

Quantum Physics 1

Class 10

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Functional vector space

Last Time :

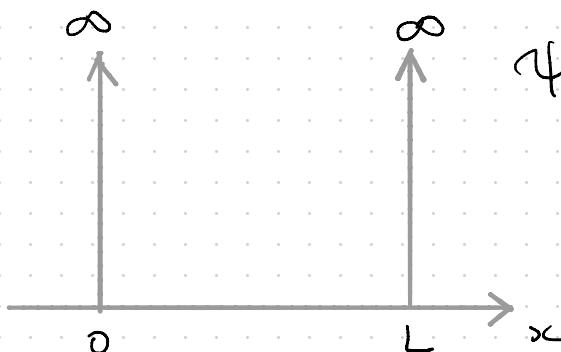
we solved S.E. using separation of variables

$$\underline{\Psi}(x, t) = \Psi(x) f(t)$$

$$\left\{ \frac{\hbar^2}{2m} \frac{\partial^2 \Psi(x)}{\partial x^2} + V(x) \Psi(x) = E \Psi(x) \dots \textcircled{1} \right.$$
$$\frac{df}{dt} = -i \frac{E}{\hbar} f(t) \dots \textcircled{2}$$
$$\Rightarrow f(t) = f(0) e^{-iEt/\hbar}$$

Infinite Square Well :

↗ normaliz.
const.



$$\Psi(x) = \sqrt{\frac{2}{L}} \sin(k_n x)$$

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$

↗ Kinetic energy

VIP points:

- (1) Energy is quantized
- (2) Ground state energy > 0
- (3) $\Psi^*\Psi$ is time independent.

In-class 10.1, 10.2,
10.3

Recall superposition:

$$\Psi(x,t) = \sum c_n \varphi_n(x) e^{-i E n t / \hbar}$$

at $t=0$ $\Psi(x,0) = \sum c_n \varphi_n(x)$ ** Fourier Series
 \hookrightarrow what is Expansion
 c_n ?

$$\Rightarrow \Psi(x,0) = \sum c_n \varphi_n(x)$$

multiply both sides by φ_m^*

$$\Rightarrow \varphi_m^* \Psi(x,0) = \sum c_n \varphi_m^* \varphi_n(x)$$

$$\Rightarrow \int (\varphi_m^* \Psi(x,0)) dx = \sum c_n \underbrace{\int \varphi_m^* \varphi_n(x) dx}_{\{}}$$

$$\underbrace{\int \varphi_m^* \varphi_n dx}$$

$$\int \sin\left(\frac{m\pi x}{L}\right) \sin\left(\frac{n\pi x}{L}\right) dx$$

$$= \begin{cases} 0 & m \neq n \\ 1 & m = n \end{cases}$$

$$= \delta_{m,n} \text{ Kronecker delta.}$$

$$\Rightarrow \int \varphi_m^* \bar{\Psi}(x, 0) dx = \sum c_n \delta_{m,n}$$

$$= c_m$$

$\therefore c_m = \int \varphi_m^* \bar{\Psi}(x, 0) dx \equiv \text{Fourier coefficient.}$

NB: $|c_m|^2 \sim \text{probability of a particle being in the } m^{\text{th}} \text{ state.}$

Now, recall for $t > 0$

$$\bar{\Psi}(x, t) = \sum_n c_n \varphi_n(x) e^{-i E_n t / \hbar}$$

Consider: $\int_0^L \bar{\Psi}^*(x,t) \bar{\Psi}(x,t) dx dt = 1$

$$= \int_0^L \sum_n c_n^* \psi_n^*(x) \sum_m c_m \psi_m(x) dx$$

$$= \sum_n \sum_m c_n^* c_m \delta_{m,n}$$

$$= \sum_n |c_n|^2 = 1$$

↳ probability of finding
particle in the
 n^{th} state

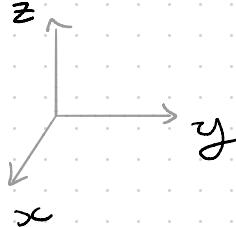
In-class 10.4

Functional Vector Space

recall Geometrical vector space-

vectors: \vec{A}, \vec{B}, \dots

in Euclidean space



$$\vec{A} = A_x \hat{i} + A_y \hat{j} + A_z \hat{k}$$

$$= \sum A_i \hat{\gamma}_i ; \hat{\gamma}_1 = \hat{i}$$

$$\hat{\gamma}_2 = \hat{j}$$

$$\hat{\gamma}_3 = \hat{k}$$

recall:

$$(i) \hat{\gamma}_i \cdot \hat{\gamma}_j = \delta_{ij}$$

$$(ii) A_i \hat{i} = \vec{A} \cdot \hat{\gamma}_i$$

$$= A_x$$

compare with:

Wavefunction vector space
with objects:

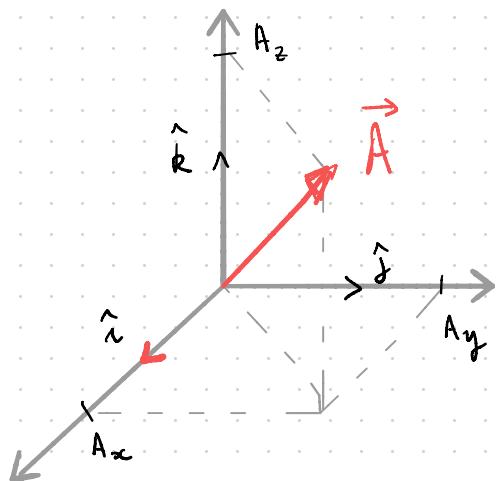
$$\Psi(x, t), \varphi_n, \varphi_n^*$$

$\varphi_n \Rightarrow$ unit vectors
(basis vectors)

