

Dynamics of Rigid Body

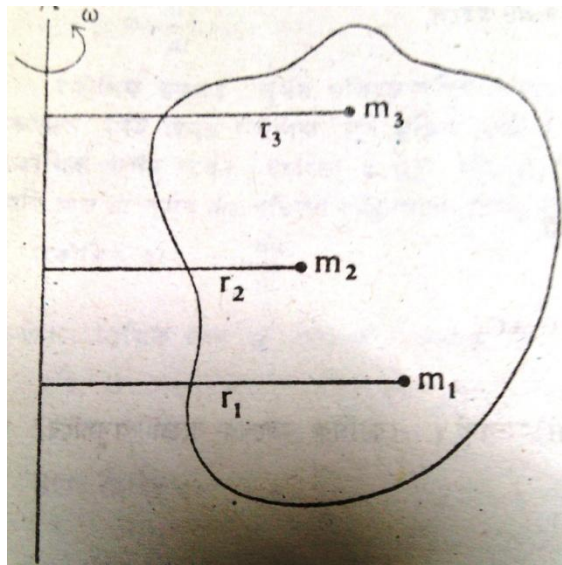
Rigid Body:

A rigid body is defined as that body which does not undergo any change in shape and volume when external forces are applied on it.

When forces are applied on a rigid body, the distance between any two particles of the body will remain unchanged, however large the force may be. In actual practice, nobody is perfectly rigid. For practical purpose solid bodies are taken as rigid bodies.

Moment of Inertia:

A body maintains the current state of motion unless acted upon by an external force." The measure of the inertia in the linear motion is the **mass** of the system and its angular counterpart is the so-called **moment of inertia**. The moment of inertia of a body is not only related to its mass but also the distribution of the mass throughout the body.



Let us consider a body of mass M and any axis YY' . Imagine the body to be composed of a large number of particles of masses m_1, m_2, m_3 , etc. at distance r_1, r_2, r_3, \dots etc. from the axis YY' . Then the moment of inertia of the particle m_1 about YY' is $m_1 r_1^2$, that of the particle m_2 is $m_2 r_2^2$ and so on. Therefore, the moment of inertia, I of the whole body, about the axis YY' is equal to the sum of $m_1 r_1^2, m_2 r_2^2, m_3 r_3^2, \dots$ etc.

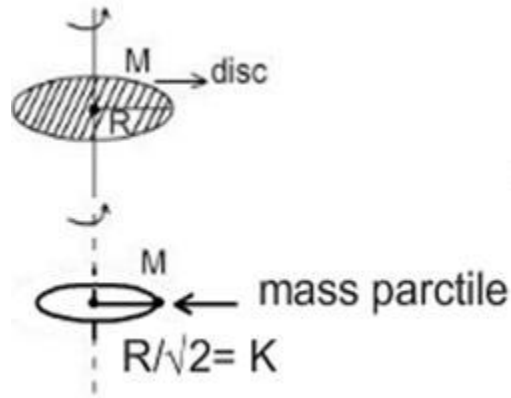
Thus,

$$I = m_1 r_1^2 + m_2 r_2^2 + m_3 r_3^2 + \dots$$
$$I = \sum_{i=1}^n m_i r_i^2$$

Radius of Gyration:

The radius of gyration of the body about an axis may be defined as the distance of a mass point from the same axis, whose mass is equal to the mass of the whole body and whose moment of Inertia is equal to the moment of inertia of the body, if rotated about the same axis.

The example is given below:



In the above diagram we have shown a disc of mass M and radius R , just below (to make you understand) we have shown a particle of mass M which is moving about the same axis in the circle of radius K .

So for disc and particle,

$$I = \frac{MR^2}{2} \dots\dots\dots(1)$$

$$I = MK^2 \dots\dots\dots(2)$$

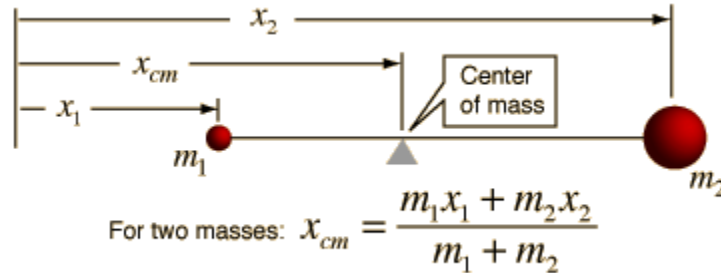
Comparing (1)&(2)

$$K = \frac{R}{\sqrt{2}} .$$

Here K is the radius of Gyration.

