

D₈

-1-

Eret Zohar, 6.3.22

$$D_8 = \{g\} = \{a^p x^q \mid p=0, \dots, 3; q=0, 1\}$$

$$|D_8| = 8$$

$$a \mapsto \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = i\sigma_3, \quad x \mapsto \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \sigma_2 \rightarrow \underbrace{\text{Nontrivial rep.}}_{j=2}.$$

$\left(\frac{\pi}{2} \text{ rotation}\right) \quad (\text{reflection})$

Trivial representation, $j=0$: $D^0(g) = 1 \quad \forall g$

First Sign representation, $j=\bar{0}$: $D^{\bar{0}}(g) = \det(D^2(g)) =$

$$= \det(a)^p \det(x)^q = 1^p (-1)^q = (-1)^q.$$

Second Sign representation: $D^1(g) = (-1)^p$

Third Sign representation: $D^{\bar{1}}(g) = (-1)^{p+q}$

$$|D_8| = \sum_j \dim(j) = 2^2 + 1 \cdot 1^2 = 8 \quad \checkmark$$

nontrivial

-1 is contained in any \checkmark representation.

$$\text{In particular } D^2(a^2) = -1_{2 \times 2}$$

From now on, we use $D(g) \equiv D^2(g)$.

Group Element States

$$|g\rangle = |p, q\rangle$$

Group Multiplication Rules:

$$\text{?} \quad (p, q) \cdot (p', q') =$$

$$= a^p x^q a^{p'} x^{q'}$$

$$xa = i\sigma_y \sigma_z = -i\sigma_z \sigma_y = -ax = a^3 x \quad (a^2 = -1)$$

~~$x^q a^{p'} = \cancel{x} \quad \cancel{a}$~~

$$\underline{q=0}: x^q a^{p'} = a^{p'} x^q$$

$$\underline{q=1}: x^q a^{p'} = a^3 x^q a^{p'-1} = a^6 x^q a^{p'-2} =$$

$$= \dots = a^{3p'} x^q$$

$$\Rightarrow x^q a^{p'} = a^{2p'q} a^{p'} x^q =$$

$$\Rightarrow (p, q) \cdot (p', q') = a^p x^q a^{p'} x^{q'} =$$

$$= a^{p+p'+2p'q} x^{q+q'}$$

$$\Rightarrow \boxed{(p, q) \cdot (p', q') = (p+p'+2p'q, q+q')}$$

$$\underline{\text{Inversion}}: (p, q)^{-1} = (a^p x^q)^{-1} = (x^{-1})^q (a^{-1})^p$$

$$x^{-1} = x$$

$$a^{-1} = a^3$$

$$= x^q a^{3p} = a^{3p+6p'q} x^q$$

~~$= \cancel{a^{9p'}} \cancel{x^q} =$~~

$$\Rightarrow \text{If } q=0: (p, 0)^{-1} = a^{3p} \quad [a^3 = 1] = \cancel{a^{3p}} \cancel{a^p x^q} = \cancel{a^p x^q}$$

$$\text{If } q=1: (p, 1)^{-1} = a^q p x^q = a^p x^q = (p, 1) \quad [a^3 p = 1]$$

Left multiplication:

$$\begin{aligned}\tilde{\Theta}(p', q') |p, q\rangle &= |(p', q')^{-1} \cdot (p, q)\rangle \\ &= |(3p' + 6p'q', q') \cdot (p, q)\rangle = \\ &= |(3p' + 6p'q' + p, + \cancel{2p(p'+q'q')} , q + q')\rangle \\ &\quad + 2pq'\end{aligned}$$

Right multiplication:

$$\begin{aligned}\Theta(p', q') |p, q\rangle &= |(p, q) \cdot (p', q')^{-1}\rangle \\ &= |(p + 3p' + 6p'q', q + 2q(p' + 6p'q'))\rangle\end{aligned}$$

To replace the fermions (2 comp. spinors)
by hard-core bosons $\{\gamma_m^+\}_{m=1,2}$,

We need to find the operator P on
each link, such that:

$$PUP = -U$$

P should correspond to either left
or right multiplication by $a^2 = -1$
(doesn't matter which).

