

Time & Freq. responses of 1st-order lowpass and highpass filters (Linear systems 1)

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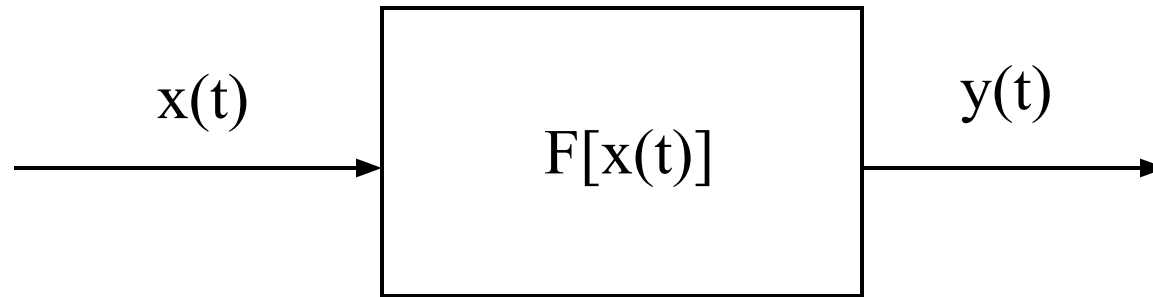
- Linear systems
 - Linear time-invariant systems
 - Convolution and impulse response
 - Complex numbers and transfer functions
- RC integrators and differentiators
 - Frequency analysis
- VLSI Integrators
 - Follower integrator
 - Translinear principle
 - Log-domain current-mode integrator

Linear Systems, Small Signal Analysis, and Integrators

(WS07-08)

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Linear Systems (I)



If a linear time-invariant system has not internally stored energy, then the output of the system $y(t)$ is the forced response due entirely to the input $x(t)$: $y(t)=F[x(t)]$.

A system is said to be linear if it obeys the two fundamental principles of **homogeneity** and **additivity**.

Homogeneity and Additivity

Homogeneity If

$$\hat{x}(t) = \alpha x(t)$$

then

$$\hat{y}(t) = F[\alpha x(t)] = \alpha F[x(t)]$$

Additivity If

$$x(t) = \sum_k a_k x_k(t)$$

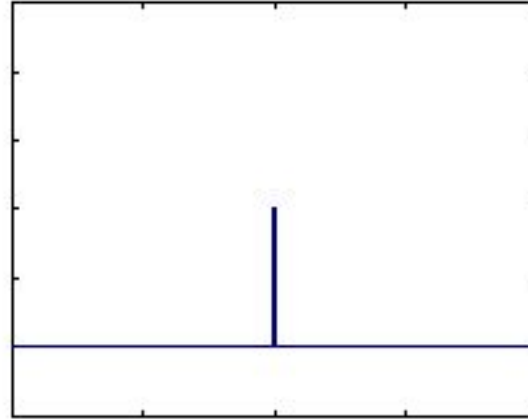
with a_k constant $\forall k$

then

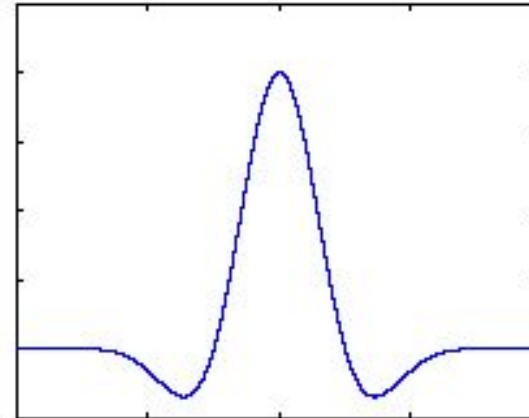
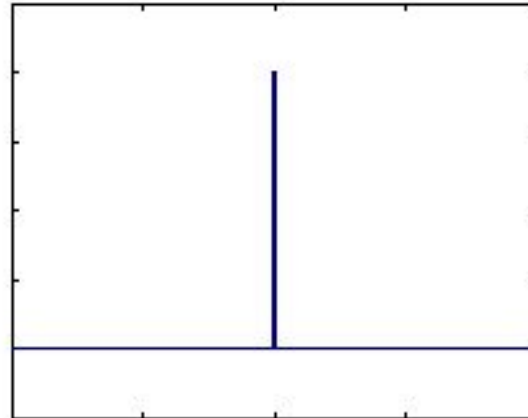
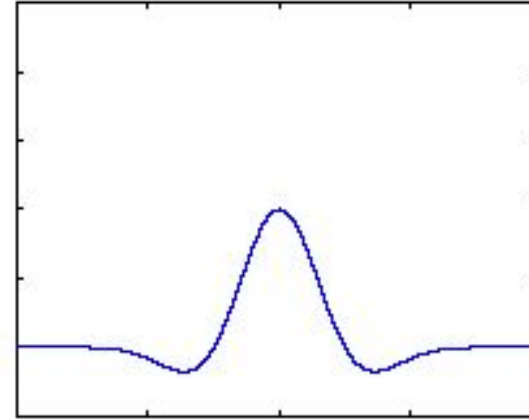
$$y(t) = \sum_k a_k F[x_k(t)]$$

Homogeneity

Input

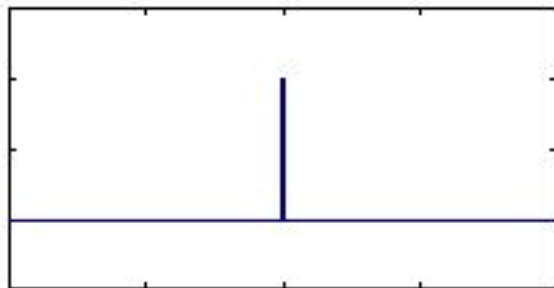


Output

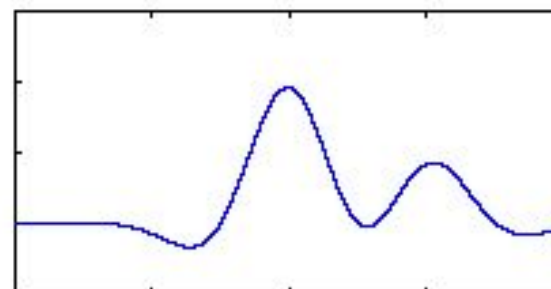
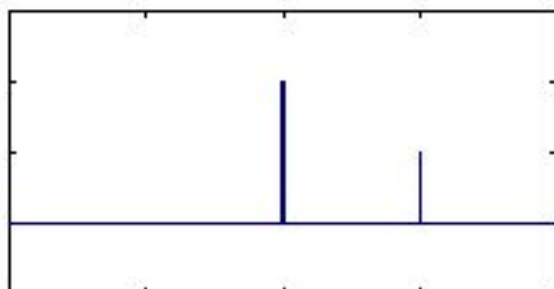
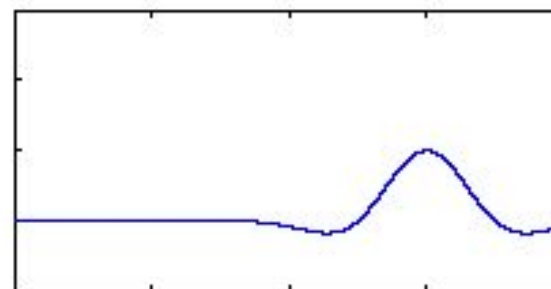
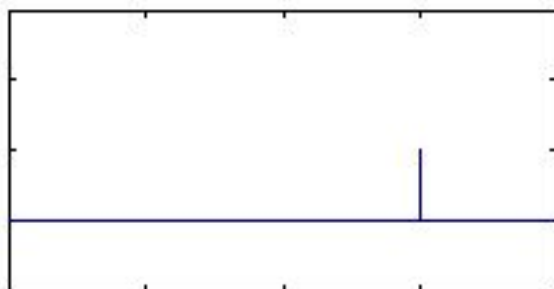
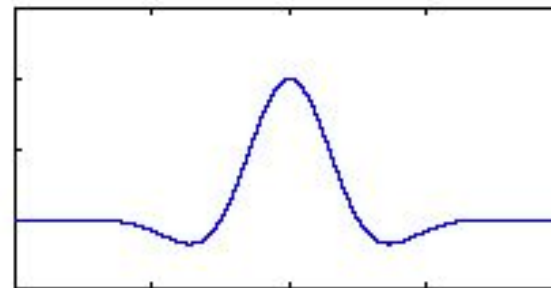


Additivity

Input



Output



Time Invariance

A time-invariant system obeys the following time-shift invariance property: If the input signal $\hat{x}(t)$ is