Center of Mass:

The terms "center of mass" and "center of gravity" are used synonymously in a uniform gravity field to represent the unique point in an object or system which can be used to describe the system's response to external forces and torques. The concept of the center of mass is that of an average of the masses factored by their distances from a reference point. In one plane, that is like the balancing of a seesaw about a pivot point with respect to the torques produced.



If you are making measurements from the center of mass point for a two-mass system then the center of mass condition can be expressed as

$$m_1 r_1 = m_2 r_2 \quad m_1 = m_2 \frac{r_2}{r_1}$$

where r_1 and r_2 locate the masses. The center of mass lies on the line connecting the two masses.

Rotational Kinetic Energy:

Let us consider a rigid body rotates with angular velocity ω about an axis XY. Let the body contain 'n' number of particles and their mass are $m_1, m_2, m_3 \dots \dots m_n$ respectively and distance from the rotational axis $r_1, r_2, r_3 \dots \dots r_n$ respectively.

Hence, if the radial velocity of the particles are $v_1, v_2, v_3 \dots \dots v_n$, then the kinetic energy of the particle of mass m_1 , is

$$k_1 = \frac{1}{2} m_1 v_1^2.$$

$$k_1 = \frac{1}{2} m_1 \omega^2 r_1^2 \dots \dots \dots (1)$$

Similarly, the kinetic energy of the particle of mass m_2 is

$$k_2 = \frac{1}{2} m_2 \omega^2 r_2^2 \dots \dots \dots (2)$$

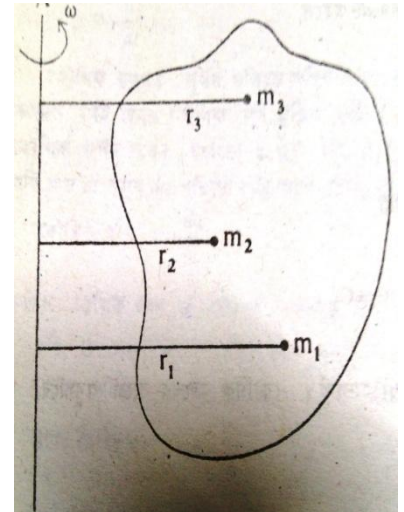
Similarly,

$$k_3 = \frac{1}{2} m_3 \omega^2 r_3^2$$

.....

.....

$$k_n = \frac{1}{2} m_n \omega^2 r_n^2.$$



Hence, the kinetic energy of the rigid body is equal to the sum of kinetic energy of the particles. i.e.,

$$\begin{aligned} k &= k_1 + k_2 + k_3 + \dots \dots \dots + k_n. \\ &= \frac{1}{2} m_1 \omega^2 r_1^2 + \frac{1}{2} m_2 \omega^2 r_2^2 + \dots \dots \dots + \frac{1}{2} m_n \omega^2 r_n^2. \\ &= \frac{1}{2} \omega^2 (m_1 r_1^2 + m_2 r_2^2 + \dots \dots \dots + m_n r_n^2) \\ &= \frac{1}{2} \omega^2 \sum_{i=1}^n m_i r_i^2 \\ &= \frac{1}{2} \omega^2 I \dots \dots \dots \because I = \sum_{i=1}^n m_i r_i^2 \\ &\therefore k = \frac{1}{2} \omega^2 I. \end{aligned}$$

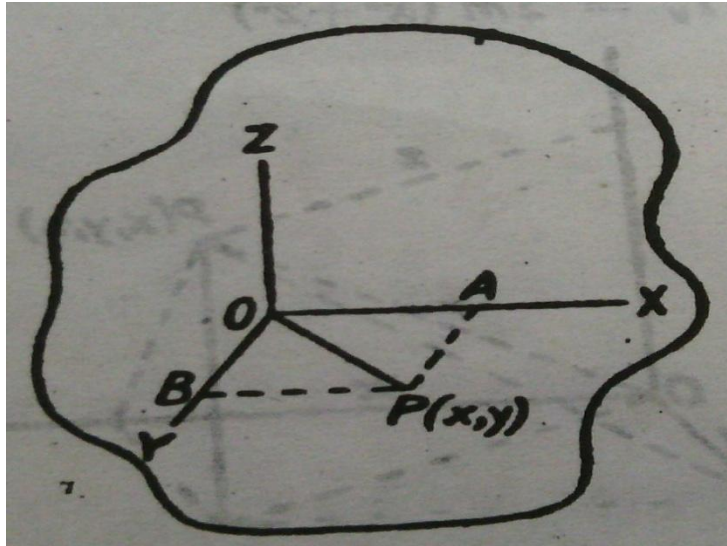
Therefore,

Kinetic energy of rotation = $\frac{1}{2} \times (\text{moment of inertia}) \times (\text{angular velocity})$.