For that, let's define the rep. basis states:

$$|0\rangle, |\bar{0}\rangle, |1\rangle, |\bar{1}\rangle, \{|_{2mn}\rangle\}_{m,n=1,2}$$

~~$\langle j'm'|\rho|jmn\rangle = \sum_g$~~

$$\langle j'm'n'|P|jmn\rangle = \sum_g \underbrace{\langle j'm'n'|g'Xg|P|gXg|jmn\rangle}_{\delta g, a^2 g} \\ \stackrel{?}{=} \frac{\text{dim}(j)}{|D_8|} \sum_g \langle j'm'n'|a^2 g X g|jmn\rangle$$

$$= \frac{\text{dim}(j)}{|D_8|} \sum_g \bar{D}_{m'n'}^{j'}(a^2 g) D_{mn}^j(g) =$$

~~$= \frac{\text{dim}(j)}{|D_8|} \bar{D}$~~

$$D^{j'}(a^2) = \begin{cases} 1 & \text{for all } j \neq 2 \\ -1 & \text{for } j=2 \end{cases}$$

$$= \frac{\text{dim}(j)}{|D_8|} D^{j'}(a^2) \sum_g \bar{D}_{m'n'}^{j'}(g) D_{mn}^j(g)$$

Or. orth.

$$\stackrel{j}{=} \frac{\text{dim}(j)}{|D_8|} D^{j'}(a^2) \frac{|D_8|}{\text{dim}(j)} \delta_{j'j} \delta_{m'm} \delta_{n'n}$$

Then,

$$= D^{j'}(a^2) \delta_{j'j} \delta_{m'm} \delta_{n'n}$$

$$\Rightarrow P = \sum_j D^j(a^2) \langle jmn | jmn \rangle$$

$$= \underbrace{\prod_{j=1}^2}_{\substack{\text{proj. onto the } j \neq 2 \\ (\text{id}) \text{ irrep}}} - \underbrace{\prod_{j=2}^2}_{\substack{\text{proj. onto the } j=2 \\ (\text{id}) \text{ irrep}}} \rightarrow \text{proj onto the } j=2 \text{ irrep}$$

\Rightarrow The transformed gauge-matter interaction:

$$\tilde{H}_{\text{int}} = \epsilon \sum_x P(x-1) \eta_m^+(x) U_{mn}(x) \eta_n(x+1) + \text{h.c.}$$

Now we proceed and remove the fermions altogether! For that, we need two indep. Gauss laws that will allow us to solve for $n_1 = \eta_1^+ \eta_1$ and $n_2 = \eta_2^+ \eta_2$ separately.

~~We can~~

The group is finite, so we can write Gauss laws only in terms of finite transformations, and we have 8 such:

$$\forall g \in D_8, \quad \tilde{\hat{O}}_g(x) \hat{O}_g^+(x+1) |\psi\rangle = \hat{O}_g(x) |\psi\rangle$$

Where $\hat{O}_g(x)$ is the matter's transformation;

$$\hat{O}_g(x) = e^{i \eta_m^+(x) q_g} \eta_m(x) \quad (\text{if we don't stagger which is pointless for a finite group}).$$

$$q_g = -i \log(D^{j^2}(g))$$

(Zohar, Burello PRD 2015)

To Date

We need two commuting, independent $g \in D_8 \rightarrow$ pick α_x^2 and x .

$$q_{\alpha^2 x} = -i \log \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} = -i \begin{pmatrix} i\pi & 0 \\ 0 & 0 \end{pmatrix} = \pi \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

$$q_x = -i \log \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = -i \begin{pmatrix} 0 & 0 \\ 0 & i\pi \end{pmatrix} = \pi \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\Rightarrow \hat{\theta}_{\alpha^2 x} = e^{i\pi(\eta_1^+ \eta_1^- + \eta_2^+ \eta_2^-)} = e^{i\pi n_1} \cancel{e^{i\pi n_2}} = (1-2n_1) \cancel{(1-2n_2)}$$

$$\left(e^{i\pi n} = 1 + \sum_{k=1}^{\infty} \frac{(i\pi)^k}{k!} n = 1 + (e^{i\pi} - 1)n = 1 - 2n \right)$$