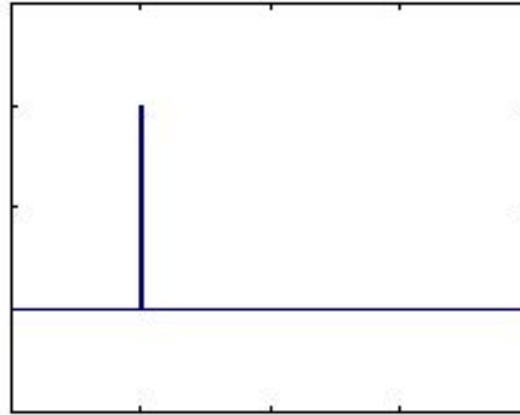
$$\hat{x}(t) = x(t - \tau)$$

then for any real constant τ ,

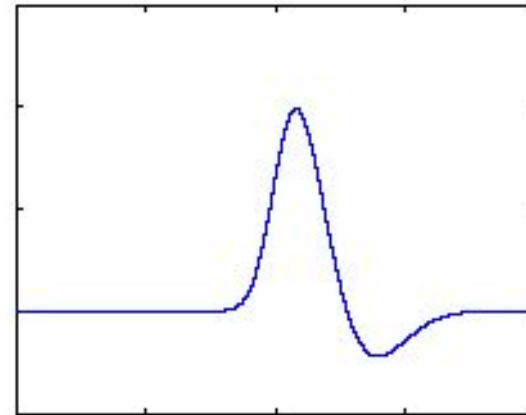
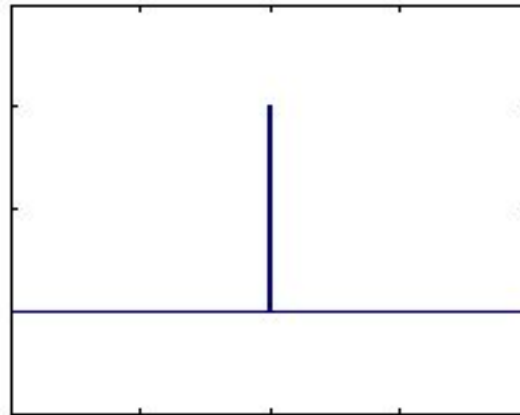
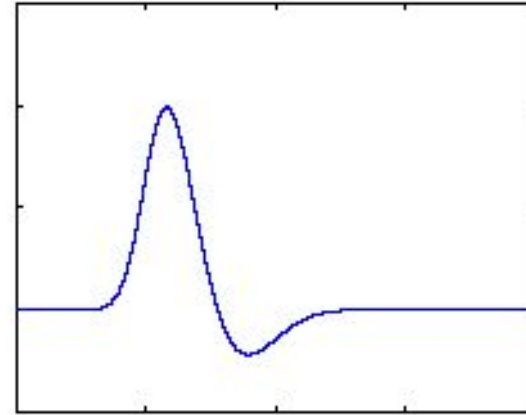
$$\hat{y}(t) = F[x(t - \tau)] = y(t - \tau)$$

Time Invariance

Input



Output



Convolution

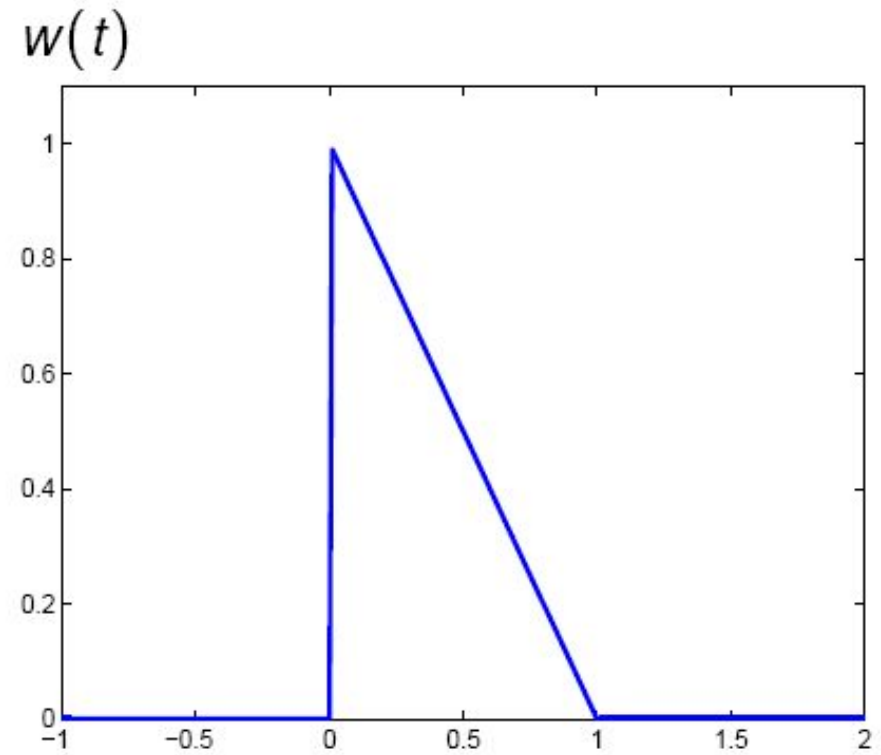
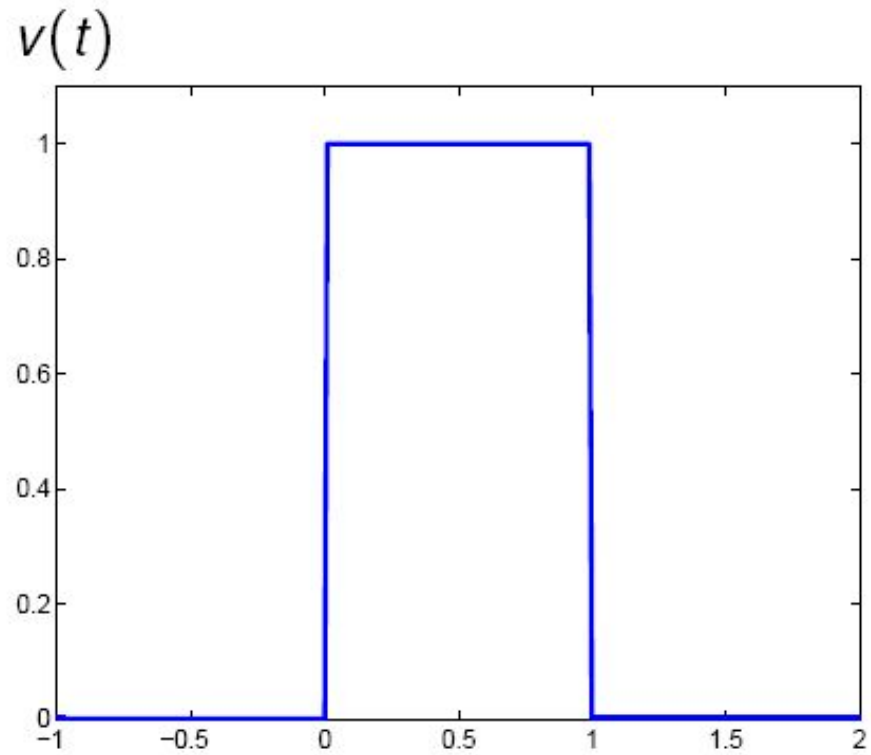
It is a useful mathematical tool for analyzing linear systems.

$$\begin{aligned}v(t) * w(t) &\doteq \int_{-\infty}^{+\infty} v(\lambda) w(t - \lambda) d\lambda \\ &= \int_{-\infty}^{+\infty} v(t - \lambda) w(\lambda) d\lambda\end{aligned}$$