Theorem of Perpendicular Axes:

(a) For a plane Lamina:

The theorem of perpendicular axes states that the moment of inertia of a plane lamina about an axis perpendicular to the plane is equal to the sum of the moments of inertia about two mutually perpendicular axes in the plane of the lamina such that the three mutually perpendicular axes have a common point of intersection.



Let us consider a plane lamina having the axes OX and OY in the plane of the lamina. The axis OZ passes through O and is perpendicular to the plane of the lamina. Let the lamina be divided into a large number of particles, each of mass m. Let a particle of mass m be at P with coordinates (x,y) and situated at a distance r from the point of intersection of the axes (Fig.-01).

$$\therefore r^2 = x^2 + y^2 \dots \dots \dots (1)$$

The moment of inertia of the particle P about the axis

$$OZ = mr^2$$

The moment of inertia of the whole lamina about the axis OZ is given by

$$I_Z = \sum mr^2 \dots \dots \dots (2)$$

The moment of inertia of the whole lamina about the axis OX is given by

$$I_X = \sum my^2 \dots \dots \dots (3)$$

Similarly,

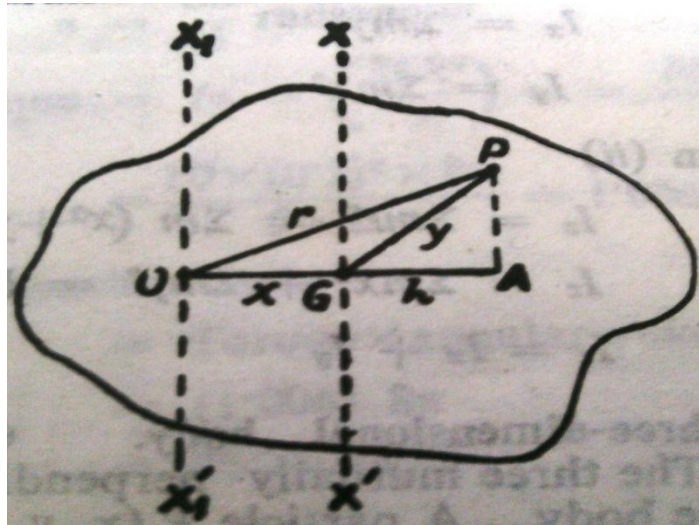
$$I_Y = \sum mx^2 \dots \dots \dots (4)$$

Therefore from equation (2)

$$\begin{aligned} I_Z &= \sum mr^2 = \sum m(x^2 + y^2) \\ I_Z &= \sum mx^2 + \sum my^2 = I_Y + I_X \\ \therefore I_Z &= I_Y + I_X \end{aligned}$$

Theorem of Parallel Axes:

The theorem of parallel axes states that the moment of inertia of a body about any axis is equal to the sum of the moment of inertia of the body about a parallel axis passing through the centre of gravity and the product of the mass of the body and square of perpendicular distance between the two parallel axes.



Let us consider a plane lamina having its centre of gravity at G. The axis XX' passes through the centre of gravity and is perpendicular to the plane of the lamina. The axis X_1X_1' passes through the point O and is parallel to the axis XX' . The distance between the two parallel axes is x , which is shown in Fig.-01.

Let the lamina be divided into large number of particles each of mass m . The moment of inertia of the particle of mass m at P about the axis X_1X_1' is equal to mr^2 . The moment of inertia of the whole lamina about the axis X_1X_1' is given by

$$I_o = \sum mr^2 \dots \dots \dots (1)$$

In the ΔOPA

$$\begin{aligned} OP^2 &= (OA)^2 + (AP)^2 \\ r^2 &= (x+h)^2 + (AP)^2 \\ r^2 &= x^2 + 2xh + h^2 + (AP)^2 \\ r^2 &= x^2 + 2xh + h^2 + y^2 - h^2 \\ r^2 &= x^2 + y^2 + 2xh \end{aligned}$$

Putting the above value in equation (1)

$$\begin{aligned} I_o &= \sum m(x^2 + y^2 + 2xh) \\ I_o &= \sum mx^2 + \sum my^2 + \sum m(2xh) \\ I_o &= Mx^2 + I_G + 2x \sum mh \end{aligned}$$

Here $\sum my^2 = I_G$ and $\sum mh = 0$

This is because the body balances about centre of mass at G. Therefore, the algebraic sum of moments of all the particles about the centre of gravity, i.e.,

$$\sum mgh = 0$$

As g is constant

$$\sum mh = 0$$

Hence, we can write

$$I_o = I_G + Mx^2$$

Conservation Theorem of Energy:

Statement: Energy may be transformed from one kind to another, but it cannot be created or destroyed; the total energy is constant. This is the principle of the conservation of energy.