$\eta^k = n \quad \forall k \in \mathbb{N}$

$$\hat{\theta}_x = e^{i\pi n_2} = (1-2n_2)$$

Define: $G_1^{(x)} = \tilde{\theta}_{\alpha^2}(x) \hat{\theta}_{\alpha^2}^+(x-1) \cancel{\tilde{\theta}_x^+}$

$$G_2^{(x)} = \tilde{\theta}_x^-(x) \hat{\theta}_x^+(x-1)$$

We have:

$$G_2^{(x)} |\psi\rangle = \hat{\theta}_x(x) |\psi\rangle = (1-2n_2^{(x)}) |\psi\rangle$$

$$\Rightarrow n_2(x) |\psi\rangle = \underbrace{\left(\frac{1-G_2(x)}{2} \right)}_{\equiv G_2(x)} |\psi\rangle$$

Note that $G_2^2(x) = (\tilde{\theta}_x^-(x) \hat{\theta}_x^+(x-1))^2$

$$= \tilde{\theta}_x^{-2}(x) \hat{\theta}_x^{+2}(x-1) = 1$$

Hence G_2 is a projector with spectrum 0,1 and we need no constraint for it →

$$\text{Similarly, } \Omega_1(x)|\psi\rangle = (1 - 2n_f)|\psi\rangle$$

$$\Rightarrow n_i(x)|\psi\rangle = \frac{1 - \Omega_1(x)}{2}|\psi\rangle$$

$$= G_i(x)|\psi\rangle$$

Note that $\Omega_i^2(x) = 1$

and thus $G_i(x)$ are projectors, ~~and~~
implying that

$G_i(x)(G_i(x) - 1) = 0 \quad \forall i, x$
holds ~~as~~ as an operator equation,
and we don't need to impose these
local conditions!

$$[G_i(x), \Omega_{mn}(x)] =$$

$$\Rightarrow H_{\text{int}} = \epsilon \sum_{x, i, m, n} \cancel{P(x-1)} \hat{P}_+^m(x) \cup_{mn} \hat{P}_+^n(x) + \text{h.c.}$$

$\hat{P}_+^m(x)$ is the projector
onto $\cancel{n_m(x)} = 1$.

AND THAT'S IT.

~~$\langle \text{KMN} | 10^j \rangle = \langle \text{JLMN} | 10^j \rangle$~~

Clebsch Gordan coefficients.

~~8~~

$$\begin{aligned} |0\rangle \otimes |0\rangle &= |0\rangle \rightarrow \langle 0|00\rangle = 1 \\ |\bar{0}\rangle \otimes |\bar{0}\rangle &= |\bar{0}\rangle \rightarrow \langle \bar{0}|\bar{0}\bar{0}\rangle = \langle \bar{0}|\bar{0}0\rangle = 1 \end{aligned}$$

|

$$\left\{ \begin{array}{l} D^0(g) D^0(g) = D^0(g) \Rightarrow \langle 0|00\rangle = 1 \\ D^0(g) D^{\bar{0}}(g) = D^{\bar{0}}(g) \Rightarrow \langle \bar{0}|0\bar{0}\rangle = \langle \bar{0}|\bar{0}0\rangle = 1 \\ D^0(g) D^1(g) = D^1(g) \Rightarrow \langle 1|10\rangle = \langle 1|01\rangle = 1 \\ D^0(g) D^{\bar{1}}(g) = D^{\bar{1}}(g) \Rightarrow \langle \bar{1}|\bar{1}0\rangle = \langle \bar{1}|0\bar{1}\rangle = 1 \end{array} \right.$$

