}{c^2} E^2 = \mu_0 \epsilon_0 E^2$

$\sim \epsilon_0 = \epsilon_0 E^2 = \epsilon_0 E_0^2 \cos^2(\omega z - ct + \delta)$

energy flux density $\vec{S} = \frac{1}{\mu_0} (\vec{E} \times \vec{B})$

$$\vec{S} = c \epsilon_0 E^2 \cos^2(\omega z - ct + \delta) \hat{z} = c \epsilon_0 \frac{1}{2} \hat{z}$$

momentum density $\vec{g} = \frac{1}{c^2} \vec{S} = \frac{1}{c} \epsilon_0 \frac{1}{2} \hat{z}$ momentum density

$$\langle \epsilon_0 \rangle = \frac{1}{2} \epsilon_0 E_0^2 \quad \langle \vec{S} \rangle = \frac{1}{2} c \epsilon_0 E_0^2 \frac{1}{2} \hat{z} \quad \langle \vec{g} \rangle = \frac{1}{2c} \epsilon_0 E_0^2 \frac{1}{2} \hat{z}$$



classical em light is a wave, em energy density,
Poynting vector, momentum, $\vec{E} = \vec{r} \times \vec{g}$

quantum light particle, energy $E = h \cdot F$, momentum
is $\vec{p} = \frac{\hbar}{2\pi} \vec{k} = \hbar \vec{k}$, conjugate momentum
 $\vec{E} = \vec{r} \times \vec{p}$