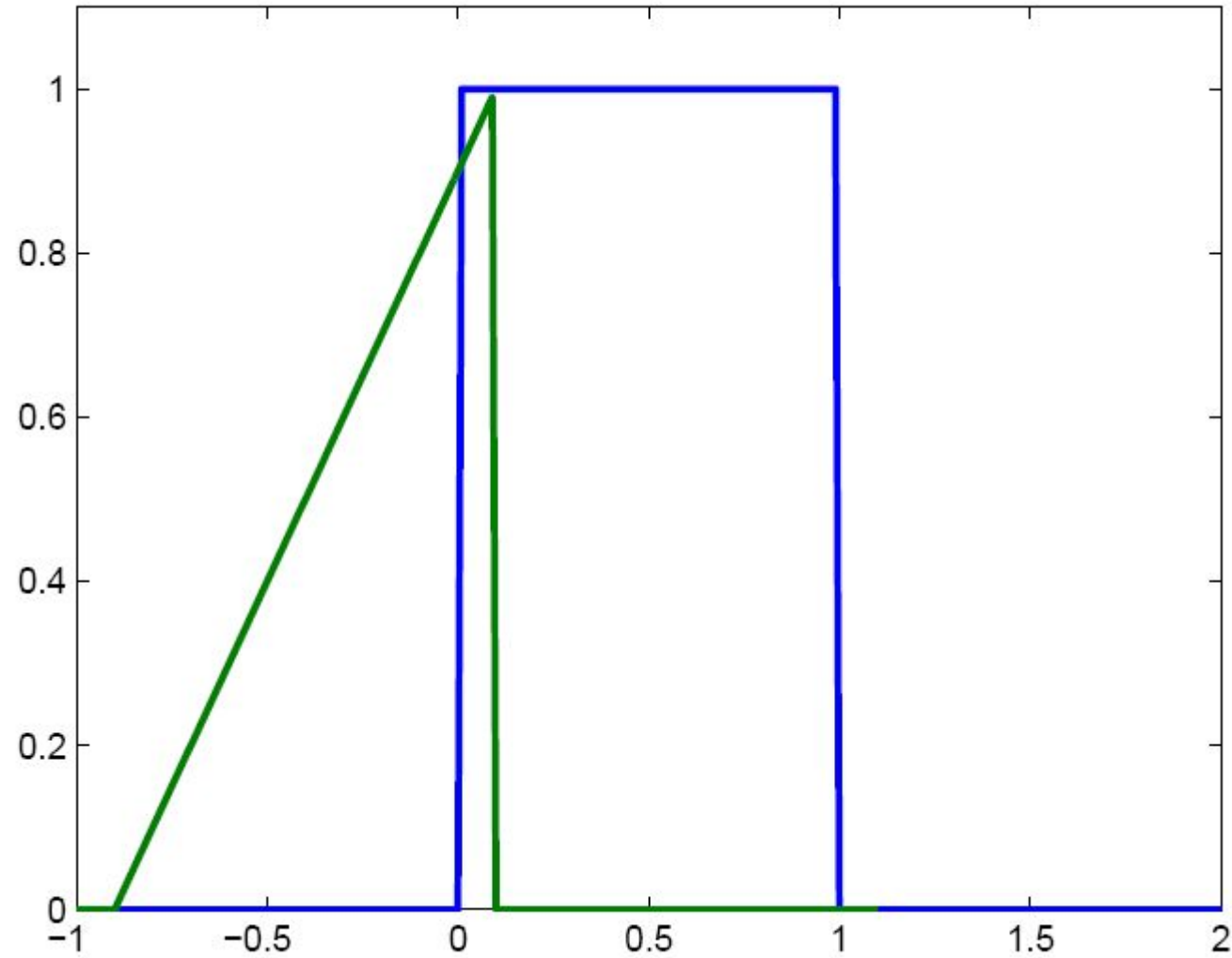
Note that in the integral t is the *independent* variable, and λ is the *integration* variable.

Insofar as the integral is concerned, t is constant.

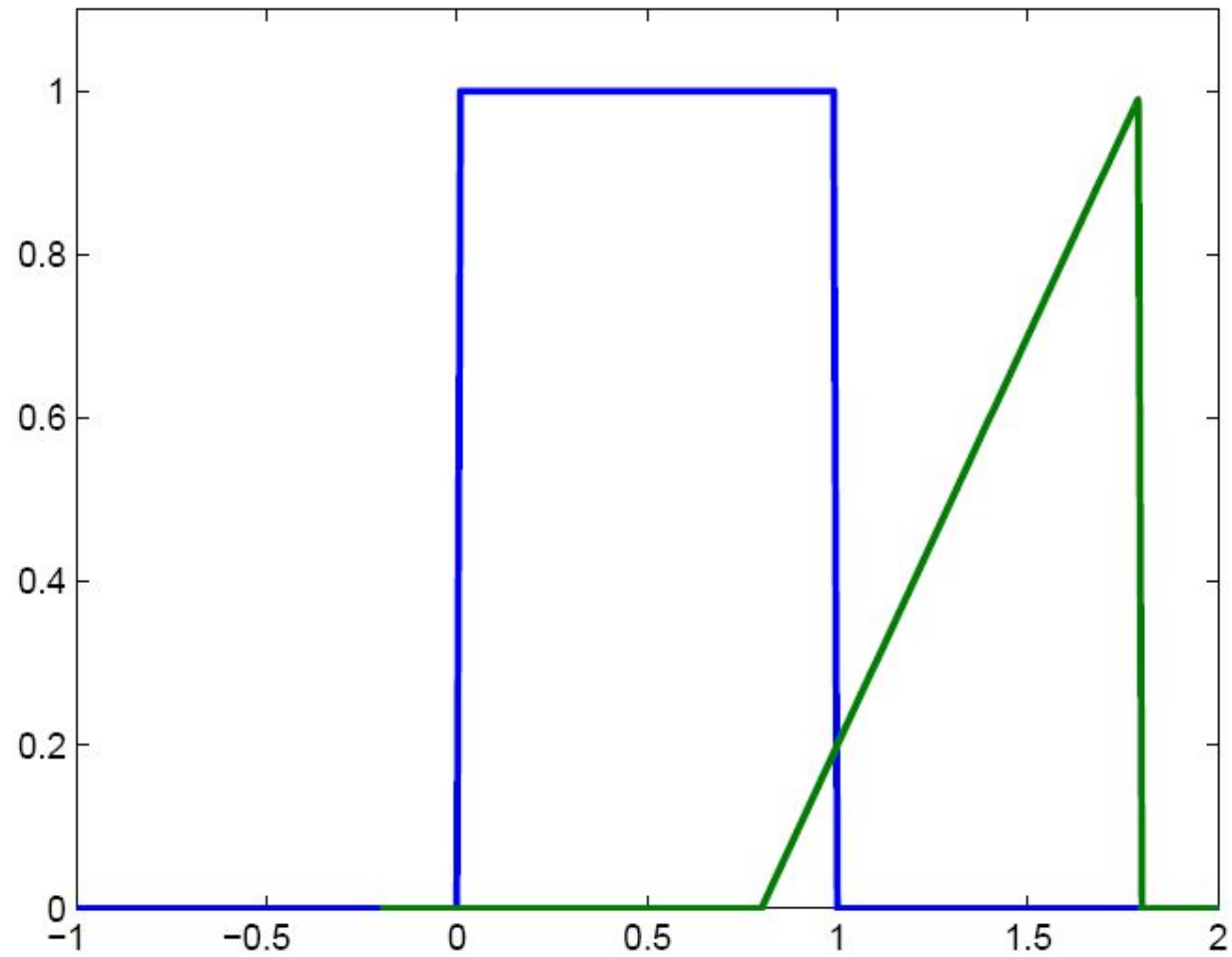
Graphical Example



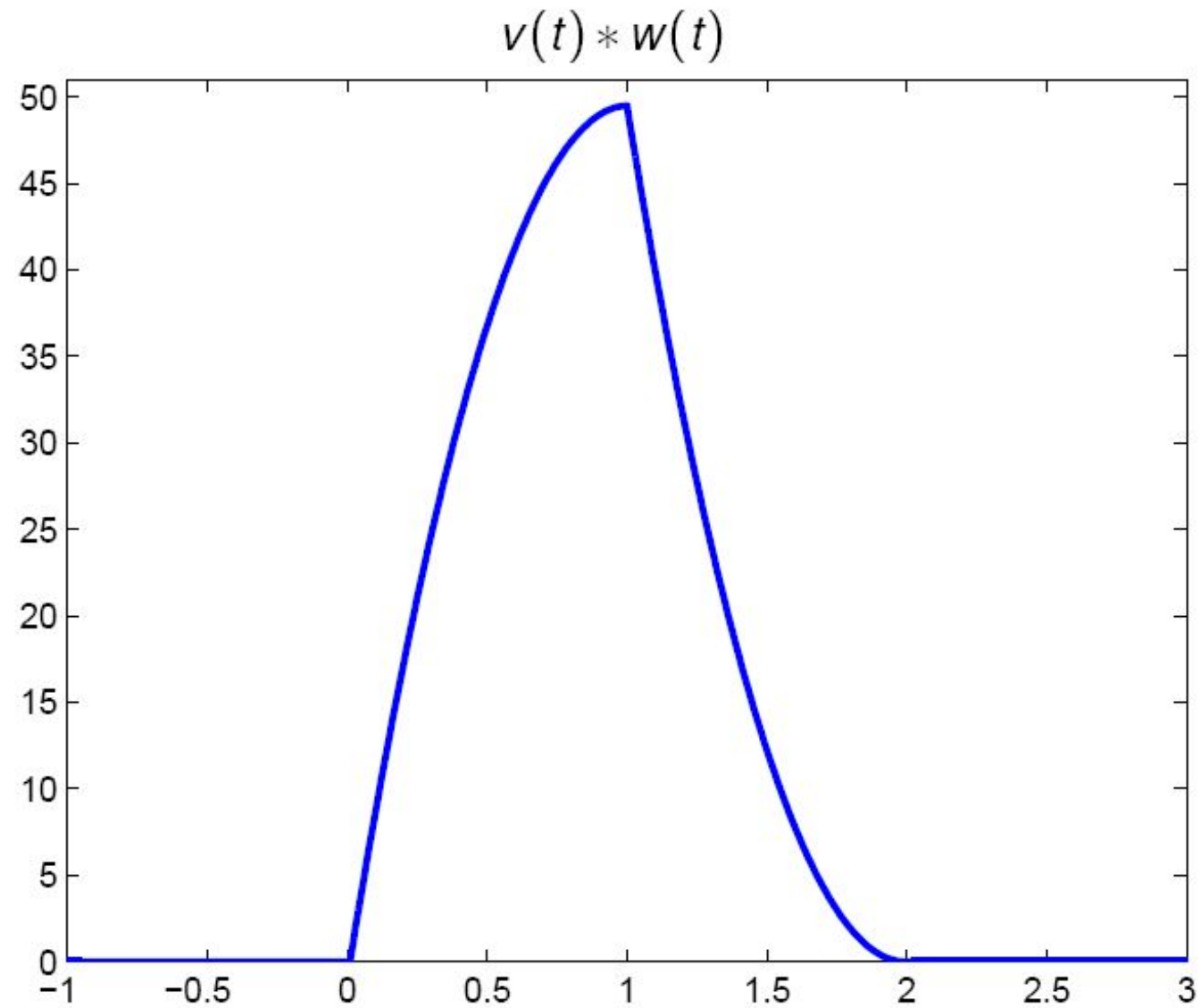
Graphical Example (II)



