

$$\{s_i\}_{i=1}^N$$

$$\langle s_i \rangle = m$$

$$\langle (s_i - m)^2 \rangle = \sigma_s^2$$

$\forall i$

and independent

$$\Delta s_i = s_i - \langle s_i \rangle$$

$$\langle (\Delta s_i)^2 \rangle = \sigma_s^2$$

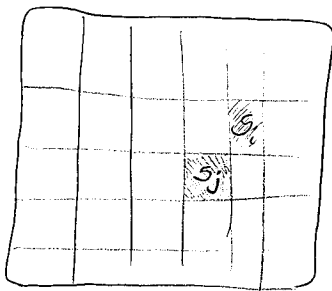
$$M = \sum_{i=1}^N s_i$$

$$1) \langle M \rangle = \left\langle \sum_{i=1}^N s_i \right\rangle = \sum_{i=1}^N \langle s_i \rangle = \sum_{i=1}^N m = Nm = N \langle s_i \rangle$$

$$\begin{aligned} 2) \langle M^2 \rangle &= \langle (M - \langle M \rangle)^2 \rangle = \left\langle \left(\sum_{i=1}^N s_i - \sum_{i=1}^N \langle s_i \rangle \right)^2 \right\rangle = \left\langle \left(\sum_{i=1}^N \underbrace{(s_i - \langle s_i \rangle)}_{\Delta s_i} \right)^2 \right\rangle \\ &= \left\langle \left(\sum_{i=1}^N \Delta s_i \right)^2 \right\rangle = \left\langle \sum_{i=1}^N \Delta s_i \sum_{j=1}^N \Delta s_j \right\rangle = \sum_{i,j} \langle \Delta s_i \Delta s_j \rangle = \\ &= \sum_{i \neq j} \underbrace{\langle \Delta s_i \Delta s_j \rangle}_{\langle \Delta s_i \rangle \langle \Delta s_j \rangle} + \sum_{i=1}^N \langle (\Delta s_i)^2 \rangle = N \langle (\Delta s_i)^2 \rangle = N \sigma_s^2 \\ &\quad \underbrace{\langle \Delta s_i \rangle}_{=0} \underbrace{\langle \Delta s_j \rangle}_{=0} \quad (i \neq j: s_i \text{ and } s_j \text{ are independent}) \end{aligned}$$

$$\sigma_M = \sqrt{N} \sigma_s$$

Stat. Mech implications: N independent subsystem



$i \neq j: \sim \text{independent}$

$$\text{relative fluctuations: } \frac{\sigma_M}{\langle M \rangle} = \frac{\sqrt{N} \sigma_s}{Nm} = \frac{1}{\sqrt{N}} \frac{\sigma_s}{m}$$

$$\propto \frac{1}{\sqrt{N}}$$

Generating function

$p(x)$ (prob. density)

$$\phi(k) = \langle e^{ikx} \rangle = \int_{-\infty}^{\infty} dx e^{ikx} p(x) \quad (\text{Fourier tr.})$$

$$p(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk e^{-ikx} \phi(k) \quad (\text{inverse Fourier tr.})$$

Properties:

$\phi(k)$: moment generating function

$$1) \quad \phi(k) \Big|_{k=0} = \int_{-\infty}^{\infty} dx p(x) = 1 \quad \text{by normalization}$$

$$2) \quad \frac{d\phi}{dk} \Big|_{k=0} = \int_{-\infty}^{\infty} dx (ix) p(x) = i \langle x \rangle$$

$$\frac{d^2\phi}{dk^2} \Big|_{k=0} = \int_{-\infty}^{\infty} dx (ix)^2 p(x) = (i)^2 \langle x^2 \rangle = -\langle x^2 \rangle$$

$$\frac{d^n\phi}{dk^n} \Big|_{k=0} = \int_{-\infty}^{\infty} dx (ix)^n p(x) = i^n \langle x^n \rangle$$

$$\langle x^n \rangle = \frac{1}{i^n} \frac{d^n\phi}{dk^n} \Big|_{k=0}$$

$$\phi(k) = \sum_{n=0}^{\infty} \frac{(ik)^n}{n!} \langle x^n \rangle$$

$$\psi(k) = \ln \phi(k)$$

↖ "free energy" ↗ partition function

$$\psi(k) = \sum_{m=1}^{\infty} \frac{(ik)^m}{m!} C_m$$

$$\psi(k) \Big|_{k=0} = 0$$

$$C_m = \frac{1}{i^m k} \frac{d^m \psi(k)}{dk^m} \Big|_{k=0}$$

$$C_0 = 0$$

$$\left. \frac{d\psi}{dk} \right|_{k=0} = \frac{1}{\phi(k)} \left. \frac{d\phi}{dk} \right|_{k=0} = i \langle x \rangle$$

$$\left. \frac{d^2\psi}{dk^2} \right|_{k=0} = \frac{\frac{d^2\phi}{dk^2} \phi(k) - \left(\frac{d\phi}{dk} \right)^2}{\phi(k)} \Big|_{k=0} = i^2 \langle x^2 \rangle - (i \langle x \rangle)^2 = -(\langle x^2 \rangle - \langle x \rangle^2)$$