$$\left\{ \begin{array}{l} D^{\bar{0}}(g) D^{\bar{0}}(g) = (-1)^{2q} = D^0(g) \Rightarrow \langle 0|\bar{0}\bar{0}\rangle = 1 \\ D^{\bar{0}}(g) D^{\bar{1}}(g) = (-1)^{p+q} = D^{\bar{1}}(g) \Rightarrow \langle \bar{1}|0\bar{1}\rangle \\ \quad = \langle \bar{1}|\bar{1}0\rangle = 1 \\ D^{\bar{0}}(g) D^{\bar{1}}(g) = (-1)^{p+2q} = D^1(g) \\ \quad \Rightarrow \langle 1|\bar{0}\bar{1}\rangle = \langle 1|\bar{1}0\rangle = 0 \end{array} \right.$$

$$\left\{ \begin{array}{l} D^1(g) D^{\bar{1}}(g) = (-1)^{2p+q} = D^{\bar{0}}(g) \\ \quad \Rightarrow \langle \bar{0}|1\bar{1}\rangle = \langle \bar{0}|\bar{1}1\rangle = 1 \end{array} \right.$$

- 9 -

$$D^0(g) D^2_{mn}(g) = D^2_{mn}(g) \Rightarrow \cancel{\langle 2mn | 2n \rangle}$$

$$\Rightarrow \langle 0, 2m | 2n \rangle = \langle 2m, 0 | 2n \rangle = \delta_{mn} \quad \begin{aligned} &\cancel{\langle 2mn | 0, 2m \rangle} \\ &= \cancel{\langle 2m n | 2m n, 0 \rangle} = 1 \end{aligned}$$

~~$D^0(g) D^2_{mn}(g) = (-1)^q D^2_{mn}(g) =$~~

~~C.G. Series~~C-G series:

~~$D^0(g)$~~ $D^{j_1(g)}_{m_1 m_1}(g) D^{j_2(g)}_{m_2 m_2} = \sum_j \langle j_1 m_1, j_2 m_2 | JM \rangle D^J_{M M'}(g) \quad \begin{aligned} &\cancel{\langle j_1 m_1 | JM' | j_2 m_2 \rangle} \\ &= \cancel{\langle JM' | j_1 m_1, j_2 m_2 \rangle} \end{aligned}$

$$D^0(g) D^2_{mn}(g) = \langle 0, 2mn | 2m'n' \rangle D^2_{m'n'}(g)$$

~~$(-1)^q A^P \times g = A^{2q+p} \times g$~~

$$(-1)^q D^2(g) = D^2(\alpha^{2q}) D^2(g) =$$

~~$= D^2(\alpha^{2q}) \alpha^P \times g$~~
 ~~$= D^2 \alpha^P$~~

 $(-D)$

$$A \in D^2(a)$$

$$X \in D^2(x)$$

$$\begin{aligned} (-1)^q A^P X^q &= \cancel{A^P} \cancel{(X^q)} \\ \det X &= A^P (\det X \cdot X^{-1})^q = \\ &\stackrel{x=x^{-1}}{=} A^P (\det X \cdot X^{-1})^q = \\ &= A^P (\epsilon X^T \epsilon^T)^q \\ &= \cancel{A^P} \stackrel{x=x^T}{=} A^P (\epsilon X \epsilon^T)^q \\ &= A^P \epsilon X^q \epsilon^T = \\ &\stackrel{?}{=} \end{aligned}$$

$$\epsilon = -i \sigma_y = -A \Rightarrow [A^P, \epsilon] = 0$$

$$\Rightarrow (-1)^q A^P X^q = \epsilon A^P X^q \epsilon^T$$

$$\Rightarrow \cancel{D^0(g)} D^2(g) D_{m'm'}^2(g) = \epsilon_{mm'} D_{m'm'}^2(g) \epsilon_{n'n}^T = \\ = \epsilon_{mm'} \epsilon_{n'n} D_{m'm'}^2(g)$$

$$\stackrel{?}{=} \langle \bar{0}, 2m | 2m' \rangle D_{m'm'}^2(g)$$

C.G. Serig $\langle \bar{0}, 2n | 2n' \rangle$

$$\Rightarrow \boxed{\langle \bar{0}, 2mn | 2m'n' \rangle = \epsilon_{mm'} \epsilon_{nn'}}$$