Graphical Example (III)



Graphical Example (IV)



Impulse

The *Dirac-delta function*, or *unit impulse* $\delta(t)$ is not a strict function in the mathematical sense. It is only defined by a set of assignment rules. If $v(t)$ is continuous at time $t = 0$,

$$\int_{t_1}^{t_2} v(t)\delta(t)dt = \begin{cases} v(0) & \text{if } t_1 < 0 < t_2 \\ 0 & \text{otherwise} \end{cases}$$

$$\Rightarrow \int_{-\infty}^{+\infty} \delta(t)dt = \int_{-\varepsilon}^{+\varepsilon} \delta(t)dt = 1$$

The *Dirac-delta function*, $\delta(t)$ has unit area at $t = 0$, is null for $t \neq 0$, and has no mathematical or physical meaning, unless it appears under the operation of integration.

Impulse Integration Properties

Replication operation

$$v(t) * \delta(t - \tau) = v(t - \tau)$$

Sampling operation

$$\int_{-\infty}^{+\infty} v(t) \delta(t - \tau) dt = v(\tau)$$

Linear System's Impulse Response

We *define* the impulse response of a system characterized by $y(t) = F[x(t)]$ as:

$$h(t) \doteq F[\delta(t)]$$

If $x(t)$ is a continuous function, $x(t) = x(t) * \delta(t)$.

$$y(t) = F[x(t)] = F \left[\int_{-\infty}^{+\infty} x(\lambda) \delta(t - \lambda) d\lambda \right]$$

If the system is linear

$$y(t) = \int_{-\infty}^{+\infty} x(\lambda) F[\delta(t - \lambda)] d\lambda$$

If the system is shift-invariant

$$y(t) = \int_{-\infty}^{+\infty} x(\lambda) h(t - \lambda) d\lambda \Rightarrow y(t) = x(t) * h(t)$$

Step and Impulse Response

If $u(t)$ is a step function such that

$$u(t) = \begin{cases} 1 & \text{if } t \geq 0 \\ 0 & \text{otherwise} \end{cases}$$

the system's *step response* is defined as $g(t) \doteq F[u(t)]$.

Exploiting the system's impulse response properties:

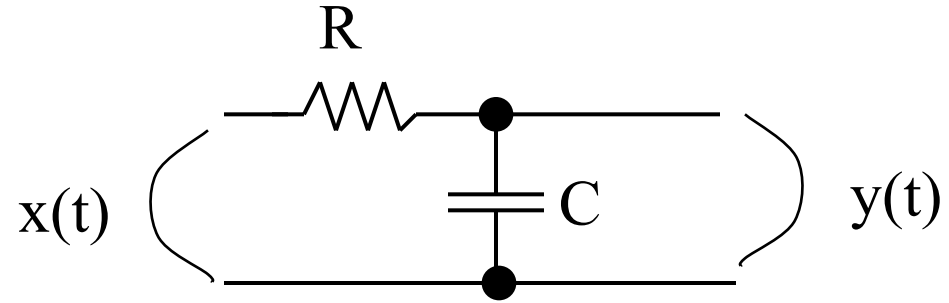
$$g(t) = h(t) * u(t)$$

$$\Rightarrow \frac{d}{dt}g(t) = h(t) * \frac{d}{dt}u(t) = h(t) * \delta(t)$$

$$h(t) = \frac{d}{dt}g(t)$$

A system's impulse response is the derivative of its step response.

Resistor-Capacitor Circuit



1st order linear differential equation:

$$RC \frac{d}{dt} y(t) + y(t) = x(t)$$

Solving for $x(t) = \delta(t)$:

$$h(t) = \frac{1}{RC} e^{-t/RC} \cdot u(t)$$

In general: $y(t) = x(t) * h(t)$

$$= \int_0^{\infty} \frac{1}{RC} e^{-t/RC} x(t - \lambda) d\lambda$$

Complex Exponentials

All solutions to linear homogeneous equations are of the form e^{st} where $s \doteq \sigma + j\omega = M\cos(\phi) + jM\sin(\phi)$, and $j = \sqrt{-1}$. σ represents the real part of the complex number, ω its imaginary part, M represents its *magnitude* and ϕ its *phase*. Magnitude and phase of a complex number obey to the following relationships:

$$\begin{aligned} M &= \sqrt{\sigma^2 + \omega^2} & e^{j\phi} &= \cos(\phi) + j\sin(\phi) \\ \phi &= \arctan\left(\frac{\omega}{\sigma}\right) & e^{-j\phi} &= \cos(\phi) - j\sin(\phi) \end{aligned}$$

It follows that s can be also written as $s = Me^{j\phi}$

Complex Exponentials

Example of second order linear differential equation:

$$\frac{d^2}{dt^2} V(t) + \alpha \frac{d}{dt} V(t) + \beta V(t) = 0$$

Substitute e^{st} for V : $s^2 e^{st} + \alpha s e^{st} + \beta e^{st} = 0$

$$\text{Solving for } s: \quad s = \frac{-\alpha \pm \sqrt{\alpha^2 - 4\beta}}{2}$$

If $\alpha^2 - 4\beta \geq 0$, s is real, otherwise s is complex

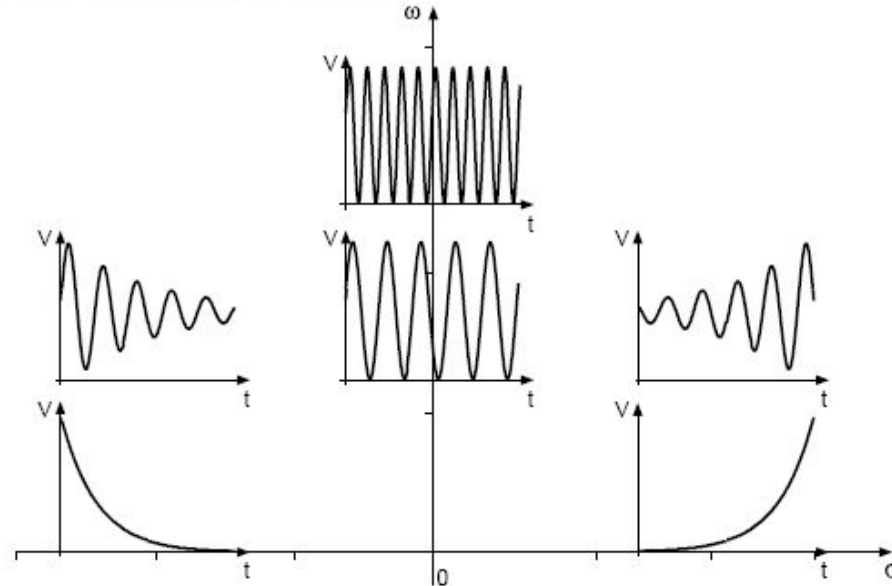
Measured Responses

In practice, we measure responses of the type $V = e^{st}$ with *real* instruments, so if e^{st} were a complex exponential, we would be able to measure only

$$\operatorname{Re}\{e^{st}\} = \operatorname{Re}\{e^{(\sigma+j\omega)t}\} = e^{\sigma t} \operatorname{Re}\{e^{j\omega t}\}$$

So the measured response of the system would be

$$V_{\text{meas}} = e^{\sigma t} \cos \omega t$$



Heaviside-Laplace Transform

Oliver Heaviside, in analyzing analog circuits made the following observation:

$$\frac{d^n}{dt^n} e^{st} = s^n e^{st}$$

we can consider s as an operator meaning derivative with respect to time. Similarly, we can view $\frac{1}{s}$ as the operator for integration with respect to time.