This theorem can be expressed as

$$\Delta K + \Delta U + Q + (\text{Change in other form of energy}) = 0$$

Here, ΔK is the change in kinetic energy.

ΔU is the change in potential energy.

Q is the heat produced due to friction.

Conservation Theorem of Momentum:

Statement:

If the vector sum of the external forces on a system is zero, the total momentum of the system is constant.

Let us consider an idealized system consisting of two particles that interact with each other but not with anything else. Each particle exerts a force on the other; according to Newton's third law, the two forces are always equal in magnitude and opposite in direction. Hence the impulses that act on the two particles will be equal and opposite, and the changes in momentum of the two particles will be equal and opposite.

Let $\vec{F}_{B \text{ on } A}$ be the force exerted by particle B on particle A and $\vec{F}_{A \text{ on } B}$ be the force exerted by A on B.

Therefore, the rate of change of momentum of the two particles is

$$\vec{F}_{B \text{ on } A} = \frac{d\vec{P}_A}{dt}, \quad \vec{F}_{A \text{ on } B} = \frac{d\vec{P}_B}{dt} \dots \dots \dots (1)$$

Where, \vec{P}_A = momentum of particle A.

\vec{P}_B = momentum of particle B.

Since $\vec{F}_{B \text{ on } A}$ and $\vec{F}_{A \text{ on } B}$ are equal and opposite, then

$$\vec{F}_{B \text{ on } A} = -\vec{F}_{A \text{ on } B}$$

$$\Rightarrow \vec{F}_{B \text{ on } A} + \vec{F}_{A \text{ on } B} = 0 \dots \dots \dots (2)$$

From equation (1) and (2), we get

$$\begin{aligned} \frac{d\vec{P}_A}{dt} + \frac{d\vec{P}_B}{dt} &= 0 \\ \Rightarrow \frac{d}{dt}(\vec{P}_A + \vec{P}_B) &= 0 \\ \Rightarrow \frac{d\vec{P}}{dt} &= 0 \quad [\vec{P} = \text{Total momentum} = \vec{P}_A + \vec{P}_B] \\ \therefore \vec{P} &= \text{Constant.} \end{aligned}$$

Collision:

In a collision a relatively large force acts on each colliding particle for a relatively short time. The basic idea of a “collision” is that the motion of the colliding particles changes rather abruptly and that we can make a relatively clean separation of times that are “ before the collision” and those that are “ after the collision”.

There are two types of collision:

- i. Elastic Collision
- ii. Inelastic Collision

Elastic Collision: A collision in which the kinetic energy of the system is the same after the collision as before.

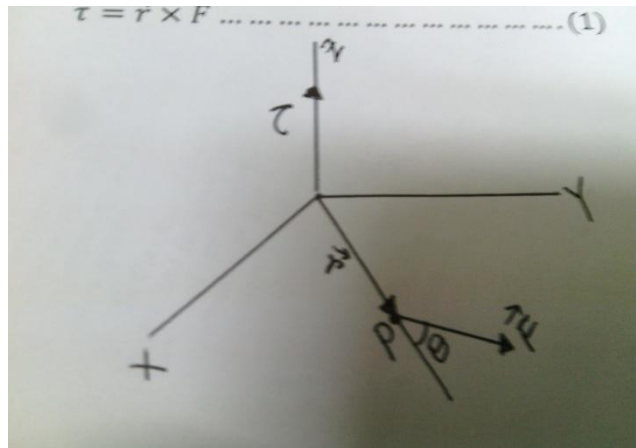
Inelastic Collision: A collision in which the total kinetic energy after the collision is less than that before the collision is called an inelastic collision.

Torque:

Torque is defined as the tendency of a force to rotate an object about an axis.

If a force \vec{F} act on a single particle at a point P whose position with respect to the origin O of the inertial reference frame is given by the displacement vector \vec{r} , the torque τ acting on the particle with respect to the origin O is defined as

$$\vec{\tau} = \vec{r} \times \vec{F} \dots \dots \dots (1)$$



Torque is a vector quantity. Its magnitude is given by

$$\tau = rF \sin\theta$$

where θ is angle between \vec{r} and \vec{F} ; its direction is normal to the plane formed by \vec{r} and \vec{F} .

Flywheel:

A flywheel is a heavy metal disc with its mass concentrated mostly in its rim. The wheel is capable of rotation about a horizontal or a vertical axis. A thick rod which is called the axle passes through the center of gravity of the wheel which rotates about the rod as axis. The rod and the wheel are rigidly connected. The wheel is on a supported horizontal axis (Fig.-01). One end of a flexible cord is fixed to a small peg on the axle. The other end of the cord, which is wrapped round the axle, carries a mass M . The length of the cord is such that it becomes detached from the axle when the mass strikes the ground.

Fig.-01