$$\boxed{\langle 2mn, \bar{0} | 2m'n' \rangle = \epsilon_{mm'} \epsilon_{nn'}}$$

$$\boxed{\langle \bar{0}, 2m | 2m' \rangle = \langle 2m, \bar{0} | 2m' \rangle} = \epsilon_{mm'} = -(\epsilon_y)_{mm'} (\epsilon_y)_{nn'}$$

$$D^1(g) D_{mn}^2(g) = (-1)^P A^P X^q = (-A)^P X^q$$

$$\begin{aligned} &= \cancel{(-A)^P X^q} = \cancel{(\bar{A}^T)^P X^q} \\ &= \cancel{(\det A \bar{\epsilon}^T A^{-1} \bar{\epsilon})^P X^q} = \\ &= -A = A^{-1} = XAX \end{aligned}$$

$$D^1(g) D_{mn}^2(g) = (XAX)^P X^q =$$

$$\begin{aligned} &= \cancel{X} \\ &= X A^P X^{q+1} = \\ &= X A^P X^q X = (X D^2(g) X)_{mn} \end{aligned}$$

$$\begin{aligned} D^1(g) D_{mn}^2(g) &= X_{mm'} D_{m'n'}^2(g) X_{n'n}^T = \Rightarrow \langle 1, 2m | 2m' \rangle \\ &= X_{mm'} X_{nn'} D_{m'n'}^2(g) \quad = \langle 2m, 1 | 2n' \rangle = (\Gamma_2)_{mn} \end{aligned}$$

$$\Rightarrow \boxed{\begin{aligned} &\langle 1, 2m | 2m' \rangle \rightarrow X_{mm'} X_{m'm} = (\Gamma_2)_{mm'} (\Gamma_2)_{m'm} \\ &\langle 2m, 1 | 2m' \rangle = X_{mm'} X_{nn'} = \end{aligned}}$$

$$\begin{aligned} D^1(g) D_{mn}^2(g) &= (-1)^{P+q} D_{mn}^2(g) = \\ &= (-1)^P \epsilon_{mm'} \epsilon_{nn'} D_{m'n'}^2(g) \\ &= \epsilon_{mm'} \epsilon_{nn'} \cancel{X_{m'm''} X_{n'n''}} D_{m'n''}^2(g) \\ &= (\epsilon X)_{mm'} (\epsilon X)_{nn'} D_{m'n'}^2(g) \end{aligned}$$

$$\Rightarrow \boxed{\begin{aligned} &\epsilon X = \cancel{i\Gamma_y \Gamma_z} \quad i\Gamma_y \Gamma_z = \Gamma_x \Rightarrow \langle 1, 2m | 2m' \rangle = \langle 2m, 1 | 2m' \rangle = (\Gamma_x)_{mm'} \\ &\langle 1, 2m | 2m' \rangle = \langle 2m, 1 | 2m' \rangle = (\Gamma_x)_{mm'} (\Gamma_x)_{nn'} \end{aligned}}$$

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-12-

$$D_{mn}^2(g) D_{m'n'}^2(g) = \underbrace{(A \otimes A)^P}_{\downarrow} (X \otimes X)^q$$

$$[A \otimes A, X \otimes X] = 0 \quad \cancel{\Rightarrow}$$

$$(A \otimes A)^2 = (X \otimes X)^2 = 1_4 \quad \Rightarrow \text{Forming a } \mathbb{Z}_2 \times \mathbb{Z}_2 \text{ group}$$

Only Abelian irreps
are included.

