

# PHYS 472L #13#18

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## 1 #13

### 1.1 Newtonian

In a Newtonian system we see that momentum must be conserved so if in the instantaneous rest frame the rocket ejects a mass  $dm$  with a momentum  $u dm$  it must adopt the momentum  $dp = -u dm$ . in other words  $mdv = -u dm$  or

$$v = \int_{m_0}^m -u/m dm = u * \ln(m) \Big|_{m_0}^m = u * \ln(m_0/m)$$

### 1.2 relativistic

(take all interactions to be along the  $\hat{x}$  axis)

In relativistic mechanics in the instantaneous rest frame of the rocket the rocket ejects the fuel of mass  $\delta m$  with 4-velocity  $(\gamma', \gamma' u)$  where  $u$  is the velocity of the fuel and 4-momentum  $\delta m(\gamma', \gamma' u)$ . Consider that the rocket must have started at rest in this frame and then has some velocity. We see by conservation of relativistic momentum  $(M, 0) = M'(\gamma, \gamma dv) + \delta m(\gamma', \gamma' u)$ . Let's calculate  $dv$ ,  $M' \gamma dv = -\delta m \gamma' u \Rightarrow \delta m = -M' \frac{\gamma dv}{\gamma' u}$ . Now  $M = M' \gamma + \delta m \gamma'$  becomes  $M = M' \gamma - M' \frac{\gamma dv}{\gamma' u} \gamma' \Rightarrow M = M' \gamma (1 - \frac{dv}{u}) \Rightarrow M \frac{\sqrt{1-dv^2}}{1-dv/u} = M'$  doing a Taylor expansion and eliminating higher order terms in  $dv$  we end up at  $dv = \frac{u}{M} dM$ .

In our lab frame this  $dv$  causes a change in the rockets velocity  $v' + dv' = \frac{v' + dv}{1 + v' dv} = (v' + dv) * (1 - v' dv + O(dv^2))$  with linearisation yields:  $dv' = (1 - v'^2) \frac{u}{M} dM$ . Now we need only integrate over our change in velocity

$$\int_{M_0}^M \frac{u}{M} dM = \int_0^{v_f} \frac{1}{(1 - v'^2)} dv'$$

After calculus

$$v_f = \tanh(-u \ln(M_0/M))$$

### 1.3 relativistic mass conservation

Consider a frame that the rocket is at rest in that watches the rocket move in the  $\hat{x}$  for a bit and sees all of the ejected fuel moving in the  $-\hat{x}$  direction. We immediately note that there is more kinetic energy then there was initially so by conservation of the first terms of the four momentums -the terms that are equivalent to  $M+T$  we see that  $\sum M = \sum M' + \sum T'$  since some  $T'$  are not 0 the sum of masses must have changed. Mass can not be conserved.

## 2 #18

In all frames  $P \bullet P = (P^1)^2 + p^2$  in this lab the energy of the particle  $P^1$  is given by the first quantity of the particles 4-momentum, witch can be obtained in a frame independent way by taking  $u \bullet P = E = P_{lab}^1$  so we can now put this in for the energy we get  $P \bullet P = (u \bullet P)^2 + p^2$  so  $p = \sqrt{P \bullet P - (u \bullet P)^2}$ .