

1. Consider a hypothetical Fermi system with N particles in volume V and with the density of states $g(\varepsilon)$ given by

$$g(\varepsilon) = \begin{cases} 0 & \text{if } \varepsilon < 0 \\ \alpha V & \text{if } \varepsilon > 0 \end{cases},$$

where α is a constant.

- Find the Fermi energy ε_F and the internal energy of the system at zero temperature.
- Using the Sommerfeld expansion, find the chemical potential, the internal energy, and the specific heat at low temperatures.

2. The dimensionality plays a fundamental role in phase transitions and critical phenomena. Its importance also shows in the peculiar phenomenon of Bose-Einstein condensation. Show that there is *no* Bose-Einstein condensation in the *one or two-dimensional* ideal Bose gas with the dispersion relation $\varepsilon(p) = p^2 / 2m$ at any nonzero temperature. (At $T = 0$, of course, all particles are in the ground state.)

3. Find the behavior of the *isotherm compressibility*, κ_T , in as the critical temperature is approached *from above*, i.e., in the $T \rightarrow T_c + 0$ limit. This means that you should obtain an expression for κ_T as a function of $(T - T_c)$ in the vicinity of the transition temperature. You should review the notes posted on LMS regarding the behavior of the Bose gas just above T_c . In particular, use

$$\mu(T) \approx -kT \left(\frac{3\zeta(3/2)}{4\sqrt{\pi}} \right)^2 \left(\frac{T - T_c}{T_c} \right)^2 \approx -kT_c \left(\frac{3\zeta(3/2)}{4\sqrt{\pi}} \right)^2 \left(\frac{T - T_c}{T_c} \right)^2$$

and

$$f_{3/2}^-(e^{-\alpha}) \approx \zeta(3/2) - 2\sqrt{\pi}\alpha^{1/2}, \text{ where } \alpha = -\frac{\mu}{kT}.$$

Also, you will need to employ the basic thermodynamic relation

$$-\left(\frac{\partial V}{\partial P} \right)_{N,T} = \frac{V^2}{N^2} \left(\frac{\partial N}{\partial \mu} \right)_{V,T}.$$

4. Consider the simple model for a one-dimensional solid consisting of N atoms of identical mass m along a chain. Atoms are connected only with nearest-neighbor atoms by “springs” with spring constant κ . The equilibrium separation between atoms (i.e., the lattice constant) is a . The springs are relaxed in equilibrium and we allow only for

longitudinal oscillations. For simplicity, use periodic boundary conditions, $u_{j+N} = u_j$, where u_j is the displacement of the j th atom, measured from its equilibrium position.

(a) Show that the equation of motion for the displacements is

$$m\ddot{u}_j = \kappa(u_{j+1} + u_{j-1} - 2u_j), \quad j = 1, 2, \dots, N.$$

(b) Solve the above set of equation by means of complex Fourier series, yielding the normal modes and the spectrum. In particular, show that the frequency – wave-number dispersion relation is given by

$$\omega(k) = 2\sqrt{\frac{\kappa}{m}} \left| \sin\left(\frac{ka}{2}\right) \right|,$$

where $k = \frac{2\pi}{Na}n$, $n = -\frac{N}{2}, \dots, \frac{N}{2}$ (assumed even N for simplicity).

(c) From the above dispersion relation it follows that for small k values we can use $\omega(k) = ck$, where c is the speed of longitudinal vibrations in this solid. Obtain the low-temperature behavior of the specific heat of this *one-dimensional* solid of size $L = Na$ in the Debye approximation.

5. Obtain an *estimate* for the

(a) *Fermi energy and Fermi temperature in copper* (assume one conduction or “free” electron per atom). For the mass density of copper use $\rho = 9 \frac{\text{g}}{\text{cm}^3}$, and the atomic mass is 63.5g/mol.

(b) critical temperature for the Bose-Einstein condensation in an ideal He^4 “gas” with density $\rho = 0.145 \frac{\text{g}}{\text{cm}^3}$. The atomic mass is 4g/mol.

(c) *Debye temperature in copper* (use the parameters given in (a) and the effective sound velocity $c \approx 4000 \frac{\text{m}}{\text{s}}$).