This observation was formalized by mathematicians, when they realized that Heaviside's method was using polynomials in s that were *Laplace transforms*. The Laplace transform is an operator that allows to link functions that operate in the time domain with functions of complex variables:

$$\bullet L[y(t)] = Y(s) = \int_{-\infty}^{\infty} y(t) e^{-st} dt$$

Transfer Function

• Transfer function is defined as $H(s) = \frac{Y(s)}{X(s)}$

If $x(t) = \delta(t)$, then

$$X(s) = \int_{-\infty}^{\infty} \delta(t) e^{-st} dt = 1$$

and

$$Y(s) = \int_{-\infty}^{\infty} y(t) e^{-st} dt = \int_{-\infty}^{\infty} h(t) e^{-st} dt$$

Therefore

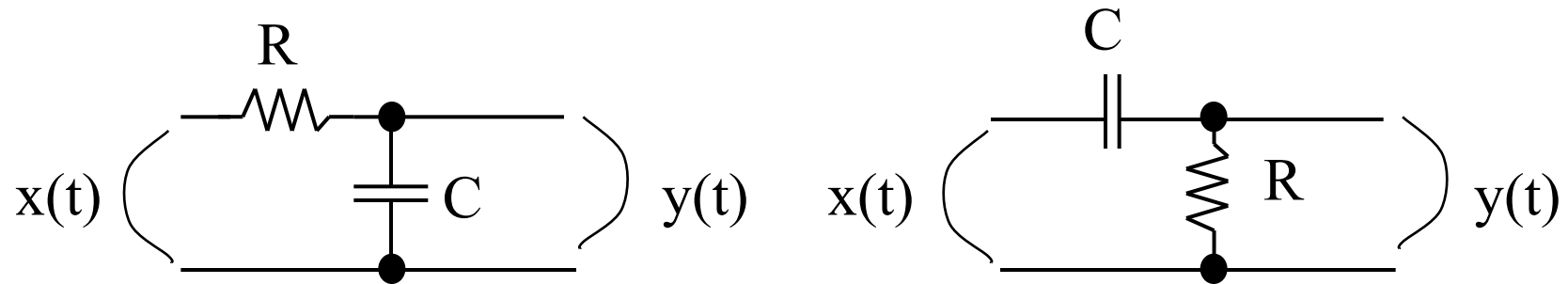
$$H(s) = \int_{-\infty}^{\infty} h(\lambda) e^{-\lambda s} d\lambda = L[h(t)]$$

Linear Systems, Small Signal Analysis, and Integrators

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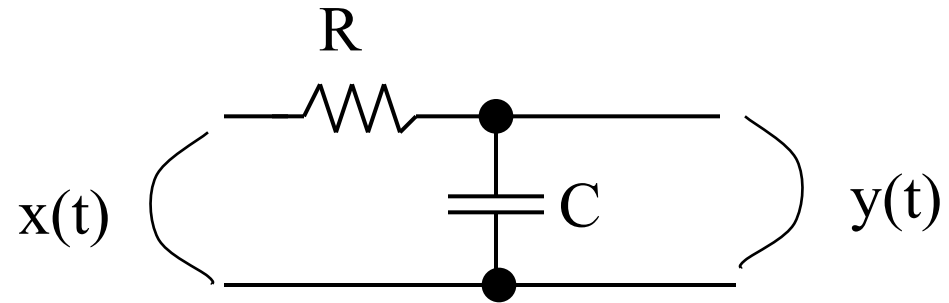
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Resistor-Capacitor Circuits



- First order, linear time-invariant systems
- Input is $x(t)$, output is $y(t)$
- RC integrator circuit
- CR differentiator circuit

R-C Integrator



Time domain:

$$RC \frac{d}{dt} y(t) + y(t) = x(t)$$

Solving for $x(t) = \delta(t)$:

$$h(t) = \frac{1}{RC} e^{-t/RC} u(t)$$

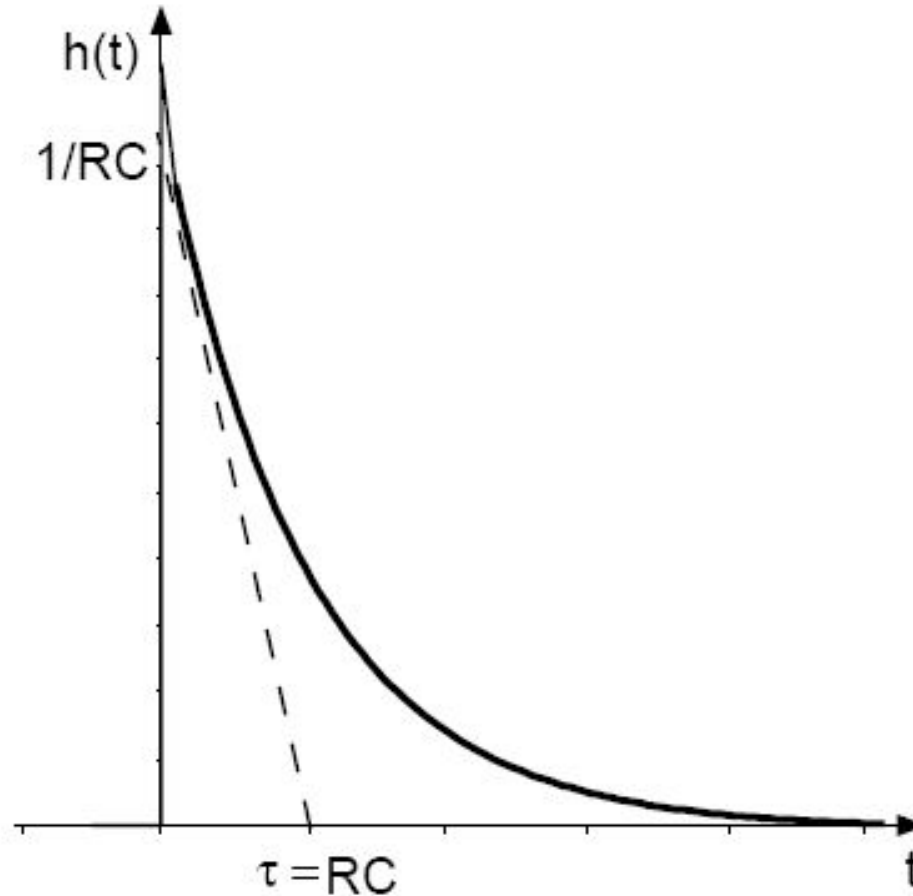
Laplace domain:

$$RCsY(s) + Y(s) = X(s)$$

Solving for $X(s) = 1$:

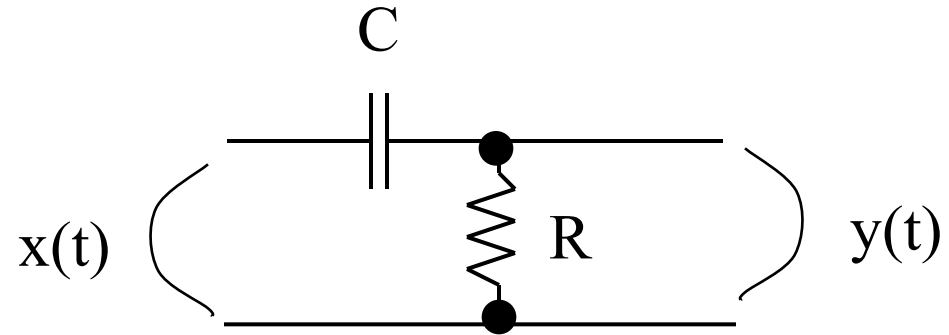
$$H(s) = \frac{1}{1 + RCs}$$

RC Integrator Impulse Response



The circuit's *time constant* $\tau = RC$ is the time required to discharge the capacitor, through the resistor, to 36.8% ($1/e$) of its final steady state value.

C-R Differentiator



Time domain:

$$\tau \frac{d}{dt} y(t) + y(t) = \tau \frac{d}{dt} x(t)$$

For a step input, the derivative of the input is an impulse at $t=0$.

Laplace domain:

$$\tau s Y(s) + Y(s) = \tau X(s)$$

Solving for $X(s) = 1$:

$$H(s) = \frac{\tau s}{1 + \tau s}$$

Frequency Domain Analysis

A system's response in the frequency domain can be determined by using pure sinusoids as input signals. For time-invariant linear systems, the amplitude and phase of the input sinusoid will be altered by the system's transfer function.

Input: $x(t) = \sin(\omega t)$

Output: $y(t) = A \sin(\omega t + \phi)$

where A and ϕ determine the amount of scaling and phase shift

RC-Integrator Frequency Domain Analysis

In this domain $s = j\omega$ and the RC-circuit's transfer function is:

$$H(j\omega) = \frac{1}{1 + j\omega\tau}$$

- $\omega\tau \ll 1 \Rightarrow Y(j\omega) \approx X(j\omega)$
- $\omega\tau \gg 1 \Rightarrow Y(j\omega) \approx \frac{1}{j\omega\tau} X(j\omega)$

The frequency $f_c = \frac{\omega}{2\pi} = \frac{1}{2\pi\tau}$ is defined to be the **cutoff frequency**.

In electronics, cutoff frequency (f_c) is the frequency either above which or below which the power output of a circuit, such as a line, amplifier, or filter, is reduced to 1/2 of the passband power; the half-power point. This is equivalent to a voltage (or amplitude) reduction to 70.7% of the passband, because voltage V^2 is proportional to power P . This happens to be close to -3 decibels, and the cutoff frequency is frequently referred to as the -3 dB point.