$$\psi(k) = \ln \phi(k) = \sum_{n=1}^{\infty} \frac{(ik)^n}{n!} C_n$$

$$C_n = \frac{1}{i^n} \left. \frac{d^n \psi}{dk^n} \right|_{k=0}$$

$$C_1 = \langle x \rangle$$

$$C_2 = \langle x^2 \rangle - \langle x \rangle^2 = \sigma^2$$

$$C_3 = \langle x^3 \rangle - 3 \langle x^2 \rangle \langle x \rangle + 2 \langle x \rangle^3$$

$$\vdots$$

$$C_n = \dots$$

The Central Limit Theorem

assume $\{x_i\}_{i=1}^N$ identically distributed, independent variables
and $\sigma = \text{finite}$ $\langle x_i \rangle = \langle x \rangle$
 $\sigma_i = \sigma \quad \forall (i.d.)$

Then, for $z = \frac{\sum_{i=1}^N x_i - N\langle x \rangle}{\sqrt{N}\sigma}$

$$\phi(z) = \frac{1}{\sqrt{2\pi}} e^{-z^2/2}$$

Proof: for simplicity: $\langle x_i \rangle = 0$ (w.l.g.)

(if not we could consider $x_i' = x_i - \langle x \rangle$ $\langle x_i' \rangle = 0$)

$$z = \frac{\sum x_i}{\sqrt{N}\sigma}$$

(normalizes $\sigma_z^2 = 1$)

$$\begin{aligned} \phi_z(k) &= \langle e^{ikz} \rangle = \left\langle e^{ik \frac{\sum x_i}{\sqrt{N}\sigma}} \right\rangle = \left\langle \prod_{i=1}^N e^{i \frac{k}{\sqrt{N}\sigma} x_i} \right\rangle = \prod_{i=1}^N \langle e^{i \frac{k}{\sqrt{N}\sigma} x_i} \rangle \\ &= \prod_{i=1}^N \phi\left(\frac{k}{\sqrt{N}\sigma}\right) = \phi^N\left(\frac{k}{\sqrt{N}\sigma}\right) \end{aligned}$$

$$\psi_z(k) = \ln \phi^N\left(\frac{k}{\sqrt{N}\sigma}\right) = N \ln \phi\left(\frac{k}{\sqrt{N}\sigma}\right) = N \psi\left(\frac{k}{\sqrt{N}\sigma}\right)$$

$$= N \sum_{n=1}^{\infty} \frac{(i \frac{k}{\sqrt{N}\sigma})^n}{n!} C_n = N \sum_{n=1}^{\infty} \frac{(ik)^n}{n!} (N\sigma^2)^{-n/2} C_n$$

$$= \sum_{n=1}^{\infty} \frac{(ik)^n}{n!} \underbrace{\sigma^{-n} N^{1-n/2}}_{C_n^z} C_n$$

$$C_n^z = \frac{C_n}{\sigma^n N^{n/2-1}}$$

for z : $C_1^z = \langle z \rangle = \frac{C_1}{\sigma/\sqrt{N}} = \frac{\langle x \rangle}{\sigma/\sqrt{N}} = 0 \quad (\text{all } N)$

$$C_2^z = \sigma_z^2 = \frac{1}{\sigma^2} C_2 = \frac{1}{\sigma^2} \sigma^2 = 1$$

$$C_n^z = \frac{C_n}{\sigma^n N^{n/2-1}} \xrightarrow{N \rightarrow \infty} 0$$

for $n > 2$
 $n=3: o\left(\frac{1}{\sqrt{N}}\right)$

Thus, for $N \rightarrow \infty$

$$\psi_2(k) = -\frac{k^2}{2}$$

$$\psi_2(k) = \ln \phi_2(k) \Rightarrow \phi_2(k) = e^{-\frac{k^2}{2}}$$

$$\rho(z) = \frac{1}{2\pi i} \int_{-\infty}^{+\infty} dk e^{-ikz} \phi_2(k) = \frac{1}{2\pi i} \int_{-\infty}^{+\infty} dk e^{-\frac{k^2}{2} - ikz} = \frac{1}{\sqrt{2\pi i}} e^{-\frac{z^2}{2}}$$

Details

$$-\frac{k^2}{2} - ikz = -\frac{1}{2}(k+iz)^2 - \frac{z^2}{2}$$

$$\Rightarrow \frac{1}{2\pi i} \int_{-\infty}^{+\infty} dk e^{-\frac{1}{2}(k+iz)^2 - \frac{z^2}{2}} = \frac{1}{2\pi i} e^{-\frac{z^2}{2}} \int_{-\infty}^{+\infty} dk e^{-\frac{1}{2}(k+iz)^2} =$$

$$= \frac{1}{2\pi i} e^{-\frac{z^2}{2}} \sqrt{2\pi i} = \frac{1}{\sqrt{2\pi i}} e^{-\frac{z^2}{2}}$$