$$A \otimes A = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \otimes \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = -\sigma_y \otimes \sigma_y$$

$$\sigma_y = U \sigma_z U^\dagger \quad U = e^{i \frac{\pi}{4} \sigma_x} \quad U = U \otimes U$$

$$A \otimes A = U (-\sigma_z \otimes \sigma_z) U^\dagger$$

$$(A \otimes A)^P = (-1)^P U (\sigma_z \otimes \sigma_z)^P U^\dagger$$

$$X \otimes X = \sigma_z \otimes \sigma_z$$

$$D^2 \otimes D^2 = (-1)^P U (\sigma_z \otimes \sigma_z)^P U^\dagger (\sigma_z \otimes \sigma_z)^q$$

$$Q = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 & 0 & 1 \\ 0 & 1 & 1 & 0 \\ 0 & 1 & -1 & 0 \\ 1 & 0 & 0 & -1 \end{pmatrix} \quad Q^T Q = 1 \quad (\text{Orthogonal})$$

$$Q (A \otimes A) Q^T = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

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-13-

$$Q(X \otimes X) Q^T = \begin{pmatrix} 1 & & & \\ & -1 & & \\ & & -1 & \\ & & & 1 \end{pmatrix}$$

$$D^2(g) D^2(g) = (A \otimes A)^P (A \otimes A)^Q =$$

$$= Q \begin{pmatrix} 1 & & & \\ & (-1)^{P+Q} & & \\ & & (-1)^Q & \\ & & & (-1)^P \end{pmatrix} Q^T$$

The four 1d irreps.

Ordered as: 0, 1, 0, 1

$$\Rightarrow \cancel{D^2(g)} \quad \cancel{\otimes \otimes}$$

$$D_{mn}^2(g) D_{m'n'}^2(g) = \sum_{\alpha} Q_{(mn), \alpha} \cdot Q_{(m'n'), \alpha} D^{\alpha}(g)$$

↓
Direct sum

Q-Clebsch-Gordan coefficients!

$$Q_{mn'}^{\alpha} = \langle \alpha | \cancel{2m m'} \rangle \quad 2m, 2m'$$

$$\cancel{\langle \alpha | 21, 21 \rangle} = \cancel{\frac{1}{\sqrt{2}}}$$

$\alpha = 0, 1, 0, \cancel{1}$

Order $(m, m') = (1, 1), (1, 2), (2, 1), (2, 2)$

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First row:

$$\langle 0 | 2m, 2m' \rangle = Q_{mm'}^0 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 & 0 & 1 \end{pmatrix}$$

Second row:

$$\langle 1 | 2m, 2m' \rangle = Q_{mm'}^1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 1 & 0 \end{pmatrix}$$

Third row:

$$\langle \bar{0} | 2m, 2m' \rangle = Q_{mm'}^{\bar{0}} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 1 & -1 & 0 \end{pmatrix}$$

Fourth row:

$$\langle \bar{1} | 2m, 2m' \rangle = Q_{mm'}^{\bar{1}} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 & 0 & -1 \end{pmatrix}$$

All C-G coefficients of D_8 :

$$\langle 0 | 00 \rangle = 1$$

$$\langle \bar{0} | \bar{0}0 \rangle = \langle \bar{0} | 0\bar{0} \rangle = 1$$

$$\langle 1 | 10 \rangle = \langle 1 | 01 \rangle = 1$$

$$\langle \bar{1} | \bar{1}0 \rangle = \langle \bar{1} | 0\bar{1} \rangle = 1$$

$$\langle 0 | 0\bar{0} \rangle = 1$$

$$\langle \bar{1} | 0\bar{1} \rangle = \langle \bar{1} | \bar{1}0 \rangle = 1$$

$$\langle 1 | \bar{0}\bar{1} \rangle = \langle 1 | \bar{1}\bar{0} \rangle = 1$$

$$\langle \bar{0} | 1\bar{1} \rangle = \langle \bar{0} | \bar{1}1 \rangle = 1$$

$$\langle \bar{0} | 2m, 2m' \rangle = \langle 2m, \bar{0} | 2m' \rangle = \delta_{mm'}$$

$$\langle 1 | 2m, 2m' \rangle = \langle 2m, 1 | 2m' \rangle = (\hat{O}_z)_{mm'}$$

$$\langle 0 | 2m, 2m' \rangle = \langle 2m, 0 | 2m' \rangle = \delta_{mm'}$$

$$\langle \bar{1} | 2m, 2m' \rangle = \langle 2m, \bar{1} | 2m' \rangle = (\hat{O}_x)_{mm'}$$

The rest - given on top of this page, or zero.