Magnitude, Phase, Bode Plots

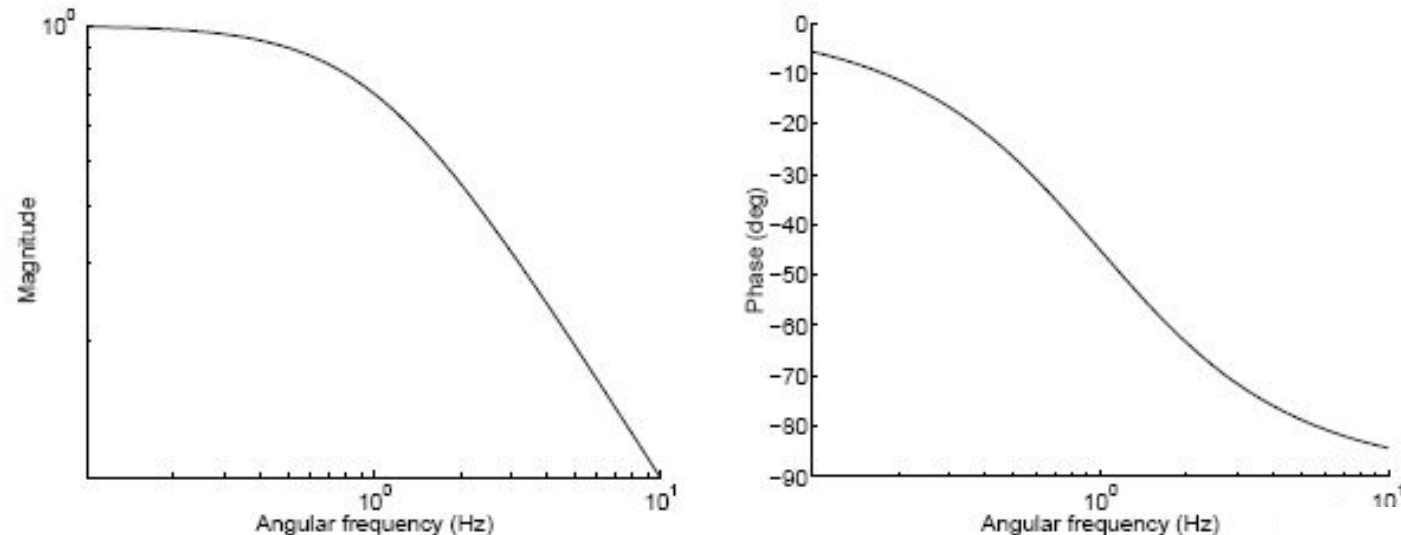
The magnitude of the RC-integrator's transfer function is:

$$|H(j\omega)| = \frac{1}{\sqrt{1 + (\omega\tau)^2}}$$

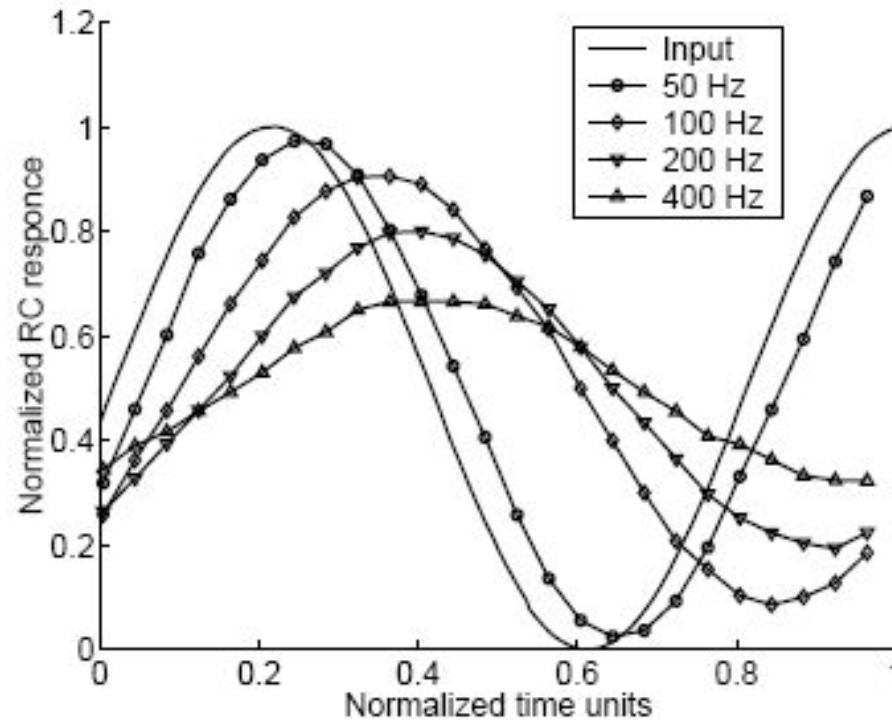
Its phase is

$$\phi = \arctan(-\omega\tau)$$

The plots of a transfer function's magnitude and phase versus input frequency are called *Bode plots* (after Hendrik Wade Bode).



RC-Integrator Frequency Response



Responses of an RC lowpass filter with $R = 10M\Omega$ and $C = 1nF$.

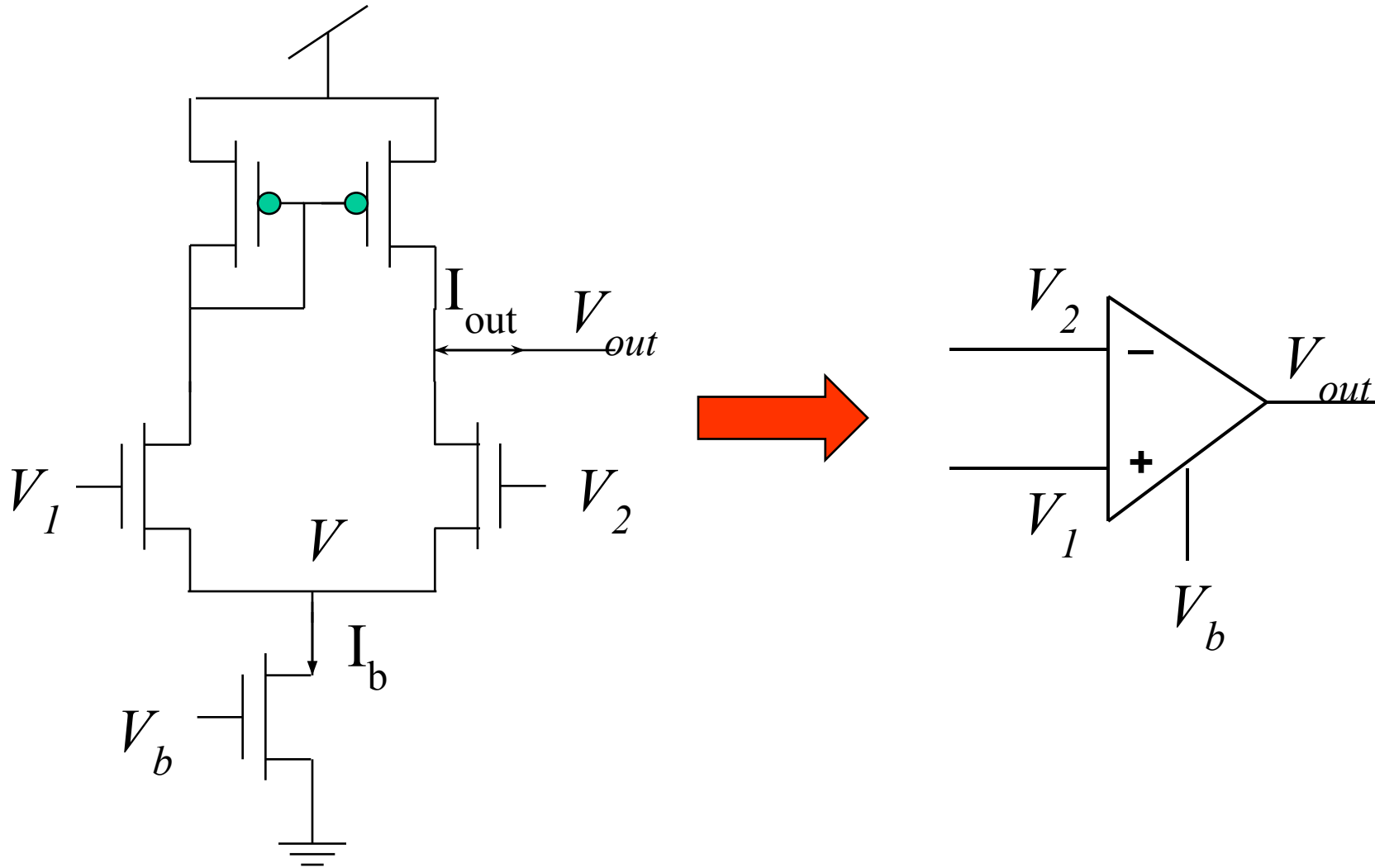
All data is plotted on a normalized scale. The signals' amplitudes have been normalized with respect to the input and the time-base has been normalized to unity.

