# PHYS 351 #9

Parker Whaley

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## 1

In class we derived the first TdS equation,  $TdS = C_V dT + T \frac{\partial P}{\partial T}|_V dV$ . Let v be the molar volume and let s be the molar entropy, that is, v = V/n and s = S/n. Let  $c_V$  be the molar specific heat,  $c_V = C_V/n$ . Then the first TdS equation can be written for molar specific quantities as  $Tds = c_V dT + T \frac{\partial P}{\partial T}|_V dv$ .

#### 1.1

The Ideal Gas equation of state is Pv = RT. Assume that  $c_V$  is constant and integrate the equation of state to show

$$s = s_0 + c_V \ln(\frac{T}{T_0}) + R \ln(\frac{v}{v_0})$$

Note that  $\frac{\partial P}{\partial T}|_V = R/v$  so  $Tds = c_V dT + T \frac{\partial P}{\partial T}|_V dv$  becomes  $ds = c_V/TdT + R/vdv$ . We immediately see that if we take the integral we get  $s - s_o = c_V (\ln(T) - \ln(T_o)) + R(\ln(v) - \ln(v_o)) \Rightarrow s = s_0 + c_V \ln(\frac{T}{T_o}) + R \ln(\frac{v}{v_o})$ .

### 1.2

The van der Waals equation of state is  $(P+a/v^2)(v-b) = RT$ . Assume that  $c_v$  is constant and integrate the equation of state to show

$$s = s_0 + c_V \ln(T/T_0) + R \ln(\frac{v-b}{v_0 - b})$$

Proceeding in the same fashion.  $P = \frac{RT}{(v-b)} - a/v^2$  so  $\frac{\partial P}{\partial T}|_V = \frac{R}{(v-b)}$  and  $ds = c_V/TdT + \frac{R}{(v-b)}dv$ . Integrating both sides and adding  $s_0$  we get  $s = s_0 + c_V \ln(T/T_0) + R \ln(\frac{v-b}{v_0-b})$ .

## $\mathbf{2}$

Consider one mole of a solid or a liquid under hydrostatic pressure; as above, v is the molar volume.

### 2.1

Develop an equation of state by expanding dv in terms of dT and dP, and then using the definitions of compressibility and expansion coefficients,  $\beta$  and  $\kappa$ . Show that, if we assume that  $\beta$  and  $\kappa$  are constant, one obtains

$$v = v_0[1 + \beta(T - T_0) - \kappa(P - P_0)]$$

for small changes in v.

We can easily get a expression for v assuming v(T,P) in the usual way  $dv = 1/n\frac{\partial V}{\partial T}|_P dT + 1/n\frac{\partial V}{\partial P}|_T dP$ . Note using the definitions of  $\kappa$  and  $\beta$  that  $dv/v = \beta dT - \kappa dP$ . So assuming those two are constant we see  $\ln(\frac{v-v_0}{v_0}+1) = \beta(T-T_0) - \kappa(P-P_0)$ . Noting that  $\frac{v-v_0}{v_0}$  is a small quantity (we also must note that it is unit-less) we may use a Taylor expansion on  $\ln(x+1)$  and take the first term x. Thus we get  $\frac{v-v_0}{v_0} = \beta(T-T_0) - \kappa(P-P_0)$  or  $v = v_0(1+\beta(T-T_0) - \kappa(P-P_0))$ 

## 2.2

In class we derived the second TdS equation,  $TdS = C_P dT - T(\frac{\partial V}{\partial T})|PdP$ ; in terms of molar specific quantities this is  $Tds = c_p dT - T(\frac{\partial v}{\partial T})|PdP$ . Now use the second TdS equation to find an expression for the change in molar entropy as T and P change. (You should assume that  $c_p$  is constant and, as above, that  $\beta$  and  $\kappa$  are constant and that changes in v are small.)

We note that the work from the above section will still be valid here and thus  $(\frac{\partial v}{\partial T})|P = v_0 \kappa$ . so our equation for entropy becomes  $ds = c_p/T dT - v_0 \kappa dP$  and  $s = s_0 + \ln(T/T_0) - v_0 \kappa (P - P_0)$ .