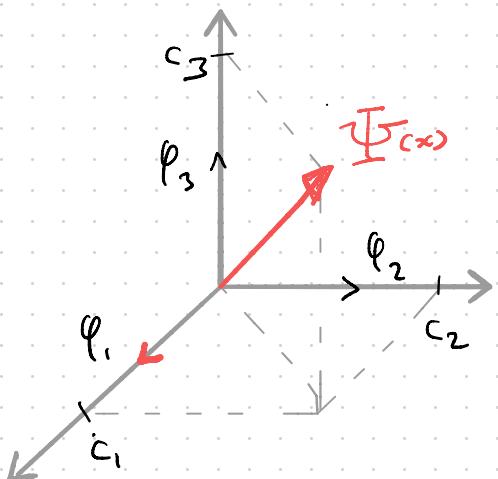
$$\underline{\Psi}(x) = \sum c_n \varphi_n(x)$$

$$c_n = \int \varphi_n^* \underline{\Psi}(x) dx$$

Geometrical vector space



Wavefunction vector space:



Now,

$$\hat{\vec{r}}_i \rightarrow \varphi_n$$

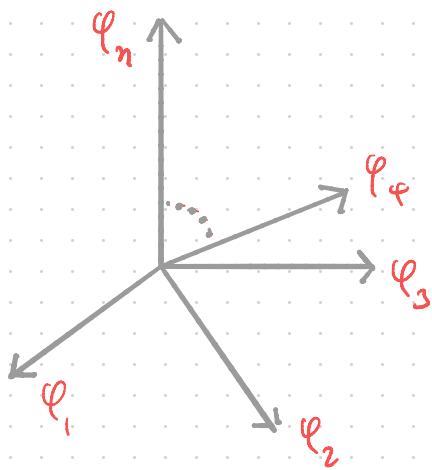
$$A_i \rightarrow c_n$$

$$\left\{ \begin{array}{l} \vec{A} \cdot \hat{\vec{r}}_i = A_x \\ c_n = \int \varphi_n^*(x) \Psi^{(\infty)}(x) dx \end{array} \right.$$

$$\left\{ \begin{array}{l} \hat{\vec{r}}_i \cdot \hat{\vec{r}}_j = \delta_{ij} \\ \int \varphi_m^* \varphi_n dx = \delta_{mn} \end{array} \right.$$

in general:

\exists a High-dimensional space



Inner product:

$$\vec{A} = \sum A_n \hat{q}_n = A_x \hat{i} + A_y \hat{j} + A_z \hat{k}$$

$$\vec{B} = \sum B_n \hat{q}_n = B_x \hat{i} + B_y \hat{j} + B_z \hat{k}$$

$$\begin{aligned}\vec{A} \cdot \vec{B} &= \left(\sum_n A_n \hat{q}_n \right) \cdot \left(\sum_m B_m \hat{q}_m \right) \\ &= A_x B_x + A_y B_y + A_z B_z\end{aligned}$$

[Dot product]

$$\underline{\psi}(x) = \sum c_n \phi_n(x)$$

$$\underline{\phi}(x) = \sum D_n \phi_n(x)$$

$$\oint \int \underline{\psi}^* \underline{\phi}$$

$$= \int \sum_n \sum_n c_n^* \phi_n^*(x) D_n \phi_n(x) dx$$

$$= \sum_n \sum_n c_n^* D_n \delta_{nn}$$

$$= \sum_n c_n^* D_n$$

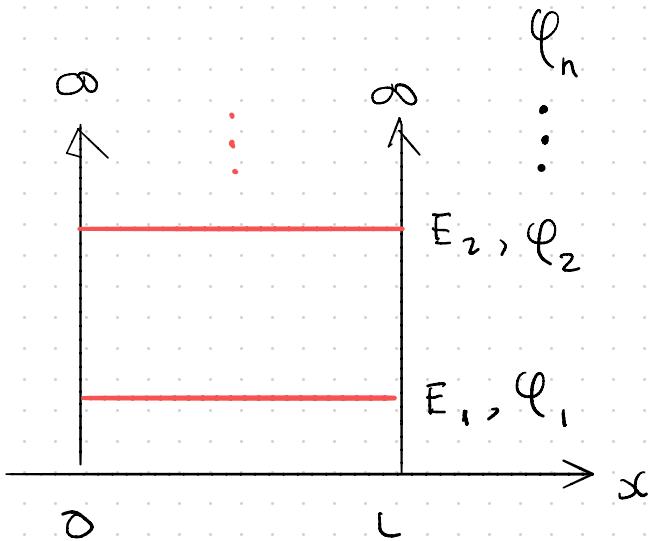
Now, reconsider the time independent Schrödinger equation of infinite square well.

$$\Rightarrow \phi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) \rightarrow \boxed{\text{basis vector}}$$

$$\Rightarrow \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right] \phi(x) = E \phi(x)$$

$\underbrace{\quad}_{\equiv \hat{H}}$

[Aside: Classically $\hat{H} = \frac{p^2}{2m} + V$]



where $\psi_n(x)$ = Eigenfunctions of
the Hamiltonian \hat{H}