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Transconductance Amplifier



Transconductance Amplifier Specs

$$I_{out} = I_b \tanh \left(\frac{\kappa}{2U_T} (V_1 - V_2) \right)$$

Transconductance:

$$g_m = \frac{\partial I_{out}}{\partial (V_1 - V_2)} = \frac{\kappa I_b}{2U_T}$$

Output conductance:

$$g_d = -\frac{\partial I_{out}}{\partial V_{out}} \approx \frac{I_b}{V_E}$$

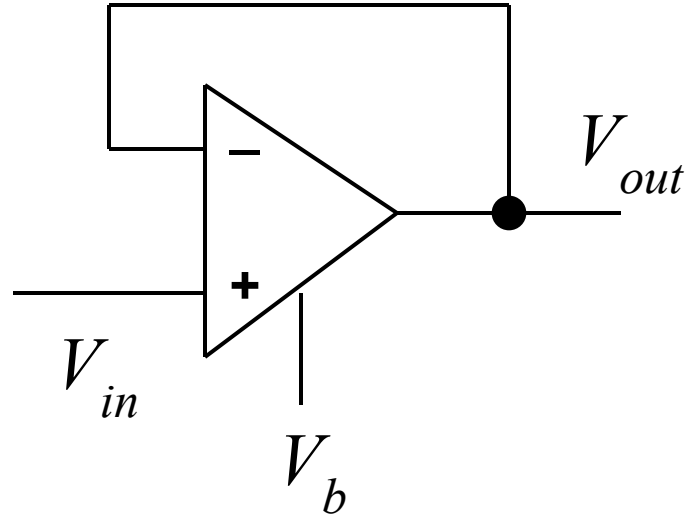
Input conductance:

$$g_{in} = \frac{\partial I_{in}}{\partial V_{out}} \approx 0$$

Voltage gain:

$$A = \frac{\partial V_{out}}{\partial (V_1 - V_2)} = \frac{g_m}{g_d}$$

Unity-Gain Follower



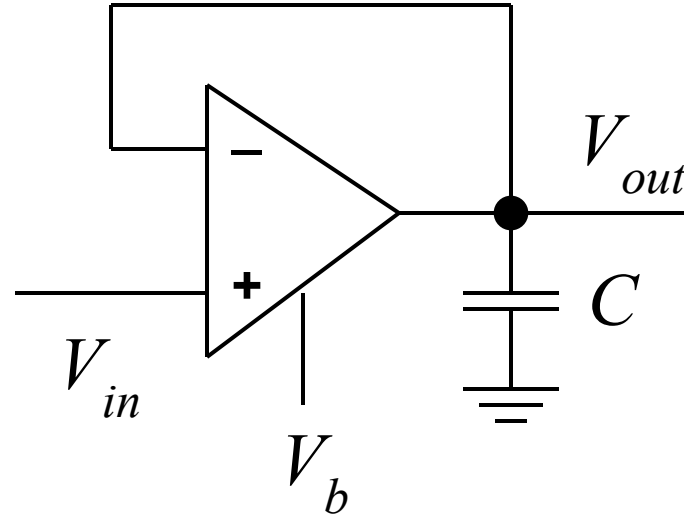
- $\partial V_{out} = A(\partial V_{in} - \partial V_{out})$

- Transfer function: $\frac{\partial V_{out}}{\partial V_{in}} = \frac{A}{1 + A}$

Input impedance: $Z_{in} = \frac{\partial V_{in}}{\partial I_{in}} \rightarrow \infty$

Output impedance: $Z_{out} = -\frac{\partial V_{out}}{\partial I_{out}} = \frac{1}{g_d} \approx \frac{V_E}{I_b}$

Follower Integrator



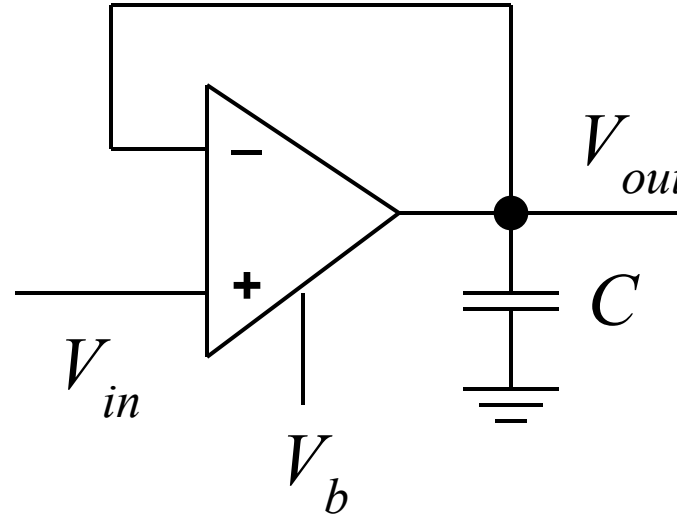
$$C \frac{dV_{out}}{dt} = I_b \tanh \left(\frac{\kappa(V_{in} - V_{out})}{2U_T} \right)$$

In the small-signal regime

$$C \frac{dV_{out}}{dt} = G(V_{in} - V_{out})$$

where $G = \frac{kI_b}{2U_T}$ is the amplifier transconductance

Follower Integrator (Laplace domain)



- $$C \frac{dV_{out}}{dt} = G(V_{in} - V_{out}) \Rightarrow sCV_{out} = G(V_{in} - V_{out})$$

Transfer function:
$$\frac{V_{out}}{V_{in}} = \frac{1}{s\tau + 1} \text{ where } \tau = \frac{C}{G}$$

Amplifier operates in linear region only if

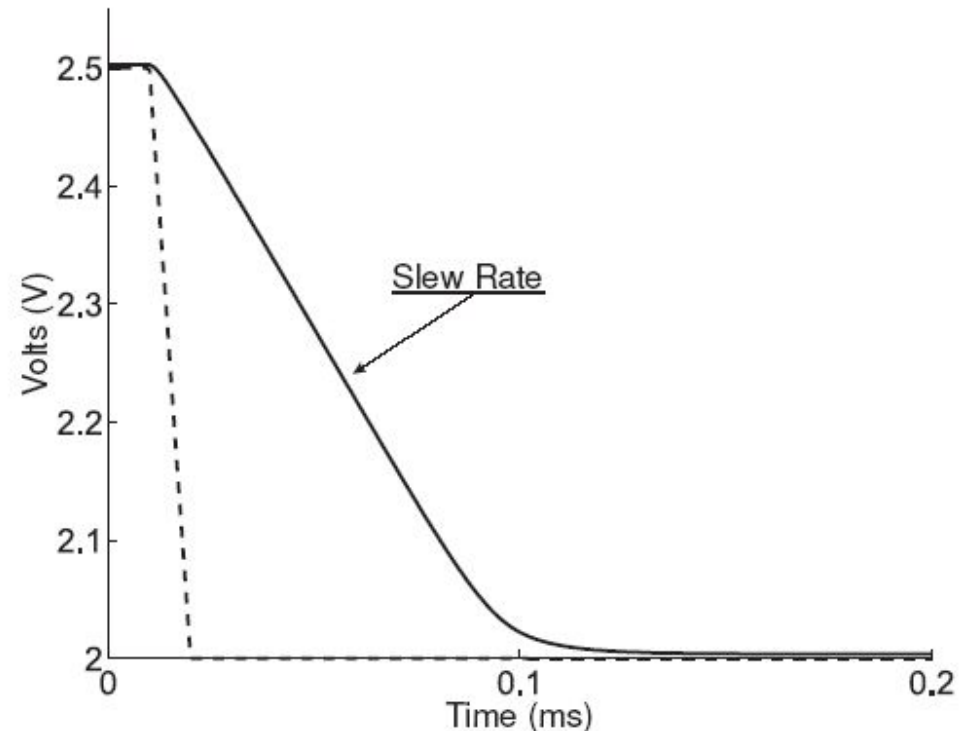
$$V_{in} - V_{out} < 4U_T$$

Follower Integrator (Large Signal Behavior)

For large variations of V_{in} , the output current of the transconductance amplifier saturates to I_b or $-I_b$.

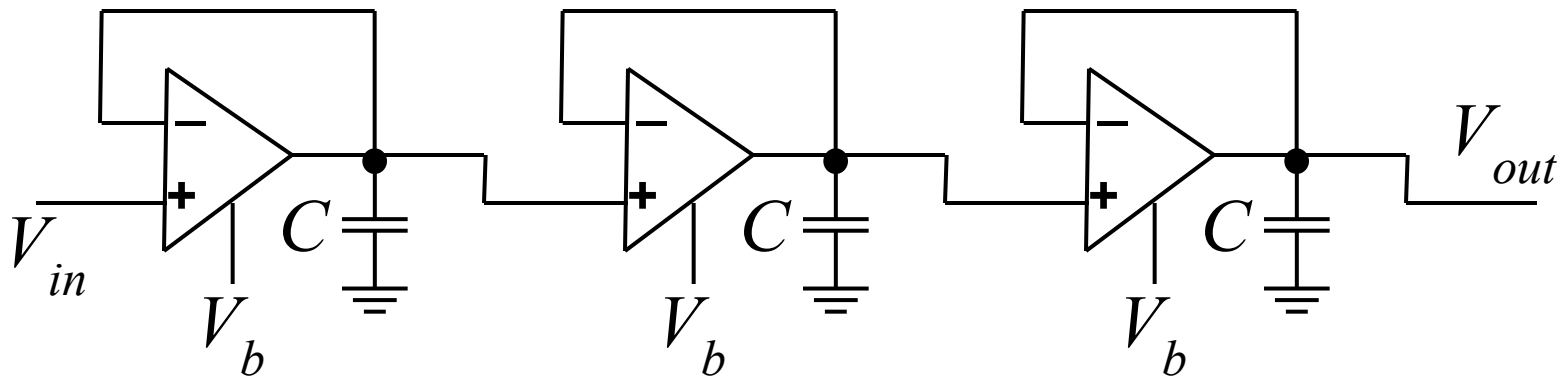
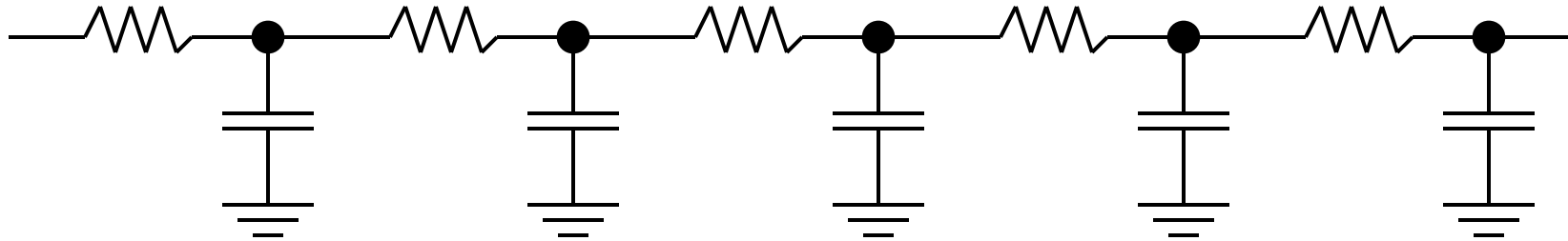
As long as $|V_{out} - V_{in}| > 4 U_T$, V_{out} changes linearly over time. This region is also called the **slew rate** of an amplifier.

If $|V_{out} - V_{in}| < 4 U_T$, the amplifier behaves like a linear conductance and V_{out} changes exponentially.



Delay Line

Cascade of RC integrators and follower-integrators



Follower-Integrator Composition Property

Transfer function of delay line with n delay elements:

$$\frac{V_{out}}{V_{in}} = \left(\frac{1}{s\tau + 1} \right)^n$$

We can write the transfer function directly in terms of

magnitude and phase valid for $(j\omega\tau) \ll 1$

$$\frac{V_{out}}{V_{in}} \approx \left(\frac{1}{1 + \frac{n}{2} (j\omega\tau)^2} \right) e^{(-jn\omega\tau)}$$

Translinear Principle

- Coined by Barrie Gilbert in 1975, *translinear* means that the bipolar junction transistor's *transconductance* is *linear* in its collector current.
- In a bipolar transistor, the collector current is exponential in the base-emitter voltage. This exponential dependence is also captured in the subthreshold domain of a MOSFET.

Translinear Principle

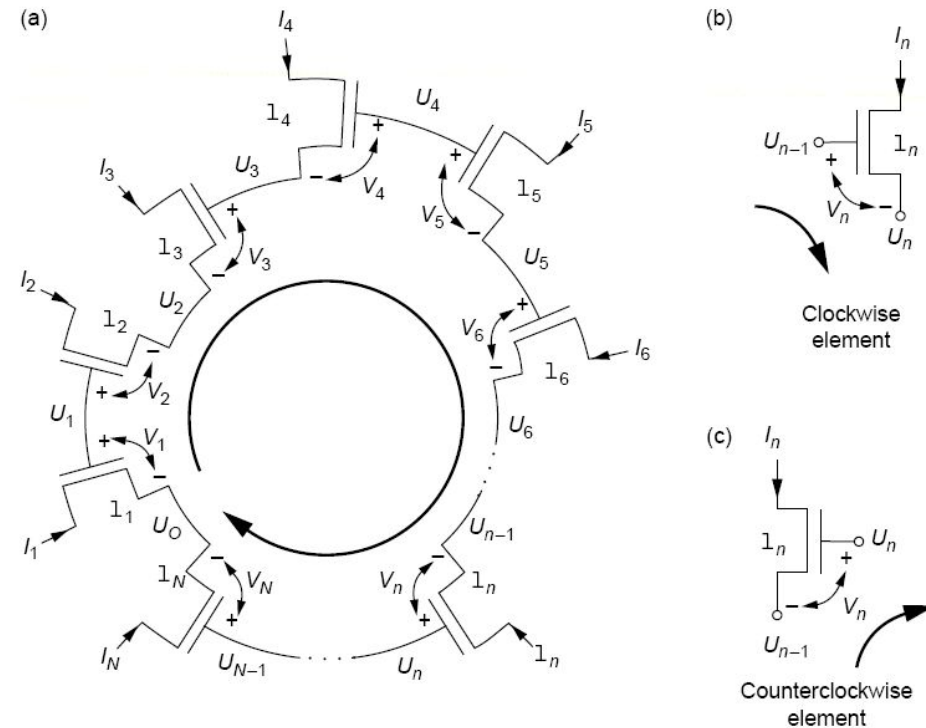
Kirchoff's Voltage Law around loop:

$$\sum_{n \in CCW} V_n = \sum_{n \in CW} V_n \quad (1)$$

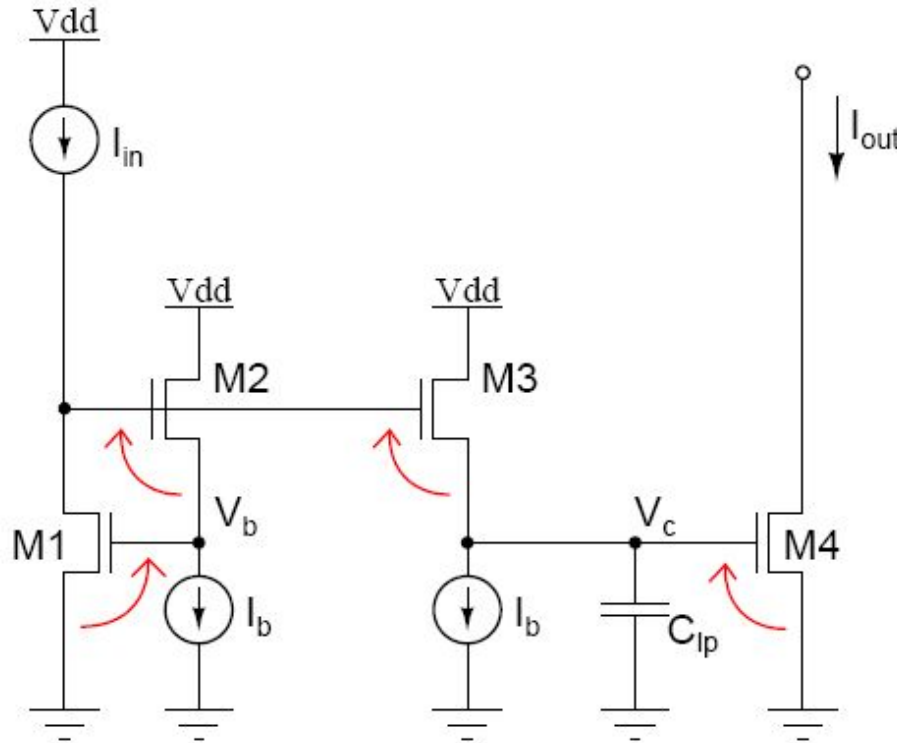
Substituting for I in (1) :

$$\sum_{n \in CCW} U_T \log\left(\frac{I_n}{I_0}\right) = \sum_{n \in CW} U_T \log\left(\frac{I_n}{I_0}\right)$$

$$\rightarrow \prod_{n \in CCW} I_n = \prod_{n \in CW} I_n$$



Current-Mode Low-Pass Filter



$$V_{gs}^{M1} + V_{gs}^{M2} - V_{gs}^{M3} - V_{gs}^{M4} = 0$$

$$I_{ds} = I_0 e^{\frac{\kappa V_{gs} - V_s}{U_T}}$$

$$I_{ds}^{M1} \cdot I_{ds}^{M2} = I_{ds}^{M3} \cdot I_{ds}^{M4}$$

$$I_{ds}^{M3} = I_b + C_{lp} \frac{d}{dt} V_c$$

$$\frac{d}{dt} I_{out} = \frac{\kappa}{U_T} I_{out} \frac{d}{dt} V_c$$

$$I_{in} \cdot I_b = \left(I_b + \frac{C_{lp} U_T}{\kappa} \frac{\dot{I}_{out}}{I_{out}} \right) \cdot I_{out}$$

$$\tau \dot{I}_{out} + I_{out} = I_{in}$$

$$\text{with } \tau = \frac{C_{lp} U_T}{\kappa I_b}$$

(Frey, 1982)

THE END

Next week: Current-mode and winner-take-all circuits